

The place of the local group in the cosmic web

Sebastian Bustamante ^{*1} Jaime E. Forero-Romero ²

¹*Instituto de Física - FCEN, Universidad de Antioquia, Calle 67 No. 53-108, Medellín, Colombia*

²*Departamento de Física, Universidad de los Andes, Cra. 1 No. 18A-10, Edificio Ip, Bogotá, Colombia*

13 October 2013

ABSTRACT

We present here a study about the influence of the environment on the local group (LG) of galaxies in the context of Λ CDM. In this study we use a large volume high resolution N-body cosmological simulation (Bolshoi) together with the most recent methods to quantify the cosmic web (T-web, V-web schemes); furthermore we propose a novel approximation, base upon the minimization of the mean density of void regions, to determinate the optimum threshold value λ_{th} , which have been treated until now as a free parameter. Following the recent work of Courtois et al. (2013), where was found that the LG is located near to a large void, we also do an extensive study of voids, applying a FOF algorithm to find void regions and performing an analysis of their shape based upon the reduced inertia tensor. Using the recent observations that constrained the tangential velocity of M31 with respect to the Milky Way (MW), the previously established radial velocity, the estimated masses of dark halos, along with some criteria to guarantee the gravitational isolation of these systems (Forero-Romero et al. 2013-1), we select a set of halos pairs as a representative sample of LG systems in the Bolshoi simulation. We look for possible bias and correlations between the environment properties of each LG system and its kinematic and formation properties. Among our main result we find [\[summarize our results here!\]](#).

Key words: Cosmology: large-scale Structure of Universe, galaxies: star formation - line: formation

1 INTRODUCTION

The spatial distribution of galaxies describes a web-like pattern, the so-called cosmic web. Today it is understood that such configuration is driven by gravitational instabilities. ...

The study of the influence of the cosmic web on galaxy properties start with the seminal work of Dressler [\[reference here\]](#) and extends to recent works using large observational surveys that look for signatures of the web into the evolution of galaxy populations. With the advent of more detailed observations and sophisticated computational models it is now within our reach to understand what physical processes dominate.

This makes that the mass assembly history of a galaxy is deeply connected with its position in the cosmic web. There is an extensive body of literature on the effects of the web environment on the observable properties of galaxies.

This environmental study is also of paramount importance to understand the formation of our Galaxy. In our local neighborhood, the observations of dwarf galaxies around

the Milky Way (MW) and the Andromeda galaxy (M31) show filamentary and disk-like patterns that can be linked to a preferential infall direction, very likely connected with the cosmic web where the Local Group (LG) of galaxies is embedded.

In this paper we quantify the velocity shear environment of DM halo pairs representative of the principal members of the Local Group (LG), the Milky Way (MW) and Andromeda galaxy (M31). We perform this study in an unconstrained cosmological simulation from random phases in the initial conditions, and unlike previous works, where were used constrained cosmological simulations which have been setup as to reproduce the large scale structure of the local universe, we use directly observational measurements of the kinematics properties of the local group [\[Reference here\]](#) in order to build faithful samples of LG-like systems.

We pay special attention to the correlation of the present velocity shear environment with the assembly and the kinematics properties of the pairs. The motivation to have that focus is that it has been previously shown that the LG present in three different realizations of the constrained simulations have assembly histories biased towards

* sbustama@pegasus.udea.edu.co

early formation times and absence of major mergers (ratio 1:10) in the last 10 Gyr. In the case of the kinematic properties, recent observational constraints to the galactocentric tangential velocity of M31 has enabled to establish how typical is the LG in a cosmological context [reference to Forero-Romero et.al 2013-1], that is why we focus here how a specific kind of host environment biases these kinematics properties.

2 THE SIMULATION

As it was previously mentioned, we use an unconstrained cosmological simulation, the Bolshoi simulation, to identify the possible large scale environment of the Local Group. This is a similar approach to the one already used by [reference here].

The Bolshoi simulation follows the non-linear evolution of a dark matter density field on a cubic volume of size $250h^{-1}\text{Mpc}$ sampled with 2048^3 particles. The cosmological parameters in the simulation are $\Omega_m = 0.27$, $\Omega_\Lambda = 0.73$, $h = 0.70$, $n = 0.95$ and $\sigma_8 = 0.82$ for the matter density, cosmological constant, dimensionless Hubble parameter, spectral index of primordial density perturbations and normalization for the power spectrum. The mass of each particle in the simulation is $m_p = 1.4 \times 10^8 h^{-1} \text{M}_\odot$.

