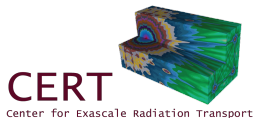


A High-Order Low-Order Algorithm with Exponentially-Convergent Monte Carlo for Thermal Radiative Transfer

Simon Bolding and Jim Morel

4 August 2016



We are interested in modeling thermal radiation transport in the high energy density physics regime

Materials under extreme conditions

Temperatures $\mathcal{O}(10^6)$ K or more

Photon radiation transports through a material

Significant **energy** may be exchanged

We want to improve efficiency of calculations

e.g., inertial confinement fusion, supernovae, et. al.

Our method has been applied to a simplified model:
the 1D grey radiative transfer equations

Energy balance equations for radiation and material.

Radiation intensity $I(x, \mu, t)$, material temperature $T(x, t)$

$$\frac{1}{c} \frac{\partial I}{\partial t} + \mu \frac{\partial I}{\partial x} + \sigma_t I(x, \mu, t) = \frac{1}{4\pi} (\sigma_a a c T^4 + \sigma_s \phi),$$
$$C_v \frac{\partial T(x, t)}{\partial t} = \sigma_a \phi(x, t) - \sigma_a a c T^4$$

Equations are **nonlinear** and may be tightly coupled

Absorption opacity (σ_a) can be a strong function of T

Implicit Monte Carlo (IMC) is the standard Monte Carlo transport method for TRT problems

The system is *linearized* over a time step $t \in [t^n, t^{n+1}]$
Opacities are evaluated with $T(t^n)$

- ▶ Produces a linear MC transport problem
with effective emission and scattering terms
- ▶ Emission source is **not** fully implicit.
Monte Carlo integration over Δt for intensity

Our high-order low-order (HOLo) method improves on several drawbacks of IMC

Standard IMC

Large **statistical noise** possible

Effective scattering can make MC very expensive

Linearization can cause **non-physical** results (maximum principle violations)

Reconstruction of linear emission shape limits artificial energy propagation

HOLo Method

ECMC is **very efficient** for TRT problems

MC solution has **no scattering**

Fully **implicit** time-discretization and LO solution **resolves nonlinearities**

Linear-discontinuous FE for $T(x)$ preserving equilibrium diffusion limit

A HOLO Algorithm for Thermal Radiative Transfer



Overview of algorithm

Derivation of the LO equations

Exponentially Convergent MC High-Order Solver

Computational Results

Ongoing Research

A HOLO Algorithm for Thermal Radiative Transfer



Overview of algorithm

Derivation of the LO equations

Exponentially Convergent MC High-Order Solver

Computational Results

Ongoing Research

Solve a non-linear low-order system with high-order angular correction from efficient MC simulations

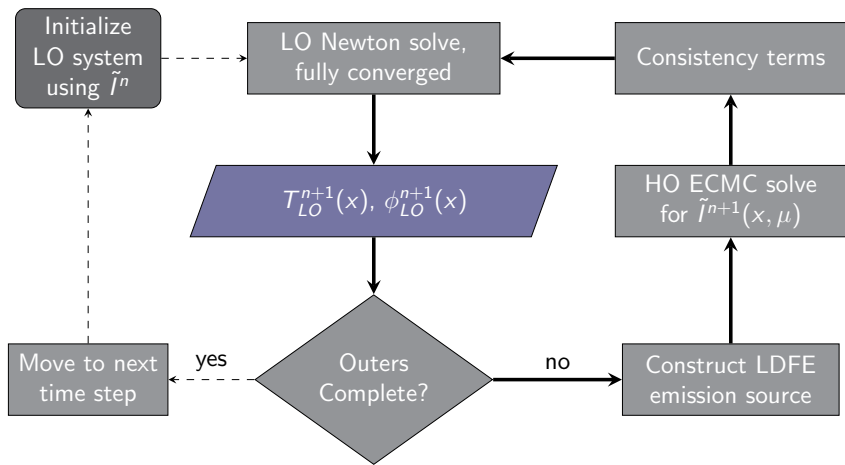
The **LO system** is space-angle moment equations, on a fixed finite-element (FE) spatial mesh

