

SOLUTION(1)

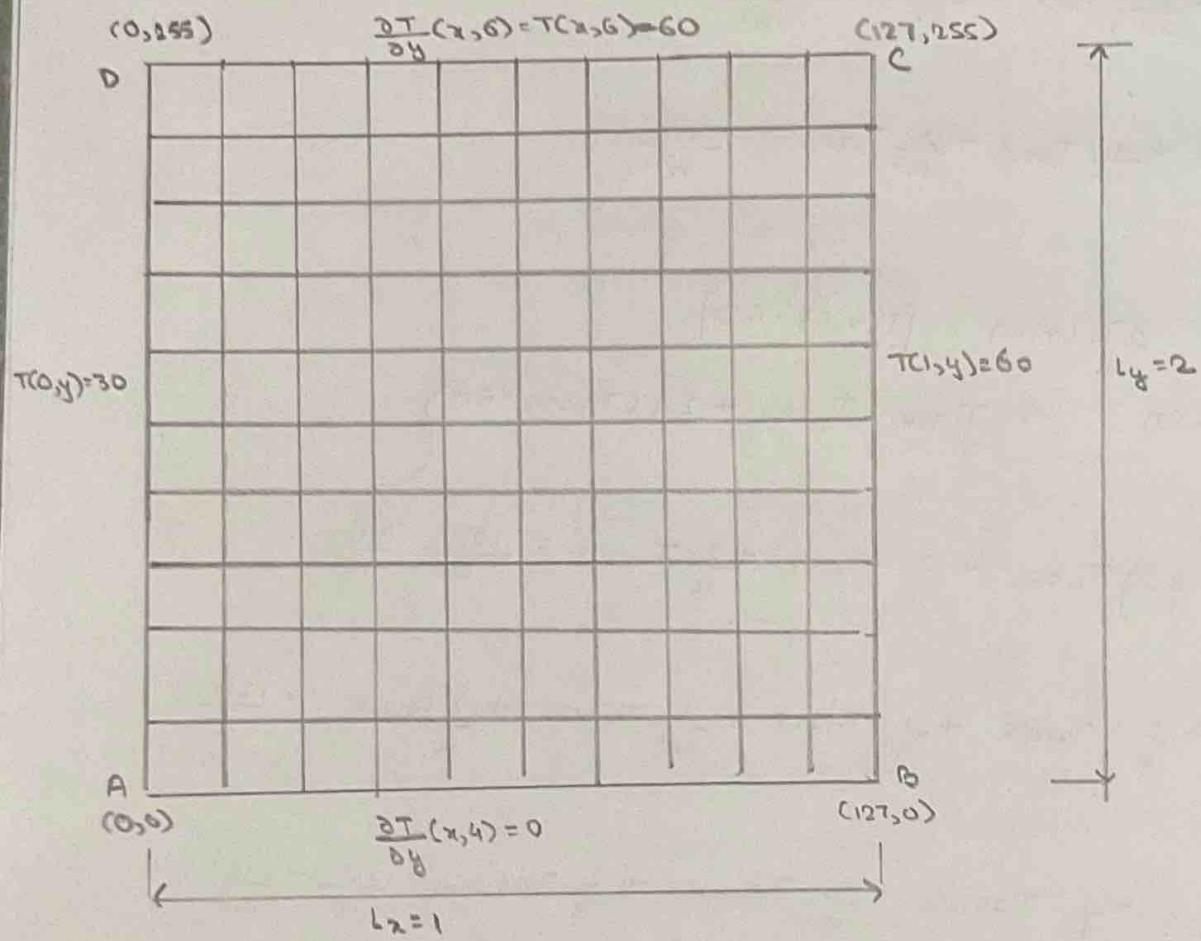
Heat conduction equation: $\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0 \quad \dots (1)$

Discretizing the domain and approximating heat conduction equation using second order central difference scheme:

$$\frac{T_{i+1,j} - 2T_{i,j} + T_{i-1,j}}{h^2} + \frac{T_{i,j+1} - 2T_{i,j} + T_{i,j-1}}{h^2} = 0 \quad \dots (2)$$

Boundary conditions:

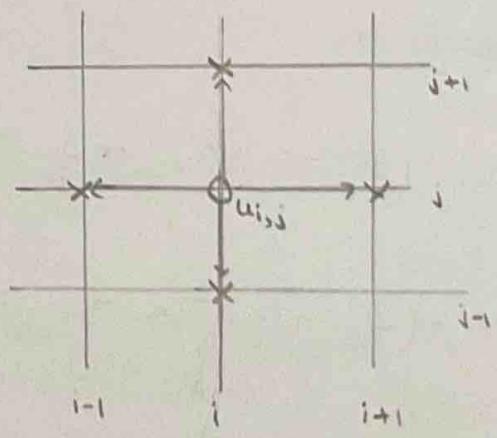
$$(1) T(2,y) = 30 \quad (2) T(3,y) = 60 \quad (3) \frac{\partial T}{\partial y}(x,4) = 0 \quad (4) \frac{\partial T}{\partial y}(x,6) = T(x,6) - 60$$



Given: $N_x = 128$, $N_y = 256$

$$\text{Then, } h = \Delta x = \frac{L_x}{N_x} = \frac{1}{128-1} = \frac{1}{127}$$

$$h_y = \Delta y = \frac{L_y}{N_y} = \frac{2}{256-1} = \frac{2}{255}$$



b) Now, finding coefficient matrix

$$\frac{1}{h^2} T_{i+1,j} - \left(\frac{2}{h^2} + \frac{2}{k^2} \right) T_{i,j} + \frac{1}{h^2} T_{i-1,j} + \frac{1}{k^2} T_{i,j+1} + \frac{1}{k^2} T_{i,j-1} = 0 \quad \dots (2)$$

Bottom surface AB:

$$BC: \frac{\partial T}{\partial y}(x, 4) = 0 \Rightarrow \frac{\partial T}{\partial y} \Big|_{(i,j,0)} = \frac{T_{i,j,1} - T_{i,j,-1}}{2k} = 0$$

$$\therefore T_{i,j,1} = T_{i,j,-1}$$

Node (1, 0):

$$\frac{1}{h^2} T_{0,0} - \left(\frac{2}{h^2} + \frac{2}{k^2} \right) T_{1,0} + \frac{2}{k^2} T_{1,1} = -\frac{30}{h^2}$$

Node (2, 0):

$$\frac{1}{h^2} T_{3,0} - \left(\frac{2}{h^2} + \frac{2}{k^2} \right) T_{2,0} + \frac{1}{h^2} T_{1,0} + \frac{2}{k^2} T_{2,1} = 0$$

Node (126, 0):

$$-\left(\frac{2}{h^2} + \frac{2}{k^2} \right) T_{126,0} + \frac{1}{h^2} T_{125,0} + \frac{2}{k^2} T_{126,1} = -\frac{60}{h^2}$$

