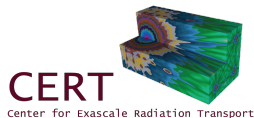


A High-Order Low-Order Algorithm with Exponentially-Convergent Monte Carlo for Thermal Radiative Transfer

Simon Bolding and Jim Morel

4 August 2016



We are interested in modeling thermal radiation transport in the high energy density physics regime

Materials under extreme conditions

Temperatures $\mathcal{O}(10^6)$ K or more

Photon radiation transports through a material

Significant **energy** may be exchanged

We want to improve efficiency of calculations

e.g., inertial confinement fusion, supernovae, et. al.

Our method has been applied to a simplified model:
the 1D grey radiative transfer equations

Energy balance equations for radiation and material.

Radiation intensity $I(x, \mu, t)$, material temperature $T(x, t)$

$$\frac{1}{c} \frac{\partial I}{\partial t} + \mu \frac{\partial I}{\partial x} + \sigma_t I(x, \mu, t) = \frac{1}{4\pi} \sigma_a a c T^4,$$
$$C_v \frac{\partial T(x, t)}{\partial t} = \sigma_a \phi(x, t) - \sigma_a a c T^4$$

Equations are **nonlinear** and may be tightly coupled

Absorption opacity (σ_a) can be a strong function of T

Implicit Monte Carlo (IMC) is the standard Monte Carlo transport method for TRT problems

The system is *linearized* over a time step $t \in [t^n, t^{n+1}]$
Opacities are evaluated with $T(t^n)$

- ▶ Produces a linear MC transport problem
with effective emission and scattering terms
- ▶ Emission source is **not** fully implicit.
Monte Carlo integration over Δt for intensity

Our high-order low-order (HOLON) method improves on several drawbacks of IMC

Standard IMC

Large **statistical noise** possible

Effective scattering can make MC very expensive

Linearization can cause **non-physical** results (maximum principle violations)

Reconstruction of linear emission shape limits artificial energy propagation

HOLON Method

ECMC is **very efficient** for TRT problems

MC solution has **no scattering**

Fully **implicit** time-discretization and LO solution **resolves nonlinearities**

Linear-discontinuous FE for $T(x)$ preserving equilibrium diffusion limit

A HOLO Algorithm for Thermal Radiative Transfer



Overview of algorithm

Derivation of the LO equations

Exponentially Convergent MC High-Order Solver

Computational Results

Ongoing Research

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Solve a non-linear low-order system with high-order angular correction from efficient MC simulations

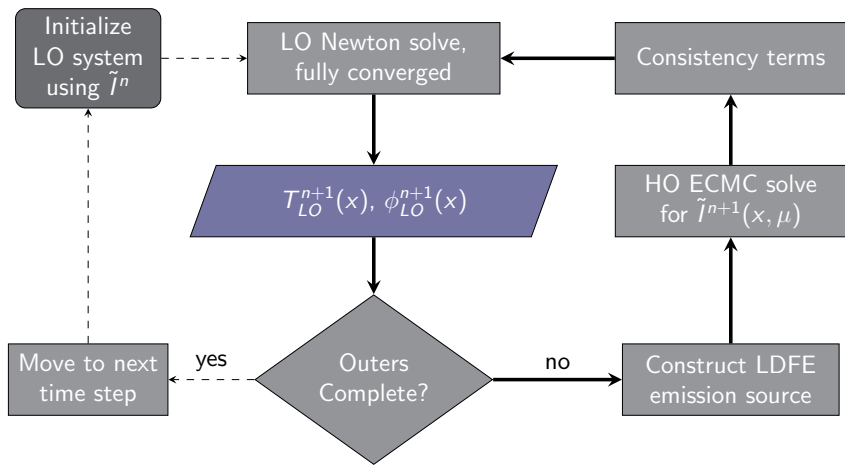
The **LO system** is space-angle moment equations, on a fixed finite-element (FE) spatial mesh

- ▶ Reduced dimensionality in angle
allows for solution with Newton's method
- ▶ **Output:** linear-discontinuous $\phi(x)$ and $T(x)$
Construct LDFE scattering and emission source

