

# Intro to Free Probability Theory

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## My Motivation

- Given two matrices :  $A \times B$  (with eigenvals/vecs)
- What can we say about the eigenvalues  
of  $A+B$  (or)  $\underline{A \cdot B}$
- Depends on the eigenvalues & eigenvectors  
of  $A \times B$
- Tools I know:
  - Trace method.
  - Weyl's ineq.

: not always useful.

$$\mathbb{E}[(A+B)^2] = \mathbb{E}[A^2] + \mathbb{E}[B^2] + 2\mathbb{E}[A]\mathbb{E}[B]$$

What if Eigen vectors are "uncorrelated".

look at  $A + Q^T B Q$        $Q$  is a random rotation matrix

eigenvecs of  $B$ :  $b_1, \dots, b_n$   
" "  $Q^T B Q$ :  $Q^T b_1, \dots, Q^T b_n$

can we say something about eigenvals of  
 $A + Q^T B Q$  now?

"Takes eigenvecs out of the question".

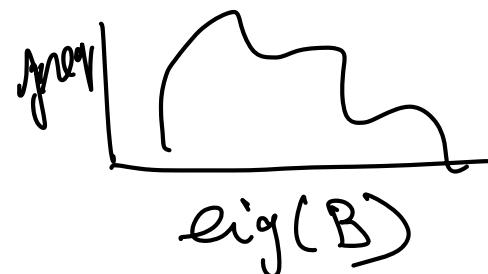
→ can reason about spectrum  
based on  $\text{eig}(A) \times \text{eig}(B)$ .

## Free Probability Result:

As size of matrix  $\rightarrow \infty$ , we have

$$\text{eigenvals}(A + \underline{Q^T B Q}) = \text{Unif}(\text{eig}(A)) \boxplus \text{Unif}(\text{eig}(B))$$

free additive convolution

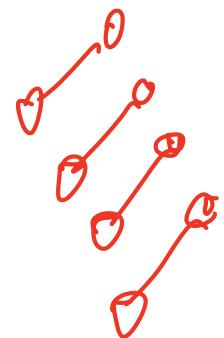


# Why Care About $A + Q^T B Q$ ?

- Let  $P$  be a random permutation matrix.  
Let's look at  $A + P^T B P$
- Let  $M$  be a perfect matching matrix:

$$A_G := P_1^T M P_1 + P_2^T M P_2 + P_3^T M P_3$$

this creates a 3-regular graph  
(maybe multi-edges, but prob. decreases as size of graph increases)



[MSS] : You can bound any one eigenvalue of  $A_G$  "nicely".

$$[\text{MSS}] \rightarrow \mathbb{E}_Q \chi_x(A + Q^T B Q) = \mathbb{E}_P \chi_x(A + P^T B P)$$

[MSS] : Find "good"  $P$  s.t.

$$\text{k-th root}(\mathbb{E} \chi_x(A + Q^T B Q)) \leq \text{k-th root}(\chi_x(A + P^T B P)) .$$
$$\max(\chi_x(A + P^T B P)) \leq \max(\mathbb{E} \chi_x(A + Q^T B Q))$$

## Free Probability Basics

if one more prob. theory

- Freeness  $\approx$  Independence
- Central limit theorem
- Some analytic tools

(We work with  
moments of  
distributions)

## Classical Prob. Thy.

Classical independence gives a way of determining mixed moments by marginals.

$$\text{eg: } \mathbb{E}[abab] = \mathbb{E}[a^2]\mathbb{E}[b^2] \quad a \bowtie b \text{ are indep.}$$

# Freeness

Def: [ Non-commutative Prob space (ncps) ]

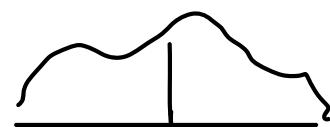
( $\mathcal{A}$ ,  $\psi$ )

- $\mathcal{A}$ : unital  $\mathbb{C}$ -algebra  
(ring  $\times \mathbb{C}$ -vector space)  $\underline{1 \in \mathcal{A}}$   $\underline{\psi(1)=1}$
- $\underline{\psi: \mathcal{A} \rightarrow \mathbb{C}}$
- $\psi$  is tracial :  $\psi(ab) = \psi(ba)$
- Elements  $a \in \mathcal{A}$  are called (nc) random var.
- "Think of  $\psi(a) = \mathbb{E}[a]$ "

