A High-Order Low-Order Algorithm with Exponentially-Convergent Monte Carlo for Thermal Radiative Transfer

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MC 2015 - April 22



Introduction



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Outline

- Introduction
- 2 Low-Order Solver
- High-Order Solver
- 4 Algorithm
- **6** Computational Results
- 6 Conclusions



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Overview

Introduction

- We are interested in modeling thermal radiation transport in the high-energy density physics regime
 - Temperatures on order of 10^6 K or more
 - Significant energy and momentum may be exchanged with material



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Overview

- We are interested in modeling thermal radiation transport in the high-energy density physics regime
 - Temperatures on order of 10⁶ K or more
 - Significant energy and momentum may be exchanged with material
- Radiative transfer simulations important in modeling:
 - Material under extreme conditions
 - Inertial confinement fusion
 - Supernovae and other astrophysical phenomena.



The Grey TRT Equations

Introduction

• The 1D, grey equations

$$\frac{1}{c}\frac{\partial I}{\partial t} + \mu \frac{\partial I}{\partial x} + \sigma_t I(x, \mu) = \frac{1}{2}\sigma_a ac T^4(x),$$

$$C_v \frac{\partial T}{\partial t} = \sigma_a \left(\phi(x) - ac T^4\right)$$

$$\phi(x) = \int_{-1}^1 I(x, \mu) d\mu.$$

• Fundamental unknowns are the radiation intensity $I(x, \mu)$ and material temperature T(x)



Conclusions

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$$\phi(x) = \int_{-1}^1 I(x, \mu) d\mu.$$

- Fundamental unknowns are the radiation intensity $I(x,\mu)$ and material temperature T(x)
- Absorption cross section (σ_a) can be a strong function of T
- Equations are nonlinear and may be tightly coupled

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The Implicit Monte Carlo Method

Introduction

- Equations often solved with Monte Carlo (MC) via the Implicit Monte Carlo (IMC) method
- The material energy equation is linearized over a time step and eliminated from the system
 - Results in linear transport equation with a discrete emission and scattering terms
 - Continuous integration of time derivative



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- The material energy equation is linearized over a time step and eliminated from the system
 - Results in linear transport equation with a discrete emission and scattering terms
 - Continuous integration of time derivative
- Drawbacks

Introduction

- Effective scattering cross section can be very large
 - Local acceleration methods based on discrete diffusion approxmiations
- Nonlinearities not converged
- Reconstruction of linear source shape in cell

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An alternative High-Order Low-Order approach

Basic Idea

Introduction

Solve a fully non-linear, low-order (LO) system that preserves a high-order (HO) solution from efficent MC simulations, for each time step



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Solve a fully non-linear, low-order (LO) system that preserves a high-order (HO) solution from efficient MC simulations, for each time step

- The LO system consists of space-angle moment equations, formed over a fixed finite-element (FE) spatial mesh
 - Lower dimensionality in angle
 - Non-linear temperature dependence converged with Newton's method
 - Output: a linear-discontinuous (LD) FE representation of scattering and emission source

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 - Output: a linear-discontinuous (LD) FE representation of scattering and emission source
- The HO system is a fixed-source, pure-absorber transport problem
 - No effective scattering events
 - Solved with ECMC for efficient reduction of statisical noise
 - Output: angular consistency terms

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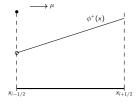
Outline

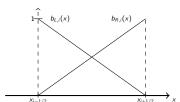
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LO Discretization & Space-Angle Moments

- The LO system is "S₂-like" equations (similar to [Wolters 2013]), with a backward Euler discretization in time
- Linear discontinuous (LD) FE in space for ϕ , T, and $T^4(x)$

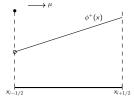


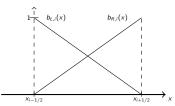


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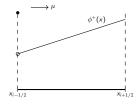


• Half-range integrals over + and $-\mu$

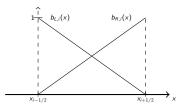
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Introduction



- Half-range integrals over + and $-\mu$
- Examples of moments:

$$\frac{\text{Spatial: left basis}}{\langle \cdot \rangle_{L,i} = \frac{2}{h_i} \int_{x_{i-1/2}}^{x_{i+1/2}} b_{L,i}(x) (\cdot) \mathrm{d}x} \qquad \frac{\text{Angular: positive flow}}{\phi^+(x) = \int_0^1 I(x,\mu) \mathrm{d}\mu}$$