2.1 Halos and Merger Trees

We identify halos with two algorithms, the Friends-of-Friends [reference here] algorithm and the Bound Density Maximum algorithm. The constructed catalogues also provide the basis for the mass aggregation history studies. We also include in the catalogues information about the substructure.

All the results presented here must be interpreted in term of host halos, without any information of the substructure. In particular the merger of two FOF halos corresponds to the epoch of first overlap, and not to the fusion and/or disruption of an accreted sub-halo with a dominant halo.

The linking length is $b = 0.17$ times the mean interparticle separation. All objects with 20 particles or more are considered a bona fide halo and are included in the construction of the merger tree, this corresponds to a minimum halo mass of $M_{\min} = 2.70 \times 10^9 h^{-1} \text{M}_\odot$ in the Bolshoi simulation.

The halo identification for the simulation was done for XX snapshots in the redshift range $0 < z < 7$ more or less evenly spaced in look-back time.

3 ALGORITHMS TO QUANTIFY THE COSMIC WEB

3.1 The tidal web (T-web)

The first algorithm we use to identify the cosmic web is based upon the diagonalization of the tidal tensor, defined as the Hessian of a normalized gravitational potential

$$T_{\alpha\beta} = \frac{\partial^2 \phi}{\partial x_\alpha \partial x_\beta} \quad (1)$$

where the physical gravitational potential has been rescaled by a factor $4\pi G\bar{\rho}$ in such a way that ϕ satisfies the following equation

$$\nabla^2 \phi = \delta, \quad (2)$$

where $\bar{\rho}$ is the average density in the Universe, G is the gravitational constant and δ is the dimensionless matter overdensity.

3.2 The velocity web (V-web)

We also use a kinematical method to define the cosmic-web environment in the simulation. The method has been thoroughly described in XXX and applied to study the shape and spin alignment in the Bolshoi simulation here XX. We refer the reader to these papers to find a detailed description of the algorithm, its limitations and capabilities. Here we summarize the most relevant points for the discussion.

The V-web method for environment finding is based on the local shear tensor calculated from the smoothed DM velocity field in the simulation. The central quantity is the following dimensionless quantity

$$\Sigma_{\alpha\beta} = -\frac{1}{2H_0} \left(\frac{\partial v_\alpha}{\partial x_\beta} + \frac{\partial v_\beta}{\partial x_\alpha} \right) \quad (3)$$

where v_α and x_α represent the α component of the comoving velocity and position, respectively. $\Sigma_{\alpha\beta}$ can be represented by a 3×3 symmetric matrix with real values, that ensures that is possible to diagonalize and obtain three real eigenvalues $\lambda_1 > \lambda_2 > \lambda_3$ whose sum (the trace of $\Sigma_{\alpha\beta}$) is proportional to the divergence of the local velocity field smoothed on the physical scale \mathcal{R} .

The relative strength of the three eigenvalues with respect to a threshold value λ_{th} allows for the local classification of the matter distribution into four web types: voids, sheets, filaments and peaks, which correspond to regions with 3, 2, 1 or 0 eigenvalues with values larger than λ_{th} . Below we shall discuss a novel approach to define an adequate threshold value based on the visual impression of void regions, furthermore we study other possible values based on other visual features of the cosmic web.

3.3 The cosmic web in Bolshoi

Both established schemes to quantify the cosmic web depend on continuous and smooth physical quantities, as the peculiar velocity field and the density field. In order to calculate the necessary tensors, a discretization of the simulation volume is performed, so all the properties are reduced to single values associated to discrete cells. According to this, we divide the overall volume into $(256)^3$ cells, so each cell has an associated comoving cubic volume of 0.98 Mpc h^{-1} . Finally, to reduce possible effects due to the discretization process, a gaussian softening is performed between neighbour cells.

Once defined the numerical details about the classification schemes, we shall analyse the dependence on the threshold value λ_{th} for each one. In the figure 1 we show the variation of the mean density parameter δ with the threshold value for cells marked in each of the adopted types of environment.

