

Quantum Mechanics

1. Black body Radiation spectrum.
2. Photo electric effect
3. X-Ray spectrum
4. Compton effect / scattering
5. Hydrogen-Line spectrum
OR. Atomic -

Domains where classical mechanics failed

- a. Where the velocity of object is very high (10^8 ms^{-1}) (Theory of Relativity)
- b. Atomic Realm (Quantum mechanics)

14 dec. 1900

→ Max. Planck produced a paper to explain Black body radiation,

$$E = h\nu$$

1905

→ Albert Einstein explain P.E.E by applying Max planck's T.O.R.

1913 →

Neil's Bohr, gave Bohr's atomic model by introducing quantum condition on Angular momentum.

1920 A.H. compton

1927

Davisson - Germer experiment

→ particles have wave-character

Heisenberg (1925)

Erwin Schrödinger (1927)

Matrix Mechanics

Wave mechanics

Postulates of quantum mechanics

Postulate 1: The state of any physical system is given by space vector $\Psi(x, t)$ in Hilbert space. This $\Psi(x, t)$ contains all the information we need. Any superposition of state vectors is a state vector.

$$\Psi(x, t) = a\Psi_1(x, t) + b\Psi_2(x, t)$$

Postulate 2: To every observable and/or dynamical variable, there corresponds a linear Hermitian operator whose eigenfunction forms a complete basis.

Postulate 3: The measurement of an observable A may be represented by the action of \hat{A} on state vector $\Psi(x, t)$ which results a $a\Psi$ where a is called eigenvalue of \hat{A} .

Postulate 4: The time evolution of the state vector $\Psi(x, t)$ is given by the time dependent Schrödinger's equation.

$$\frac{-\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V\Psi = i\hbar \frac{\partial \Psi}{\partial t}$$

$$\frac{-\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial t^2} + V\Psi = i\hbar \frac{\partial \Psi}{\partial x}$$

Failure of classical (EM) Theory

- C.P could not explain the phenomena related to atomic domain.
- C.P could not explain the Hydrogen - spectra.
- This theory was not able to explain the stability of atom.
- The enormous range of conductivity of metals compared to Insulators also could not be explained.
- The range of magnetic susceptibility is also very large, but it couldn't explained.
- Classical theory could not explained the emission of α - particle from a nucleus despite of the fact that it has insufficient energy.
- The specific heat variation with temperature was not explain by the classical theory.
- Black body radiation, PEC, Raman effect, Compton effect was not successfully explained by classical Theory. Dual nature of particle etc.

Photoelectric Effect :-

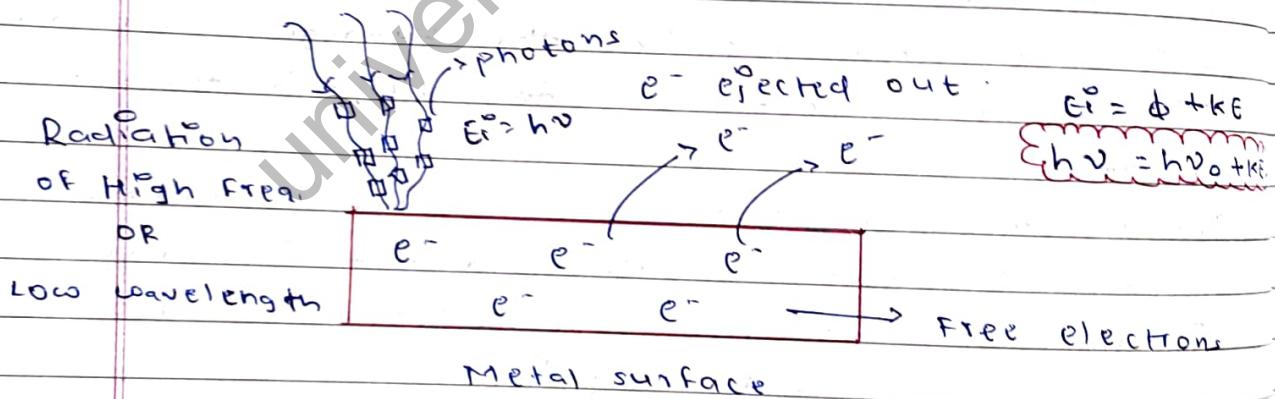
Random light → Wave nature
 Random light → Particle Nature

Particle Nature of Light.

Light → 1675 → Newton → Light → Particle

1801 → Young → Diffraction → Light wave
 1801 → Young → Polarisation → Light wave

1905 → Photo electric } → Light → Particle



- ① Light consists of particle called photons.
 'wave packets'
 'energy packets'
 'quanta'

- ② Mass of photon is not defined in sense of Newtonian mechanics. The Rest mass of photon zero

3. All photons travel with $299,792,458$ m/s in air & vacuum $\approx 3 \times 10^8$ m/s.
4. Photons are neutral in nature.
5. Each photon has a definite Energy & definite Momentum.

6.

$$E = \frac{hc}{\lambda}$$

$h \rightarrow$ Planck's constant

$$h = 6.62 \times 10^{-34} \text{ Js}$$

$$E = 1240 \text{ eV}$$

λ (in nm)

$$E = hF$$

$$P = \frac{h}{\lambda}$$

$\nwarrow N = F \lambda$

$\nwarrow f = N / \lambda$

$\nwarrow F = c / \lambda$

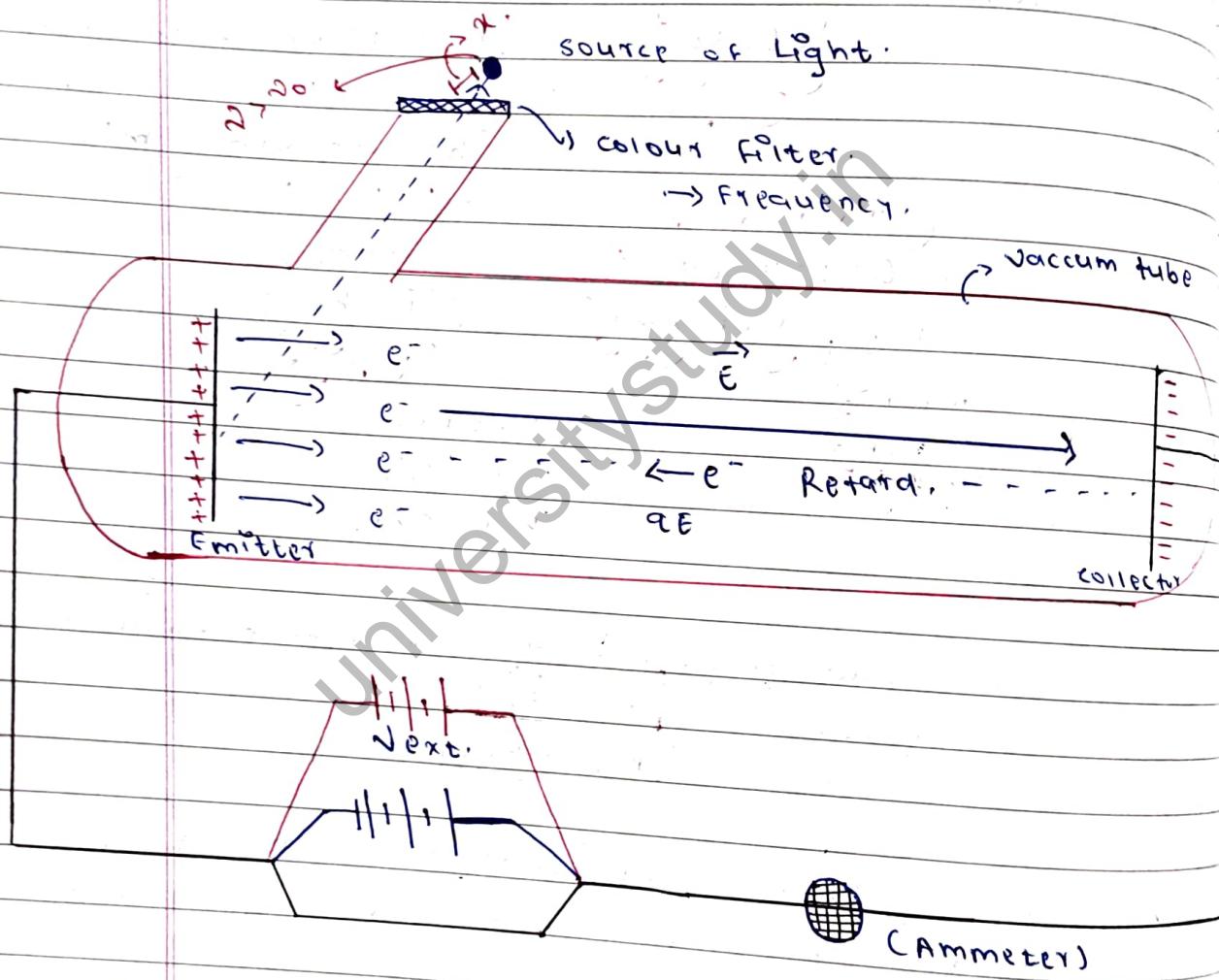
6. All collisions b/w a photon $\xrightarrow{\Sigma}$ a matter particle are perfectly elastic.

$$\begin{array}{|c|c|} \hline \xrightarrow{\Sigma} & \xrightarrow{\Sigma} \\ P_i = P_f & \\ \hline K \cdot E_i = K \cdot E_f & \\ \hline \end{array}$$

7. On increasing Intensity of Radiation (Light), the no. of photons increased but the energy of each photon does not change.