Angular: positive flow

Conclusions

Applying moments to the TRT equations yields 4 radiation equations, and 2 material energy equations, per cell



Introduction

- Applying moments to the TRT equations yields 4 radiation equations, and 2 material energy equations, per cell
- ② The resulting radiation equations are manipulated to produce weighted averages, which we refer to as consistency terms, e.g.,

$$\{\mu\}_{L,i}^{+} := \frac{\int\limits_{0}^{1} \int\limits_{x_{i-1/2}}^{x_{i+1/2}} \mu \, b_{L,i}(x) I^{n+1}(x,\mu) \, \mathrm{d}x \mathrm{d}\mu}{\int\limits_{0}^{1} \int\limits_{x_{i-1/2}}^{x_{i+1/2}} b_{L,i}(x) I^{n+1}(x,\mu) \, \mathrm{d}x \mathrm{d}\mu}$$

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3 Use $\tilde{l}_{HO}^{n+1}(x,\mu)$ & LD spatial closure to close the system

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- **3** Use $\tilde{l}_{HO}^{n+1}(x,\mu)$ & LD spatial closure to close the system
 - Cell unknowns: $\langle \phi \rangle_{L,i}^{n+1,+}$, $\langle \phi \rangle_{R,i}^{n+1,+}$, $\langle \phi \rangle_{L,i}^{n+1,-}$, $\langle \phi \rangle_{R,i}^{n+1,-}$, $T_{L,i}$, $T_{R,i}$
- Non-linear, fully discrete system, solved with approximate Newton's method

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Overview of Exponentially Convergent Monte Carlo

Introduction

- Each batch tallies the error in current estimate of solution
 - Can reduce statistical error globally $\propto e^{-\alpha N}$
 - In TRT problems, old angular intensity provides a *very good* guess of new solution, significantly reducing the required number of histories



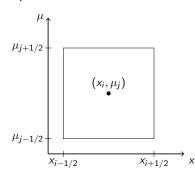
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Overview of Exponentially Convergent Monte Carlo

- Each batch tallies the error in current estimate of solution
 - Can reduce statistical error globally $\propto e^{-\alpha N}$
 - In TRT problems, old angular intensity provides a *very good* guess of new solution, significantly reducing the required number of histories
- Requires a discretized form of the angular intensity $\tilde{I}(x, \mu)$

Introduction

- Use <u>projection</u> of the solution onto a space-angle FE mesh
- LD FE $\tilde{I}^{HO}(x,\mu)$ allows for direct computation of consistency terms



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High Order System and ECMC Algorithm

Introduction

• Pure absorber transport problem because we have LD representation of T^4 from LO solution

$$\left[\mu \frac{\partial}{\partial x} + \left(\sigma_t + \frac{1}{c\Delta t}\right)\right] I^{n+1}(x,\mu) = \frac{I^n}{c\Delta t} + \left[\frac{1}{2}\sigma_a a c T_{LO}^{n+1,4}\right]$$

$$\mathbf{L}I^{n+1} = q_{LO}$$

High Order System and ECMC Algorithm

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For each batch:

Introduction

• Residual Equation: $\mathbf{L}\tilde{\epsilon}^{(m)} = \tilde{r}^{(m)} = q - \mathbf{L}\tilde{I}^{n+1,(m)}$

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- Residual Equation: $\mathbf{L}\tilde{\epsilon}^{(m)} = \tilde{r}^{(m)} = q \mathbf{L}\tilde{I}^{n+1,(m)}$
- Compute $\tilde{\epsilon}^{(m)} = \mathbf{L}^{-1}\tilde{r}^{(m)}$ with MC
 - Particles are allowed to stream along path s, $w(s) = w_0 e^{-\sigma_t s}$
 - Use cell-wise systematic sampling for $r^{(m)}$ source

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High Order System and ECMC Algorithm