The **HO system** is a pure-absorber transport problem

- ▶ Solved with exponentially-convergent MC (ECMC)
for *efficient* reduction of statistical noise
- ▶ **Output:** consistency terms

Iterations between the HO and LO systems are performed each time step



A HOLO Algorithm for Thermal Radiative Transfer



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Exponentially Convergent MC High-Order Solver

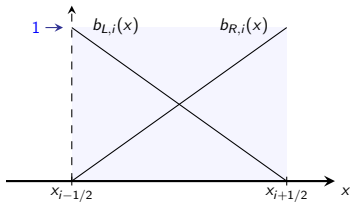
Computational Results

Ongoing Research

The LO equations are formed as *consistently* as possible with spatial and angular moments of TRT equations

The time discretization is backward Euler
for both the HO and LO equations

FE basis functions weight spatial moments:



$$\langle \cdot \rangle_{L,i} = \frac{2}{h_i} \int_{x_{i-1/2}}^{x_{i+1/2}} b_{L,i}(x)(\cdot) dx$$

Half-range integrals reduce angular dimensionality

$$\phi^+(x) = 2\pi \int_0^1 I(x, \mu) d\mu$$

Apply moments to the TRT equations
and manipulate to form **angular consistency terms**

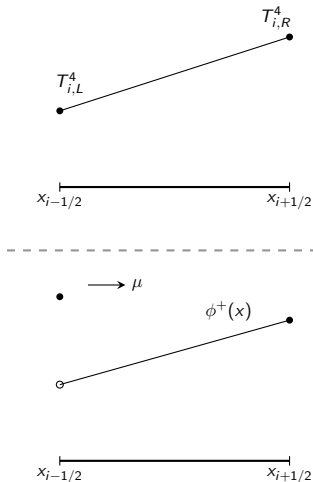
Ultimately, we get 4 radiation and 2 material equations
for each spatial element i

The streaming terms are algebraically manipulated
to form consistency terms (**no approximation here**)

For example, apply $\langle \cdot \rangle_{L,i}^+$ to a streaming term:

$$\begin{aligned} \left\langle \mu \frac{\partial I}{\partial x} \right\rangle_L^+ &= -\frac{2}{h_i} \{ \mu I \}_{i-1/2}^+ + \langle \mu I \rangle_{L,i}^+ + \langle \mu I \rangle_{R,i}^+ \\ &= -\frac{2}{h_i} \frac{\{ \mu I \}_{i-1/2}^+}{\phi_{i-1/2}^+} \phi_{i-1/2}^+ + \frac{\langle \mu I \rangle_{L,i}^+}{\langle \phi \rangle_{L,i}^+} \langle \phi \rangle_{L,i}^+ + \frac{\langle \mu I \rangle_{R,i}^+}{\langle \phi \rangle_{R,i}^+} \langle \phi \rangle_{R,i}^+ \end{aligned}$$

We can close equations with HO angular information and a linear-discontinuous (LD) spatial discretization



1. Assume $T(x)$ and $T^4(x)$ are LD

2. A lagged $\tilde{T}_{\text{HO}}^{n+1}$ is used to evaluate consistency terms

3. Eliminate $\phi_{i+1/2}^\pm$ with LD closure preserving equi. diff. limit

$$\phi_{i+1/2}^+ = 2\langle\phi\rangle_{R,i}^+ - \langle\phi\rangle_{L,i}^+$$

4. Global system solved with Newton's method and lagged **implicit opacities**
Energy is always conserved

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Computational Results

Ongoing Research

Exponentially Convergent Monte Carlo can efficiently reduce noise globally

Each MC batch tallies the **error** in the solution

- ▶ standard MC particle transport,
but a **complex** source
- ▶ ECMC requires a **functional** representation of $I(x, \mu)$