Ejection of e^- from Metal surface \rightarrow Work f.

When light of particular freq. ν is incident on a surface e^- are emitted from it. These e^- are called photoelectrons. The effect is called photoelectric effect.



* $\nu < \nu_0$ NO PEE will occur.

* E.F will set up in our tube

* e^- will be trapped in E.F & will accelerate

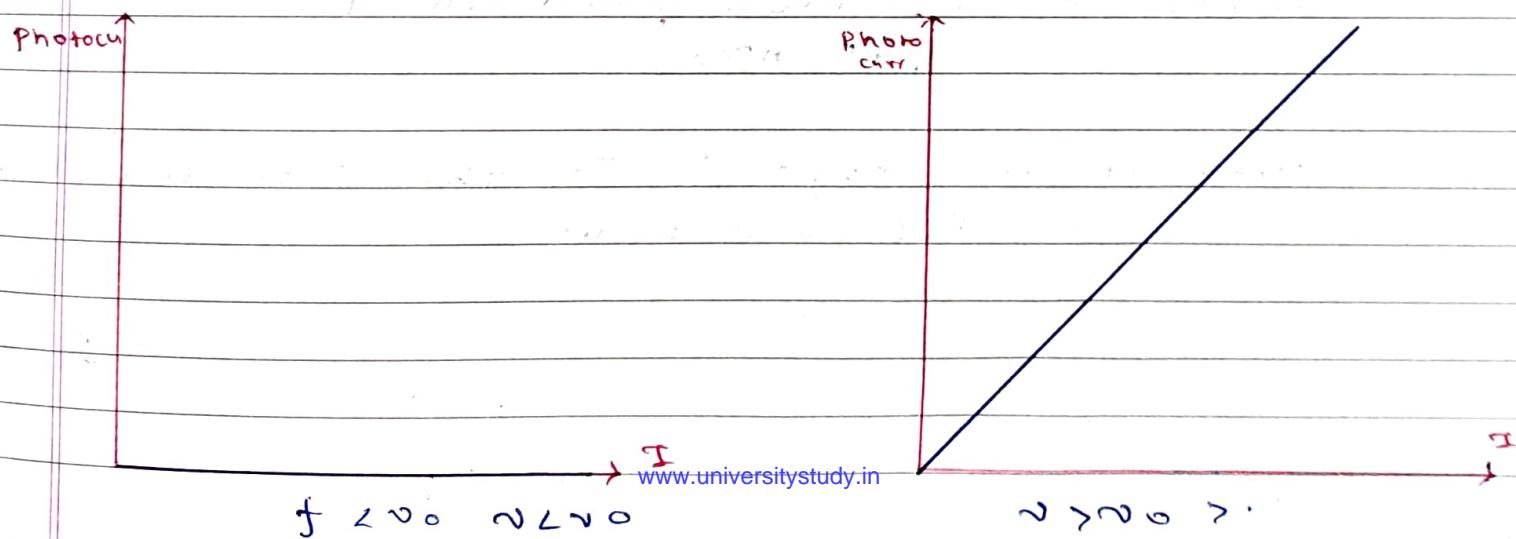
- Stopping potential \rightarrow -ve voltage of collector at which photocurrent falls to zero.
- -ve voltage at which e^- with highest KE also stops.

Analysis :-

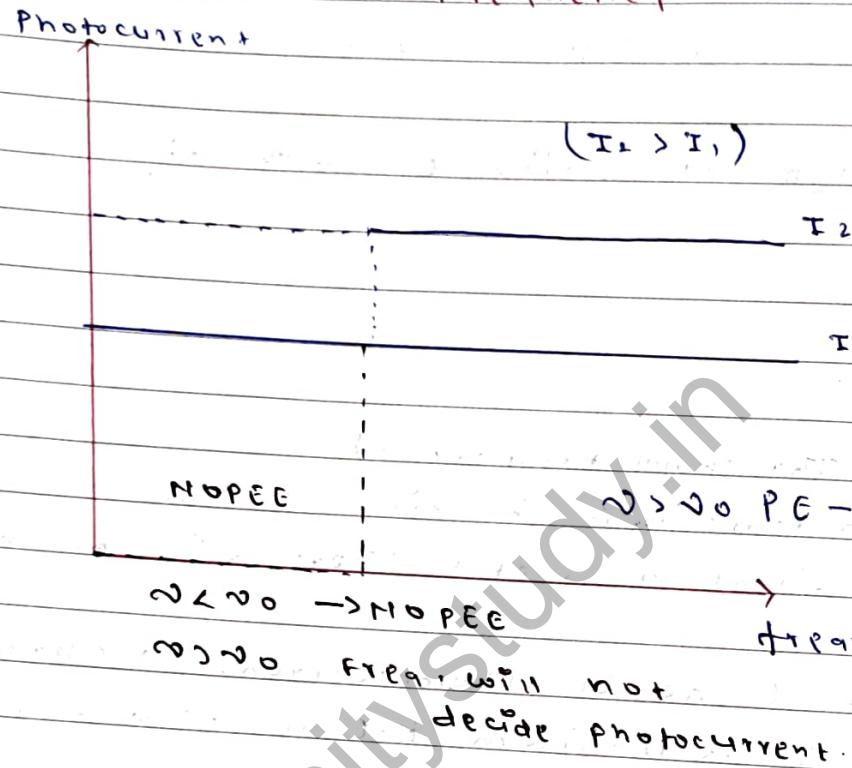
- $\nu > \nu_0$ (Threshold frequency)
- $\nu < \nu_0$ whatever may be time for which light is incident on metal there will be no PEE.
- I of light \propto no. of photoelectrons \propto current.
- F of light \propto maximum K-E of e^- \propto stopping pot.
- If $\nu = \nu_0$ e^- will come out with zero KE due to which metal will be positively charged as it absorbs back the e^- .

Observation :-

1. Intensity Vs Photocurrent.



2. Photocurrent vs Frequency :-



Photocurrent depends on Intensity of light, & is independent of frequency.

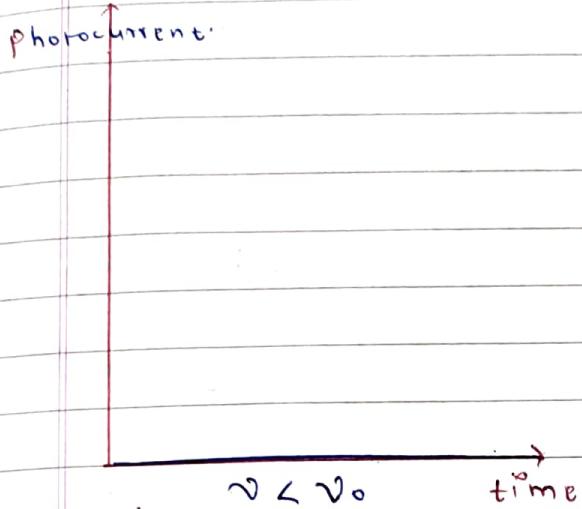
Threshold Frequency :- Min. Frequency at which PEE starts occurring (n_0)

Work Function :- Min. energy Required to remove an e^- from metal in PEE.

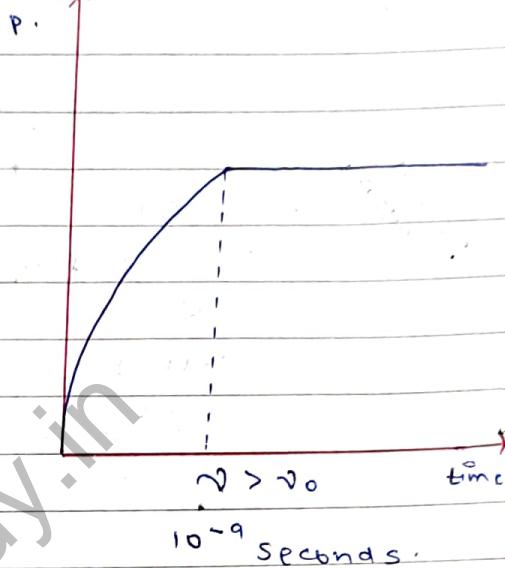
$$\Phi_0 = h\nu_0$$

Cut-off-wavelength :- Is the Max wavelength upto which PEE will occur.

3. Photocurrent vs Time.



photocurrent.