- ▶ Reduced dimensionality in angle
allows for solution with Newton's method
- ▶ **Output:** linear-discontinuous $\phi(x)$ and $T(x)$
Construct LDFE scattering and emission source

The **HO system** is a pure-absorber transport problem

- ▶ Solved with exponentially-convergent MC (ECMC)
for *efficient* reduction of statistical noise
- ▶ **Output:** consistency terms

Iterations between the HO and LO systems are performed each time step



A HOLO Algorithm for Thermal Radiative Transfer



Overview of algorithm

Derivation of the LO equations

Exponentially Convergent MC High-Order Solver

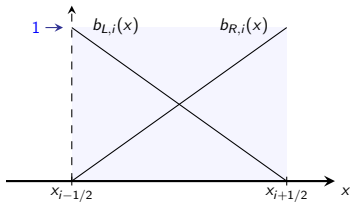
Computational Results

Ongoing Research

The LO equations are formed as *consistently* as possible with spatial and angular moments of TRT equations

The time discretization is backward Euler
for both the HO and LO equations

FE basis functions weight spatial moments:



$$\langle \cdot \rangle_{L,i} = \frac{2}{h_i} \int_{x_{i-1/2}}^{x_{i+1/2}} b_{L,i}(x)(\cdot) dx$$

Half-range integrals reduce angular dimensionality

$$\phi^+(x) = 2\pi \int_0^1 I(x, \mu) d\mu$$

Apply moments to the TRT equations
and manipulate to form **angular consistency terms**

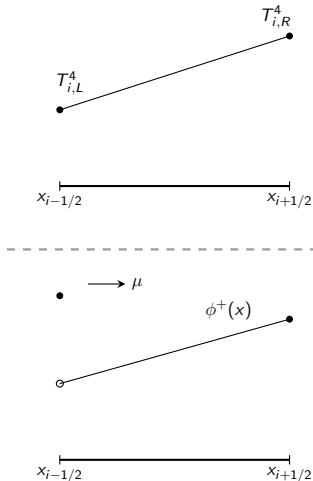
Ultimately, we get 4 radiation and 2 material equations
for each spatial element i

The streaming terms are algebraically manipulated
to form consistency terms (**no approximation here**)

For example, apply $\langle \cdot \rangle_{L,i}^+$ to a streaming term:

$$\begin{aligned} \left\langle \mu \frac{\partial I}{\partial x} \right\rangle_L^+ &= -\frac{2}{h_i} \{ \mu I \}_{i-1/2}^+ + \langle \mu I \rangle_{L,i}^+ + \langle \mu I \rangle_{R,i}^+ \\ &= -\frac{2}{h_i} \frac{\{ \mu I \}_{i-1/2}^+}{\phi_{i-1/2}^+} \phi_{i-1/2}^+ + \frac{\langle \mu I \rangle_{L,i}^+}{\langle \phi \rangle_{L,i}^+} \langle \phi \rangle_{L,i}^+ + \frac{\langle \mu I \rangle_{R,i}^+}{\langle \phi \rangle_{R,i}^+} \langle \phi \rangle_{R,i}^+ \end{aligned}$$

We can close equations with HO angular information and a linear-discontinuous (LD) spatial discretization



1. Assume $T(x)$ and $T^4(x)$ are LD

2. A lagged $\tilde{l}_{\text{HO}}^{n+1}$ is used to evaluate consistency terms

3. Eliminate $\phi_{i+1/2}^\pm$ with LD closure preserving equi. diff. limit

$$\phi_{i+1/2}^+ = 2\langle\phi\rangle_{R,i}^+ - \langle\phi\rangle_{L,i}^+$$

4. Global system solved with Newton's method and lagged **implicit opacities**
Energy is always conserved

Use source-iteration with diffusion-synthetic acceleration (DSA) to solve LO equations