Can reduce solution error **globally** $\propto e^{-\alpha N}$

Adaptive h -refinement can help represent error

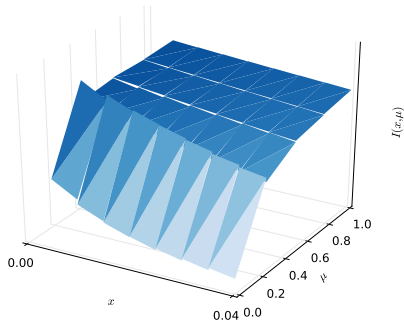
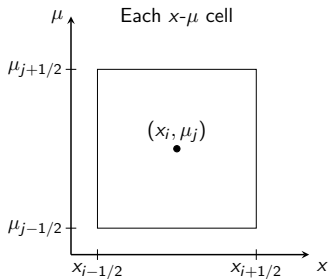
$I^n(x, \mu)$ often provides an **excellent** estimate of $I^{n+1}(x, \mu)$

No MC sampling from thermal equilibrium regions

We use a **projection** $\tilde{I}(x, \mu)$ of the angular intensity onto a LDFFE space-angle mesh

local volumetric tallies

$$\tilde{I}_{ij}(x, \mu) = I_a + \frac{2}{h_x} I_x (x - x_i) + \frac{2}{h_\mu} I_\mu (\mu - \mu_i)$$



We apply the ECMC algorithm to the **pure-absorber** HO transport equation

$$\left[\mu \frac{\partial}{\partial x} + \left(\sigma_t + \frac{1}{c \Delta t} \right) \right] I^{n+1} = \frac{1}{4\pi} \left[\sigma_a a c (T_{LO}^{n+1})^4 + \sigma_s \phi_{LO}^{n+1} \right] + \frac{\tilde{I}^n}{c \Delta t}$$
$$\mathbf{L} I^{n+1} = q$$

For each batch m :

- ▶ Evaluate residual source: $r^{(m)} = q - \mathbf{L} \tilde{I}^{n+1,(m)}$
- ▶ Estimate $\epsilon^{(m)} = \mathbf{L}^{-1} r^{(m)}$ via **MC simulation**
- ▶ Update solution: $\tilde{I}^{n+1,(m+1)} = \tilde{I}^{n+1,(m)} + \tilde{\epsilon}^{(m)}$

Our HO system allows for simple variance reduction methods

Histories stream without collision

along path s , weight reduces as $w(s) = w_0 e^{-\sigma_t s}$

Use cell-wise systematic sampling for $|r^{(m)}|$ source
Particularly effective in thick cells

- ▶ Particles in each $x-\mu$ cell $\propto |r^{(m)}|$ in cell
- ▶ Set minimum n for cells
except for cells in thermal equilibrium

A HOLO Algorithm for Thermal Radiative Transfer



Overview of algorithm

Derivation of the LO equations

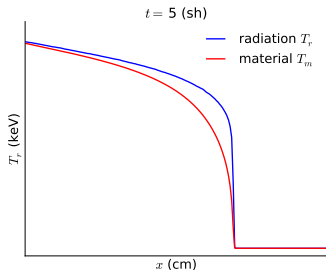
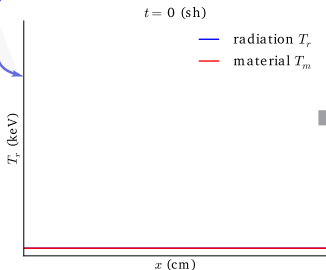
Exponentially Convergent MC High-Order Solver

Computational Results

Ongoing Research

We will test our method with several standard **Marshak Wave** problems

constant radiation
boundary source



Results show radiation temperature $T_r = \sqrt[4]{\phi/ac}$

Several specifics for most results shown:

Initial Δt of 0.001 sh (0.01 ns), linearly increased to 0.01 sh by 15% per step

One HO solve per time step (predictor-corrector)

- ▶ *each HO* solve has 3 ECMC batches
no mesh refinement

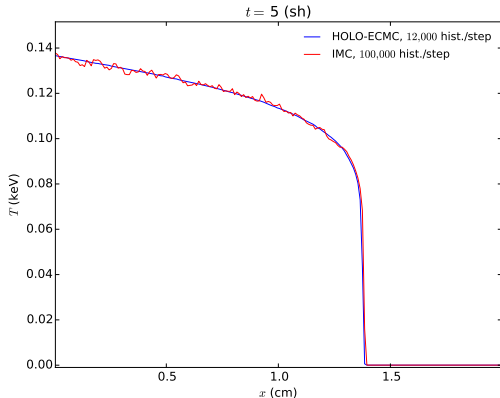
$$\text{FOM} = \frac{1}{\|s\|^2 N_{\text{total}}}$$

normalized to IMC result

The HOLO method produces significantly less noise than IMC in a Marshak Wave test problem