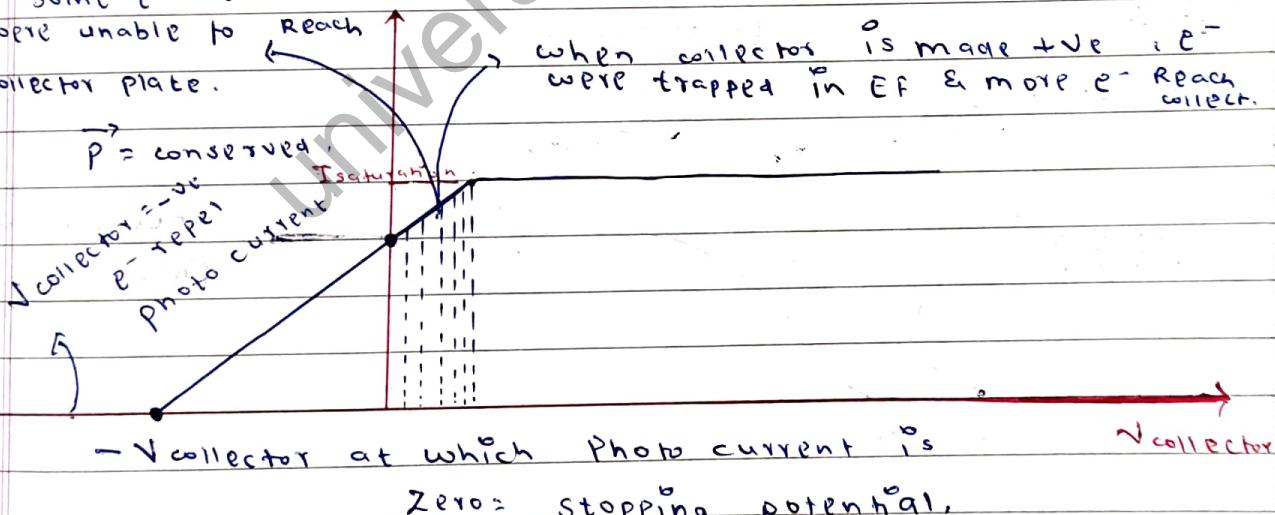


4. Photocurrent vs Voltage.

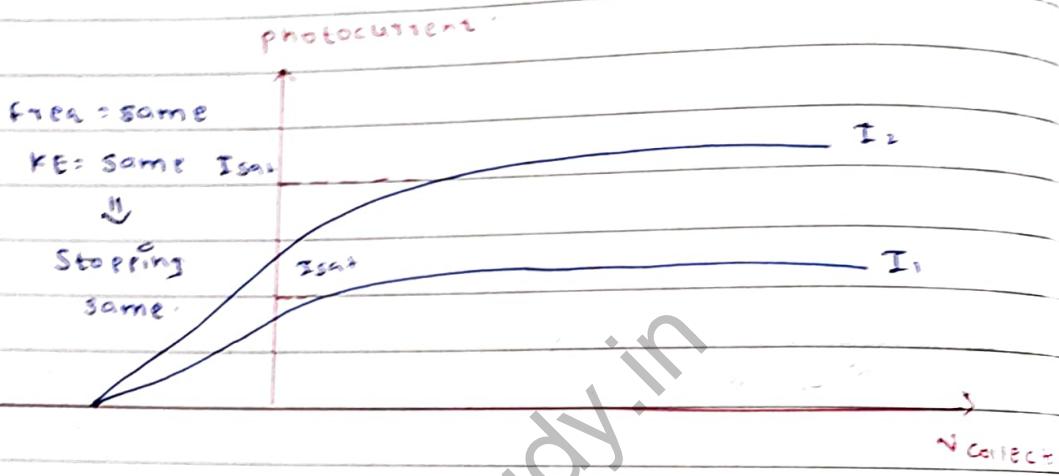
 $(V > V_0)$

Some e^- were scattered. Photocurrent
& were unable to reach collector plate.

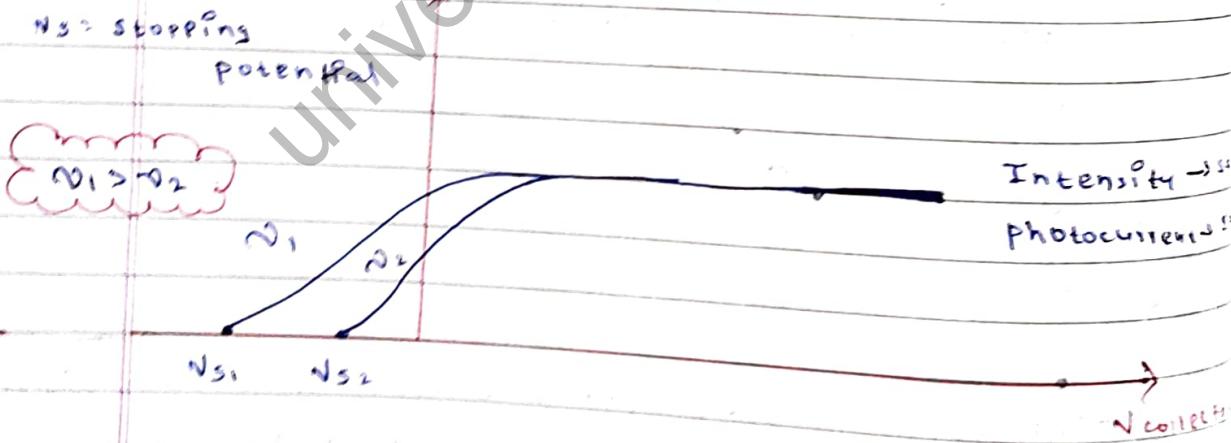
when collector is made +ve i.e. e^-
were trapped in EF & more e^- reach collect.



case 1: Different Intensity
 → It will impact photocurrent
 → no change of stopping pot. (V_{stop})

B.TECH
Subject

case 2: Different Frequency (Intensity constant)
 → saturated p.c.
 → const.

B.TECH
Subject

Einstein Explanation of photoelectric effect:

Quantum of lights \rightarrow Packet of Energy.

- Radiation observed by metal are in form of quantum (photon). Energy of each photons depends on Frequency.

$$\text{Energy of light} = h\nu$$

Photon Energy

$$E = h\nu$$

If n no. of

$$\nu$$

Photons are

incident $E = nh\nu$

- One photon interact with one electron only and interaction is instantaneous process.
(one-to-one interaction) (instantaneous process.)
 $\sim 10^{-9}$ s.
- In the interaction the photon transfer complete energy to the electric.
- During the interaction the energy and momentum is conserved. \rightarrow (collision process b/w $1P_2, 1e^-$)
(Momentum is conserved)
- The I of light is directly proportional to no. of photons hence If intensity increases the no. of photon incident increases which leads to more emission of electrons (rise of photocurrent).

$$\left\{ \begin{array}{l} I = E \\ tA \end{array} \right\}$$

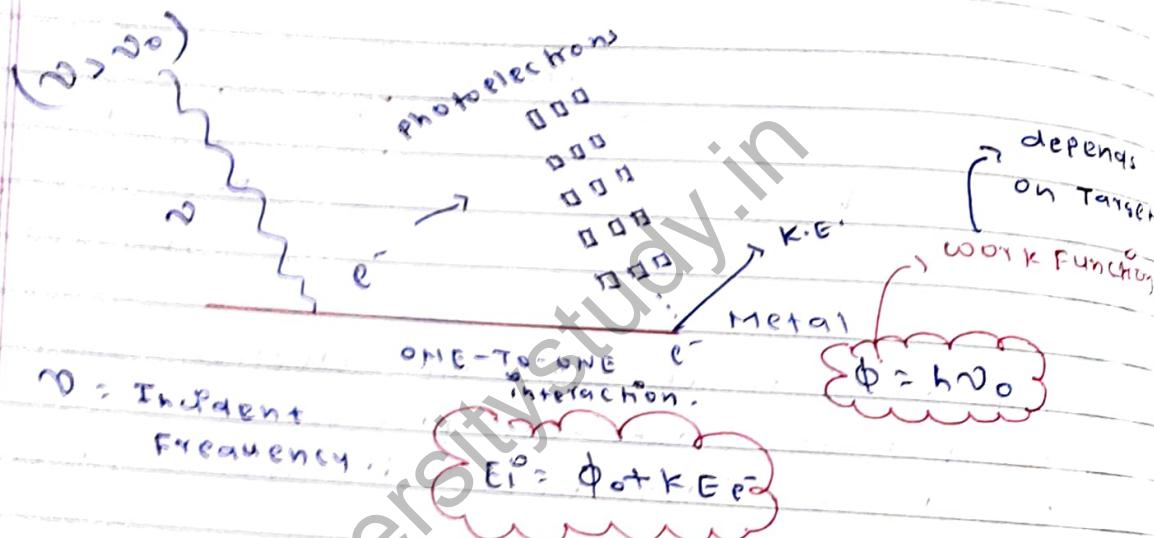
Acc. to Einstein Energy = $n h \nu$

n = Total no. of photons

$I \propto n$

($n \propto I^{\frac{1}{2}}$, $I \propto$ no. of collisions,
no. of photo e⁻)