Eg:

$\mathcal{A}$ :  $n \times n$  matrices

$\psi$ : normalized trace



$$\underline{\psi(a)} = \frac{1}{n} \sum_{i=1}^n a_{ii} = \underbrace{\frac{1}{n} \sum_{i=1}^n}_{\lambda_i \text{ vs spectral}} \lambda_i(a) \xrightarrow{\mathbb{E}[\lambda]} \psi(a).$$

$$\psi(I) = 1$$

NCPS: Also equipped with \* operation

$$\begin{aligned} * : A &\rightarrow A \\ a &\mapsto a^* \end{aligned} \quad \text{s.t.}$$

$$\left( \begin{array}{l} \cdot (a^*)^* = a \\ \cdot (ab)^* = b^* a^* \\ \cdot (a+b)^* = a^* + b^* \\ \cdot \psi(a^* a) \geq 0 \quad \text{with equality iff } a=0. \end{array} \right)$$

\* With this, we can define a to be :

- Self adjoint :  $a = a^*$
- Unitary :  $a^* a = a a^* = I$
- Normal :  $a a^* = a^* a$

eg: We can prove we get inner prod  $\therefore$  prove things like C-S ineq  
 $|\psi(b^* a)|^2 = \psi(a^* a) \psi(b^* b)$ .

- Recall:
- $\Phi$  encodes moments.
  - Algebra encodes relation b/w r.v.s

Analytic Dist:

Let  $(A, \Phi)$  be a ncps,  $a \in A$  normal. If  $\exists$  cpt. supp dist  $\mu$  on  $\mathbb{C}$  s.t.

$$\int z^k \bar{z}^l d\mu(z) = \Phi(a^k (a^*)^l) \quad \forall k, l \in \mathbb{N},$$

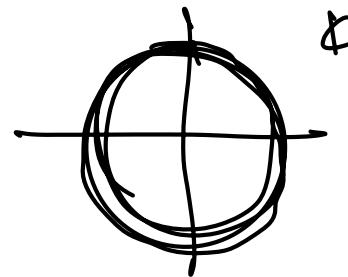
then  $\mu$  is uniquely determined & is called analytic dist. of  $a$ .

Eg : Def: Haar Unitary :

Let  $(A, \psi)$  be a ncfs.  $\forall k \in \mathbb{Z}$  is Haar unitary if it is unitary and

$$\psi(u^k) = \begin{cases} 1 & \text{if } k=0 \\ 0 & \text{o.w.} \end{cases}$$

The corresponding analytic dist. is the unif dist over unit circle  $\mathbb{C}$ . (verify)



Def: Let  $(A, \varphi)$  be a ncfs.  $a, b \in A$  are free if all centered alternating mixed moments vanish.

$$\text{Eq: } \varphi(\bar{a} \bar{b}) = 0 \quad \underbrace{a - \varphi(a)}_{\sim} = \bar{a}$$

$$\varphi(\underbrace{\bar{a}^2}_{\sim} \underbrace{\bar{b}^*}_{\sim} \underbrace{\bar{a}}_{\sim} \underbrace{\bar{b}^3}_{\sim}) = 0$$

Eg: 1

$$0 = \varphi(\underbrace{(a - \varphi(a))}_{\sim} \underbrace{(b - \varphi(b))}_{\sim})$$

$$= \varphi(\underbrace{ab}_{\sim} - \underbrace{\varphi(a)b}_{\sim} - \underbrace{a\varphi(b)}_{\sim} + \underbrace{\varphi(a)\varphi(b)}_{\sim})$$

$$= \varphi(ab) - \varphi(a)\varphi(b) - \varphi(a)\varphi(b) + \varphi(a)\varphi(b)$$

$$\Rightarrow \varphi(ab) = \varphi(a)\varphi(b) \quad \Rightarrow \text{like in classical case.}$$

$$2: \underline{\psi(\bar{a}\bar{b}\bar{a}\bar{b}) = 0} \Rightarrow$$

$$\underline{\psi((a-\psi(a))(b-\psi(b))(a-\psi(a))(b-\psi(b))) \dots}}$$

We get:

$$\underline{\psi(abab) = \psi(a^2)\psi(b^2) + \psi(a)^2\psi(b^2) - \psi(a)^2\psi(b)^2}.$$

In classical indef:

$$\underline{\psi(abab) = \psi(a^2)\psi(b^2)}.$$

Why haven't we seen freeness:

- Commuting r.v. are free only if one of them is a constant.
- Freeness is an  $n$ -dim. phenomenon.