As was previously established by [Hoffman et al. \(2012\)](#) and as can be seen in the figure 1, the behaviour of the V-web scheme is significantly more sensible to variations of the λ_{th} value compared with the T-web scheme; nevertheless, the behaviour of the mean density parameter for voids, sheets and filaments, are qualitatively quite similar for the two schemes, reaching extreme values in the marked λ_{th} critical values respectively. For instance, voids reach a minimum value at some critical threshold, increasing for higher threshold values, while sheets and filaments reach a maximum value at other critical threshold, decreasing for higher threshold values. Although we shall focus our analysis on voids because they are completely dominant in the visual impression of the cosmic web, an analogous analysis might be performed for other type of environments.

On cosmic scales, the presence of highly non-linear structures implies the existence of very vast regions with density lower than the mean cosmological value due to the mass conservation. That is why the visual impression of the cosmic web must be necessarily dominated by these under-dense regions. Keeping that in mind, our novel proposal is based upon the correct quantification of these regions, so the optimum threshold value must be chosen such that: sheet regions do not invade real void regions (in such case, the mean density parameter of sheet regions would become negative) and void regions do not invade real sheet and filament regions (in such case, the mean density parameter of void regions would increase due to the contribution of over-dense regions). Thus, the optimum value is simply where the mean density parameter of void regions is minimized. According to this, we obtain the next threshold values for the T-web and the V-web respectively, $\lambda_{th}^T = 0.61$ and $\lambda_{th}^V = 0.26$. To verify our analysis, we show in the figure 2 the visual impression for each defined critical value, and as can be seen, the chosen values reproduce properly the expected impression according to the density field.

Our classification scheme may be thought as a refinement of the recent schemes, where the threshold value is used to taking as a free parameter, based on the classic methods, where the classification is performed based on a cut off of the density field directly [\[references here\]](#). So we make use of the objectivity achieved by the analysis of the mean density, but keeping all the environmental information provided by the tensorial schemes instead of the poor description provided by the density field.

3.4 Method to find void regions

Following the recent work of [Courtois et al. 2013](#), we use a method based on a FOF algorithm to find extended regions of voids in order to select halo systems according to the proximity to those regions. To achieve this, we build the input catalogue of the FOF method with the positions of the center coordinate of every cell marked as void according to the web scheme adopted; furthermore we set an adequate linking length to connect even diagonal neighbour cells.

Following the work of [Forero-Romero et al. 2008](#), we also perform a percolation analysis in order to select the best threshold parameter that reduce percolation in cells, thereby accounting for physical void regions. In the figure 3 we show the obtained result of our percolation analysis for both web schemes. In both cases it can be noted the volume of the

largest void region is minimized and the volume distribution of voids is relatively flat at $\lambda_{th} = 0.0$, that means percolation is completely reduced for this threshold value. So despite of the previously established λ_{th} optimum values for each scheme, we shall use $\lambda_{th} = 0.0$ just for the detection of void regions. Besides, due to the domination of the large scale visual impression by voids, it is inevitable the presence of percolation phenomenon, so the current chosen threshold value for percolation is justified because in spite of voids are necessarily connected, we are just interested in detecting bulk regions.

Next, we shall calculate the reduced inertia tensor of each void region in order to determinate their principal directions of inertia and analyse the size-shape distribution of voids.

$$\tau_{ij} = \sum_l \frac{x_{l,i}x_{l,j}}{R_l^2} \quad (4)$$

where l is a index associated to each cell of the current region, i and j indexes run over each spatial direction and finally R_l is defined as $R_l^2 = x_{l,1}^2 + x_{l,2}^2 + x_{l,3}^2$, all positions are measured from the respective center of mass of the region.

The eigenvalues of the reduced inertia tensor, i.e. the principal moments of inertia, are used to quantify the shape of void regions. They are denoted as τ_1 , τ_2 and τ_3 such that $\tau_1 \leq \tau_2 \leq \tau_3$. In the figure 4 we show the computed distributions for τ_1/τ_2 and τ_2/τ_3 , where we rather calculate histograms for these ratio quantities instead of each value in order to avoid using an arbitrary normalization. For both schemes, it can be noticed that the shape-distribution is completely spread out, thereby indicating a non-preferred geometry of void regions, which is in agreement with the well-established anisotropic flow of matter associated to this type of region [\[reference here\]](#). Because of that, we shall look for possible alignments between the plane of rotation of halo pairs and the principal directions of inertia of the nearest void regions.