Einstein's Equation of P.E.E



$$T_F \approx \nu < \nu_0 \quad E_i^0 < \phi_0$$

Forms

$$E_i^0 = \phi_0 + K.E$$

$$\hbar\nu = h\nu_0 + \frac{1}{2}mv^2$$

$$\frac{hc}{\lambda} = \frac{hc}{\lambda_0} + \frac{1}{2}mv^2$$

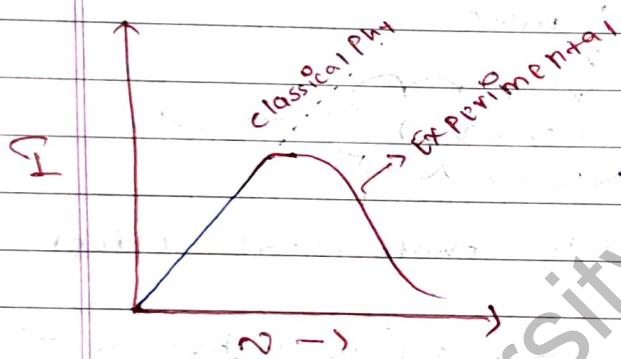
Black Body Radiation

$$I(\nu) = \frac{8\pi h\nu^3}{c^3} \frac{1}{e^{h\nu/kT} - 1}$$

h = Planck's constant

c = speed of light

ν = Frequency.



Einstein Photoelectric Equation

- has explained various law of PEE on the basis of Planck's law,

When a photon of energy $h\nu$ strikes on a metal surface a part of its energy known as work function (ϕ) is used to liberate the electrons from the surface, whereas the remaining energy is spent for its kinetic energy.

Thus, the energy of photon = the sum of work function and kinetic energy of the electron, that is $h\nu = \phi + \frac{1}{2}mv^2$ — ①

Dual nature of a particle / Matter.

→ De Broglie suggestion was based on the reasoning that nature is symmetrical & universe is composed of matter & radiations, thus these two physical entities matter and radiations should have symmetry. The wave associated with a moving particles is known as de-Broglie wave, & the wave-length associated with this particle is known as de-Broglie wave-length it is denoted by symbol (λ).

De - Broglie wave Equation.

Acc. to Planck's quantum theory energy of photon with frequency ν is given as,

$$E = h\nu \quad \text{--- (1)}$$

The Einstein's Mass-Energy relation is given as,

$$E = mc^2 \quad \text{--- (2)}$$

From eq (1) & (2),

$$mc^2 = h\nu$$

$$m = \frac{h\nu}{c^2} \quad \text{--- (3)}$$

We know, $P = \text{Mass} \times \text{Velocity}$
 $P = mc. \quad \text{--- (4)}$

→ particle is moving at speed of light

when a min. threshold frequency is provided in that case, electron will just eject from the surface & velocity will become zero.

In that case $\nu = \nu_0$, $v = 0$.

Eqn ① will be modified as $\{ h\nu_0 = \omega_0 \} \quad \text{--- } ②$

Eq ① & ②.

$$h\nu = h\nu_0 + \frac{1}{2}mv^2$$

$$\therefore \left\{ \frac{1}{2}mv^2 = h(\nu - \nu_0) \right\} \quad \text{--- } ③$$

Eq. ③ is the Einstein's photoelectric effect.

$\nu < \nu_0 \rightarrow \text{K.E will become -ve.}$

which is not possible

Hence, NO PE emission

is possible, if frequency of incident photon is less

than, threshold frequency.

Explanation of de-broglie wave-equation.

- If mass or velocity is kept large, in that case, de-broglie wavelength will be small.
- De-broglie wavelength is independent on charge.
- If velocity = 0, in that case $\lambda \rightarrow \infty$. Such waves can't be visualized hence, the de-broglie wavelength is associated only with the moving particle.

De-broglie wavelength for an electron.

- Consider an electron with mass m , charge e & potential V so $E = eV$ — ①
- Consider the K.E. of the electron is given as $\frac{1}{2}mv^2$ — ②.

On multiplying & dividing eq. ② (R.H.S) with mass.

$$E = \frac{1}{2} \frac{mv^2 \cdot m}{m}$$

$$= \frac{1}{2} (mv)^2 \quad [\text{where } P = mv].$$

$$E = \frac{P^2}{2m}$$

$$P = \sqrt{2mE} \quad - \textcircled{3}.$$

From Eq. ① & ④

$$P = \frac{h\nu}{c\lambda} \quad \text{--- ⑤}$$

$$\boxed{P = \frac{h\nu}{c\lambda}} \quad \text{--- ⑤}$$

We know $c = \nu\lambda$ --- ⑥

From ⑤ & ⑥

$$P = \frac{h\nu}{m\nu\lambda}$$

$$P = \frac{h}{\lambda}$$

$$\boxed{\frac{h}{\lambda} = \frac{h}{\lambda}} \quad \text{--- ⑦}$$

$$\lambda = \frac{h}{P} \quad \text{--- ⑧}$$

When a particle is moving with a velocity v , then momentum $P = mv$.

From Eq. ① & ⑦

$$\lambda = \frac{h}{mv}$$

Above Equation is www.universitystudy.in - broglie's wave equation.

From Equation ① & ③ -

$$p = \sqrt{2mE} \quad \text{--- (4)}$$

We know de-broglie $\lambda = \frac{h}{p}$ \rightarrow ⑤

From ④ & ⑤,

$\lambda = h / \sqrt{2mE}$

de-broglie wavelength for an electron,

Heisenberg Uncertainty Principle.

According to U.P., it is impossible to measure simultaneously the exact position & momentum.

Mathematically,

$$\Delta x \cdot \Delta p \geq \frac{\hbar}{4\pi}$$

or

$$\Delta x \cdot \Delta p \geq \frac{\hbar}{2}$$

$$\hbar = \frac{h}{2\pi}$$

In terms of Energy & time the uncertainty is given as

$$\Delta E \cdot \Delta t \geq \frac{\hbar}{4\pi}$$

or

$$\Delta E \cdot \Delta t \geq \frac{\hbar}{2}$$

Phase Velocity & Group Velocity.

(v_p) (v_g)

Phase Velocity (v_p)

→ It is the velocity with which planes of equal phase travel through a medium, the phase velocity is associated with a single wave.

The equation of propagation of a wave is given as

$$y = a \sin(\omega t - kx) \quad \text{--- (1)}$$

Where a represent amplitude,

ω is angular frequency

t is time

k is propagation constant

x is displacement.

Where $(\omega t - kx)$ is the phase of the wave

If the phase is constant then $\omega t - kx$ will be equal to constant.

$\omega t - kx$ is phase

$$\omega t - kx = \text{constant} \quad \text{--- (2)}$$

on diff Eq. 2

$$\frac{\omega t - kx}{dt} = 0$$

at

$$k \frac{dx}{dt} = \omega$$

$$v_p = \frac{\omega}{k}$$

$$\frac{dx}{dt} = \frac{\omega}{k}$$

Above expression represent the phase velocity of the wave.

where

$$\omega = 2\pi\nu$$

$$\& \kappa = \frac{2\pi}{\lambda}$$

Relationship b/w phase velocity & particle vel.

Since, phase velocity $v_p = \frac{\omega}{\kappa} \quad \rightarrow \textcircled{1}$

Also $\omega = 2\pi\nu \quad \& \kappa = \frac{2\pi}{\lambda} \quad \rightarrow \textcircled{3}$
 $\rightarrow \textcircled{2}$

$$v_p = \frac{2\pi\nu}{2\pi}$$

$$v_p = \lambda\nu \quad \rightarrow \textcircled{4}$$

The energy of photon is given as $E = h\nu$
 or $\nu = \frac{E}{h} \quad \rightarrow \textcircled{5}$

De-Broglie wave-equation is given
 as $\lambda = \frac{h}{mv} \quad \rightarrow \textcircled{6}$

$$v_p = \frac{\nu \cdot E}{mv}$$

$$v_p = \frac{E}{mv} \quad \rightarrow \textcircled{7}$$

From Einstein's Equation $E=mc^2 \rightarrow \textcircled{8}$

From $\textcircled{7}$ & $\textcircled{8}$

$$v_p = \frac{mc^2}{\rho h v}$$

$$v_p = \frac{c^2}{v}$$

∴ A phase velocity is not associated with a particle in a motion.

Group Velocity (v_g).

→ The wave packet/bundle/group moves on its own velocity known as group velocity.