Upper surface CD:

$$\text{Boundary condition: } \frac{\partial T}{\partial y}(x, 6) = [T(x, 6) - 60]$$

$$\frac{\partial T}{\partial y} \Big|_{(i,255)} = T_{i,255} - 60 \Rightarrow T_{i,256} = T_{i,254} + 2k(T_{i,255} - 60)$$

Node (1, 255):

$$\frac{1}{h^2} T_{2,255} - \left(\frac{2}{h^2} + \frac{2}{k^2} \right) T_{1,255} + \frac{2}{k^2} T_{1,254} + \frac{2}{k} T_{1,255} = \frac{120}{k} - \frac{30}{h^2}$$

Node (2, 255):

$$\frac{1}{h^2} T_{3,255} - \left(\frac{2}{h^2} + \frac{2}{k^2} \right) T_{2,255} + \frac{1}{h^2} T_{1,255} + \frac{2}{k^2} T_{2,254} + \frac{2}{k} T_{2,255} = \frac{120}{k}$$

Node (126, 255):

$$-\left(\frac{2}{h^2} + \frac{2}{k^2} \right) T_{126,255} + \frac{1}{h^2} T_{125,255} + \frac{2}{k^2} T_{126,254} + \frac{2}{k} T_{126,255} = \frac{120}{k} - \frac{60}{h^2}$$

Interior points

Node (1, 1):

$$\frac{1}{h^2} T_{2,1} - \left(\frac{2}{h^2} + \frac{2}{k^2} \right) T_{1,1} + \frac{1}{k^2} T_{1,2} + \frac{1}{k^2} T_{1,0} = -\frac{30}{h^2}$$

Node (1, 2):

$$\frac{1}{h^2} T_{2,1} - \left(\frac{2}{h^2} + \frac{2}{k^2} \right) T_{1,2} + \frac{1}{k^2} T_{1,3} + \frac{1}{k^2} T_{1,1} = -\frac{30}{h^2}$$

Node (2, 2):

$$\frac{1}{h^2} T_{3,2} - \left(\frac{2}{h^2} + \frac{2}{k^2} \right) T_{2,2} + \frac{1}{h^2} T_{1,2} + \frac{1}{k^2} T_{2,3} + \frac{1}{k^2} T_{2,1} = 0$$

Node (2,1)

$$\frac{1}{h^2} T_{2,11} - \left(\frac{2}{h^2} + \frac{2}{R^2} \right) T_{2,12} + \frac{1}{h^2} T_{1,11} + \frac{1}{R^2} T_{2,12} + \frac{1}{R^2} T_{2,00} = 0$$

Using all the equations we got and arranging them in coefficient matrix

$$\begin{bmatrix} a & \frac{1}{h^2} & 0 \\ \frac{1}{h^2} & a & \frac{1}{h^2} \\ 0 & \frac{1}{h^2} & a \end{bmatrix} \begin{bmatrix} \frac{2}{R^2} \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} T_{1,0} \\ T_{2,0} \\ T_{125,0} \\ T_{126,0} \\ T_{1,11} \\ T_{2,11} \\ T_{125,11} \\ T_{126,11} \end{bmatrix} = \begin{bmatrix} -\frac{30}{h^2} \\ 0 \\ -\frac{30}{h^2} \\ 0 \\ -\frac{60}{h^2} \\ \vdots \\ -\frac{60}{h^2} \end{bmatrix}$$

$$\begin{bmatrix} \frac{2}{R^2} & 0 & 0 \\ 0 & a & \frac{1}{h^2} \\ 0 & \frac{1}{h^2} & a \end{bmatrix} \begin{bmatrix} 0 \\ \frac{2}{R^2} \\ 0 \end{bmatrix} = \begin{bmatrix} T_{1,255} \\ T_{2,255} \\ T_{125,255} \\ T_{126,255} \end{bmatrix} = \begin{bmatrix} \frac{120}{R} - \frac{30}{h^2} \\ \frac{120}{R} \\ \frac{120}{R} \\ \frac{120}{R} - \frac{60}{h^2} \end{bmatrix}$$

$$\text{where, } a = -\left(\frac{2}{h^2} + \frac{2}{R^2}\right) \rightarrow b = -\left(\frac{2}{h^2} + \frac{2}{R^2}\right) + \frac{2}{R}$$

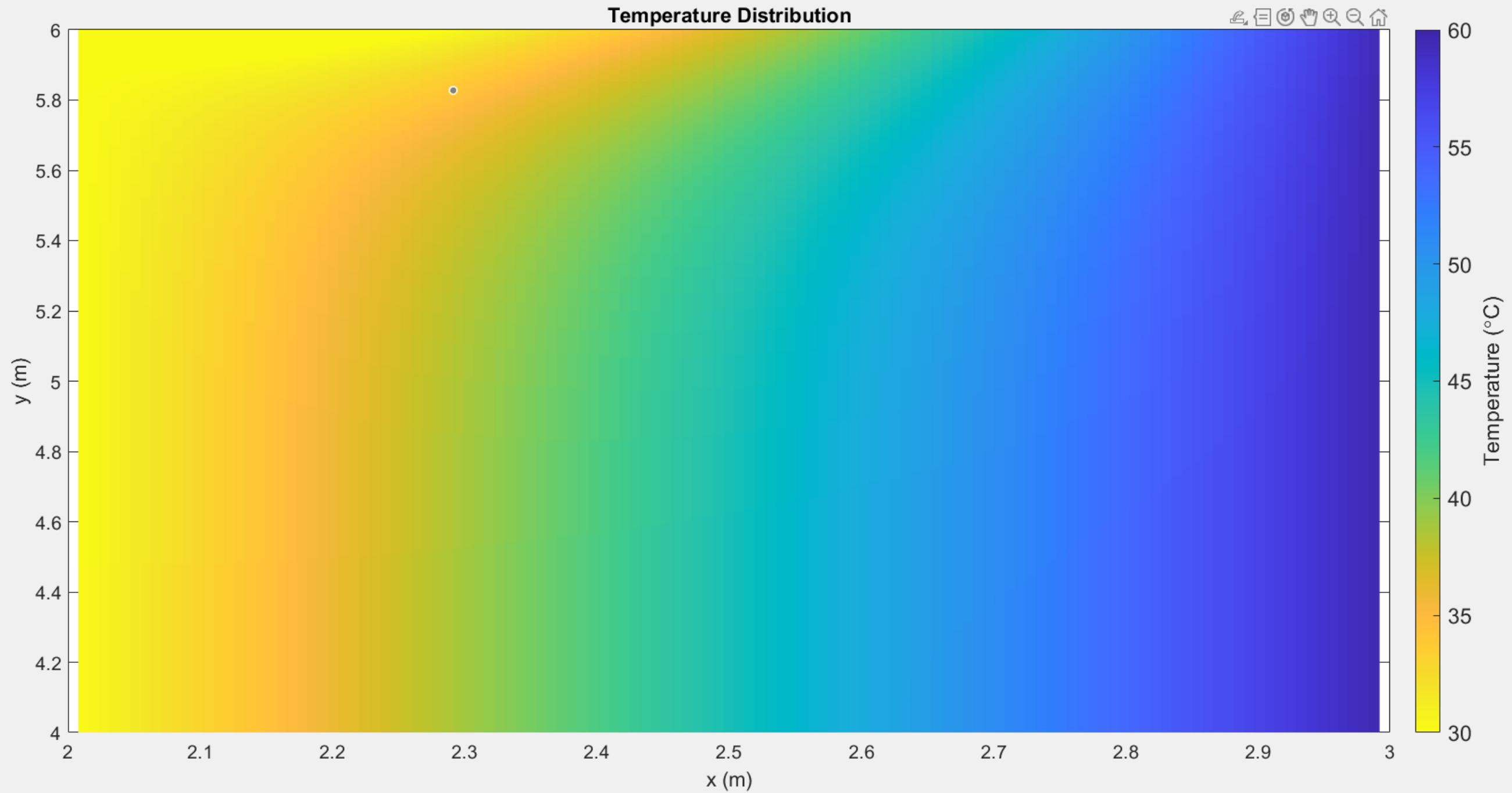
Above matrix can be written in terms of Block Tridiagonal matrix :