• Pure absorber transport problem because we have LD representation of \mathcal{T}^4 from LO solution

$$\left[\mu \frac{\partial}{\partial x} + \left(\sigma_t + \frac{1}{c\Delta t}\right)\right] I^{n+1}(x,\mu) = \frac{I^n}{c\Delta t} + \left[\frac{1}{2}\sigma_a ac T_{LO}^{n+1,4}\right]$$

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For each batch:

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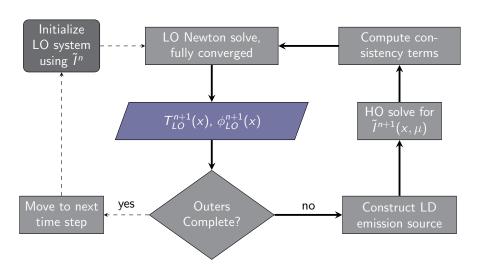
- Residual Equation: $L\tilde{\epsilon}^{(m)} = \tilde{r}^{(m)} = q L\tilde{I}^{n+1,(m)}$
- Compute $\tilde{\epsilon}^{(m)} = \mathbf{L}^{-1}\tilde{r}^{(m)}$ with MC
 - Particles are allowed to stream along path s, $w(s) = w_0 e^{-\sigma_t s}$
 - Use cell-wise systematic sampling for $r^{(m)}$ source
- Update: $\tilde{I}^{n+1,(m)} = \tilde{I}^{n+1,(m-1)} + \tilde{\epsilon}^{(m)}$
 - ullet If $ilde{\epsilon}$ is reduced each batch, exponential convergence achieved

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High-Order Low-Order Algorithm



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Relevant Implementation Specifics

Introduction

- Difficulties in resolving the solution near the wave-front
 - The LD representation for I results in negativite values within a cell
 - In these results, no correction is applied to the HO solution, and the LO solution uses lumped LD and S₂ equivalent terms in bad cells

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Relevant Implementation Specifics

- Difficulties in resolving the solution near the wave-front
 - The LD representation for I results in negativite values within a cell
 - In these results, no correction is applied to the HO solution, and the LO solution uses lumped LD and S₂ equivalent terms in bad cells
- For all results

Introduction

- LO Newton iterations are is fully converged each solve
- Fixed Δt of 0.001 sh
- One HO solve per time step

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Marshak Wave Test Problem

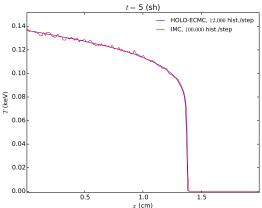
• From equilibrium, a radiation source is applied at the left boundary at t=0. With $\sigma_a \propto T^{-3}$, energy slowly moves across the system.



Marshak Wave Test Problem

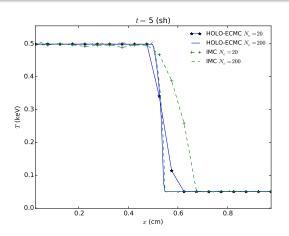
Introduction

- From equilibrium, a radiation source is applied at the left boundary at t=0. With $\sigma_a \propto T^{-3}$, energy slowly moves across the system.
- Transient solution after 5 shakes plotted as $T_r = \sqrt[4]{\phi/ac}$, with 200 x cells



Two Material Problem, Comparison of Spatial Convergence

 Problem the same as Marshak Wave, but with constant opacities and an optically thin (left) and optically thick (right) region

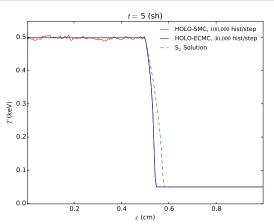


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Comparison of statistical noise for standard and ECMC HO solvers

Introduction

 One HO solve, with a fixed number of histories per time step, for two different HO solvers: a comparison of ECMC with 3 batches and standard MC (SMC), as well as S₂ solution



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Current & Future Development

Introduction

- Can accurately reproduce IMC results with HOLO method
 - ECMC requires significantly less particles
 - LO solver determines non-linear Material temperature distribution
 - Linear shape within a cell mitigates teleportation error
 - Very efficient for diffusive problems



Current & Future Development

- Can accurately reproduce IMC results with HOLO method
 - ECMC requires significantly less particles
 - LO solver determines non-linear Material temperature distribution
 - Linear shape within a cell mitigates teleportation error
 - Very efficient for diffusive problems
- Currently developing strategies for dealing with unresolvable solutions
- Next step will be to implement in 2 spatial dimensions

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Questions?

