

1 Expanding a power-sum MacMahon symmetric function in the monomial basis

Fix a positive integer r . Throughout, we will work with MacMahon symmetric functions in r alphabets. The complete set of variables is

$$\{x_{i,j} \mid i \in \mathbb{P}, j \in [r]\}.$$

Fix vector partitions $\lambda = (\lambda_1, \dots, \lambda_\ell)$ and $\mu = (\mu_1, \dots, \mu_m)$ of the same vector, where

$$\begin{aligned}\lambda_i &= (\lambda_{i,1}, \dots, \lambda_{i,r}) \in \mathbb{P}^r \setminus \{\vec{0}\}, \\ \mu_i &= (\mu_{i,1}, \dots, \mu_{i,r}) \in \mathbb{P}^r \setminus \{\vec{0}\}.\end{aligned}$$

The monomial MacMahon symmetric function indexed by μ is the sum of all monomials in the diagonal \mathfrak{S}_n -orbit of the monomial

$$x_\mu = \prod_{i=1}^m \prod_{j=1}^r x_{i,j}^{\mu_{i,j}}.$$

The power-sum MacMahon symmetric function indexed by λ is

$$\begin{aligned}P_\lambda &= \prod_{i=1}^\ell P_{\lambda_i} \\ &= \prod_{i=1}^\ell \sum_{k \in \mathbb{P}} x_{k,1}^{\lambda_{i,1}} \cdots x_{k,r}^{\lambda_{i,r}} \\ &= \sum_{(k_1, \dots, k_\ell) \in \mathbb{P}^\ell} \prod_{i=1}^\ell x_{k_i,1}^{\lambda_{i,1}} \cdots x_{k_i,r}^{\lambda_{i,r}}.\end{aligned}$$

Thus the coefficient of x_μ in P_λ is

$$\# \left\{ (k_1, \dots, k_\ell) \in \mathbb{P}^\ell \mid \sum_{i: k_i=h} \lambda_i = \mu_h \quad \forall h \in [m] \right\}.$$

Note that we can replace \mathbb{P}^ℓ by m^ℓ .

For example, fix $r = 2$. Let $\mu = (1, 1)(1, 2)$, so that

$$x_\mu = x_{1,1}^1 x_{1,2}^1 x_{2,1}^1 x_{2,2}^2$$

and let $\lambda = (1, 1)(1, 0)(0, 1)(0, 1)$ so that $\ell = 4$ and

$$\begin{aligned}P_\lambda &= \sum_{(k_1, \dots, k_4)} \prod_{i=1}^4 x_{k_i,1}^{\lambda_{i,1}} x_{k_i,2}^{\lambda_{i,2}} \\ &= \sum_{(k_1, \dots, k_4)} x_{k_1,1}^{\lambda_{1,1}} x_{k_1,2}^{\lambda_{1,2}} x_{k_2,2}^{\lambda_{2,1}} x_{k_2,2}^{\lambda_{2,2}} x_{k_3,1}^{\lambda_{3,1}} x_{k_3,2}^{\lambda_{3,2}} x_{k_4,1}^{\lambda_{4,1}} x_{k_4,2}^{\lambda_{4,2}} \\ &= \sum_{(k_1, \dots, k_4)} x_{k_1,1}^1 x_{k_1,2}^1 x_{k_2,2}^1 x_{k_2,2}^0 x_{k_3,1}^0 x_{k_3,2}^1 x_{k_4,1}^0 x_{k_4,2}^1\end{aligned}$$

so in order to make this monomial equal to x_μ , we need

$$\sum_{i: k_i=1} \lambda_i = \mu_1 = (1, 1), \quad \sum_{i: k_i=2} \lambda_i = \mu_2 = (1, 2).$$