Consider a group of wave consisting two sine waves with same amplitude "a" and slightly different angular frequency " ω_1 & ω_2 " & prop. constant " k_1 & k_2 ". These waves are represented as:

$$y_1 = a \sin(\omega_1 t - k_1 x) \quad \text{--- (1)}$$

$$y_2 = a \sin(\omega_2 t - k_2 x) \quad \text{--- (2)}$$

The superposition of these two waves will give

$$Y = y_1 + y_2. \quad \text{--- (3)}$$

From Eq. (1) + (2) + (3)

$$Y = a \sin(\omega_1 t - k_1 x) + a \sin(\omega_2 t - k_2 x)$$

$$\sin A + \sin B = 2 \sin \left(\frac{A+B}{2} \right) \cos \left(\frac{A-B}{2} \right)$$

$$y = 2a \sin \left(\frac{\omega_1 t - k_1 x}{2} + \frac{\omega_2 t - k_2 x}{2} \right) \cos \left(\frac{\omega_1 t - k_1 x}{2} + \frac{\omega_2 t - k_2 x}{2} \right)$$

$$y = 2a \sin \left(\frac{\omega_1 t - k_1 x}{2} + \frac{\omega_2 t - k_2 x}{2} \right) \cdot \cos \left(\frac{\omega_1 t - k_1 x}{2} + \frac{\omega_2 t - k_2 x}{2} \right)$$

$$y = 2a \sin \left[\frac{(\omega_1 + \omega_2)t}{2} - \left(\frac{k_1 + k_2}{2} \right)x \right] \cos \left[\frac{(\omega_1 - \omega_2)t}{2} - \left(\frac{k_1 - k_2}{2} \right)x \right] \quad - (4)$$

$$\frac{\omega_1 + \omega_2}{2} = \omega, \quad \frac{k_1 + k_2}{2} = K.$$

$$y = 2a \sin [\omega t - Kx] \cdot \cos \left[\frac{t \Delta \omega}{2} - \frac{x \Delta K}{2} \right] \quad - (5)$$

The general wave equation is written as

$$y = a \sin (\omega t - Kx) \quad - (6)$$

On comparing Eq. (5) & (6)

$$\text{Amplitude} = 2a \cos \left[\frac{t \Delta \omega}{2} - \frac{x \Delta K}{2} \right]$$

$$2\omega \cos \left[\frac{\Delta k}{2} \left[\frac{\Delta \omega}{\Delta k} + t - \alpha \right] \right]$$

where $v_g = \frac{\Delta \omega}{\Delta k}$

The above expression represents the group velocity associated with the particle.

R/s b/w v_p & v_g .

since, $v_p = \frac{\omega}{k}$

or $\omega = k v_p$ — ①.

on diff. both side w.r.t to k in eq ①

ωdk

$$\frac{d\omega}{dk} = v_p + k \frac{dv_p}{dk}$$

$$v_g = v_p + k \frac{dv_p}{dk} \quad [\because v_g = \frac{d\omega}{k}] - ②$$

$$k = \frac{2\pi}{\lambda}$$

We know, $k = \frac{2\pi}{\lambda}$ — ③.

on diff. both side w.r.t ' λ '

$$\frac{dk}{d\lambda} = -\frac{2\pi}{\lambda^2} \cdot d\lambda \quad - ④$$

on substituting ③ & ④ in eq ② -

$$\nu_g = \nu_p + \frac{2\pi}{\lambda} \cdot \frac{d\nu_p}{dx}$$

$$\nu_g = \nu_p + \frac{2\pi \cdot \lambda \cdot d\nu_p}{-2\pi}$$

$$\nu_g = \nu_p - \lambda \frac{d\nu_p}{d\lambda}$$

The above equation represent RL b/w Group Velocity & phase velocity.

where, λ is the dispersion in the medium

case ① :- when there is no dispersion, then $\lambda = 0$,

in that case $\nu_g = \nu_p$.

case ② :- For normal dispersive medium

$$\nu_p > \nu_g,$$

case ③ :- For anomalous dispersive medium

$$\nu_p < \nu_g,$$

RHS below ν_g & Particle Velocity (v).

Consider, a particle moving with velocity (v).

The K.E is given as $E = \frac{1}{2}mv^2$

In terms of Momentum $E = \frac{P^2}{2m} \quad [\because P = mv]$

$$\hookrightarrow \textcircled{1} \rightarrow \frac{\hbar}{\lambda} \times$$

$$\text{Since } E = \hbar\omega \quad \text{--- (2)}$$

$$P = \hbar k \quad \text{--- (3)}$$

$$E = \hbar\nu \quad \rightarrow \omega$$

$$E = \frac{\hbar}{2\pi} \cdot 2\pi\nu$$

$\hookrightarrow \hbar$

on substituting, Eq. (2) & (3)

in Eq. (1)

$$E = \hbar\omega$$

$$\hbar\omega = \frac{\hbar^2 k^2}{2m}$$

$$\lambda = \frac{\hbar}{P}$$

$$\omega = \frac{\hbar k^2}{2m} \quad \text{--- (4)}$$

$$P = \frac{\hbar}{\lambda}$$

on diff. Eq. (4) w.r.t k .

$$P = \frac{\hbar}{2\pi} \cdot \frac{2\pi}{\lambda} \cdot k \quad \hookrightarrow \hbar$$

$$\frac{d\omega}{dk} = \frac{2\hbar k}{2m}$$

→ $\textcircled{5}$

$$\frac{d\omega}{dk} = \frac{\hbar k}{m}$$

$$\nu_g = \frac{\hbar k}{m}$$

$$[\because \nu_g = \frac{d\omega}{dk}]$$

From Eq. (3) (4) (5)

$$\nu_g = \frac{P}{m}$$

Since $P = mv$, so.

$$Vg = \frac{Pv}{m}$$

$$\boxed{Vg = V}$$

Wave function (Ψ)

The quantity whose variation make up the matter waves is known as wave function. Matter wave is associated with the moving particle. Thus the wave function is the function which describes the matter wave associated with a moving particle. The wave function is a function of both space & time,

The wave function $\boxed{\Psi = A + iB}$

$A \& B \rightarrow \text{constant.}$

The Ψ can be expressed as

$$\boxed{\Psi = A e^{i(kx - \omega t)}}$$

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$$\begin{aligned}|4\vec{v}|^2 &= (2 - 3i)(2 + 3i) \\&= (2)^2 - (3i)^2 \\&= 4 + 9 = 13\end{aligned}$$

classmate

Date _____

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Physical Significance of Ψ .

$$|\Psi|^2 = \Psi^* \Psi$$

The wave function Ψ has no direct physical significance, the probability of Ψ gives the physical significance and the probability is written as $| \Psi |^2 = \Psi^* \Psi$, where the probability should have real & true values.

Normalization of Ψ

→ As the particle must be found somewhere,
 so one can say that the probability of
 finding the particle ⁱⁿ whole space is unity
 that is ∞ :

$$\int_{-\infty}^{\infty} |\Psi(x, t)|^2 = 1$$

Schrodinger's Time dependent & Independent Wave Equation.

→ Time dependent wave equation

$$-\frac{i\hbar}{\partial t} \partial \Psi = -\frac{\hbar}{2m} \frac{\partial^2 \Psi}{\partial x^2}$$

→ Time independent wave equation

$$\nabla^2 \Psi + \frac{2m(E-V)}{\hbar^2} \Psi = 0$$

Operators

1. Energy operators [E]

$$\Psi = A e^{i(kx - \omega t)} \quad \text{--- (1)}$$

$$p = \hbar k \Rightarrow k = \frac{p}{\hbar} \quad \text{--- (2)}$$

$$E = \hbar \omega \quad \omega = \frac{E}{\hbar} \quad \text{--- (3)}$$

From (1) (2) & (3)

$$\Psi = A e^{i(CRt - Et)} \quad \text{--- (4)}$$

$$\frac{\partial \Psi}{\partial t} = -\frac{i}{\hbar} E A e^{i(CRt - Et)} \quad \text{--- (5)}$$

From (4) & (5)

$$\frac{\partial \Psi}{\partial t} = -\frac{i}{\hbar} e \Psi$$

$$e \Psi = -\frac{\hbar}{i} \frac{\partial \Psi}{\partial t}$$

$$E \Psi = -\frac{e\hbar}{i \cdot i} \frac{\partial \Psi}{\partial t}$$

$$E \Psi = \frac{i \hbar \partial \Psi}{i \cdot i \partial t}$$

$$E = i \hbar \frac{\partial}{\partial t}$$

2. Momentum operator [\hat{P}]

$$\Psi = A e^{i \frac{\omega}{\hbar} (kx - \omega t)} \quad \text{--- (1)}$$

(2)

(3)

(4)

$$\frac{\partial \Psi}{\partial x} = \frac{i}{\hbar} p A e^{i \frac{\omega}{\hbar} (p x - E t)}$$

$$\frac{\partial \Psi}{\partial x} = \frac{i}{\hbar} p \Psi$$

$$p\Psi = \frac{\hbar}{i} \frac{\partial \Psi}{\partial x}$$

$$p\Psi = \frac{i \hbar}{i \cdot i} \frac{\partial \Psi}{\partial x}$$

$$p\Psi = i \hbar \frac{\partial \Psi}{\partial x}$$

$p = i \hbar \frac{\partial}{\partial x}$

Kinetic Energy operator (\hat{T})

$$K.E = T = \frac{1}{2} m v^2$$

or

$$T = \frac{p^2}{2m}$$

$$[\hat{T}] = [\frac{p^2}{2m}]$$

$$\hat{P} = -i \hbar \frac{\partial}{\partial x}$$

$$\hat{T} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$$

Hamiltonian operator (\hat{H})

$$\hat{H} = -\frac{\hbar^2}{2m} (\nabla^2) + V$$

Position operator (\hat{x})

$$\hat{x} = x$$

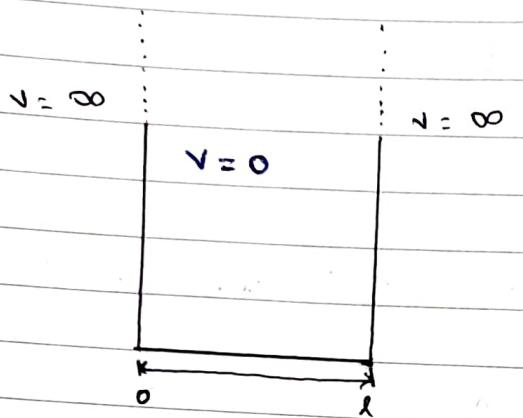
Potential operator (\hat{V})

$$\hat{V} = V$$

Q. Light of wavelength 5000 \AA falls on a metal surface of work function (ϕ) = 1.9 eV. Find

- (i) The energy of photons in eV
- (ii) The K.E. of photo-electrons
- (iii) Stopping Potential.

