# A High-Order Low-Order Algorithm with Exponentially-Convergent Monte Carlo for Thermal Radiative Transfer

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Introduction



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- Introduction
- 2 Low-Order Solver
- High-Order Solver
- 4 Algorithm
- **6** Computational Results
- 6 Conclusions



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#### Overview

Introduction

- We are interested in modeling thermal radiation transport in the high-energy density physics regime
  - Temperatures on order of  $10^6$  K or more
  - Significant energy and momentum may be exchanged with material



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#### Overview

- We are interested in modeling thermal radiation transport in the high-energy density physics regime
  - Temperatures on order of 10<sup>6</sup> K or more
  - Significant energy and momentum may be exchanged with material
- Radiative transfer simulations important in modeling:
  - Material under extreme conditions
  - Inertial confinement fusion
  - Supernovae and other astrophysical phenomena.



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## The Grey TRT Equations

Introduction

• The 1D, grey equations

$$\frac{1}{c}\frac{\partial I}{\partial t} + \mu \frac{\partial I}{\partial x} + \sigma_t I(x, \mu) = \frac{1}{2}\sigma_a ac T^4(x),$$

$$C_v \frac{\partial T}{\partial t} = \sigma_a \left(\phi(x) - ac T^4\right)$$

$$\phi(x) = \int_{-1}^1 I(x, \mu) d\mu.$$

• Fundamental unknowns are the radiation intensity  $I(x, \mu)$  and material temperature T(x)



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- Fundamental unknowns are the radiation intensity  $I(x, \mu)$  and material temperature T(x)
- Absorption cross section  $(\sigma_a)$  can be a strong function of T
- Equations are nonlinear and may be tightly coupled

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#### The Implicit Monte Carlo Method

Introduction

- Equations often solved with Monte Carlo (MC) via the Implicit Monte Carlo (IMC) method
- The material energy equation is linearized over a time step and eliminated from the system
  - Results in linear transport equation with a discrete emission and scattering terms
  - Continuous integration of time derivative



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- Drawbacks

Introduction

- Effective scattering cross section can be very large
  - Local acceleration methods based on discrete diffusion approxmiations
- Nonlinearities not converged
- Reconstruction of linear source shape in cell

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### An alternative High-Order Low-Order approach

#### Basic Idea

Introduction

Solve a fully non-linear, low-order (LO) system, that preserves a high-order (HO) solution from efficent MC simulations

• The LO system consists of space-angle moment equations, formed over a fixed finite-element (FE) spatial mesh

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  - Lower dimensionality in angle
  - Non-linear temperature dependence converged with Newton's method
  - Output: a linear-discontinuous (LD) FE representation of scattering and emission source

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- The HO system is a fixed-source, pure-absorber transport problem
  - No effective scattering events
  - Solved with FCMC for efficient reduction of statistical noise
  - Output: angular consistency terms

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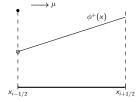
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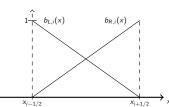
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### LO Discretization & Space-Angle Moments

- The LO system is "S<sub>2</sub>-like" equations (similar to [Wolters 2013]), with a backward Euler discretization in time
- Linear discontinuous (LD) FE in space for  $\phi$ , T, and  $T^4(x)$



Introduction

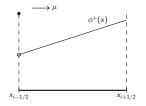


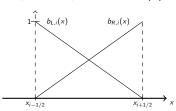
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- Half-range integrals over + and  $-\mu$
- Examples of moments:

Spatial: left basis
$$\langle \cdot \rangle_{L,i} = \frac{2}{h_i} \int_{x_{i-1/2}}^{x_{i+1/2}} b_{L,i}(x)(\cdot) dx$$
Angular: positive flow
$$\phi^+(x) = \int_0^1 I(x,\mu) d\mu$$

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#### Forming the LO System

Introduction

- Applying moments to the TRT equations yields 4 radiation equations, and 2 material energy equations, per cell
- The resulting radiation equations are manipulated to produce weighted averages, which we refer to as consistency terms, e.g.,

$$\{\mu\}_{L,i}^{+} := \frac{\int\limits_{0}^{1} \int\limits_{x_{i-1/2}}^{x_{i+1/2}} \mu \, b_{L,i}(x) I^{n+1}(x,\mu) \, \mathrm{d}\mu \mathrm{d}x}{\int\limits_{0}^{1} \int\limits_{x_{i-1/2}}^{x_{i+1/2}} b_{L,i}(x) I^{n+1}(x,\mu) \, \mathrm{d}\mu \mathrm{d}x}$$

**1** Use HO  $\tilde{I}^{n+1}(x,\mu)$  and the LD spatial closure to close the system

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- **③** Use HO  $\tilde{I}^{n+1}(x,\mu)$  and the LD spatial closure to close the system
- Cell unknowns:  $\langle \phi \rangle_{L,i}^{n+1,+}$ ,  $\langle \phi \rangle_{R,i}^{n+1,+}$ ,  $\langle \phi \rangle_{L,i}^{n+1,-}$ ,  $\langle \phi \rangle_{R,i}^{n+1,-}$ ,  $T_{L,i}$ ,  $T_{R,i}$

