

# A High-Order Low-Order Algorithm with Exponentially-Convergent Monte Carlo for Thermal Radiative Transfer

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# Outline

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# Overview

Thermal radiative transfer is relevant in very high temperature heat transfer scenarios, representing a complicated interaction between matter and a radiation field. The radiation field heats the material and moves energy through photon streaming and scatter. The material conducts energy and emits radiation proportional to  $T^4$ . Time dependent, radiative transfer simulations are important in:

- Inertial confinement fusion
- High energy plasma calculations
- Astrophysics calculations

# The thermal radiative transfer equations

- The 1D, frequency-integrated (grey) equations

$$\frac{1}{c} \frac{\partial I}{\partial t} + \mu \frac{\partial I}{\partial x} + \sigma_t I(x, \mu) = \frac{1}{2} \sigma_a a c T^4 \quad (1)$$

$$C_v \frac{\partial T}{\partial t} = \sigma_a \phi - \sigma_a a c T^4. \quad (2)$$

$$\phi(x) = \int_{-1}^1 I(x, \mu) d\mu \quad (3)$$

- Fundamental unknowns are the radiation intensity  $I(x, \mu)$  and material temperature  $T$
- Cross sections ( $\sigma$ ) can be a strong function of  $T$
- Equations are tightly coupled through source terms and contain the non-linear Planckian emission source  $a c T^4$

# The implicit Monte Carlo method

- These equations are typically solved with Monte Carlo (MC) via the Implicit Monte Carlo (IMC) method
- In IMC, the material energy equation is linearized **over a time step** and eliminated from the system
  - Results in an effective source and scattering cross section for a **linear** transport problem
  - This linear transport equation is solved with standard MC particle transport algorithms
  - Continuous integration in time
- Drawbacks
  - Effective scattering cross section can be very large, resulting in many expensive MC scattering events
  - System is not truly implicit in time, because of linearization
  - Artificial reconstruction of linear source shape in cell

# An alternative High-Order Low-Order approach

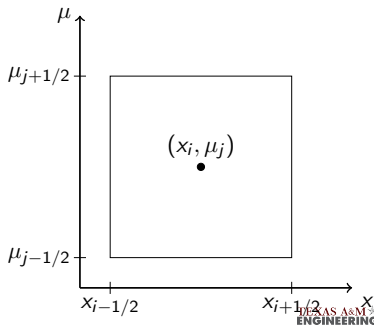
## Basic Idea

Build a low-order (LO) system that can be efficiently solved, with high-order (HO) correction from simpler MC simulations

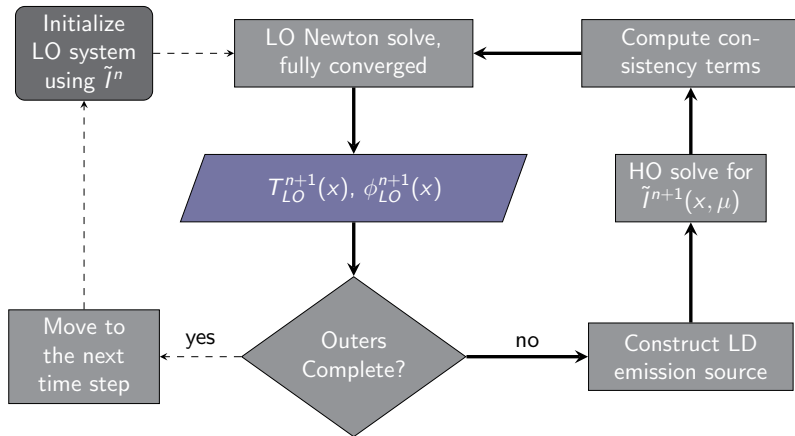
- Useful when global solution of particle density needed
- The LO solution preserves moments of the HO solution upon convergence
- The LO solver resolves non-linear temperature dependence with Newtons method
  - Lower dimensional problem in angular variable
  - Produces a linear-discontinuous (LD) finite element representation of sources
- The HO system is a fixed-source, pure-absorber transport problem, which we solve with ECMC
  - No effective scattering events