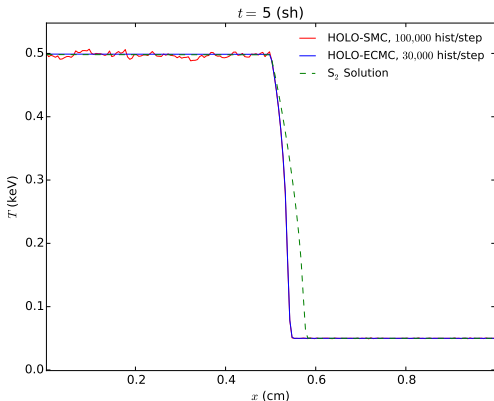
► $\sigma_a \propto T^{-3}$

- Transient solution after 5 shakes
200 x cells and 4 μ cells (for ECMC)



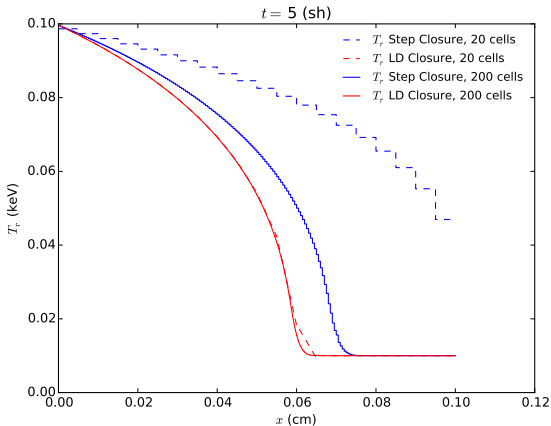
ECMC is more efficient than standard MC (SMC) as a HO solver

- ▶ Left half is optically thin ($\sigma=0.2 \text{ cm}^{-1}$), right half is thick ($\sigma_a=2000 \text{ cm}^{-1}$). 8 μ cells
- ▶ Different HO solvers: ECMC with 3 batches, standard MC (SMC), and an S_2 solution



The LDFE discretization for the LO equations preserves the equilibrium diffusion limit

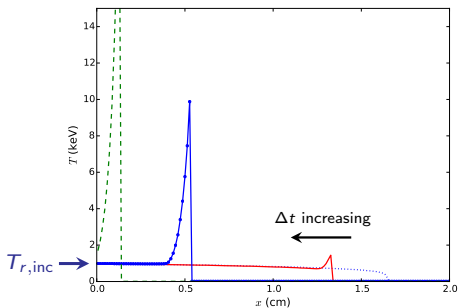
- Large, constant σ_a and small c_v



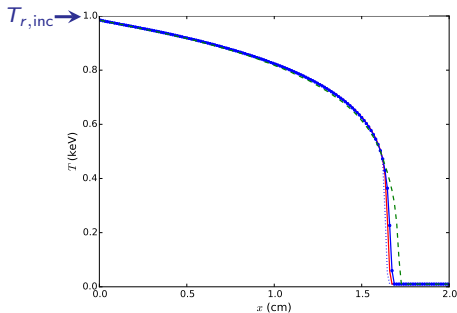
Our HOLO method preserves the maximum principle with sufficient nonlinear convergence

- ▶ **Material temperatures** plotted; all simulations end at $t = 0.1$ sh
 $\sigma_a \propto T^{-3}$, c_v small, $\Delta t \in [10^{-4}, 10^{-2}]$ sh
- ▶ LO Newton iterations required damping

IMC T_m



HOLO T_m



A HOLO Algorithm for Thermal Radiative Transfer



Overview of algorithm

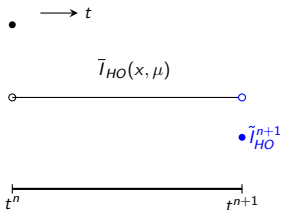
Derivation of the LO equations

Exponentially Convergent MC High-Order Solver

Computational Results

Ongoing Research

A doubly-discontinuous trial space in time allows for a MC temporal closure



Include $\frac{1}{c} \frac{\partial}{\partial t} (\cdot)$ in transport operator \mathbf{L}
with $T(x)$ still implicit