In higher dimensions, the scattering terms cannot be directly inverted easily

Use source iteration with WLA-D
for (effective) scattering source of each Newton step

1. Sweep for a new ϕ^\pm
with a lagged scattering source
2. Solve approximate **spatially continuous** diffusion equation for error in scattering iterations
3. Update with local balance equations over elements

Recast iterations as DSA-preconditioned Krylov method to resolve convergence difficulties with inconsistencies

A HOLO Algorithm for Thermal Radiative Transfer



Overview of algorithm

Derivation of the LO equations

Exponentially Convergent MC High-Order Solver

Computational Results

Ongoing Research

Exponentially Convergent Monte Carlo can efficiently reduce noise globally

Each MC batch tallies the **error** in the solution

- ▶ standard MC particle transport,
but a **complex** source
- ▶ ECMC requires a **functional** representation of $I(x, \mu)$

Can reduce solution error **globally** $\propto e^{-\alpha N}$

Adaptive h -refinement can help represent error

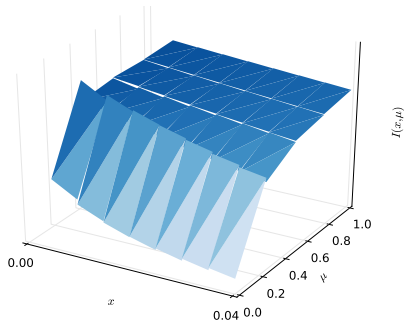
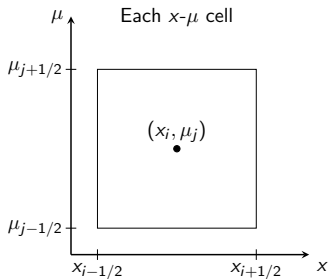
$I^n(x, \mu)$ often provides an *excellent* estimate of $I^{n+1}(x, \mu)$

No MC sampling for equilibrium regions

We use a **projection** $\tilde{I}(x, \mu)$ of the angular intensity onto a LDFFE space-angle mesh

local volumetric tallies

$$\tilde{I}_{ij}(x, \mu) = I_a + \frac{2}{h_x} I_x (x - x_i) + \frac{2}{h_\mu} I_\mu (\mu - \mu_i)$$



We apply the ECMC algorithm to the **pure-absorber** HO transport equation

$$\left[\mu \frac{\partial}{\partial x} + \left(\sigma_t + \frac{1}{c \Delta t} \right) \right] I^{n+1} = \frac{1}{4\pi} \left[\sigma_{aac} (T_{LO}^{n+1})^4 + \sigma_s \phi_{LO}^{n+1} \right] + \frac{\tilde{I}^n}{c \Delta t}$$
$$\mathbf{L} I^{n+1} = q$$

For each batch m :

- ▶ Evaluate residual source: $r^{(m)} = q - \mathbf{L} \tilde{I}^{n+1,(m)}$
- ▶ Estimate $\epsilon^{(m)} = \mathbf{L}^{-1} r^{(m)}$ via **Monte Carlo simulation**
- ▶ Update solution:

$$\begin{aligned} \tilde{I}^{n+1,(m+1)} &= \tilde{I}^{n+1,(m)} + \tilde{\epsilon}^{(m)} \\ &= \tilde{I}^{n+1,(m)} + \mathbf{L}^{-1} q - \mathbf{L}^{-1} \mathbf{L} \tilde{I}^{n+1,(m)} \end{aligned}$$

Our HO system allows for simple variance reduction methods

Histories stream without collision

along path s , weight reduces as $w(s) = w_0 e^{-\sigma_t s}$

Use cell-wise systematic sampling for $|r^{(m)}|$ source
Particularly effective in thick cells