# Overview of Exponentially Convergent Monte Carlo

- Iterative form of residual Monte Carlo
  - Each batch tallies the **error** in current estimate of solution, which is a transport problem with a reduced source
  - Can reduce statistical error **globally**  $\propto e^{-\alpha N}$
  - In TRT problems, old angular intensity provides a very good guess of new solution, significantly reducing the required number of histories
- Requires a **discretized** form of the angular intensity  $\tilde{I}(x, \mu)$ 
  - Use **projection** of the solution onto a space-angle FE mesh
  - Projection computed using path-length estimators of moments of  $I$



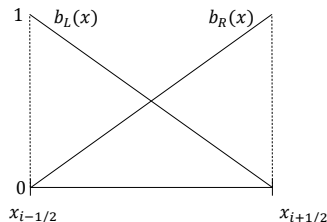
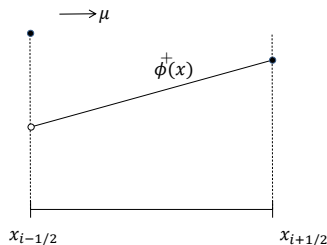
# High-Order Low-Order Algorithm





# LO Discretization & Space-Angle Moments

- Backward Euler discretization in time
- Linear discontinuous (LD) FE in space for  $\phi$  and  $T$ , and half-range averages in angle



- *Examples of moments:*

Spatial: left basis

$$\langle \cdot \rangle_{L,i} = \frac{2}{h_i} \int_{x_{i-1/2}}^{x_{i+1/2}} b_{L,i}(x)(\cdot) dx$$

Angular: positive flow

$$\phi^+(x) = 2\pi \int_0^1 \psi(x, \mu) d\mu$$

## Forming the LO System

- We take spatial and angular moments of the equations to reduce dimensionality.
- Use a backward Euler discretization in time
- Taking moments of T.E. yields 4 radiation equations, and 2 material energy equations, per cell
- Conserves the total energy in the system
- Cell unknowns:  $\langle \phi \rangle_{L,i}^{n+1,+}$ ,  $\langle \phi \rangle_{R,i}^{n+1,+}$ ,  $\langle \phi \rangle_{L,i}^{n+1,-}$ ,  $\langle \phi \rangle_{R,i}^{n+1,-}$ ,  $T_{L,i}$ ,  $T_{R,i}$
- The global system for radiation, **fully implicit** in time, is

$$\mathbf{D}^{n+1} \underline{\Phi}^{n+1} = \underline{B}(T^{n+1}) + \frac{\Phi^n}{c\Delta t}$$

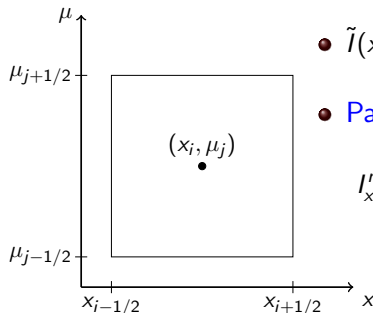
# Computing LO Consistency terms from HO Solution

- The streaming operator  $D$  contains unknown weighted angular averages, called consistency terms
- For  $\mu > 0$ ,  $L$  moment

$$\langle \mu \rangle_{L,i}^+ \simeq \frac{\frac{2}{h_i} \int_0^1 \int_{x_{i-1/2}}^{x_{i+1/2}} \mu b_{L,i}(x) \tilde{\psi}^{HO}(x, \mu) d\mu dx}{\frac{2}{h_i} \int_0^1 \int_{x_{i-1/2}}^{x_{i+1/2}} b_{L,i}(x) \tilde{\psi}^{HO}(x, \mu) d\mu dx} \quad (4)$$

- ECMC gives LDfE  $\tilde{\psi}^{HO}(x, \mu)$  for computing terms directly
- To get  $S_2$ , set  $\langle \mu \rangle^\pm = \pm \frac{1}{\sqrt{3}}$

# Space-Angle Mesh and MC Implementation Details



- $\tilde{l}(x, \mu) = l_{a,i} + l_{x,i} \frac{2}{h_x} (x - x_i) + l_{\mu,i} \frac{2}{h_\mu} (\mu - \mu_i)$

- Path-length estimators of moments, e.g.

$$l_{x,i}^{n+1} = \frac{6}{h_x^2 h_\mu} \iint_{\mathcal{D}} (x_{c,j} - x_i) l^{n+1}(x, \mu) dx d\mu$$