4 DEFINING SAMPLES

In order to establish an adequate set of criteria to define a LG-like sample in unconstrained simulations, we proceed from the the general dark halo catalogues, constructed using the FOF scheme with a linking length of $b = 0.17$ and the Bound Density Maxima (BDM) scheme for halos with an associated overdensity value of 200 times the critical density ([\[references here\]](#)). These samples are defined over all the Bolshoi simulation and will be referred as General Halos samples, GH_{FOF} and GH_{BDM} respectively¹. To be consistent with the mass range observationally determined of halos that host disk-form galaxies, we select halos in the mass range $7 \times 10^{11} < M_h/h^{-1}M_{\odot} < 7 \times 10^{12}$, referred here as Individual Halos (IH) samples.

As a primal approach to define gravitational bounded halo pairs we select halos within the IH sample that satisfied to be the closest to each other. They constitute the Pairs (P) sample. To maintain the consistency with previous

¹ For each sample defined, there will be a FOF sample and a BDM sample in order to have robust results

Sample	FOF	BDM
General Halos (<i>GH</i>)	432000	404432
Individual Halos (<i>IH</i>)	65109	61165
Pairs (<i>P</i>)	17115	15981
Isolated Pairs (<i>IP</i>)	1249	1647
Reduced Isolated Pairs (<i>RIP</i>)	138	116

Table 1. Size of each sample defined in the Bolshoi simulation and for each of the two schemes used to detect halos (FOF and BDM).

works regarding algorithms used for LG selection ([\[references here\]](#)), we define here the next list of conditions that have to fulfil a pair system in order to belong to the Isolated Pairs (*IP*) sample. All these considerations are based upon the relative dynamics of the Milky Way and M31, and its isolation from massive structures:

- (i) With respect to each halo, there cannot be any other halo within the mass range $7 \times 10^{11} < M_h/h^{-1}M_\odot < 7 \times 10^{12}$ closer than its partner. It means that there cannot be ambiguity on the identity of the pair members.
- (ii) The relative radial velocity between the two halos must be negative.
- (iii) The distance between the center of the two halos must be less than $1.0h^{-1}\text{Mpc}$.
- (iv) The distance to any halo more massive than any of the pair members cannot be less than $3h^{-1}\text{Mpc}$.
- (v) The distance to cluster-like halos with masses larger than $7 \times 10^{13} h^{-1}M_\odot$ must be larger than $4h^{-1}\text{Mpc}$.

Following the previous criteria established by Forero-Romero et al. 2013-1, we define a more reduced sample, the Reduced Isolated Pairs sample (*RIP*), as pairs that fulfil the next two extra criteria :

- (i) The separation between the center of mass of the halos is in the range $700 - 800\text{kpc}$.
- (ii) The total mass of the two halos is in the range $1 - 4 \times 10^{12}/h^{-1}M_\odot$.

Finally in the table 1 we show the sizes of all defined samples.

5 FINDING THE LOCAL GROUP ENVIRONMENT

In this section we prepare some numerical experiments with the halos samples and their environment, with special emphasis on the isolated pairs and LG samples. All this in order to find for environment correlations and common properties between LG systems.

5.1 Defining the LG environment

With the aim of constructing a sample of LG systems in Bolshoi simulations, we use the eigenvalues of the three LG systems in each CLUES simulation and choose a range for each one based in the extreme values of the six halos, this with the expectation of reproducing the LG specific conditions in constrained simulations. This criteria can be thought as a

first approximation to establish a more faithful sub-sample into the isolated pairs. In table ?? we show the used extreme values for each eigenvalue of local shear tensor.

As a self-consistency test, we apply this criteria to CLUES simulations and find, on average, three LG systems in each one. But to avoid confusion, we keep defining the LG sample as the three initial pairs. Finally, we construct the LG sample of Bolshoi, where the sample size is illustrated in Table ??.

Once we have established the equivalence of samples in each simulation and have defined the LG samples, we proceed to calculate correlations between halo samples and their environment. As was mentioned in section 3.2, we do not use an specific value of eigenvalue threshold λ_{th} , instead of this, we explore a relatively wide range of this parameter (i.e. $0 \leq \lambda_{th} \leq 1$) and calculate distributions respect to each eigenvalue individually.