- ▶ Particles in each $x-\mu$ cell $\propto |r^{(m)}|$ in cell
- ▶ Set minimum n for cells
except for cells in thermal equilibrium

A HOLO Algorithm for Thermal Radiative Transfer



Overview of algorithm

Derivation of the LO equations

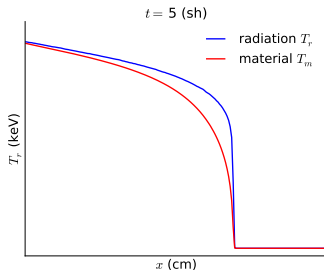
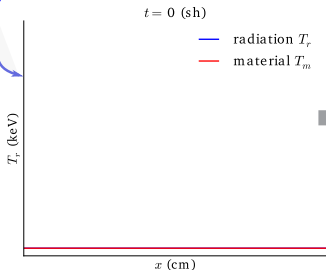
Exponentially Convergent MC High-Order Solver

Computational Results

Ongoing Research

We will test our method with several standard **Marshak Wave** problems

constant radiation
boundary source

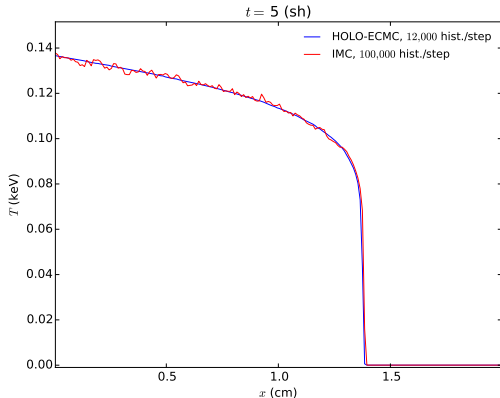


Results show radiation temperature $T_r = \sqrt[4]{\phi/ac}$

The HOLO method produces significantly less noise than IMC in a Marshak Wave test problem

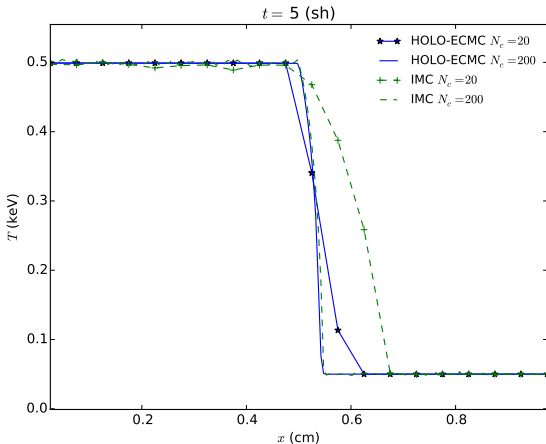
► $\sigma_a \propto T^{-3}$

- Transient solution after 5 shakes
200 x cells and 4 μ cells (for ECMC)

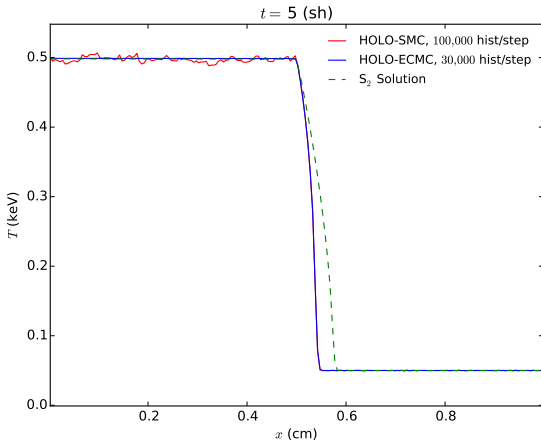


The LDFE representation has higher spatial accuracy than IMC source tilting for two-material problem

Left half is optically thin ($\sigma=0.2 \text{ cm}^{-1}$) and right half is thick ($\sigma_a=2000 \text{ cm}^{-1}$). ECMC uses 8μ cells



