0.1 Estimating the Spatial Closure from the HO Solution

0.2 MC Estimation of the Spatial Closure

This sections describes an alternative spatial closure to the LO equations based on a parametric estimate from the HO solution. After algebraically manipulating the time-discretized moment equations to form angular consistency parameters, there is still more unknowns than equations, for each cell, due to the . A half-range balance equation for $\mu > 0$ is formed by adding the exact L and R moment equations given by Eq. (??) and (??), i.e.,

$$\overline{\mu}_{i+1/2}^+ \phi_{i+1/2}^+ - \overline{\mu}_{i-1/2}^+ \phi_{i-1/2}^+ + \sigma_{a,i} h_i \phi_i^+ = \frac{h_i}{2} q_i, \tag{1}$$

where q is a general, isotropic source. The angular consistency terms can be estimated with the HO solution and the inflow term $\phi_{i-1/2}^+$ is upwinded from the previous cell. An additional equation is needed to eliminate $\phi_{i+1/2}^+$ to produce an equation for a single unknown ϕ_i^+ . The outflow can be eliminated as a parametric combination of the other variables, i.e.,

$$\phi_{i+1/2}^+ = f(\gamma_i^{HO}, \phi_i^+, \phi_x^+, \phi_{i-1/2}^+), \tag{2}$$

where γ_i^{HO} is some constant estimated with the HO solution and f is some function of some number of input variables. If the problem is linear, i.e., q is known a priori, then application of this closure will ensure that the HO and LO equations exactly preserve the same moments. As the problem is typically non-linear (e.g., scattering or thermal emission are included in q), then this will not be strictly true.

We will explore two different closure options: a scaled slope, i.e.,

$$\phi_{i\pm 1/2}^{\pm} = \phi_i^+ \pm \gamma_i \phi_x^+ \tag{3}$$

and a scaled average

$$\phi_{i\pm 1/2}^{\pm} = \gamma_i \phi_i^+ \pm \phi_x^+, \tag{4}$$

where a value of $\gamma_i = 1$ produces the standard linear discontinuous expressions for the extrapolated outflows. The closures are only needed to eliminate the extrapolated face values, not the inflow values for the particular direction equation.

We now use the HO solution to estimate γ_i . For example,

$$\gamma_i^{+,HO} = \frac{\phi_{i+1/2}^+ - \phi_x^+}{\phi_i^+} \tag{5}$$

in the scaled slope case. For this closure, as the slope goes to zero this expression becomes undefined. In cells where the slope is $O(10^{-13}\psi_i)$, we use $\gamma_i = 1$. No problems have been observed with the fact that at relatively modest slopes γ becomes very large because the solution is changing minimally in such sections. The main reason for using this closure is it allows for values of γ that are equivalent to step and lumping.

The expression for the outflow face term is substituted in each equation, using the γ estimated from a HO solution. For instance, the positive balance equation becomes

$$\overline{\mu}_{i+1/2}^{+} \left(\gamma_i^{+,HO} \phi_i^{+} + \phi_x^{+} \right) - \overline{\mu}_{i-1/2}^{+} \phi_{i+1/2}^{+} + \frac{\sigma_{a,i} h_i}{2} \phi_i^{+} = \frac{h_i}{2} q_i, \tag{6}$$

noting that ϕ_i^+ and ϕ_x^+ remain as unknowns. The MC solution can provides the face values.

Our solution is in terms of the L and R moments, however.