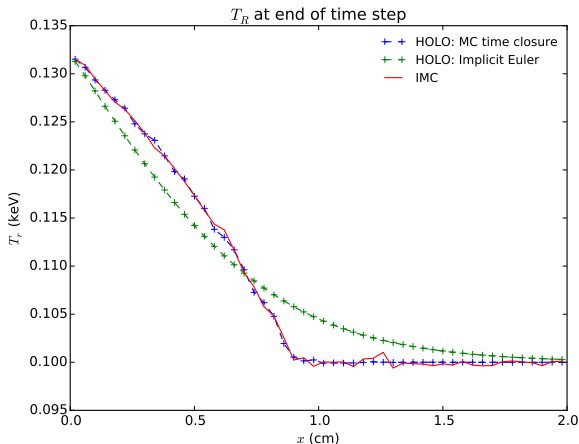
Sample and track in time
Tally error projection at t^{n+1}

In **LO equations** parameterize ϕ_{LO}^{n+1}
in terms of **time-averaged** unknowns,
e.g.,

$$\langle \phi \rangle_{L,i}^{n+1} = 2 \gamma_{L,i}^{HO} \overline{\langle \phi \rangle}_{L,i} - \langle \phi \rangle_{L,i}^n$$

The time-closure parameters preserve the accuracy of MC time integration

- Material is near-void, so temperature uncouples
take 3 large time steps and compare T_R^{n+1}



A few other topics are under investigation

Negative intensities can occur in optically thick cells
Investigating artificial source with alternate trial space

Using HO solution to estimate spatial closure
by including face tallies

Source iteration with diffusion synthetic acceleration
for the LO equations

Conclusions

ECMC is very efficient for TRT simulations
and fits well in HOLO context

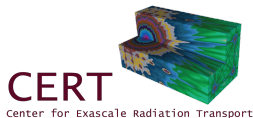
The LO system can resolve nonlinearities
with bounded angular consistency terms

The next step is to extend to higher dimensions
main hurdle to overcome is infrastructure

Simulations of Neutron Multiplicity Experiments with Perturbations to Nuclear Data

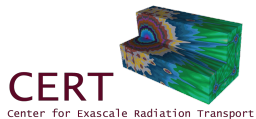
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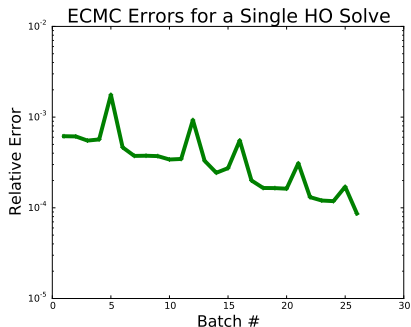
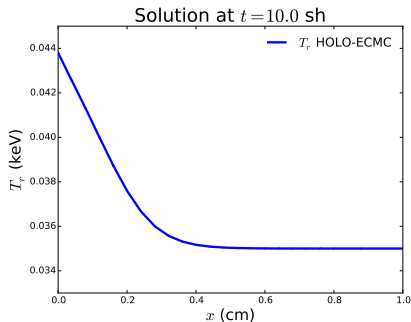


Backup Slides

S.R. Bolding and C.J. Solomon



Exponential convergence can be maintained if the LDFF mesh resolves the solution reasonably



We need a way to resolve issues when the LDFF representation of the intensity is negative

Negative intensities can occur in optically thick cells
Mesh refinement is of minimal use

$\tilde{I}_{HO}(x, \mu)$ must be positive for consistency terms
to produce a physical, stable LO solution

Independent fix up for LO solution

E.g., lumping or preserving balance with floored $\phi(x)$

Can add source δ to produce a positive projection \tilde{l}_{pos} such that \tilde{l}_{pos} satisfies the latest residual equation