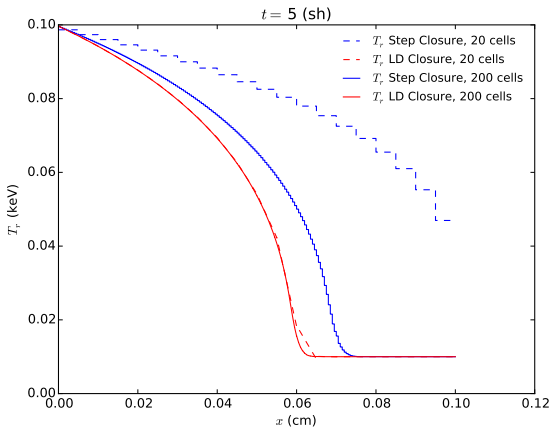
ECMC is more efficient than
standard MC (SMC) as a HO solver



Different HO solvers: **ECMC** with 3 batches,
standard MC (**SMC**), and an S_2 solution

The LDFE discretization for the LO equations preserves the equilibrium diffusion limit

Large constant σ_a and small c_v

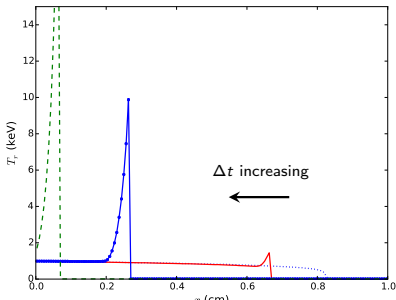


Our HOLO method preserves the maximum principle with sufficient nonlinear convergence

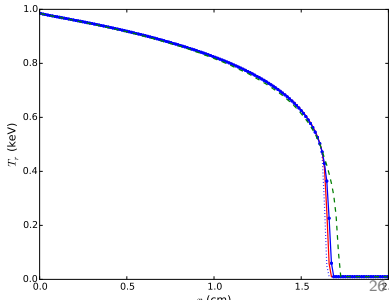
Comparison of material temperatures for different Δt
(different axis scales)

- ▶ All simulations end at $t = 0.1$ sh
 $\sigma_a \propto T^{-3}$, c_V is small
- ▶ LO Newton iterations required damping

IMC

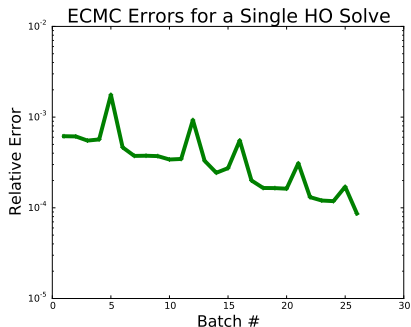
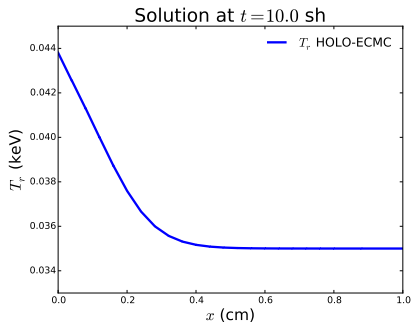


HOLO



DSA calculations

Exponential convergence can be maintained if the LDFF mesh resolves the solution reasonably



A HOLO Algorithm for Thermal Radiative Transfer



Overview of algorithm

Derivation of the LO equations

Exponentially Convergent MC High-Order Solver

Computational Results

Ongoing Research

We need a way to resolve issues when the LDFF representation of the intensity is negative

Negative intensities can occur in optically thick cells
Mesh refinement is of minimal use

$\tilde{I}_{HO}(x, \mu)$ must be positive for consistency terms
to produce a physical, stable LO solution

Independent fix up for LO solution
E.g., lumping or preserving balance with floored $\phi(x)$