- Discrete in time, resulting in fixed-source problem
- Particles are allowed to stream,  $w(s) = w_0 e^{-\Sigma_t s}$
- LDFE and upwinding eliminates surface tallies
- Cell-wise, global representation allows for stratified sampling
  - $N_{i,j} \propto |r_i(x, \mu)|$

# High Order System and ECMC Algorithm

- **Pure absorber** transport problem because we know RHS from LO solution

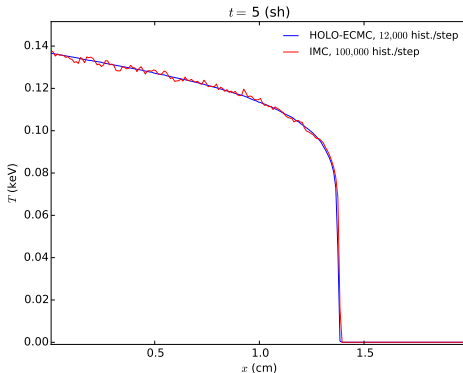
$$\left[ \mu \frac{\partial}{\partial x} + \left( \sigma_t + \frac{1}{c \Delta t} \right) \right] I^{n+1}(x, \mu) = \boxed{\frac{I^n}{c \Delta t} + \frac{1}{2} \sigma_{aac} T_{LO}^{n+1,4}}$$
$$\mathbf{L} I^{n+1} = q$$

## ECMC Algorithm

- Residual Equation:  $\mathbf{L} \tilde{\epsilon}^{(m)} = \tilde{r}^{(m)} = q - \mathbf{L} \tilde{j}^{n+1,(m)}$
- Compute  $\tilde{\epsilon}^{(m)} = \mathbf{L}^{-1} \tilde{r}^{(m)}$  with MC, **projecting** the solution
- Update:  $\tilde{j}^{n+1,(m)} = \tilde{j}^{n+1,(m-1)} + \tilde{\epsilon}^{(m)}$ 
  - If  $\tilde{\epsilon}$  is reduced each batch, **exponential convergence achieved**
  - Must sufficiently reduce noise in  $\epsilon$  tallies, each batch
  - Issues when solution cannot be represented within a cell
- Repeat for fixed number of batches

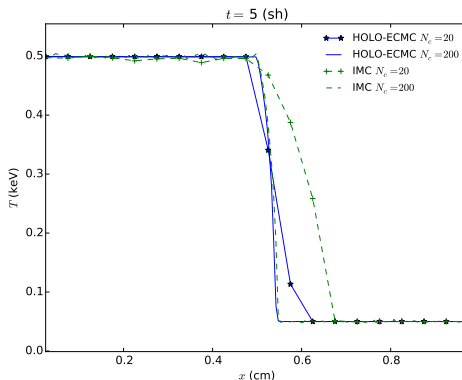
# Marshak Wave Test Problem

- From equilibrium, a radiation source is applied at the left boundary,  $\sigma \propto T^{-3}$ .
- Plot of transient solution for  $T_r = \sqrt[4]{\phi/ac}$  after 5 shakes, 200 x cells



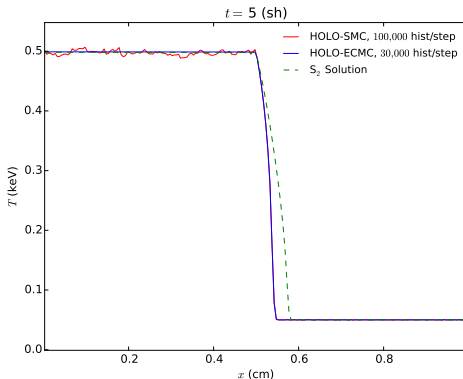
## Two Material Problem, Comparison of Spatial Convergence

- Same as Marshak Wave, but with constant opacities and an optically thin (left) and optically thick (right) region
- Convergence of spatial mesh:



# Comparison of statistical noise for standard and ECMC HO solvers

- Compare statistical accuracy of HO solvers
- One HO solve, with a *fixed number of histories* per time step. Comparison of **ECMC** with 3 batches and standard MC (**SMC**), as well as  $S_2$  solution