The solutions for (k_1, k_2, k_3, k_4) and the corresponding expressions for x_μ are

$$\begin{array}{ll} (1, 2, 2, 2) & x_{1,1}^1 x_{1,2}^1 x_{2,1}^1 x_{2,2}^2 = (x_{1,1}^1 x_{1,2}^1) (x_{2,2}^1 x_{2,2}^0) (x_{2,1}^0 x_{2,2}^1) (x_{2,1}^0 x_{2,2}^1) \\ (2, 1, 1, 2) & x_{1,1}^1 x_{1,2}^1 x_{2,1}^1 x_{2,2}^2 = (x_{2,1}^1 x_{2,2}^1) (x_{1,2}^1 x_{1,2}^0) (x_{1,1}^0 x_{1,2}^1) (x_{2,1}^0 x_{2,2}^1) \\ (2, 1, 2, 1) & x_{1,1}^1 x_{1,2}^1 x_{2,1}^1 x_{2,2}^2 = (x_{2,1}^1 x_{2,2}^1) (x_{1,2}^1 x_{1,2}^0) (x_{2,1}^0 x_{2,2}^1) (x_{1,1}^0 x_{1,2}^1) \end{array}$$

So the coefficient is 3. This could certainly be calculated algorithmically.

2 General linear algebra

Let V be a vector space of finite dimension n over ground field \mathbb{k} (it doesn't matter what it is). Suppose that we have an inner product $\langle u, v \rangle$ on V such that $B = \{v_1, \dots, v_n\}$ and $B^* = \{v_1^*, \dots, v_n^*\}$ are dual bases, i.e., $\langle v_i, v_j^* \rangle = \delta_{ij}$ for all i, j (this is the Kronecker delta). Then for every $u \in V$, we can expand u with respect to basis B as

$$u = \sum_{i=1}^n \langle v_i^*, u \rangle v_i \tag{1}$$

because, for each $j \in [n]$,

$$\left\langle v_j^*, \sum_{i=1}^n \langle v_i^*, u \rangle v_i \right\rangle = \sum_{i=1}^n \langle v_i^*, u \rangle \langle v_j^*, v_i \rangle = \sum_{i=1}^n \langle v_i^*, u \rangle \delta_{ij} = \langle v_j^*, u \rangle$$

which says that every element of B^* has the same inner product with the left and right-hand sides of (1), which is necessary and sufficient for equality.

3 Transitions between bases, from Rosas 2001

Notation:

- \mathbb{N} = nonnegative integers;
- For $u \in \mathbb{N}^k$, a **vector partition** $\lambda \vdash u$ is an unordered sequence of vectors (**parts**) summing to u ; zero vectors can be ignored.
- Weight of a vector partition = sum of all entries in all vectors in it (ex.: $\text{wt}(\{120, 100, 100, 013\}) = 9$)
- \mathfrak{M}_u = MacMahon symmetric functions of homogeneous multidegree u
- $\text{sign}(\lambda) = (-1)^{\text{number of parts with even sum}}$
- $m_\lambda, e_\lambda, h_\lambda, p_\lambda, f_\lambda$: monomial, elementary, homogeneous, power-sum, and forgotten MacMahon symmetric functions (Rosas 327–328)
- For a vector partition λ in which each part v_i occurs with multiplicity m_i , define

$$|\lambda| = \prod_i m_i! \quad \text{and} \quad \lambda! = \prod_i \prod_{x \in v_i} (x!)^{m_i}.$$

```

def bars(mu):
    Mu = list(mu)
    Done = []
    output = 1
    for part in Mu:
        if not part in Done:
            Done.append(part)
            output = output * factorial(Mu.count(part))
    return output

facto = lambda mu: prod(factorial(x) for part in mu for x in part)

```

A vector partition is **unitary** if it is a partition of the all-ones vector $\mathbf{1} \in \mathbb{N}^k$. There is an obvious bijection between unitary vector partitions and set partitions of $[k]$. There is a whole theory of MacMahon symmetric functions of unitary partitions, due to Doubilet. These are the graded pieces $\mathfrak{M}_{(1)^k}$.

```

## convert a set partition to a vector partition
## e.g.,
## sage: pi = SetPartition([[1,5,7],[2,3],[4,8],[6]])
## sage: SetPartitionToVectorPartition(pi)
## [[0,0,0,0,0,1,0,0],[0,0,0,1,0,0,0,1],[0,1,1,0,0,0,0,0],[1,0,0,0,1,0,1,0]]
def SetPToVecP(pi):
    n = pi.base_set_cardinality()
    return VectorPartition([[1 if i in B else 0 for i in range(1,n+1)] for B in pi])