(Random matrices are "asymptotically" free.)

Eg: Let  $(\mathcal{A}, \gamma)$  be ncfs. let  $a, b \in \mathcal{A}$  s.t.  
 $u$  is Haar unitary that is free from  $\{a, b\}$   
(think as random rotation)

Then  $a$  and  $u^*bu$  are free.

# Central Limit Theorem (both for independence & free r.v.)

CLT asks about:

$$\lim_{N \rightarrow \infty} \frac{a_1 + \dots + a_N}{\sqrt{N}}$$

$\varphi(a_i) = 0$   
 $\varphi(a_i^2) = 1$

all  $a_i$  have  
same distribution

where convergence is in terms of moments (weak avg.):

$$a_N \rightarrow a \text{ means } \lim_m \underbrace{\lim_{N \rightarrow \infty} P_N(a_N^m)}_{= \dots} = \varphi(a^m).$$

- We want to focus on the moments.

Fix  $m, N \geq 1$ ,  $m$  is finite,  $N$  will go to infinity later

Look at:

$$\cdot \underbrace{\varphi((a_1 + \dots + a_N)^m)}_{= \dots}$$

Analyzing  $\Psi((a_1 + \dots + a_N)^m)$

$$\cdot \Psi(\underbrace{(a_1 + \dots + a_N)^m}_{\text{m times}}) = \sum_{\gamma: [m] \rightarrow [N]} \Psi(\underbrace{a_{\gamma(1)} \dots a_{\gamma(m)}}_{\text{m times}})$$

$(a_1+a_2)(a_1+a_2)(a_1+a_2)$

But note that:

$$\Psi(a_1 \underbrace{a_2 a_2}_{\text{circled}}, \underbrace{a_3}_{\text{circled}}, \underbrace{a_1 a_2}_{\text{circled}}) = \Psi(a_1 \underbrace{a_2 a_1}_{\text{circled}}, \underbrace{a_3}_{\text{circled}}, \underbrace{a_1 a_2}_{\text{circled}})$$

We encode this information as follows:

Let  $\pi = \{\nu_1, \dots, \nu_m\}$  be a partition of  $[m]$

where

$$\forall i, j \quad \nu(i) = \nu(j) \iff \exists l \quad i, j \in \nu_l$$

Eg:

$$\begin{matrix} 1 & 2 & 3 & 4 \\ a_1 & a_2 & a_2 & \underline{a_3} \end{matrix} \quad \begin{matrix} 5 & 6 \\ a_1 & a_2 \end{matrix}$$

$$\begin{matrix} a_4 & a_1 & a_1 & a_5 & a_9 & a_1 \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \end{matrix}$$

corresponds  
to

$$\left\{ \begin{matrix} \downarrow & \downarrow & \downarrow \\ \{1, 5\}, \{2, 3, 6\}, \{4\} \end{matrix} \right\}$$

$$\frac{N}{|\pi|} \pi!$$

$$\cdot \Psi((a_1 + \dots + a_N)^m) = \sum_{\pi: [m] \rightarrow [N]} \Phi(a_{\pi(1)}, \dots, a_{\pi(m)})$$

with this, we get:

$$\Psi((a_1 + \dots + a_N)^m) = \sum_{\pi: \text{partition of } [m]} K_\pi \underbrace{A_\pi^{(N)}}_{\substack{\text{common value of} \\ \Phi(a_{\pi(1)}, \dots, a_{\pi(m)}) \forall \pi \\ \text{associated with } \pi}} \xrightarrow{\substack{\text{number of } \pi \text{'s associated} \\ \text{with } \pi}} \text{only } K_\pi \text{ depends on } N$$

Obv:  $\pi$  containing singletons does not contribute:

$$K_\pi = \Psi(\sim a \sim) = \underbrace{\Psi(a)}_{=0} \Psi(\sim \sim)$$

(both in  
classical &  
free prob.)

$\therefore$  We can restrict to  $\pi$ 's with  $|\pi| \leq m/2$ .