At first place, we calculate the environment of each LG system in CLUES simulations, for this we use the λ_{th} scheme to classify it in void, sheet, filament or knot (see 3.2), with λ_{th} into the threshold range. Figure ?? is obtained.

At first place, we calculate the integrated distribution of each eigenvalue for individual halos, isolated pairs and LG samples. For this, we associate to each halo a value of environment (set of eigenvalues) according to its center of mass position in a smoothed grid (256^3 cells for Bolshoi and 64^3 for CLUES, or equivalently a resolution of $1.0\text{Mpc } h^{-1}$ per cell, according to the physical size of the pairs.)

6 RESULTS

6.1 Bias induced on kinematic and dynamics properties

... Total Mass
 ... Radial vs. tangential velocities
 ... Angular Momentum
 ... Total Mechanical energy.
 ... Reduced Spin

6.2 Bias induced on the Mass Assembly Histories

... Last major merger. Formation Time. Assembly Time.

6.3 Pair Alignment with the Cosmic Web

... Separation.
 ... Relative velocity.
 ... Angular Momentum.

7 CONCLUSIONS

ACKNOWLEDGMENTS

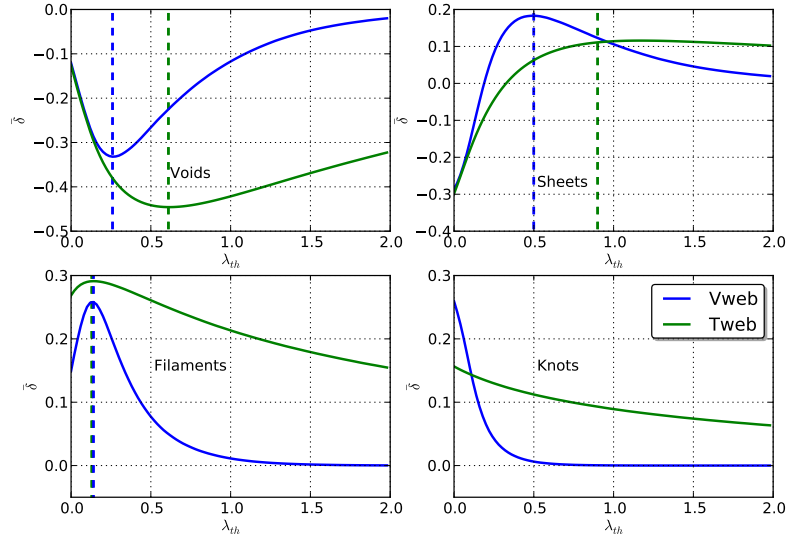


Figure 1. Mean density parameter for each of the defined environments according to the chosen λ_{th} value and for both classification schemes.