$$\begin{bmatrix} [A] & [B] & F \\ [B] & [A] & [B] \\ [B] & [B] & [A] \end{bmatrix} \begin{bmatrix} \bar{T}_0 \\ \bar{T}_1 \\ \bar{T}_2 \\ \vdots \\ \bar{T}_{254} \\ \bar{T}_{255} \end{bmatrix} = \begin{bmatrix} \bar{F}_0 \\ \bar{F}_1 \\ \bar{F}_2 \\ \vdots \\ \bar{F}_{254} \\ \bar{F}_{255} \end{bmatrix}$$

where, $[A] = \begin{bmatrix} a & \frac{1}{h^2} & 0 \\ \frac{1}{h^2} & a & \frac{1}{h^2} \\ 0 & \frac{1}{h^2} & a \end{bmatrix}$

$$[B] = \begin{bmatrix} \frac{2}{R^2} & 0 & 0 \\ 0 & a & \frac{1}{h^2} \\ 0 & \frac{1}{h^2} & a \end{bmatrix}$$

$$[C] = \begin{bmatrix} b & \frac{1}{h^2} & 0 \\ \frac{1}{h^2} & b & \frac{1}{h^2} \\ 0 & \frac{1}{h^2} & b \end{bmatrix}$$



SOLUTION (2)

Brank-Nicholson - Θ scheme:

$$u_t = \alpha u_{xx} \Rightarrow (u_t)_j^m = \alpha [(1-\theta) (u_{xx})_j^m + \theta (u_{xx})_j^{m+1}] \quad \text{where } 0 \leq \theta \leq 1$$

$$\Rightarrow \frac{u_j^{m+1} - u_j^m}{\Delta t} = \alpha(1-\theta) \left[\frac{u_{j+1}^m - 2u_j^m + u_{j-1}^m}{(\Delta x)^2} \right] + \alpha\theta \left[\frac{u_{j+1}^{m+1} - 2u_j^{m+1} + u_{j-1}^{m+1}}{(\Delta x)^2} \right] \quad \dots (1)$$

Let $\sigma = \frac{\alpha \Delta t}{(\Delta x)^2}$ $\Rightarrow \sigma$ is always positive

Using Von Neumann's stability analysis: $u_{j,n} = V^n e^{i\beta ph}$
 where, $\beta = \text{wave no.}$
 $n = \Delta x$

Substituting Ansatz in Equation (1):

$$V^{m+1} e^{i\beta ph} - V^m e^{i\beta ph} = \sigma(1-\theta) \left[V^m e^{i\beta(p+1)h} - 2V^m e^{i\beta ph} + V^m e^{i\beta(p-1)h} \right] + \sigma\theta \left[V^{m+1} e^{i\beta(p+1)h} - 2V^{m+1} e^{i\beta ph} + V^{m+1} e^{i\beta(p-1)h} \right]$$

$$\Rightarrow V^{-1} = \sigma(1-\theta) [e^{i\beta h} - 2 + e^{-i\beta h}] + \sigma\theta [V e^{i\beta h} - 2V + V e^{-i\beta h}]$$

$$\Rightarrow V = 1 + \sigma [V\theta (e^{i\beta h} + e^{-i\beta h} - 2) + (1-\theta) (e^{i\beta h} + e^{-i\beta h} - 2)]$$

$$\Rightarrow V = 1 + \sigma (e^{i\beta h} + e^{-i\beta h} - 2) (V\theta + 1 - \theta)$$

$$\Rightarrow V = \frac{1 - 4\sigma(1-\theta) \sin^2(\frac{\beta h}{2})}{1 + 4\sigma\theta \sin^2(\frac{\beta h}{2})}$$

Stability condition: $|V| \leq 1 \Rightarrow -1 \leq V \leq 1$

Solving Right side inequality:

$$\frac{1 - 4\sigma(1-\theta) \sin^2(\frac{\beta h}{2})}{1 + 4\sigma\theta \sin^2(\frac{\beta h}{2})} \leq 1$$

The above inequality is always satisfied as terms are cancelling.
 NOTE: j is denoted as p in the Von Neumann stability method.

This shows right side inequality is always satisfied

Left ~~and~~ side inequality:

$$-1 \leq \frac{1 - 4\sigma(1-\theta) \sin^2(\frac{\beta h}{2})}{1 + 4\sigma\theta \sin^2(\frac{\beta h}{2})}$$

$$\Rightarrow 2\sigma \sin^2(\frac{\beta h}{2}) (1-2\theta) \leq 1$$

$$\Rightarrow \sigma(1-2\theta) \leq \frac{1}{2 \sin^2(\frac{\beta h}{2})}$$

$$\sigma(1-2\theta) \leq \frac{1}{2 \times 1}$$

{ $\because \sigma(1-2\theta)$ should be less than minimum of RHS
also y}

$$\Rightarrow \sigma(1-2\theta) \leq \frac{1}{2} \quad \dots (2)$$

CASE 1 for $0 \leq \theta \leq \frac{1}{2}$

Then stability condition : $\sigma(1-2\theta) \leq \frac{1}{2}$ (Conditionally stable)

CASE 2 $\frac{1}{2} \leq \theta \leq 1$

Since : $\sigma > 0$ (which is always true)

This shows 2 level scheme with $\theta > 0.5$ is unconditionally stable for time derivative and second order central difference for space derivative.