0.2.1 LO Equations Trial Space

The linear shape on the interior of the cell is still used to represent the temperature unknowns, scattering source, and emission source. The choice of this trial space should allow for higher consistency between the HO and LO equations by using MC to compute outflows from space-angle cells, rather than an LD extrapolation of the solution. For the HO solution, a discontinuous outflow for a space-angle cell should produce greater angular accuracy on faces. In optically thick cells, this will mitigate observed issues with the spatial slope being artificially high to account for the discrepancy in angular shape between the internal and face solutions. The outflow will be estimated using a face-based tally of the MC solution. The error equation can be modified to estimate the outflow for a given batch. For the LO equations, a parametric relation between the outflow and internal moments, as estimated by the latest HO solution, will be used to eliminate the extra unknowns. This will produce a consistent spatial closure between the HO and LO solutions upon convergence. The closure relation will be averaged over each half-range, similar to the angular consistency parameters. Poor statistics for the face tallies may result in this trial space producing less accurate results compared to the standard LDFE solution, at least for sufficiently fine meshes where LD can accurately represent the solution. Although the closure will be applied everywhere, we expect the greatest improvement in accuracy for cells where the LDFE trial space produces a negative solution.

By introducing a LDD trial space into the LO equations, we introduce an additional unknown in each half-range equation. This extra unknown is eliminated as a function of the interior moments. However, in the HOLO context, the HO solution used a lagged source q_i .

In theory, if the problem were linear, or the nonlinear problem was fully converged, then the HO and LO solutions would produce exactly the same moments. There are several issues with ECMC that cause this to not be true, even for a linear problem. With ECMC, global energy balance is not preserved, but local energy balance is also not preserved. There are source biasing techniques for standard MC (systematic sampling) that would exactly preserve the zeroth moment of the source [21]). It is a requirement that the HO solution satisfies the zeroth moment equation. If the closure relation also uses the first moment, then it must also satisfy the first spatial moment equation. However, because we have to reconstruct the bilinear moment of x and μ , the consistency terms do not exactly preserve the first moment equation. One final reason is that the analog treatment of absorption (at low weights as discussed in Sec. ??) results in the fact that the product of ϕ_i (as estimated via path-length estimators) and the removal cross section flux does not exactly equal the absorption density of weight within that cell, due to statistical noise.

REFERENCES

- [1] S.R. Bolding and J.E. Morel. A High-Order Low-Order Algorithm with Exponentially-Convergent Monte Carlo for k-Eigenvalue problems. ANS Winter Meeting. Anaheim, CA, 2014.
- [2] Subrahmanyan Chandrasekhar. Radiative Transfer. Courier Corporation, 2013.
- [3] Jeffery D. Densmore. Asymptotic analysis of the spatial discretization of radiation absorption and re-emission in implicit monte carlo. *Journal of Computational Physics*, 230(4):1116 1133, 2011.
- [4] Jeffery D. Densmore and Edward W. Larsen. Asymptotic equilibrium diffusion analysis of time-dependent monte carlo methods for grey radiative transfer. *Journal of Computational Physics*, 199(1):175 – 204, 2004.
- [5] Jeffery D Densmore, Kelly G Thompson, and Todd J Urbatsch. A hybrid transport-diffusion monte carlo method for frequency-dependent radiative-transfer simulations. *Journal of Computational Physics*, 231(20):6924–6934, 2012.
- [6] Paul Nelson Edward W. Larsen. Finite-difference approximations and superconvergence for the discrete-ordinate equations in slab geometry. SIAM Journal on Numerical Analysis, 19(2):334–348, 1982.
- [7] Edward W. Larsen and Jim E. Morel. Advances in discrete-ordinates methodology. *Nuclear Computational Science. Springer Netherlands*, pages 1–84, 2010.
- [8] J. A. Fleck, Jr. and J. D. Cummings, Jr. An implicit monte carlo scheme for calculating time and frequency dependent nonlinear radiation transport. J. Comput. Phys., 8(3):313–342, December 1971.