Can add source δ to produce a positive projection \tilde{l}_{pos} such that \tilde{l}_{pos} satisfies the latest residual equation

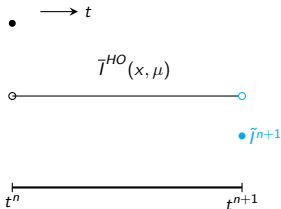
Produce \tilde{l}_{pos} by scaling $x - \mu$ moments equally,
to estimate source for the next iteration

$$\begin{array}{l} \mathbf{L} \tilde{l}^{(m)} = q - r^{(m)} \\ \mathbf{L} \tilde{l}_{pos}^{(m)} = q - r^{(m)} + \delta^{(m+1)} \end{array} \quad \longrightarrow \quad \begin{array}{l} \delta^{(m+1)} = \mathbf{L} \left(\tilde{l}^{(m)} - \tilde{l}_{pos}^{(m)} \right) \\ q \rightarrow q + \delta^{(m+1)} \end{array}$$

Could add source to LO equation
but it would affect energy conservation

Alternatively, add $\delta = -r_{pos}^{(m)}(x, \mu)$ to negative cells
 l_{pos} must be recomputing based on balance

A doubly-discontinuous trial space in time allows for a MC temporal closure



Include $\frac{1}{c} \frac{\partial}{\partial t} (\cdot)$ in transport operator \mathbf{L}
with $T(x)$ still implicit

Sample and track in time
Tally $\tilde{\epsilon}(x, \mu)$ at t^{n+1}

In **LO equations** parameterize ϕ_{LO}^{n+1}
in terms of time-averaged unknowns, e.g.,

$$\phi_{LO}^{n+1} = 2 \gamma_{HO} \bar{\phi}_{i,LO} - \phi_{LO}^n$$

Discuss briefly resolving issues with negativities, probably in smooth slide

TRY TO GET PARTICLE HISTORY COUNTS FOR SMOOTH PROBLEM

The time-closure parameters preserve the accuracy of MC time-treatment

- ▶ Problem is near void, so temperature uncouples
- ▶ 3 large time steps

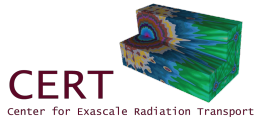
Backup slide with timing results

There are short term and long term goals left to investigate

1. Resolving issues with negative intensities
 - ▶ Accuracy of added source method
 - ▶ Consistency with LO solution
2. Using HO solution to estimate spatial closure
3. Source iterations with DSA for LO system
4. Implement damped Newton method to demonstrate maximum principle preservation in extreme problems
5. (Stretch goal) MC integration in time with consistent LO equations

Backup Slides

Simon Bolding and Jim Morel



Implementation specifics for results in the computational results section

- ▶ The LD representation of $I(x, \mu)$ is negative near the wave-front
 - ▶ Here, no correction is applied to the HO solution, and the LO solution uses lumped LD and S_2 equivalent terms in negative elements
- ▶ For all results
 1. Initial Δt of 0.001 sh (0.01 ns), linearly increased to 0.01 sh by 15% per step
 2. One HO solve per time step (predictor-corrector)
 - ▶ *each HO* solve has 3 ECMC batches
 3. $\sigma_s = 0$
 4. No mesh refinement in ECMC

Solving LO System with Newton's Method

- Linearization:

$$\underline{B}(T^{n+1}) = \underline{B}(T^*) + (T^{n+1} - T^*) \left. \frac{\partial \underline{B}}{\partial t} \right|_{t^*}$$