Produce \tilde{l}_{pos} by scaling $x - \mu$ moments equally,
to estimate source for the next iteration

$$\begin{array}{lcl} \mathbf{L}\tilde{l}^{(m)} = q - r^{(m)} & \longrightarrow & \delta^{(m+1)} = \mathbf{L} \left(\tilde{l}^{(m)} - \tilde{l}_{pos}^{(m)} \right) \\ \mathbf{L}\tilde{l}_{pos}^{(m)} = q - r^{(m)} + \delta^{(m+1)} & & q \rightarrow q + \delta^{(m+1)} \end{array}$$

We can delay error stagnation

Investigating alternative positive projection of l

Solving LO System with Newton's Method

- Linearization:

$$\underline{B}(T^{n+1}) = \underline{B}(T^*) + (T^{n+1} - T^*) \left. \frac{\partial \underline{B}}{\partial t} \right|_{t^*}$$

- Modified system

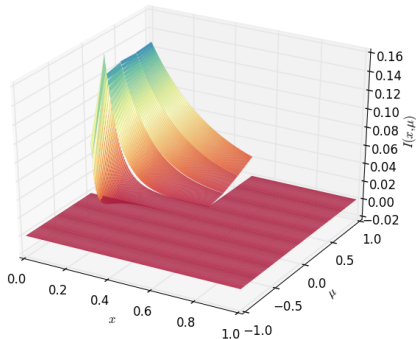
$$[\mathbf{D}(\mu^\pm) - \sigma_a^*(1 - f^*)] \underline{\Phi}^{n+1} = f^* \underline{B}(T^*) + \frac{\underline{\Phi}^n}{c\Delta t}$$

$$\hat{\mathbf{D}} \underline{\Phi}^{n+1} = \underline{Q}$$

$$f = \left(1 + \sigma_a^* c \Delta t \frac{4aT^{*3}}{\rho c_v} \right)^{-1} \quad T_i^* = \frac{T_{L,i}^* + T_{R,i}^*}{2}$$

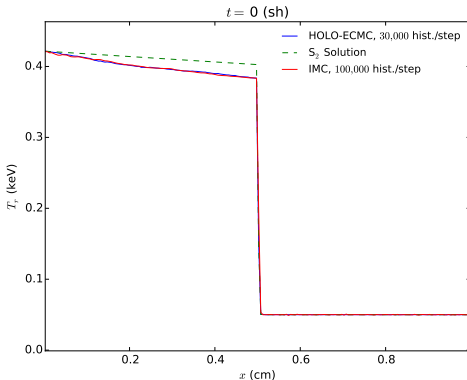
- Equation for T^{n+1} based on linearization that is conservative
- Converge T^{n+1} and $\langle \phi \rangle$ with Newton Iterations

The angular flux for the two material problem is difficult to resolve near $\mu = 0$



Two Material Problem, comparison in optically thin region

- Plot of radiation temperature after 10 time steps



Backup slide with timing results

Derivation of LO System

- Taking moments of TE yields 4 equations, per cell i , e.g.

$$\begin{aligned}
 & -2\mu_{i-1/2}^{n+1,+} \phi_{i-1/2}^{n+1,+} + \{\mu\}_{L,i}^{n+1,+} \langle \phi \rangle_{L,i}^{n+1,+} + \{\mu\}_{R,i}^{n+1,+} \langle \phi \rangle_{R,i}^{n+1,+} + \\
 & \left(\sigma_t^{n+1} + \frac{1}{c\Delta t} \right) h_i \langle \phi \rangle_{L,i}^{n+1,+} - \frac{\sigma_s h_i}{2} \left(\langle \phi \rangle_{L,i}^{n+1,+} + \langle \phi \rangle_{L,i}^{n+1,-} \right) \\
 & = \frac{h_i}{2} \langle \sigma_a^{n+1} a c T^{n+1,4} \rangle_{L,i} + \frac{h_i}{c\Delta t} \langle \phi \rangle_{L,i}^{n,+}, \quad (1)
 \end{aligned}$$

- Cell unknowns: $\langle \phi \rangle_{L,i}^+$, $\langle \phi \rangle_{R,i}^+$, $\langle \phi \rangle_{L,i}^-$, $\langle \phi \rangle_{R,i}^-$, T_L , T_R
- Need angular consistency terms and spatial closure (LD)