SOLUTION (3)

Diffusion equation : $u_t = \alpha u_{xx}$

Finite difference approximation of given PDE using explicit Euler and central order central difference in space :

$$\frac{u_i^{n+1} - u_i^n}{\Delta t} = \alpha \left[\frac{u_{i-1}^n - 2u_i^n + u_{i+1}^n}{(\Delta x)^2} \right] \quad \dots (1)$$

$$\text{Let } \sigma_\alpha = \frac{\alpha \Delta t}{(\Delta x)^2}$$

$$\text{Then, } \frac{u_i^{n+1} - u_i^n}{\Delta t} = \sigma_\alpha \left[u_{i-1}^n - 2u_i^n + u_{i+1}^n \right] \Rightarrow u_i^{n+1} = u_i^n + \sigma_\alpha [u_{i+1}^n - 2u_i^n + u_{i-1}^n] \quad \dots (2)$$

Taylor expanding the terms about u_i^n :

$$u_i^{n+1} = u_i^n + h \frac{\partial u}{\partial t} \Big|_{(x_i, t_n)} + \frac{h^2}{2} \frac{\partial^2 u}{\partial t^2} \Big|_{(x_i, t_n)} + \frac{h^3}{6} \frac{\partial^3 u}{\partial t^3} \Big|_{(x_i, t_n)} + \dots$$

$$u_{i+1}^n = u_i^n + h \frac{\partial u}{\partial x} \Big|_{(x_i, t_n)} + \frac{h^2}{2} \frac{\partial^2 u}{\partial x^2} \Big|_{(x_i, t_n)} + \frac{h^3}{6} \frac{\partial^3 u}{\partial x^3} \Big|_{(x_i, t_n)} + \frac{h^4}{24} \frac{\partial^4 u}{\partial x^4} \Big|_{(x_i, t_n)} + \dots$$

$$u_{i-1}^n = u_i^n - h \frac{\partial u}{\partial x} \Big|_{(x_i, t_n)} + \frac{h^2}{2} \frac{\partial^2 u}{\partial x^2} \Big|_{(x_i, t_n)} - \frac{h^3}{6} \frac{\partial^3 u}{\partial x^3} \Big|_{(x_i, t_n)} + \frac{h^4}{24} \frac{\partial^4 u}{\partial x^4} \Big|_{(x_i, t_n)} + \dots$$

Putting values of u_i^{n+1} , u_{i+1}^n & u_{i-1}^n in equation (2):

$$u_i^n + \left[h \frac{\partial u}{\partial t} \Big|_{(x_i, t_n)} + \frac{h^2}{2} \frac{\partial^2 u}{\partial t^2} \Big|_{(x_i, t_n)} + \frac{h^3}{6} \frac{\partial^3 u}{\partial t^3} \Big|_{(x_i, t_n)} + \frac{h^4}{24} \frac{\partial^4 u}{\partial t^4} \Big|_{(x_i, t_n)} + \dots \right] - 2u_i^n + \sigma_\alpha \left[u_i^n + h \frac{\partial u}{\partial x} \Big|_{(x_i, t_n)} + \frac{h^2}{2} \frac{\partial^2 u}{\partial x^2} \Big|_{(x_i, t_n)} + \frac{h^3}{6} \frac{\partial^3 u}{\partial x^3} \Big|_{(x_i, t_n)} + \frac{h^4}{24} \frac{\partial^4 u}{\partial x^4} \Big|_{(x_i, t_n)} + \dots \right] + u_i^n - h \frac{\partial u}{\partial x} \Big|_{(x_i, t_n)} + \frac{h^2}{2} \frac{\partial^2 u}{\partial x^2} \Big|_{(x_i, t_n)} - \frac{h^3}{6} \frac{\partial^3 u}{\partial x^3} \Big|_{(x_i, t_n)} + \frac{h^4}{24} \frac{\partial^4 u}{\partial x^4} \Big|_{(x_i, t_n)} - \dots = 0$$

$$\text{Putting } \sigma_\alpha = \frac{\alpha \Delta t}{(\Delta x)^2} = \frac{\alpha h}{h^2}$$

$$\left[\frac{\partial u}{\partial t} + \frac{h}{2} \frac{\partial^2 u}{\partial x^2} + \frac{h^2}{6} \frac{\partial^3 u}{\partial x^3} + \dots \right] = \frac{\alpha}{h^2} \left[\frac{h^2 \partial^2 u}{\partial x^2} + \frac{h^4 \partial^4 u}{\partial x^4} + \frac{h^6 \partial^6 u}{\partial x^6} + \dots \right] \Big|_{(x_i, t_n)}$$

$$(u_2)_i^n - \alpha (u_{xx})_i^n = \alpha \left[\frac{h^2}{12} \frac{\partial^4 u}{\partial x^4} + \frac{h^4}{360} \frac{\partial^6 u}{\partial x^6} + O(h^6) \right] - \left[\frac{h}{2} \left(\frac{\partial^2 u}{\partial x^2} \right) + \frac{h^2}{6} \left(\frac{\partial^3 u}{\partial x^3} \right) + \dots \right] \Big|_{(x_i, t_n)}$$

$$(u_2)_i^n - \alpha (u_{xx})_i^n = -\frac{h}{2} \left(\frac{\partial^2 u}{\partial x^2} \right) \Big|_{(x_i, t_n)} + O(h^2) + \frac{\alpha h^2}{12} \left(\frac{\partial^4 u}{\partial x^4} \right) \Big|_{(x_i, t_n)} + O(h^4)$$

We know that, $u_2 = \alpha u_{xx} \Rightarrow u_{xx} = \alpha u_{xxx} = \alpha^2 u_{xxxx}$

$$\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2} \Rightarrow \frac{\partial}{\partial t} \left(\frac{\partial u}{\partial x} \right) = \alpha \frac{\partial}{\partial t} \left(\frac{\partial^2 u}{\partial x^2} \right) = \alpha^2 \frac{\partial^2}{\partial x^2} \left(\frac{\partial u}{\partial x} \right) = \alpha^2 \frac{\partial^4 u}{\partial x^4}$$

$$(u_2)_i^n - \alpha (u_{xx})_i^n = -\frac{\alpha^2 h}{2} \left(\frac{\partial^4 u}{\partial x^4} \right) \Big|_{(x_i, t_n)} + \frac{\alpha h^2}{12} \left(\frac{\partial^4 u}{\partial x^4} \right) \Big|_{(x_i, t_n)} + O(h^2 + h^4)$$

$$(u_2)_i^n - \alpha (u_{xx})_i^n = \left(-\frac{\alpha^2 h}{2} + \frac{\alpha h^2}{12} \right) (u_{xxxx})_i^n + O(h^2 + h^4) \quad (\text{MODIFIED EQUATION})$$

We know that, if leading order term has even derivative in space, then the nature of error is DISSIPATIVE.