```

The **type** of a set partition $\pi = B_1 | \cdots | B_\ell$, with respect to a vector $u = (u_1, \dots, u_r)$ of weight n , is the vector partition $\text{type}_u(\pi) = \lambda = \lambda_1 \cdots \lambda_\ell$ where λ_k is the vector in \mathbb{N}^r whose i th coordinate is

$$\#\{j \in B_k \mid u_1 + \cdots + u_{i-1} < j \leq i + \cdots + u_i\}.$$

the number of elements of B_k such that in the i th equivalence class.

```

## Input: u is an integer vector of weight n; pi is a set partition of [n]
## Output: type_u(pi), as in Rojas p.328
def Type(u,pi):
    r = len(u)
    psum = [sum(u[:k]) for k in range(r+1)] ## partial sums
    ell = len(pi)
    n = sum(len(B) for B in pi)
    lam = [[len([j for j in B if psum[i]<j<=psum[i+1]]) for i in range(r)] for B in pi]
    return VectorPartition(lam)

```

For a vector u of weight n and a vector partition λ of weight n , the number of set partitions π of type $\text{type}_u(\pi)$ is

$$\binom{u}{\lambda} := \frac{u!}{\lambda! |\lambda|}.$$

```

Choose = lambda u,mu: facto([u])/facto(mu)/bars(mu)

```

Define a scalar product on \mathfrak{M} :

$$\langle h_\lambda, m_\mu \rangle = \delta_{\lambda\mu}.$$

The idea is to use this for basis transitions, along the lines discussed earlier. Specifically, formula (1) says that for any $F \in \mathfrak{M}_u$, we have

$$F = \sum_{\lambda \vdash u} \langle h_\lambda, F \rangle m_\lambda \quad \text{and} \quad F = \sum_{\lambda \vdash u} \langle m_\lambda, F \rangle h_\lambda. \quad (2)$$

Thus we have to know how to compute inner products.

There is a “lifting map” $\hat{\rho}$ sending MacMahon functions to unitary (Doubilet) symmetric functions. Specifically:

$$\binom{u}{\lambda} |\lambda| m_\lambda \xrightarrow{\rho} \sum_{\pi: \text{type}_u(\pi)=\lambda} m_\pi \quad (3a)$$

$$\binom{u}{\lambda} \lambda! e_\lambda \xrightarrow{\rho} \sum_{\pi: \text{type}_u(\pi)=\lambda} e_\pi \quad (3b)$$

$$\binom{u}{\lambda} \lambda! h_\lambda \xrightarrow{\rho} \sum_{\pi: \text{type}_u(\pi)=\lambda} h_\pi \quad (3c)$$

$$\binom{u}{\lambda} p_\lambda \xrightarrow{\rho} \sum_{\pi: \text{type}_u(\pi)=\lambda} p_\pi \quad (3d)$$

$$\binom{u}{\lambda} |\lambda| f_\lambda \xrightarrow{\rho} \sum_{\pi: \text{type}_u(\pi)=\lambda} f_\pi \quad (3e)$$

```

m = MacMahonSymmetricFunctions(QQ).Monomial()
p = MacMahonSymmetricFunctions(QQ).Powersum()
h = MacMahonSymmetricFunctions(QQ).Homogeneous()
e = MacMahonSymmetricFunctions(QQ).Elementary()
f = MacMahonSymmetricFunctions(QQ).Forgotten() # not yet implemented
def LiftM(mu):
    u = sum(mu)
    return 1/Choose(u,mu)/Bars(mu)*sum(m[SetPToVecP(pi)] \
        for pi in SetPartitions(sum(u)) if Type(u,pi)==mu)
def LiftP(mu):
    u = sum(mu)
    return 1/Choose(u,mu)*sum(p[SetPToVecP(pi)] \
        for pi in SetPartitions(sum(u)) if Type(u,pi)==mu)
def LiftH(mu):
    u = sum(mu)
    return 1/Choose(u,mu)/Facto(mu)*sum(h[SetPToVecP(pi)] \
        for pi in SetPartitions(sum(u)) if Type(u,pi)==mu)
def LiftE(mu):
    u = sum(mu)
    return 1/Choose(u,mu)/Facto(mu)*sum(e[SetPToVecP(pi)] \
        for pi in SetPartitions(sum(u)) if Type(u,pi)==mu)
def LiftF(mu):
    u = sum(mu)