- Modified system

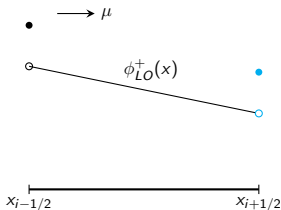
$$[\mathbf{D}(\mu^\pm) - \sigma_a^*(1 - f^*)] \underline{\Phi}^{n+1} = f^* \underline{B}(T^*) + \frac{\underline{\Phi}^n}{c\Delta t}$$

$$\hat{\mathbf{D}} \underline{\Phi}^{n+1} = \underline{Q}$$

$$f = \left(1 + \sigma_a^* c \Delta t \frac{4aT^{*3}}{\rho c_v} \right)^{-1} \quad T_i^* = \frac{T_{L,i}^* + T_{R,i}^*}{2}$$

- Equation for T^{n+1} based on linearization that is conservative
- Converge T^{n+1} and $\langle \phi \rangle$ with Newton Iterations

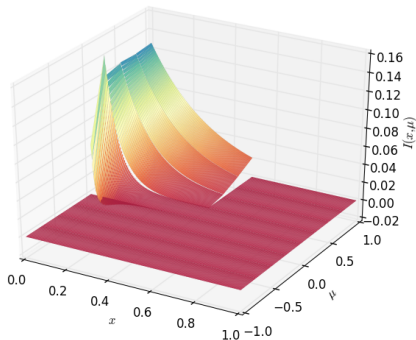
We will use a linear doubly-discontinuous (LDD) trial space to allow for a HO spatial closure



- ▶ $\tilde{I}_{HO}(x, \mu)$ will be linear in μ along face
Error estimated with MC face tallies
- ▶ In LO equations use LDD for ϕ^\pm
The linear interior preserves the EDL
- ▶ Parameterize LO spatial closure to eliminate outflow:

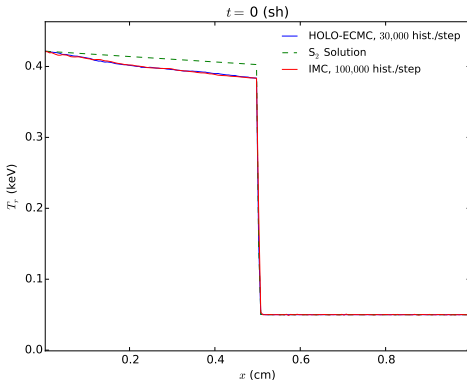
$$\phi_{i+1/2}^+ = \frac{3 + \gamma_{i,HO}^+}{2} \langle \phi \rangle_{R,i}^+ + \frac{\gamma_{i,HO}^+ - 3}{2} \langle \phi \rangle_{L,i}^+$$

The angular flux for the two material problem is difficult to resolve near $\mu = 0$



Two Material Problem, comparison in optically thin region

- Plot of radiation temperature after 10 time steps



Derivation of LO System

- Taking moments of TE yields 4 equations, per cell i , e.g.

$$\begin{aligned}
 & -2\mu_{i-1/2}^{n+1,+} \phi_{i-1/2}^{n+1,+} + \{\mu\}_{L,i}^{n+1,+} \langle \phi \rangle_{L,i}^{n+1,+} + \{\mu\}_{R,i}^{n+1,+} \langle \phi \rangle_{R,i}^{n+1,+} + \\
 & \left(\sigma_t^{n+1} + \frac{1}{c\Delta t} \right) h_i \langle \phi \rangle_{L,i}^{n+1,+} - \frac{\sigma_s h_i}{2} \left(\langle \phi \rangle_{L,i}^{n+1,+} + \langle \phi \rangle_{L,i}^{n+1,-} \right) \\
 & = \frac{h_i}{2} \langle \sigma_a^{n+1} ac T^{n+1,4} \rangle_{L,i} + \frac{h_i}{c\Delta t} \langle \phi \rangle_{L,i}^{n,+}, \quad (1)
 \end{aligned}$$

- Cell unknowns: $\langle \phi \rangle_{L,i}^+$, $\langle \phi \rangle_{R,i}^+$, $\langle \phi \rangle_{L,i}^-$, $\langle \phi \rangle_{R,i}^-$, T_L , T_R
- Need angular consistency terms and spatial closure (LD)