Particle Inside 1-D box :-



Consider a particle of mass M , that is confined to move along x -axis, box of length l . The potential energy $V = 0$, inside the box is infinite at the boundaries of the box, that is

$$V = 0 \quad 0 < x < l$$

$$V = \infty \quad x = 0, x = l.$$

Schrodinger 1-D wave equation (time independent) is given as

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{2m}{\hbar^2} (E - V) \Psi = 0.$$

Inside the box, the potential energy $V = 0$.

∴ time independent wave equation is given as

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{2mE}{\hbar^2} \Psi = 0$$

$$\frac{\partial^2 \psi}{\partial x^2} + k^2 \psi = 0 \quad \text{--- (1)}$$

Where $k^2 = \frac{2mE}{\hbar^2}$ --- (2)

$$E = \frac{1}{2} mv^2$$

$$E = p^2 / 2m$$

$$p = \hbar k$$

$$E = \frac{\hbar^2 k^2}{2m}$$

$$E =$$

$$k^2 = \frac{2mE}{\hbar^2}$$

The general solution of Eq. (1) is given as

$$\psi = A \sin kx + B \cos kx \quad \text{--- (3)}$$

A & B are constant.

To find the value of A & B we have to apply the boundary condition:

B.C I : When $x = 0$, $\psi = 0$.

on applying the B.C - I on Equation - (3)
 $x = 0$, $\psi = 0$.

$$0 = A \sin k(0) + B \cos(k(0))$$

$$B = 0 \quad \text{--- (4)}$$

From Equation (3) & (4)

$$\left\{ \begin{array}{l} \psi = A \sin kx \\ \end{array} \right. \quad \text{--- (5)}$$

B.C - II.

When $x = l$, $\psi = 0$.

on applying B.C - II, on equation (5)

$$A \sin kl = 0$$

$$\sin kl = 0$$

$$\sin kl = \sin n\pi$$

$$kl = n\pi$$

$n = 1, 2, 3, \dots, \infty$

$$k = \frac{n\pi}{l} \quad \text{--- (6)}$$

On substituting Eq. 6, k, to Eq. I (3)

$$\psi = A \sin \frac{n\pi}{l} \cdot x \quad \text{--- (7)}$$

Since, there is only one particle & there is no chance of its being located outside the box, \therefore the probability of finding the particle inside the box is 1.
Thus, the normalization condition is given as

$$\int_0^l |\psi|^2 dx = 1$$

From Eq. ⑦

$$A^2 \int_0^L \left[\sin^2 \left(\frac{n\pi x}{l} \right) \right] dx = 1$$

$$\frac{A^2}{2} \cdot \int_0^L \left[2 \sin^2 \left(\frac{n\pi x}{l} \right) \right] dx = 1$$

$$\frac{A^2}{2} \int_0^L \left[1 - \cos \frac{2n\pi x}{l} \right] dx = 1$$

$$\frac{A^2}{2} \left[x - \frac{\sin \left(\frac{2n\pi x}{l} \right)}{\frac{2n\pi}{l}} \right]_0^L = 1$$

$$\frac{A^2}{2} \left[x - \frac{l}{2n\pi} \sin \left(\frac{2n\pi x}{l} \right) \right]_0^L = 1$$

$$\frac{A^2}{2} \left[l - \frac{l}{2n\pi} \sin (2n\pi) \right] = 1$$

$$\frac{A^2}{2} \left[\frac{2n\pi l - l}{2n\pi} \right] = 0$$

$$\frac{A^2 l}{2} = 1$$

$A = \sqrt{\frac{2}{l}}$

Eq. (8) to Eq. (2)

on substituting a value from

$$\Psi = \sqrt{\frac{2}{l}} \sin\left(\frac{n\pi x}{l}\right)$$

Remember.

From Equation (2) & (6)

$$\frac{2mE}{\hbar^2} = \frac{n^2\pi^2}{l^2}$$

$$E = \frac{n^2\pi^2\hbar^2}{2ml^2}$$

The above expression represent the energy of a particle confined in one-dimension box, it signifies that the energy is always discrete

If $n=1$ $n=2$

$$E = \frac{1}{2} \frac{\pi^2 \hbar^2}{ml^2}, \quad E = \frac{4}{2} \frac{\pi^2 \hbar^2}{ml^2}$$

 $n=3$

$$E = \frac{9}{2} \frac{\pi^2 \hbar^2}{ml^2}$$

 $n=4$

$$E = \frac{16}{2} \frac{\pi^2 \hbar^2}{ml^2}$$

UNIT-V - SOLID STATE PHYSICS

Free Electron Theory :- (classical approach)

Drude & Lorentz has proposed free electron theory on the basis of following assumption:

- In conductors (Metals), the free electrons or valence electrons are free to move inside a conductor just like a gas molecule.
- Since electrons behave like gas molecule, therefore the kinetic theory of gas will be applicable on electrons.
- In the absence of E.F, the electrons move randomly, whereas the collisions will be elastic (no loss of energy).
- When the E.F is applied, the free electrons will start accelerating opposite to the applied field. This is known as drift current.
- There is no electrostatic force of attraction b/w + ions & electrons.

V-I Law

Electrical conductivity & drift current

- The electrical conductivity denoted by (σ) is the measure of the ability of a material to conduct an electric current from it. Its units are S/m .
- consider an electron, which is in motion in conductor under the influence of applied electric field. If e is the charge, m is the mass of electron, E is the E.F. the force on electron is given as

$$F = eE \quad \text{--- (1)}$$

Also, $F = ma \quad \text{--- (2)}$

a = acceleration

From (1) & (2)

$$eE = ma$$

$$a = \frac{eE}{m}$$

Rate of change of velocity w.r.t time is acceleration,

$$\therefore \frac{dv}{dt} = \frac{eE}{m}$$

$$dV = \frac{eE}{m} dt$$

on integrating both sides

$$\int dV = \frac{eE}{m} \int dt$$

$$V = \frac{eEt}{m} + C \quad \text{--- (3)}$$

When, $t = 0, V = 0$

From Equation (3)

$$C = 0 \quad \text{--- (4)}$$

From Equation (3) & (4)

$$V = \frac{eEt}{m}$$

When the transfer time is the collision time
that is relaxation time (τ)

$$\therefore \text{Therefore } V = \frac{eE\tau}{m}$$

When the velocity of electron is drift velocity
this drift velocity is the velocity of electrons,
when e^- starts moving opposite to the
applied field. The N_d is the ang. velocity.

so, the drift velocity is given as

$$v_d = \frac{eE}{2m} \quad \text{--- (5)}$$

A.C.T Ohm's law, the current density
 $\uparrow \downarrow J = \sigma E \quad \text{--- (6)}$

Since, the current density is the current
 $\uparrow \downarrow J = \frac{I}{A} \quad \text{--- (7)}$

From equation (6) & (7)

$$\frac{I}{A} = \sigma E$$

$$\sigma = \frac{I}{AE} \quad \text{--- (8)}$$

The drift current is given as

$$I = v_d n A \quad \text{--- (9)}$$

where n = no. of free electrons.

on substituting value of (9) in (8)

$$\sigma = \frac{e v_d n A}{E A}$$

$$\sigma = N d e n$$

On substituting the value of ν_d from equation ⑩ to eq. ⑤.

$$\sigma = \frac{ne^2}{2m} \tau - ⑪$$

Now, the mean free path (λ) - [The avg. dist. travelled by the electrons b/w successive collisions is known as mean free path]

$$\lambda = \tau v.$$

$$\tau = \frac{\lambda}{v} - ⑫$$

On substituting the value of τ from ⑫ to ⑪.

$$\sigma = \frac{ne^2 \lambda}{2m} - ⑬$$

The above expression represent electrical conductivity of free electrons.