SOLUTION (4)

Diffusion equation: $u_t = \alpha u_{xx}$, $\alpha = 1/5$



Initial condition: $u(x, 0) = \sin(4x) + \sin(x)$

Boundary conditions (Periodic):

$$u(0, t) = u(2\pi, t) \text{ and } u_x(0, t) = u_x(2\pi, t)$$

Using separation of variables: $u(x, t) = X(x) T(t)$

Plugging $u = X T$ into the heat equation:

$$\begin{aligned} X \frac{dT}{dt} &= \alpha T \frac{d^2X}{dx^2} \Rightarrow \frac{dT}{\alpha T} = \frac{1}{X} \frac{d^2X}{dx^2} = -\lambda \\ \Rightarrow \frac{\dot{T}}{\alpha T} &= \frac{\alpha X''}{X} = -\lambda \end{aligned}$$

$$\text{Solving for } T: \frac{\dot{T}}{\alpha T} = -\lambda \Rightarrow \int \frac{\dot{T}}{\alpha T} dt = \int -\lambda dx \Rightarrow T(t) = A e^{-\lambda t}$$

Solving second ODE:

$$X'' + \frac{\lambda}{\alpha} X = 0 \Rightarrow X'' + \gamma^2 X = 0 \quad \left\{ \text{let } \gamma^2 = \frac{\lambda}{\alpha} \right\}$$

Solution of ODE: $X(x) = C \cos(\gamma x) + D \sin(\gamma x)$

Using periodic boundary conditions:

$$u(0, t) = u(2\pi, t) \Rightarrow X(0) = X(2\pi) \quad \boxed{\gamma = n\pi} \quad \text{where, } n = 0, \pm 1, \pm 2, \pm 3, \dots$$

$$u_x(0, t) = u_x(2\pi, t) \Rightarrow X'(0) = X'(2\pi)$$

$$X_n(x) = C_n \cos(n\pi x) + D_n \sin(n\pi x)$$

$$u_n(x, t) = [C_n \cos(n\pi x) + D_n \sin(n\pi x)] e^{-n\pi t}$$

Principle of superposition shows that if u_1, u_2, \dots, u_m are solution linear homogeneous problem, then any linear combination, $c_1 u_1 + c_2 u_2 + \dots + c_m u_m$ is also the solutions.

$$u(x, t) = \sum_{n=0}^{\infty} [A_n \cos(n\pi x) + B_n \sin(n\pi x)] e^{-n\pi t}$$

Now using initial condition:

$$u(x, 0) = \sin(4x) + \sin(x) = \sum_{n=0}^{\infty} A_n \cos(n\pi x) + B_n \sin(n\pi x) \quad \dots (2)$$

Now, we know the orthogonality of sines & cosines:

$$\int_0^{2\pi} \sin\left(\frac{n\pi x}{L}\right) \sin\left(\frac{m\pi x}{L}\right) dx = \begin{cases} 0 & , m \neq n \\ \frac{4}{\pi} & , m = n \end{cases}$$

$$\int_0^{2\pi} \sin\left(\frac{n\pi x}{L}\right) \cos\left(\frac{m\pi x}{L}\right) dx = 0 \quad \text{always}$$

Multiplying equation (2) with $\sin(4x)$ and integrate :

$$\int_0^{2\pi} \sin^2(4x) dx + \int_0^{2\pi} \sin x \sin 4x dx = \sum_{n=-\infty}^{\infty} \int_0^{2\pi} A_n \cos(nx) \sin 4x dx + \sum_{n=-\infty}^{\infty} \int_0^{2\pi} B_n \sin(nx) \sin 4x dx$$

Using orthogonality of sines & cosines : Only term remaining on RHS is B_4 :

$$\int_0^{2\pi} \sin^2(4x) dx = 0 + B_4(\pi)$$

$$\Rightarrow B_4 = 1$$

Similarly, multiplying equation with $\sin x$ and integrate :

$$\int_0^{2\pi} \sin 4x \sin x dx + \int_0^{2\pi} \sin^2 x dx = \sum_{n=-\infty}^{\infty} \int_0^{2\pi} A_n \cos(nx) \sin x dx + \sum_{n=-\infty}^{\infty} \int_0^{2\pi} B_n \sin(nx) \sin x dx$$

$$0 + \int_0^{2\pi} \sin^2 x dx = B_1 \int_0^{2\pi} \sin x dx$$

$$\Rightarrow B_1 = 1$$

Multiplying equation (2) with $\cos(nx)$ gives $A_m = 0$

$$\therefore u(x,t) = \sum_{m=1,4} e^{-\alpha^2 t} \sin(mx) = e^{-\alpha t} \sin x + e^{-16\alpha t} \sin(4x)$$

$$u(x,t) = e^{-1.5t} \sin x + e^{-24t} \sin(4x) \quad \Rightarrow \text{ANALYTICAL SOLUTION}$$

And discretized equations used in the code is same as we obtained in ~~solution~~ solution (3) :

$$\frac{u_i^{m+1} - u_i^m}{\Delta t} = \alpha \left[\frac{u_{i-1}^m - 2u_i^m + u_{i+1}^m}{(\Delta x)^2} \right]$$

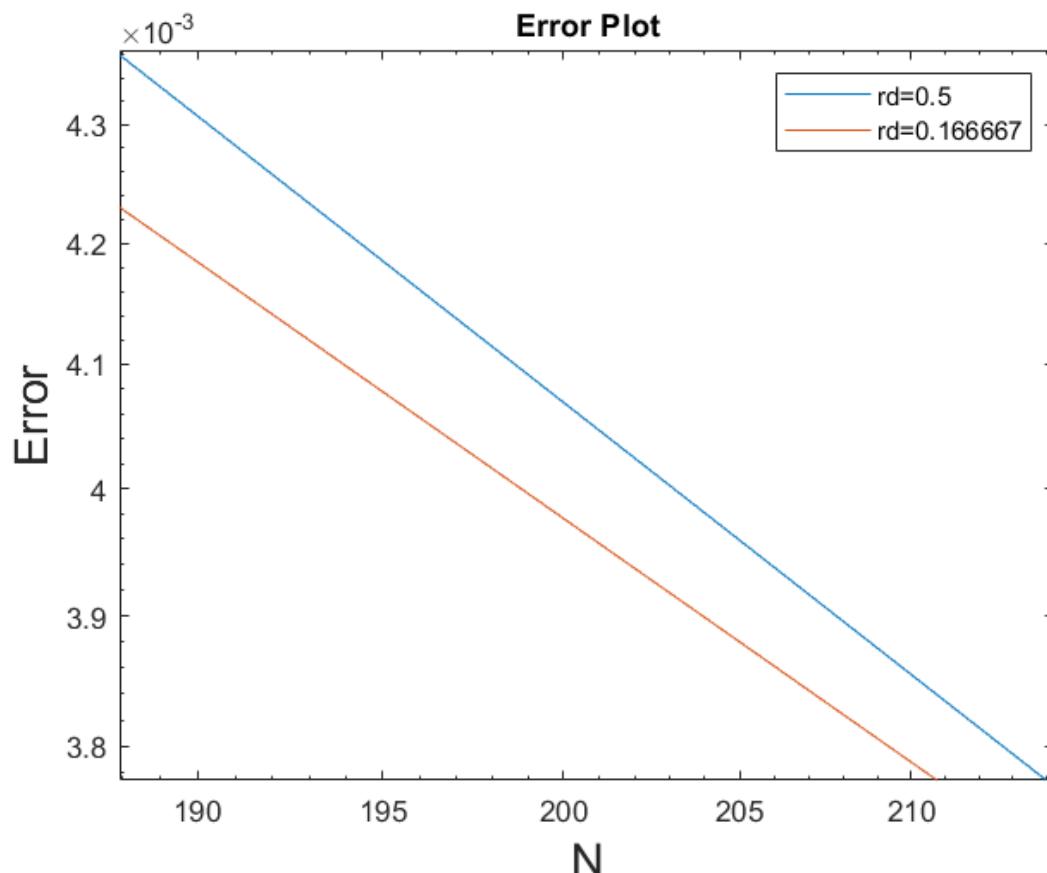
SOLUTION 4

(c)

	rd	N	E(N)
	Number	▼ Number	▼ Number
1	0.5	32	0.0309577
2	0.5	64	0.0118341
3	0.5	128	0.00667739
4	0.5	256	0.00309388
5	0.166667	32	0.0265206
6	0.166667	64	0.0122321
7	0.166667	128	0.00619624
8	0.166667	256	0.00311103