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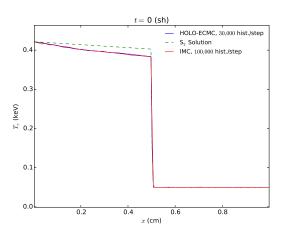




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Two Material Problem, comparison in optically thin region

• Plot of radiation temperature after 10 time steps



ECMC procedure

Algorithm

- Intialize $\tilde{I}^{n+1,(0)} := \tilde{I}^n$
- **2** Using Monte Carlo, solve $\epsilon^{(m)} = \mathbf{L}^{-1} \tilde{r}^{(m)}$
 - Use volumetric tallies, weighted with x and μ basis moments time ψ to construct LD $\tilde{\epsilon}^{(m)}(x,\mu)$ over the current space-angle mesh
- IF error stagnation:
 - ullet Refine mesh based on relative jump error in $ilde{\psi}(x,\mu)$
- **5** Repeat 2-4 until $\|\tilde{\epsilon}\|_2 < \text{tol} \times \|\psi\|_2$

HOLO Algorithm

Algorithm

- **1** Initialize $\langle \mu \rangle^{\pm}$ parameters to S_2
- Solve LO system using power iteration
- $\ \ \$ Build ${\it q}^{LD}$ for HO solver, and set $\tilde{\psi}$ to latest HO estimate on coarsest $x{-}\mu$ mesh
- Solve $\tilde{\psi}(\mathbf{x}, \mu) = \mathbf{L}^{-1} q^{LO}$ using ECMC
- **5** Compute new $\langle \mu \rangle^{\pm}$ parameters using $\tilde{\psi}^{HO}$ over LO mesh
- **6** Repeat 2-5 until Φ^{LO} is converged
 - Use adaptive convergence criteria

LO System

• Taking moments of TE yields 4 equations, per cell i, e.g.

$$-2\mu_{i-1/2}^{n+1,+}\phi_{i-1/2}^{n+1,+} + \{\mu\}_{L,i}^{n+1,+} \langle\phi\rangle_{L,i}^{n+1,+} + \{\mu\}_{R,i}^{n+1,+} \langle\phi\rangle_{R,i}^{n+1,+} + \left(\sigma_{t}^{n+1} + \frac{1}{c\Delta t}\right) h_{i}\langle\phi\rangle_{L,i}^{n+1,+} - \frac{\sigma_{s}h_{i}}{2} \left(\langle\phi\rangle_{L,i}^{n+1,+} + \langle\phi\rangle_{L,i}^{n+1,-}\right)$$

$$= \frac{h_{i}}{2} \langle\sigma_{a}^{n+1} acT^{n+1,4}\rangle_{L,i} + \frac{h_{i}}{c\Delta t}\langle\phi\rangle_{L,i}^{n,+}, \quad (1)$$

- Cell unknowns: $\langle \phi \rangle_{L,i}^+$, $\langle \phi \rangle_{R,i}^+$, $\langle \phi \rangle_{L,i}^-$, $\langle \phi \rangle_{R,i}^-$, T_L , T_R
- Need angular consistency terms and spatial closure (LD)

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Backup Slides

Solving LO System with Newton's Method

- Linearization: $\underline{B}(T^{n+1}) = \underline{B}(T^*) + (T^{n+1} T^*) \frac{\partial \underline{B}}{\partial t} \Big|_{t^*}$
- Modified system

$$\left[\mathbf{D}(\mu^{\pm}) - \sigma_a^*(1 - f^*)\right] \underline{\Phi}^{n+1} = f^* \underline{B}(T^*) + \frac{\underline{\Phi}^n}{c\Delta t}$$
$$\hat{\mathbf{D}}\Phi^{n+1} = \underline{Q}$$

$$f = \left(1 + \sigma_a^* c \Delta t \frac{4aT^{*3}}{\rho c_v}\right)^{-1} \qquad T_i^* = \frac{T_{L,i}^* + T_{R,i}^*}{2}$$

- Equation for T^{n+1} based on linearization that is conservative
- Converge T^{n+1} and $\langle \phi \rangle$ with Newton Iterations
 - Spatial representation can result in negative temperatures

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