Temperature dependency of electric conductivity.

→ Lorentz has applied the K.T.G for free electrons to explain the temp dependency of σ .

$$\frac{1}{2}mv^2 = \frac{3}{2}KT$$

$$mv^2 = 3KT$$

$$mv = \frac{3KT}{v} \quad - \quad (14)$$

From Eq. (13) (15) (14)

$\sigma = \frac{n e^2 \lambda v}{6KT}$

The above expression shows that the electric conductivity, with inc. in temperature or wi decrease.

With the help of above relation, we can measure conductivity & Resistivity of the solids.

Wiedemann - Franz Law.

- This Law gives the ratio of Thermal conductivity & Electrical conductivity.

$$\frac{K}{\sigma} = \frac{3}{2} \left(\frac{k_B}{e} \right)^2 T$$

Boltzmann's constant.

Temperature.

Thermal conductivity

Charge.

Electrical conductivity.

Diffusion current.

- The diffusion current is the current in which charge moves from higher concn to lower concn under the effect of electric field.
- The diffusion current is caused due to unequal distribution inside solids, b/w electrons & holes.

Fermi - Dirac - Distribution.

• Fermi energy & Fermi levels

- The energy of the highest occupied level at 0°C absolute energy is known as Fermi-energy & the energy level is represented as Fermi level.

The Fermi - Dirac function is given as

$$f(E) = \frac{1}{e^{(E-E_f)/KT} + 1}$$

$\therefore E_f$ = Fermi Energy.

k = Boltzmann's constant.

Case 1 :- At absolute temperature $T = 0$

$$\xrightarrow{\frac{E-E_f}{KT}} \text{where } T=0,$$

$$\xrightarrow{\frac{E-E_f}{KT}} \begin{cases} +\infty & E > E_f \\ -\infty & E < E_f. \end{cases}$$

$$f(E) = \frac{1}{e^{-\infty} + 1} = \frac{1}{0+1} = 1, \quad [e^{-\infty} = 0]$$

$$f(E) = \frac{1}{e^{\infty} + 1} = \frac{1}{\infty + 1} = 0.$$

Case 2 :- At some temperature T , $E = E_f$

$$f(E) = \frac{1}{e^{(E-E_f)KT} + 1}$$

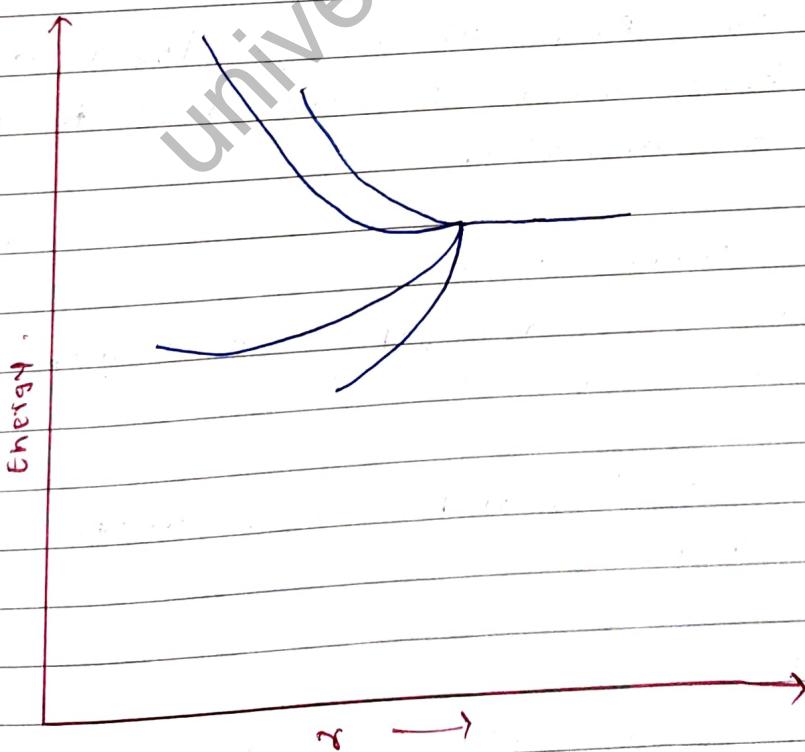
$$f(E) = \frac{1}{2} \quad (e^0 = 1)$$

* So, the Fermi level is the energy level at which the half probability of finding electron is $\frac{1}{2}$. Half

Band Theory of Solids, Formation of Allowed & Forbidden Energy bands.

Electrons are assumed to be tightly bound to the individual atoms, when atoms are bring together to form solids by interaction between the neighboring atoms causes the electrons energy level of individual atoms to be split, this is how the bands are forming.

The band theory follows Pauli's exclusion principle.



(Interatomic distance)

As they decreases the energy level get split into energy band.

Allowed Energy band.

- Allowed energy bands are range of energy levels within a material where electrons are allowed to exist. It is further categorizes into two types that is valence band & conduction band.
- Valence band.
- Range of energy level where electrons are most likely to be found it is highest occupied energy level that is fully occupied at 0 absolute temp.
- Conduction band
- It is the range of energy level above valence band where e^- are free to move within the material.

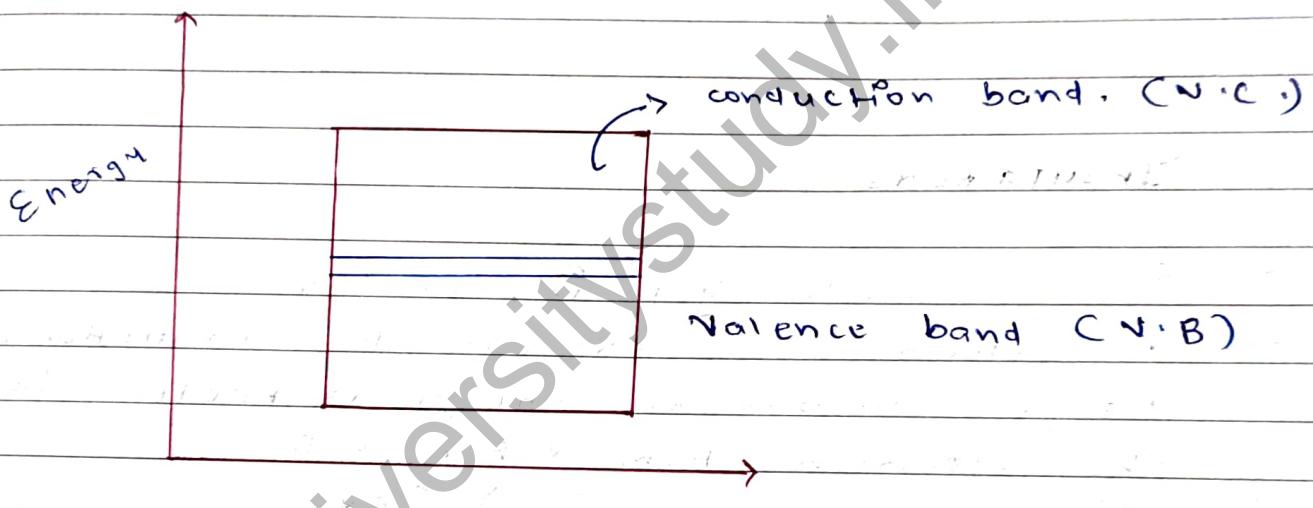
Forbidden Energy bands

- Also known as Energy gap & band gap where e^- are not allowed to exist.

Insulators, semiconductors & conductors in term of Band Gap.

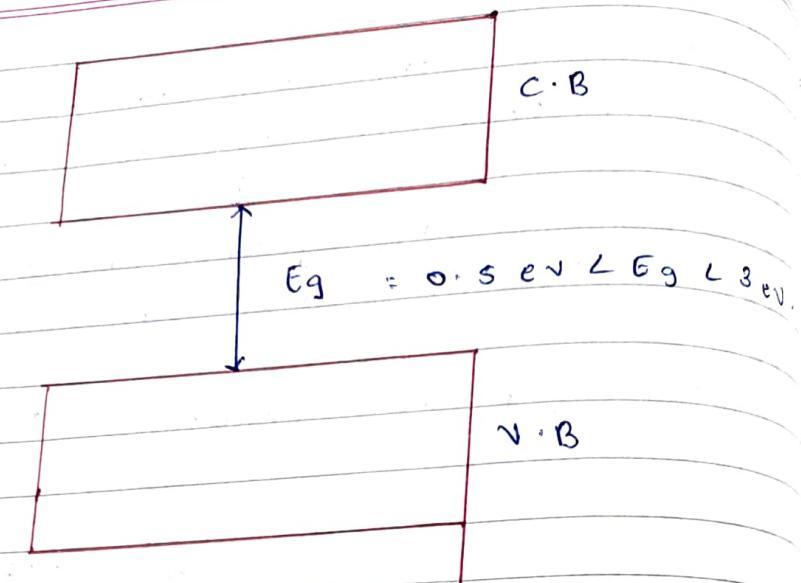
- Conductors.
- In such materials there is no forbidden gap b/w V.B & C.B