→ reason why  
 $\sqrt{N}$  appears

Calculating  $A_{\pi}^{(N)}$ :

$A_{\pi}^{(N)}$ : how many  $\tau$ 's are associated with a  $\pi$ .

$$= N(N-1)\dots(N-|\pi|+1)$$

because: let  $\pi = \{\{abc\}, \{def\}, \{f\}\}$

partition corresponds to an idx. lets say j appearing at pos. a, b, & c.

$$\xrightarrow{N \rightarrow \infty} N^{|\pi|}$$

$$\Rightarrow \lim_{N \rightarrow \infty} \varphi \left( \left( \frac{a_1 + \dots + a_N}{\sqrt{N}} \right)^m \right) = \sum_{\pi} K_{\pi} \frac{A_{\pi}^{(N)}}{N^{m/2}}$$

$$= \sum_{\pi} K_{\pi} \lim_{N \rightarrow \infty} \frac{A_{\pi}^{(N)}}{N^{m/2}}$$

$$= \sum_{\pi} K_{\pi} \lim_{N \rightarrow \infty} N^{|\pi|-m/2} \quad . \quad \begin{pmatrix} \text{If } |\pi| < \frac{m}{2} \\ \text{this goes to 0} \end{pmatrix}$$

$$\lim_{N \rightarrow \infty} \Psi\left(\left(\frac{a_1 + \dots + a_N}{\sqrt{N}}\right)^m\right) = \sum_{\pi} k_{\pi} \lim_{N \rightarrow \infty} N^{|\pi| - m \frac{r_2}{2}} \quad - (\star)$$

- We had  $|\pi| \leq m \frac{r_2}{2}$  because singletons correspond to 0.
- If  $|\pi| < m \frac{r_2}{2}$ , then  $(\star)$  goes to 0.
- Only pairings survive.

$$\lim_{N \rightarrow \infty} \Psi\left(\left(\frac{a_1 + \dots + a_N}{\sqrt{N}}\right)^m\right) = \sum_{\substack{\pi: \text{pairings} \\ g \in [m]}} k_{\pi}.$$

## Classical CLT

$$\# \text{ pairings} \quad K_n = \varphi(a_1^2) \cdot \varphi(a_2^2) \cdots = 1$$

$$\# \text{ pairings of } [m] := (m-1) \cdot (m-3) \cdots 1.$$

- pick any elem.
- $(m-1)$  choices for its partner
- choose smallest elem in remaining  $\rightarrow$  pair with any  $(m-3)$  of them.

Then,

$$\frac{1}{\sqrt{2\pi}} \int t^m e^{-t^2/2} dt = \begin{cases} (m-1)(m-3) \cdots & m \text{ even} \\ 0 & m \text{ odd} \end{cases}$$



# FREE CLT

$$\lim_{N \rightarrow \infty} \Psi\left(\left(\frac{a_1 + \dots + a_N}{\sqrt{N}}\right)^m\right) = \sum_{\substack{\pi: \text{ pairings} \\ g \in [m]}} k_\pi.$$

Play with freeness:

$$\cdot \pi = \{\{1, 2\}, \{3, 4\}\}$$



$$\Psi(aabb) \\ = \Psi(a^2 b^2)$$

$$\Psi(ab) = \Psi(a)\Psi(b)$$

$$\Psi(a^2 b^2) = \underbrace{\Psi(a^2)}_{\text{ }} \underbrace{\Psi(b^2)}_{\text{ }}$$

$$\cdot \pi = \{\{1, 3\}, \{2, 4\}\}$$



$$\Psi(abab) = 0$$

$$\cdot \pi = \{\{1, 4\}, \{2, 3\}\}$$



$$\Psi(a^b b^a) = \Psi(a^2) \Psi(b^2)$$

$$\begin{aligned} \Psi(a^b b^a) &= \Psi(a^2 b^2) \\ &= \Psi(a^2) \Psi(b^2) \end{aligned}$$

- Def: A pairing  $\pi$  is non crossing if  $\frac{\{p_1, p_2\}}{\{q_1, q_2\}} \in \pi$

$p_1 < q_1 < p_2 < q_2 \rightarrow$  doesn't happen

Fact: # of such pairings of  $[2m]$  is  $C_m$ : the  $m$ -th Catalan number

Fact: The analytic dist. whose (even) moments are Catalan numbers is the semicircular dist.

$$\frac{1}{2\pi} \sqrt{4-x^2}$$