(d)

Logarithmic Scale: Plot N vs E(N)



The order of accuracy for $rd = 1/2$ is 2 in space and 1 in time

The order of accuracy for $rd = 1/6$ is 4 in space and 2 in time

(e)

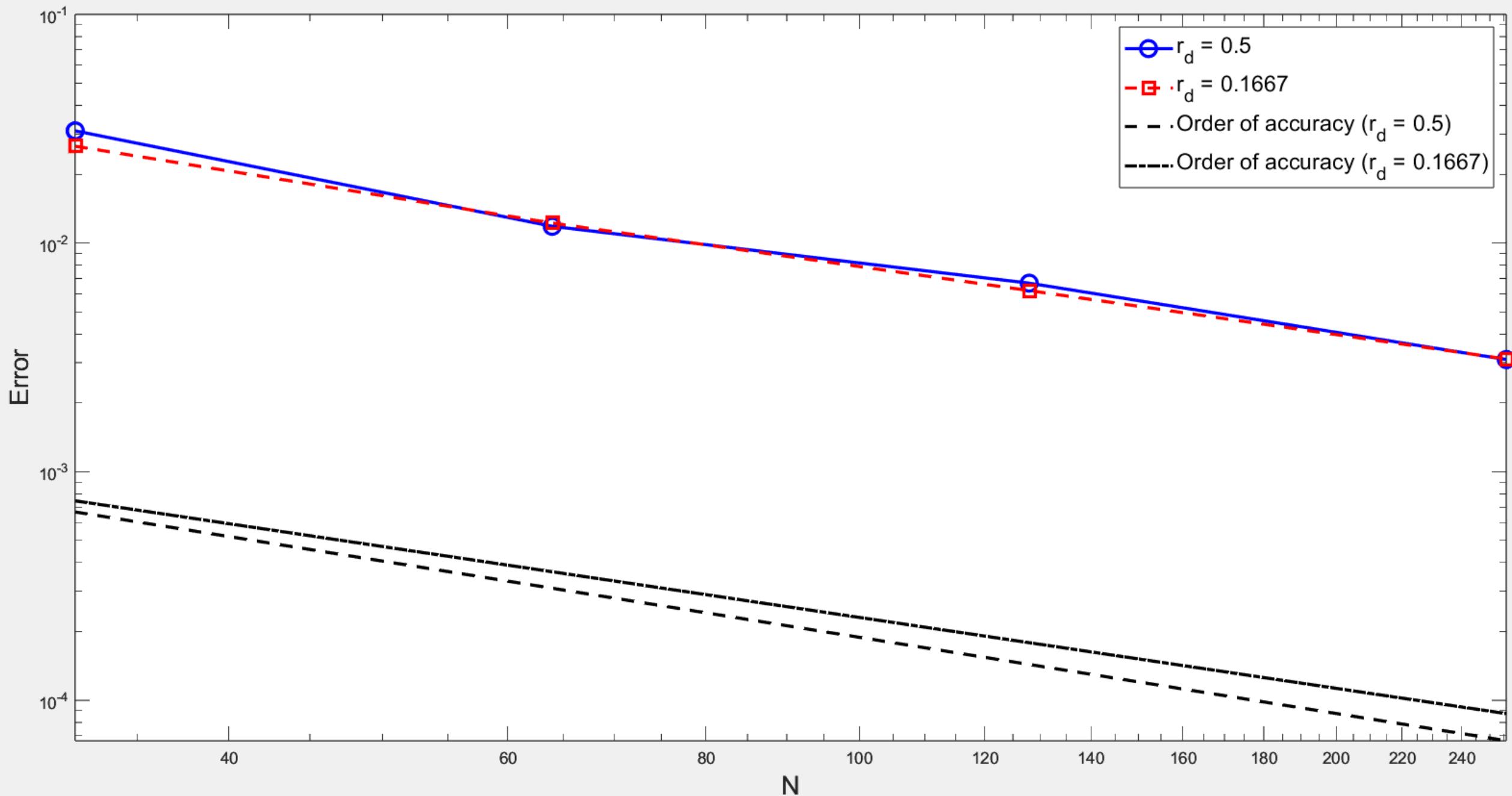
e) Show derived order of accuracy modified equation in solution (3).

$$(u_t)_i^n - \alpha (u_{xx})_i^n = \left(-\alpha^2 \frac{h}{2} + \frac{\alpha h^2}{12} \right) (u_{xxxx})_i^n + O(h^2 + h^4)$$

$$\text{If } -\frac{\alpha^2 h}{2} + \frac{\alpha h^2}{12} = 0 \Rightarrow \frac{\alpha h^2}{12} = \frac{\alpha^2 h}{2} \Rightarrow \alpha h = \frac{h^2}{6} \Rightarrow \frac{\alpha h}{h^2} = \frac{1}{6}$$

i.e., $\sigma_d = \frac{1}{6}$

Then, Order of accuracy is $O(h^2 + h^4) \Rightarrow$ Justifies abnormal order of accuracy observed in $\sigma_d = 1/6$ case.



SOLUTION (5)

Diffusion equation : $u_t = \alpha u_{xx}$

Using implicit Euler and second order central difference :

$$\frac{u_i^{n+1} - u_i^n}{\Delta t} = \alpha \left[\frac{u_{i-1}^{n+1} - 2u_i^{n+1} + u_{i+1}^{n+1}}{(\Delta x)^2} \right]$$

$$\text{Let } \sigma_\alpha = \frac{\alpha \Delta t}{(\Delta x)^2}$$

$$u_i^{n+1} - u_i^n = \sigma_\alpha [u_{i-1}^{n+1} - 2u_i^{n+1} + u_{i+1}^{n+1}] \quad \rightarrow \text{Discretized equation}$$