- [9] N. A. Gentile and Ben C. Yee. Iterative implicit monte carlo. *Journal of Computational and Theoretical Transport*, 0(0):1–31, 0.
- [10] NA Gentile. Implicit monte carlo diffusion: An acceleration method for monte carlo time-dependent radiative transfer simulations. *Journal of Computational Physics*, 172(2):543–571, 2001.
- [11] Eugene D. Brooks III, Michael Scott McKinley, Frank Daffin, and Abraham Szke. Symbolic implicit monte carlo radiation transport in the difference formulation: a piecewise constant discretization. *Journal of Computational Physics*, 205(2):737 – 754, 2005.
- [12] Eugene D. Brooks III, Abraham Szke, and Jayson D.L. Peterson. "piecewise linear discretization of symbolic implicit monte carlo radiation transport in the difference formulation". "Journal of Computational Physics", "220"(1):471 – 497, 2006.
- [13] Rong Kong and Jerome Spanier. A new proof of geometric convergence for general transport problems based on sequential correlated sampling methods. *Journal of Computational Physics*, 227:9762 – 9777, 2008.
- [14] Edward W. Larsen, Akansha Kumar, and Jim E. Morel. Properties of the implicitly time-differenced equations of thermal radiation transport. J. Comput. Phys., 238:82–96, April 2013.
- [15] E.W. Larsen, G.C. Pomraning, and V.C. Badham. Asymptotic analysis of radiative transfer problems. *Journal of Quantitative Spectroscopy and Radiative Transfer*, 29(4):285 – 310, 1983.
- [16] Elmer Eugene Lewis and Warren F Miller. Computational methods of neutron transport. John Wiley and Sons, Inc., New York, NY, 1984.

- [17] Michael Scott McKinley, Eugene D Brooks III, and Abraham Szoke. Comparison of implicit and symbolic implicit monte carlo line transport with frequency weight vector extension. *Journal of Computational Physics*, 189(1):330–349, 2003.
- [18] J.E. Morel, T.A. Wareing, and K. Smith. Linear-Discontinuous Spatial Differencing Scheme for S_n Radiative Transfer Calculations. Journal of Computational Physics, 128:445–462, 1996.
- [19] H. Park, J.D. Densmore, A.B. Wollaber, D.A. Knoll, and R.M. Ramenzahn. Monte Carlo Solution Methods in a Moment-Based Scale-Bridging Algorithm For Thermal Radiative Transfer Problems. M&C. Sun Valley, ID, 2013.
- [20] J.R. Peterson. Exponentially Convergent Monte Carlo for the 1-d Transport Equation. Master's thesis, Texas A&M, 2014.
- [21] J.K. Shultis and W.L. Dunn. Exploring Monte Carlo Methods. Academic Press, Burlington, MA 01803, 2012.
- [22] "Weston M. Stacey". "Nuclear Reactor Physics". Wiley, 2007.
- [23] T.A. Wareing. Asymptotic diffusion accelerated discontinuous finite element methods for transport problems. PhD thesis, Michigan, 1991.
- [24] T.A. Wareing, E.W. Larsen, and M.L. Adams. Asymptotic Diffusion Accelerated Discontinuous Finite Element Schemes for the S_N Equations in Slab and X-Y Geometries. In *International Topical Meeting on Advances in Mathematics*, Computations, Reactor Physics, volume 3, Pittsburgh, PA, 1991.
- [25] J. Willert, C.T. Kelly, D.A. Knoll, and H. Park. A Hybrid Approach to the Neutron Transport k-Eigenvalue Problem using NDA-based Algorithms. M&C. Sun Valley, ID, 2013.

- [26] Jeffrey Willert and H. Park. Residual monte carlo high-order solver for moment-based accelerated thermal radiative transfer equations. *Journal of Computational Physics*, 276:405 421, 2014.
- [27] Allan B Wollaber, Edward W Larsen, and Jeffery D Densmore. A discrete maximum principle for the implicit monte carlo equations. *Nuclear Science and Engineering*, 173(3):259–275, 2013.
- [28] E.R. Wolters. Hybrid Monte Carlo Deterministic Neutron Transport Methods
 Using Nonlinear Functionals. PhD thesis, Michigan, 2011.

APPENDIX A

FIRST APPENDIX

Text for the Appendix follows.



Figure A.1: TAMU figure

APPENDIX B

SECOND APPENDIX WITH A LONGER TITLE - MUCH LONGER IN FACT

Text for the Appendix follows.



Figure B.1: TAMU figure

B.1 Appendix Section