- The e^- from valence band to conduction band can freely enter due to overlapping of bands.
- Because of overlapping very low potential diff. can cause continuous flow of current.
- The band gap for conductor is $E_g = 0 \text{ eV}$



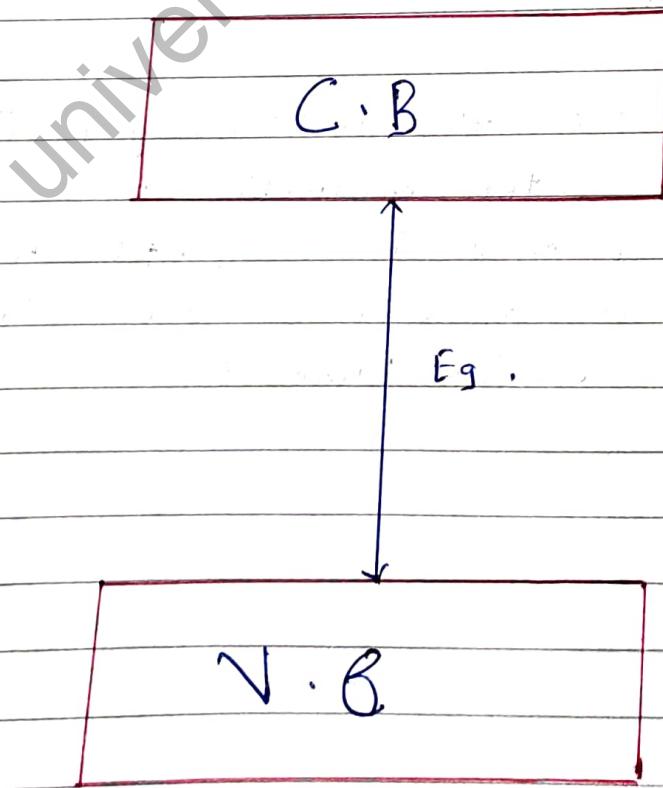
Semiconductors

- In case of S.C. the forbidden gap is very small.
- At OK temp, the conduction band is empty & the valence band is completely filled.
- When a small amount of energy is supplied the electrons can easily jump from V.B. to C.B.
- The energy gap in case of S.C. is $0.05 \text{ eV} < E_g < 3 \text{ eV}$.



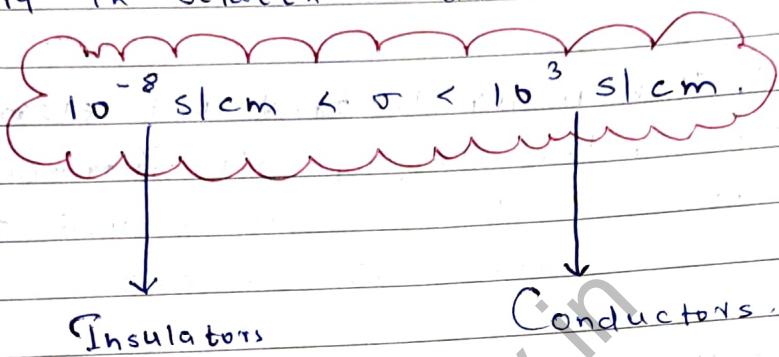
Insulators

- The forbidden gap is very large.
- No electrons are available for conduction.
- Large amount of energy is needed to move the electron from V.B to C.B.
- The energy gap is $> 3\text{ eV}$



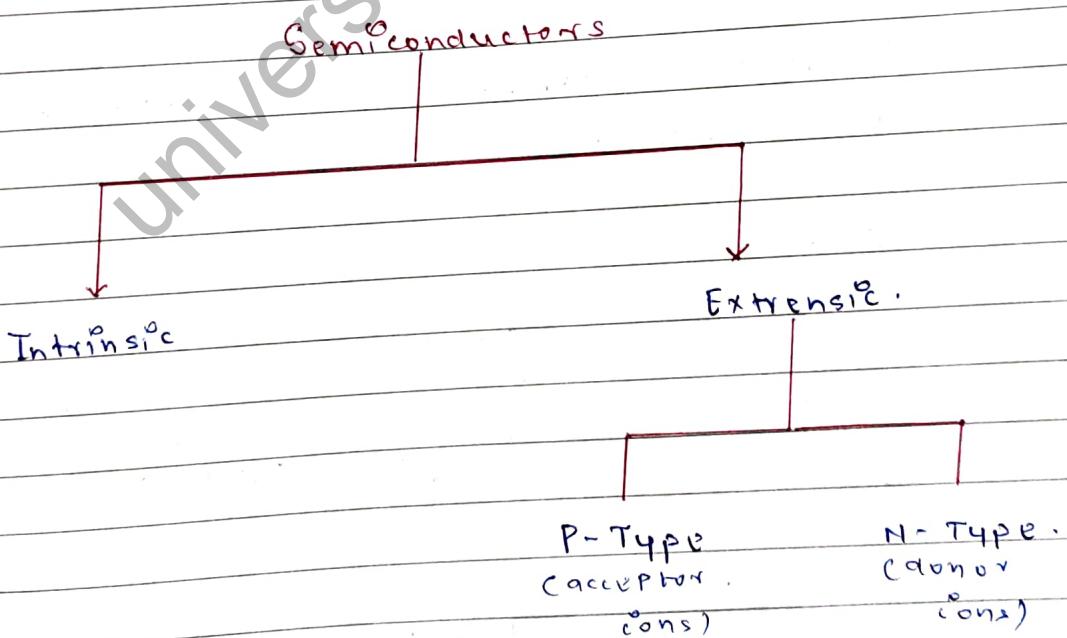
Semiconductors & its Types.

Semiconductors are the materials which have conductivity in between conductors & insulators



Types of semiconductors

The semiconductors are categorized as follow



Intrinsic semiconductors

→ such types of materials are pure semiconductors for e.g., Si & Ge that are not doped with external impurity.

Extrinsic S.C.

→ In such type of materials some specific impurities are doped to modify their electrical properties.

P-Type

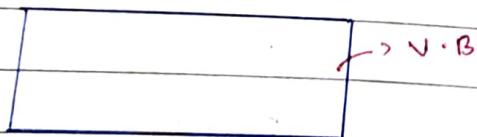
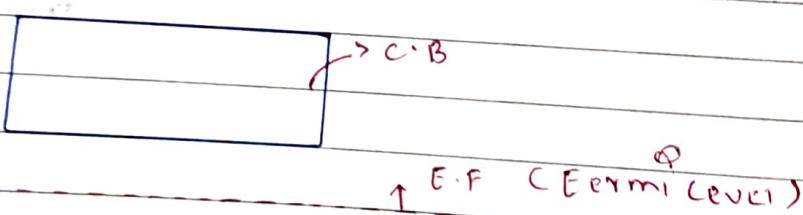
→ In such materials trivalent elements are doped, the doping will introduce excess holes, and the overall conductivity will decrease. For e.g., Group 13 elements.

N-Type

→ In such type of S.C. the pentavalent impurity is doped. The doping will introduce excess e^- , which will enhance the electrical conductivity.

Intrinsic & Extrinsic S.C. on the basis of Fermi Level.

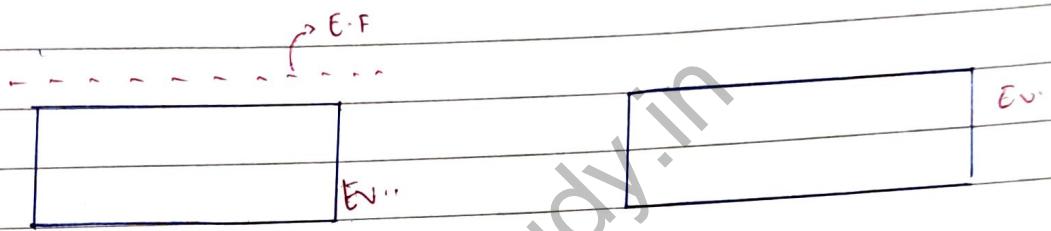
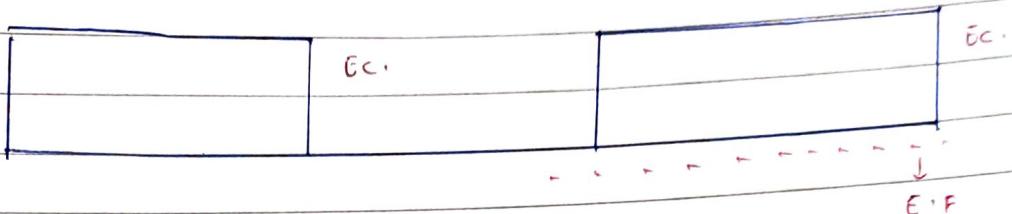
Intrinsic.



Extrinsic.

P-TYPE

N-TYPE:



- The Fermi level in terms of T.s.c will lies below ν_B & C_B
 - In p-type the Fermi level lies to near to the ν_B
 - $N = T M P C$