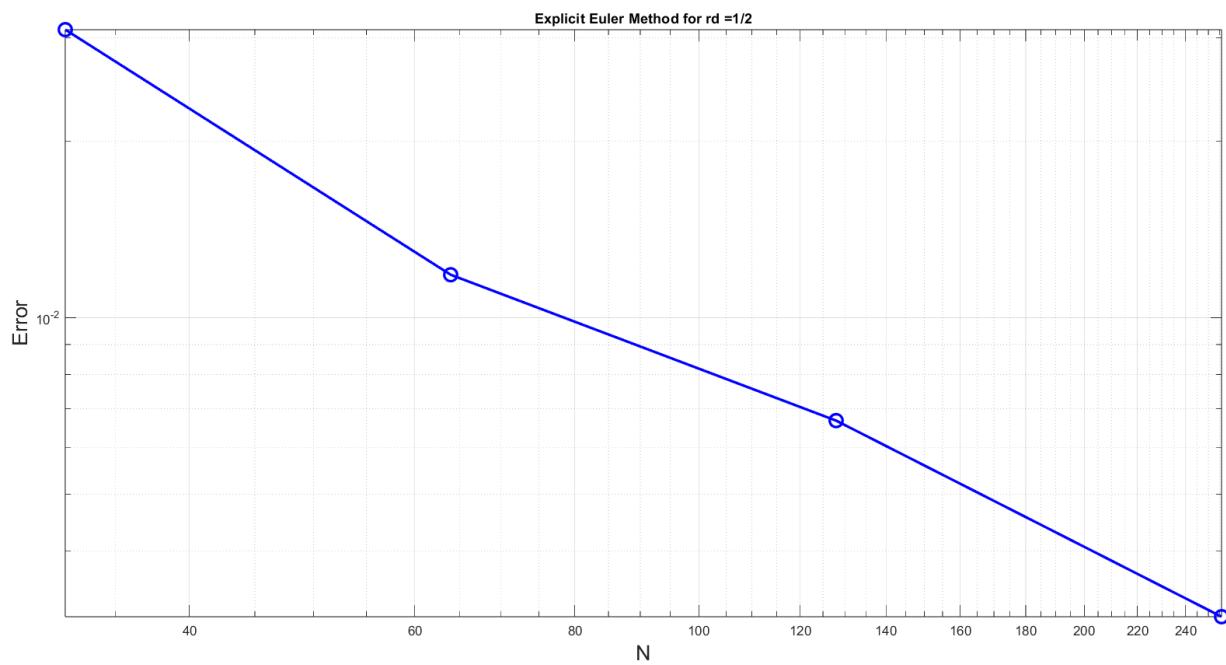
Initial condition

$$(u_0)^{100} = 5 + (u_1)^{100} = 5 = (u_2)^{100} = 5 = \dots = (u_{100})^{100} = 5$$

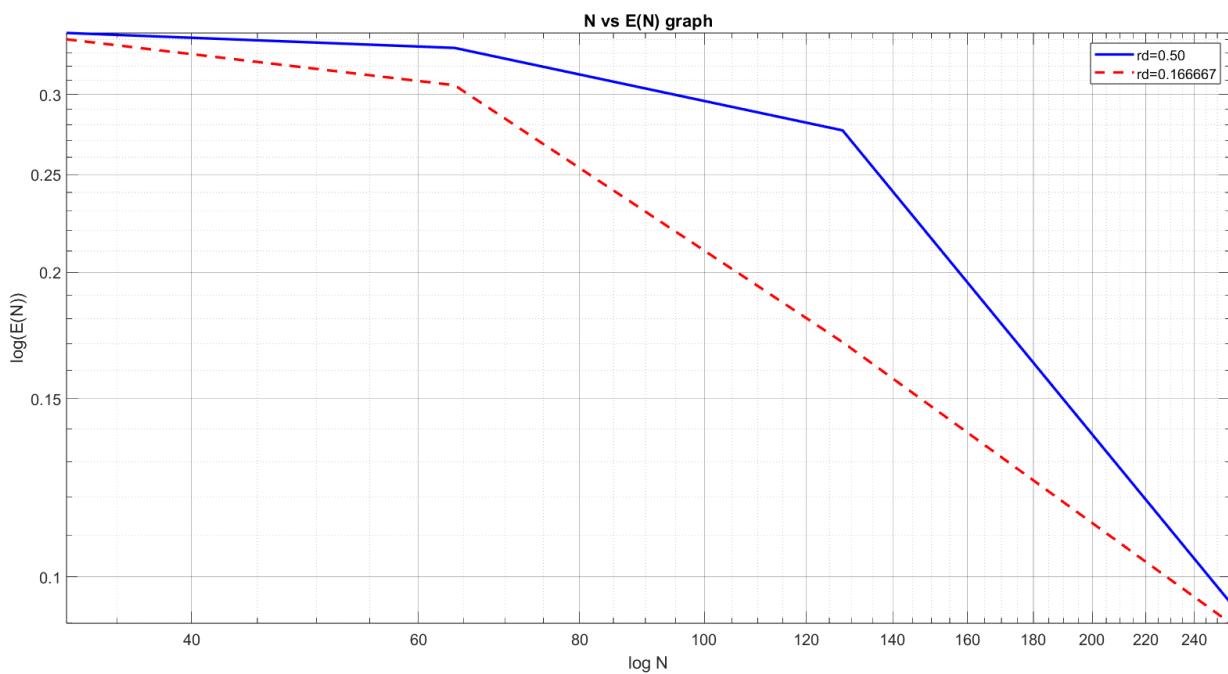
$$\left[\frac{(u_1)^{100} - 5}{\sigma_\alpha} + (u_2)^{100} - 5 = 0 \right] \dots$$