```

```

return 1/Choose(u,mu)/Bars(mu)*sum(f[SetPToVecP(pi)] \
for pi in SetPartitions(sum(u)) if Type(u,pi)==mu)

```

Moreover, for all $f, g \in \mathfrak{M}_u$ we have

$$\langle f, g \rangle = u! \langle \hat{\rho}(f), \hat{\rho}(g) \rangle \quad (4)$$

[Rosas, Prop. 7]. For instance, Doubilet's Formula #1 says that for set partitions π, σ of $[n]$ we have

$$\langle m_\pi, h_\sigma \rangle = n! \delta_{\pi\sigma} \quad (5)$$

Therefore, if λ and μ are vector partitions of u of weight n , then

$$\begin{aligned}
\langle m_\lambda, h_\mu \rangle &= u! \langle \hat{\rho}(m_\lambda), \hat{\rho}(h_\mu) \rangle && \text{(by (4))} \\
&= u! \frac{1}{\binom{u}{\lambda} |\lambda| \binom{u}{\mu} \mu!} \left\langle \sum_{\pi: \text{type}_u(\pi) = \lambda} m_\pi, \sum_{\sigma: \text{type}_u(\sigma) = \mu} h_\sigma \right\rangle && \text{(by (3a) and (3c))} \\
&= \delta_{\lambda\mu} \frac{u! n! \#\{\pi: \text{type}_u(\pi) = \lambda\}}{\binom{u}{\lambda}^2 |\lambda| \lambda!} && \text{(by linearity and (5))}
\end{aligned}$$

Since Doubilet (Appendix 2) calculated the scalar products for all five families of unitary symmetric functions, we can lift them to MacMahon functions. (We will need the lattice operations and Möbius and zeta functions in the set partition lattice Π_n

So the code could look something like this:

```

## calculate the inner product of the unitary symmetric functions m[pi] and h[sigma]
## using Doubilet formula #1: <m[pi], h[sigma]> = n! delta(pi, sigma)
## although, note, the below code just works on basis elements --- we need to extend it lin

MHPProduct = lambda pi,sigma: 0 if pi != sigma else factorial(pi.base_set_cardinality())
HMPProduct = HMPProduct

## calculate the inner product of the unitary symmetric functions p[pi] and p[sigma]
## using Doubilet formula #2: <p[pi], p[sigma]> = n! delta(pi, sigma)
## although, note, the below code just works on basis elements --- we need to extend it lin

def PPPProduct(pi,sigma):
    if pi != sigma:
        return 0
    else:
        n = pi.base_set_cardinality()
        P = posets.SetPartitions(n)
        mu = P.moebius_function
        return factorial(n) / mu(P.bottom(),pi)

## calculate the inner product of the unitary symmetric functions m[pi] and p[sigma]
## using Doubilet formula #10
## although, note, the below code just works on basis elements --- we need to extend it lin

```

```

def MPPProduct(pi, sigma):
    if not sigma in pi.coarsenings(): ## i.e., if  $\zeta(\pi, \sigma) = 0$ 
        return 0
    else:
        n = pi.base_set_cardinality()
        mu = posets.SetPartitions(n).moebius_function
        return factorial(n) * mu(pi, sigma) / abs(mu(P.bottom(), sigma))
MPPProduct = lambda sigma, pi: MPPProduct(pi, sigma)

## calculate the inner product of the unitary symmetric functions m[pi] and m[sigma]
## using Doubilet formula #7
## although, note, the below code just works on basis elements --- we need to extend it later

def MMPProduct(pi, sigma):
    n = pi.base_set_cardinality()
    if not n == sigma.base_set_cardinality():
        return 0
    else:
        P = posets.SetPartitions(n)
        mu = posets.SetPartitions(n).moebius_function
        return factorial(n) * sum( mu(pi, tau) * mu(sigma, tau) / abs(mu(P.bottom(), tau)) \
            for tau in pi.coarsenings() if tau in sigma.coarsenings())

## expand m[mu] in the h-basis
## this code will not work until
def MtoH(mu):
    sum( facto(u) * MHPProduct(LiftM(mu), LiftH(nu)) * h[nu] for nu in VectorPartitions(sum(nu)))

```