## Current & Future Development

- Can accurately reproduce IMC results with HOLO method
  - Pure absorber histories are more efficient than standard MC simulations
  - ECMC requires significantly less
  - LO solver determines non-linear Material temperature distribution
  - Linear shape within a cell mitigates teleportation error
  - Very efficient for diffusive problems
- Future work will be to implement in 2 spatial dimensions to demonstrate the elimination of ray effects

# Questions?

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## Backup Slides



## Algorithm

- 1 Initialize  $\tilde{I}^{n+1,(0)} := \tilde{I}^n$
- 2 Using Monte Carlo, solve  $\epsilon^{(m)} = \mathbf{L}^{-1} \tilde{r}^{(m)}$ 
  - Use volumetric tallies, weighted with  $x$  and  $\mu$  basis moments time  $\psi$  to construct LD  $\tilde{\epsilon}^{(m)}(x, \mu)$  over the current space-angle mesh
- 3  $\tilde{\psi}^{(j+1)} = \tilde{\psi}^{(m)} + \tilde{\epsilon}^{(m)}$
- 4 **IF** error stagnation:
  - Refine mesh based on relative jump error in  $\tilde{\psi}(x, \mu)$
- 5 Repeat 2-4 until  $\|\tilde{\epsilon}\|_2 < \text{tol} \times \|\psi\|_2$

## Algorithm

- ➊ Initialize  $\langle \mu \rangle^\pm$  parameters to  $S_2$
- ➋ Solve LO system using power iteration
- ➌ Build  $q^{LD}$  for HO solver, and set  $\tilde{\psi}$  to latest HO estimate on coarsest  $x-\mu$  mesh
- ➍ Solve  $\tilde{\psi}(x, \mu) = \mathbf{L}^{-1} q^{LO}$  using ECMC
- ➎ Compute new  $\langle \mu \rangle^\pm$  parameters using  $\tilde{\psi}^{HO}$  over LO mesh
- ➏ Repeat 2-5 until  $\underline{\phi}^{LO}$  is converged
  - Use adaptive convergence criteria

- Taking moments of TE yields 4 equations, per cell  $i$ , e.g.

$$\begin{aligned}
 & -2\mu_{i-1/2}^{n+1,+} \phi_{i-1/2}^{n+1,+} + \{\mu\}_{L,i}^{n+1,+} \langle \phi \rangle_{L,i}^{n+1,+} + \{\mu\}_{R,i}^{n+1,+} \langle \phi \rangle_{R,i}^{n+1,+} + \left( \sigma_t^{n+1} \right. \\
 & \left. - \frac{\sigma_s h_i}{2} \left( \langle \phi \rangle_{L,i}^{n+1,+} + \langle \phi \rangle_{L,i}^{n+1,-} \right) \right) = \frac{h_i}{2} \langle \sigma_a^{n+1} a_c T^{n+1,4} \rangle_{L,i} + \frac{h_i}{c \Delta t} \langle \phi \rangle_{L,i}^{n,+}, \quad (5)
 \end{aligned}$$

- Cell unknowns:  $\langle \phi \rangle_{L,i}^+$ ,  $\langle \phi \rangle_{R,i}^+$ ,  $\langle \phi \rangle_{L,i}^-$ ,  $\langle \phi \rangle_{R,i}^-$ ,  $T_L$ ,  $T_R$
- Need angular consistency terms and spatial closure (LD)

## Solving LO System with Newton's Method

- Linearization:  $\underline{B}(T^{n+1}) = \underline{B}(T^*) + (T^{n+1} - T^*) \left. \frac{\partial \underline{B}}{\partial t} \right|_{t^*}$
- Modified system

$$[\mathbf{D}(\mu^\pm) - \sigma_a^*(1 - f^*)] \underline{\Phi}^{n+1} = f^* \underline{B}(T^*) + \frac{\Phi^n}{c\Delta t}$$

$$\hat{\mathbf{D}} \underline{\Phi}^{n+1} = \underline{Q}$$

$$f = \left( 1 + \sigma_a^* c \Delta t \frac{4aT^{*3}}{\rho c_v} \right)^{-1} \quad T_i^* = \frac{T_{L,i}^* + T_{R,i}^*}{2}$$

- Equation for  $T^{n+1}$  based on linearization that is conservative
- Converge  $T^{n+1}$  and  $\langle \phi \rangle$  with Newton Iterations
  - Spatial representation can result in negative temperatures