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- Non-linear, fully discrete system, solved with approximate Newton's method

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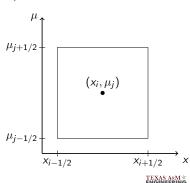


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#### Overview of Exponentially Convergent Monte Carlo

 Each batch tallies the error in current estimate of solution, which is a transport problem with a reduced source

- Can reduce statistical error globally  $\propto e^{-\alpha N}$
- In TRT problems, old angular intensity provides a very good guess of new solution, significantly reducing the required number of histories
- Requires a discretized form of the angular intensity  $\tilde{I}(x,\mu)$ 
  - Use projection of the solution onto a space-angle FE mesh
  - LD FE  $\tilde{I}^{HO}(x,\mu)$  allows for direct computation of consistency terms



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Conclusions

### High Order System and ECMC Algorithm

• Pure absorber transport problem because we have LD representation of  $\mathcal{T}^4$  from LO solution

$$\left[\mu \frac{\partial}{\partial x} + \left(\sigma_t + \frac{1}{c\Delta t}\right)\right] I^{n+1}(x,\mu) = \boxed{\frac{I^n}{c\Delta t} + \frac{1}{2}\sigma_a a c T_{LO}^{n+1,4}}$$
$$\mathbf{L} I^{n+1} = q_{LO}$$

#### For each batch:

Introduction

- Residual Equation:  $L\tilde{\epsilon}^{(m)} = \tilde{r}^{(m)} = q L\tilde{I}^{n+1,(m)}$
- Compute  $\tilde{\epsilon}^{(m)} = \mathbf{L}^{-1}\tilde{r}^{(m)}$  with MC
  - Particles are allowed to stream along path s,  $w(s) = w_0 e^{-\sigma_t s}$
  - Cell-wise, global representation allows for easy stratified sampling
- Update:  $\tilde{I}^{n+1,(m)} = \tilde{I}^{n+1,(m-1)} + \tilde{\epsilon}^{(m)}$ 
  - ullet If  $ilde{\epsilon}$  is reduced each batch, exponential convergence achieved

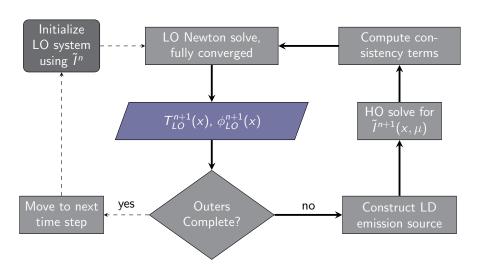
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### High-Order Low-Order Algorithm



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### Relevant Implementation Specifics

Introduction

- LO Newton iterations are is fully converged each solve
- Difficulties in resolving the solution near the wave-front
  - The LD representation for I results in negativite values within a cell
  - In these results, no correction is applied to the HO solution, and the LO solution uses lumped LD and S<sub>2</sub> equivalent terms in bad cells
- For all results, one HO solve per time step



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#### Marshak Wave Test Problem

• From equilibrium, a radiation source is applied at the left boundary at t=0. With  $\sigma_a \propto T^{-3}$ , energy slowly moves across the system.

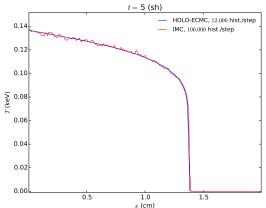


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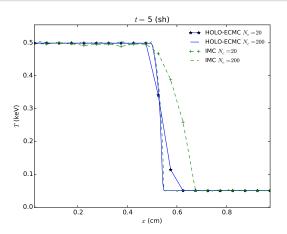
- From equilibrium, a radiation source is applied at the left boundary at t=0. With  $\sigma_a \propto T^{-3}$ , energy slowly moves across the system.
- Transient solution after 5 shakes plotted as  $T_r = \sqrt[4]{\phi/ac}$ , with 200 x cells



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## Two Material Problem, Comparison of Spatial Convergence

• Problem the same as Marshak Wave, but with constant opacities and an optically thin (left) and optically thick (right) region



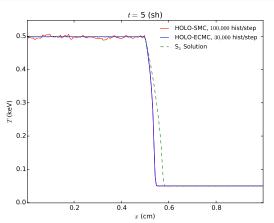
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### Comparison of statistical noise for standard and ECMC HO solvers

Introduction

 One HO solve, with a fixed number of histories per time step, for two different HO solvers: a comparison of ECMC with 3 batches and standard MC (SMC), as well as S<sub>2</sub> solution



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#### Current & Future Development

- Can accurately reproduce IMC results with HOLO method
  - ECMC requires significantly less particles
  - LO solver determines non-linear Material temperature distribution
  - Linear shape within a cell mitigates teleportation error
  - Very efficient for diffusive problems
- Currently developing strategies for dealing with unresolvable solutions
- Next step will be to implement in 2 spatial dimensions

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# Questions?

# A High-Order Low-Order Algorithm with Exponentially-Convergent Monte Carlo for Thermal Radiative Transfer

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