SOLUTION 5

Explicit Euler

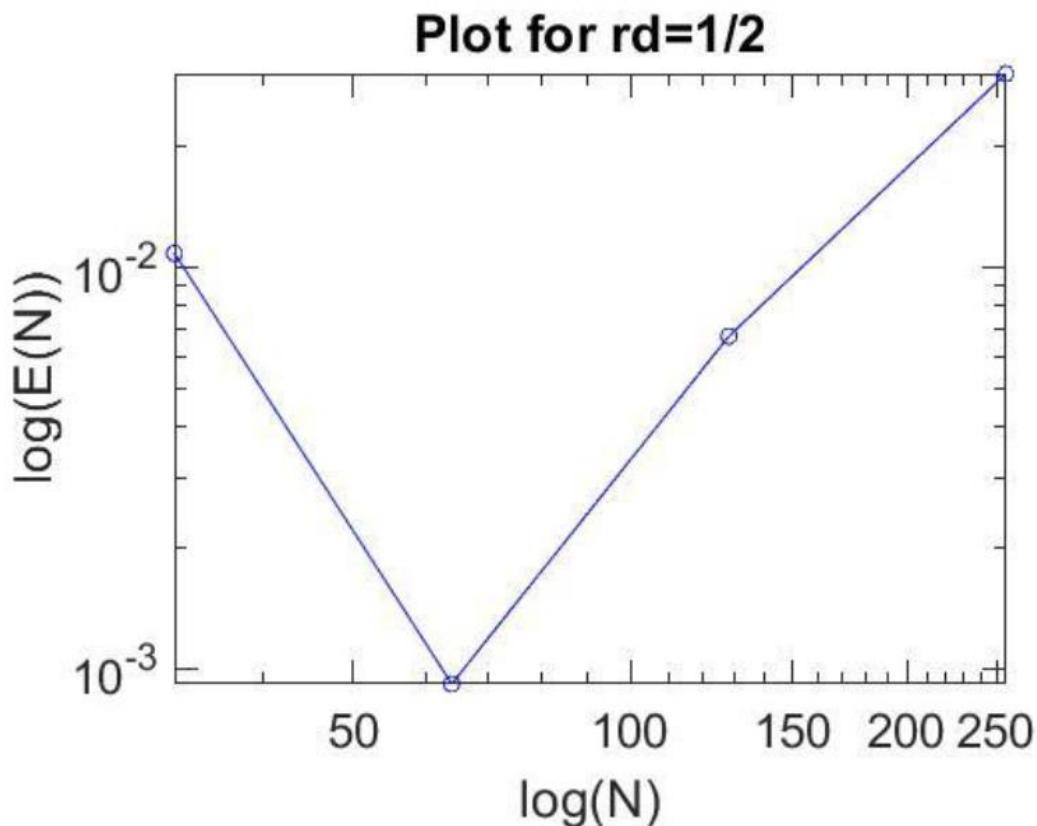


Implicit Euler



(b) *Observations*

Plot for $rd = 1/2$



- 1) As we can see, the implicit method is unconditionally stable.
- 2) Explicit method is stable only when $rd < 0.25$
- 3) Implicit method is better than the explicit one, as it requires fewer steps.
- 4) Implicit method is more accurate and stable

SOLUTION (6)

$$\text{Transient 1D Heat Conduction : } \frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2} \quad \dots (1)$$

$$\text{Given discretized equation : } \frac{T_i^{n+1} - T_i^n}{\Delta t} = \alpha \left[\frac{T_{i+1}^{n-h} - 2T_i^n + T_{i-1}^h}{(\Delta x)^2} \right] \quad \dots (2)$$

Expanding terms by Taylor series :

$$\text{Truncation error, } \gamma_i^n = \frac{\partial T}{\partial t} - \alpha \frac{\partial^2 T}{\partial x^2} = \frac{T_i^{n+1} - T_i^n}{\Delta t} \quad \text{implies} \quad -\alpha \left[\frac{T_{i+1}^{n-h} - 2T_i^n + T_{i-1}^h}{(\Delta x)^2} \right] \quad \dots (3)$$

Solving LHS of equation (2) :

$$\begin{aligned} \frac{T_i^{n+1} - T_i^n}{\Delta t} &= \frac{1}{\Delta t} \left[T_i^n + \Delta t \frac{\partial T}{\partial t} + \frac{(\Delta t)^2}{2!} \frac{\partial^2 T}{\partial t^2} + \frac{(\Delta t)^3}{3!} \frac{\partial^3 T}{\partial t^3} + \dots - T_i^n \right]_i \\ &= \frac{\partial T}{\partial t} + \left[\frac{(\Delta t)}{2!} \frac{\partial^2 T}{\partial t^2} + \frac{(\Delta t)^2}{3!} \frac{\partial^3 T}{\partial t^3} \right]_i + O(\Delta t)^3 \end{aligned}$$

$$\begin{aligned} \text{Solving: } T_{i+1}^{n-h} - 2T_i^n + T_{i-1}^h &= \left[T + (\Delta x) \frac{\partial T}{\partial x} - k \Delta t \frac{\partial T}{\partial t} + \frac{(\Delta x)^2}{2} \frac{\partial^2 T}{\partial x^2} + (\Delta x)(-k \Delta t) \frac{\partial^2 T}{\partial x \partial t} \right. \\ &\quad + \frac{(-k \Delta t)^2}{2} \frac{\partial^2 T}{\partial t^2} + \frac{(\Delta x)^3}{6} \frac{\partial^3 T}{\partial x^3} + \frac{(\Delta x)^2 (-k \Delta t)}{2} \frac{\partial^3 T}{\partial x^2 \partial t} \\ &\quad + \frac{(\Delta x)(-k \Delta t)^2}{3} \frac{\partial^3 T}{\partial x \partial t^2} + \frac{(-k \Delta t)^3}{6} \frac{\partial^3 T}{\partial t^3} + \frac{(\Delta x)^4}{24} \frac{\partial^4 T}{\partial x^4} - 2T \\ &\quad \left. + T - \Delta x \frac{\partial T}{\partial x} + \frac{(\Delta x)^2}{2} \frac{\partial^2 T}{\partial x^2} - \frac{(\Delta x)^3}{6} \frac{\partial^3 T}{\partial x^3} + \frac{(\Delta x)^4}{24} \frac{\partial^4 T}{\partial x^4} + \dots \right]_i \\ &= \left[-k \Delta t \frac{\partial T}{\partial t} + \frac{(\Delta x)^2}{2} \frac{\partial^2 T}{\partial x^2} - k \Delta t \Delta x \frac{\partial^2 T}{\partial x \partial t} - \frac{(k \Delta t)^3}{2} \frac{\partial^2 T}{\partial t^2} - \frac{k \Delta t (\Delta x)^2}{2} \frac{\partial^3 T}{\partial x \partial t} + \frac{k^2 \Delta x \Delta t^2}{2} \frac{\partial^3 T}{\partial x \partial t^2} \right. \\ &\quad \left. - \frac{k^2 \Delta t^3}{6} \frac{\partial^3 T}{\partial t^3} + \frac{\Delta x^4}{12} \frac{\partial^4 T}{\partial x^4} \right]_i \end{aligned}$$

$$\begin{aligned} \gamma_i^n &= \frac{T_i^{n+1} - T_i^n}{\Delta t} = \alpha \left[\frac{T_{i+1}^{n-h} - 2T_i^n + T_{i-1}^h}{(\Delta x)^2} \right] \\ &= \underbrace{\frac{\partial T}{\partial t} - \alpha \frac{\partial^2 T}{\partial x^2}}_{=0} + \left[\frac{\Delta t}{2} \frac{\partial^2 T}{\partial t^2} + \frac{(\Delta t)^2}{6} \frac{\partial^3 T}{\partial t^3} + \frac{\alpha k \Delta t}{(\Delta x)^2} \frac{\partial T}{\partial x} + \frac{\alpha k \Delta t}{\Delta x} \frac{\partial^2 T}{\partial x \partial t} + \frac{\alpha k^2 (\Delta t)^2}{2(\Delta x)^2} \frac{\partial^2 T}{\partial t^2} \right. \\ &\quad \left. + \frac{\alpha k x \Delta t}{2} \frac{\partial^3 T}{\partial x^2 \partial t} - \frac{\alpha k^2 (\Delta t)^2}{2 \Delta x} \frac{\partial^3 T}{\partial x \partial t^2} + \frac{\alpha k^3 (\Delta t)^3}{6 (\Delta x)^2} \frac{\partial^3 T}{\partial t^3} + \dots \right]_i \end{aligned}$$

$$\gamma_i^n = O(\Delta t, (\Delta x)^2, \frac{k \Delta t}{\Delta x}, \frac{k \Delta t}{(\Delta x)^2}) = \gamma(x, t)$$

An method is said to be consistent if $\gamma(x, t) \rightarrow 0$ as $k, h \rightarrow 0$

But as the denominator contains $\Delta x, \Delta t$. Then the expression will tend to infinity.

\therefore Given finite difference approximation is NOT CONSISTENT.