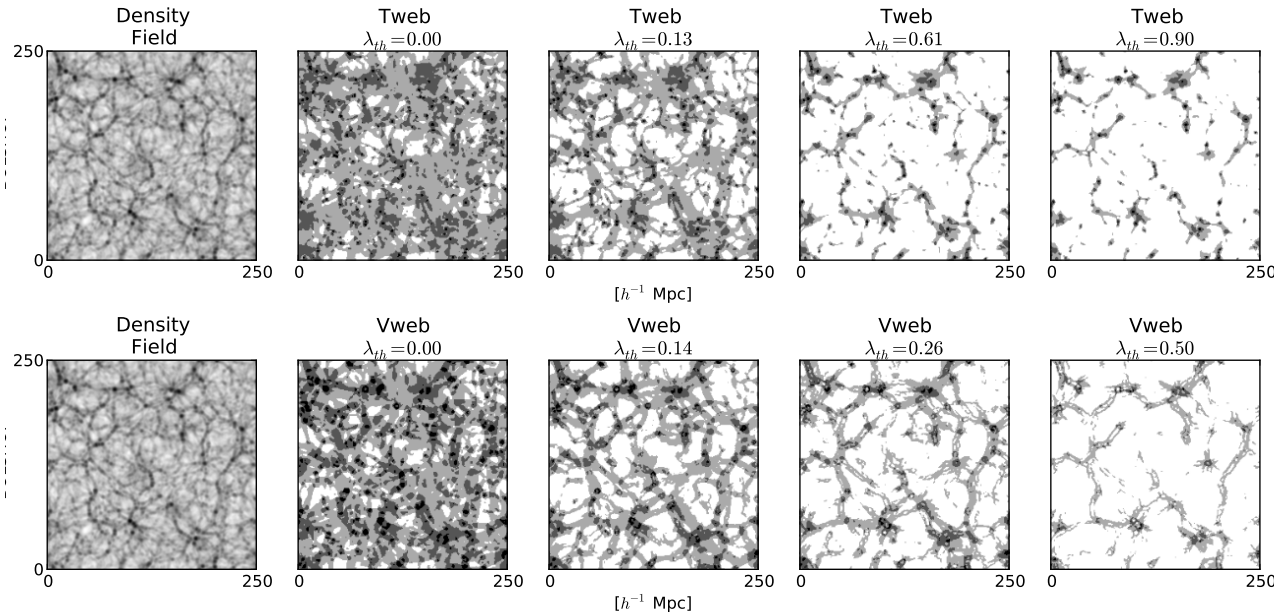


Figure 2. Visual impression of the density field (left), and of each classification scheme with the λ_{th} values obtained by our criteria (others). Our color convention for each environment is (white) - void, (light gray) - sheet, (gray) - filament, (black) - knot.

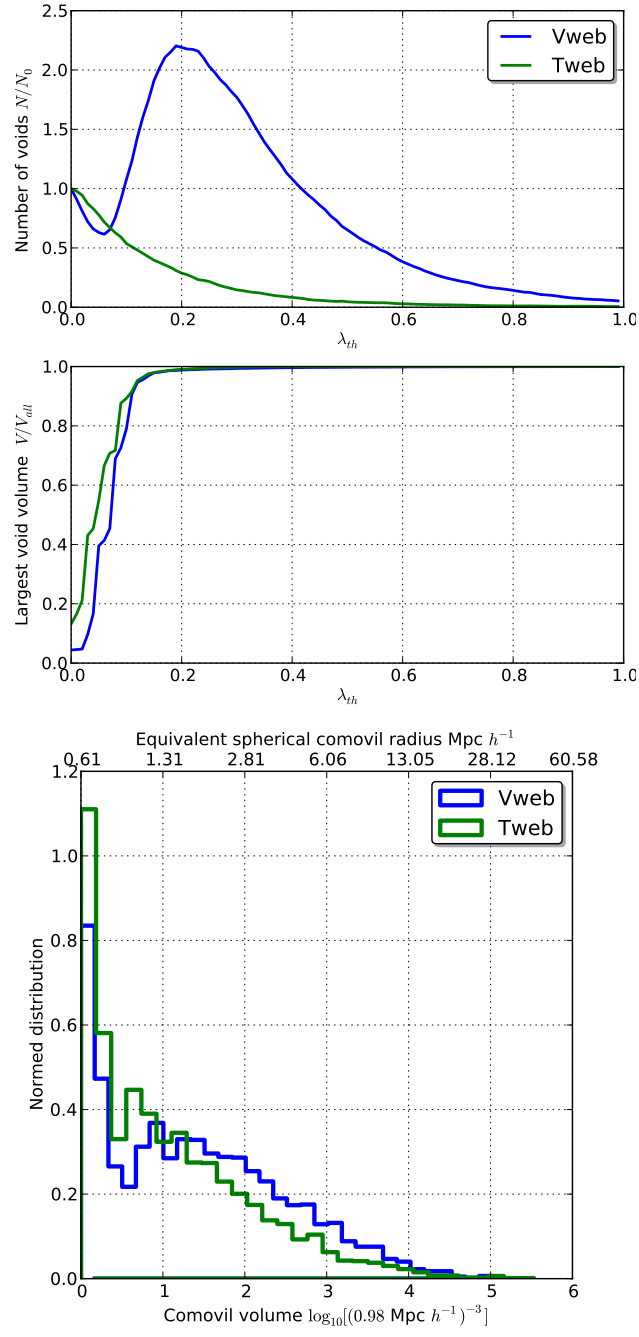


Figure 3. Percolation analysis of void regions for different λ_{th} values and for both defined classification schemes.

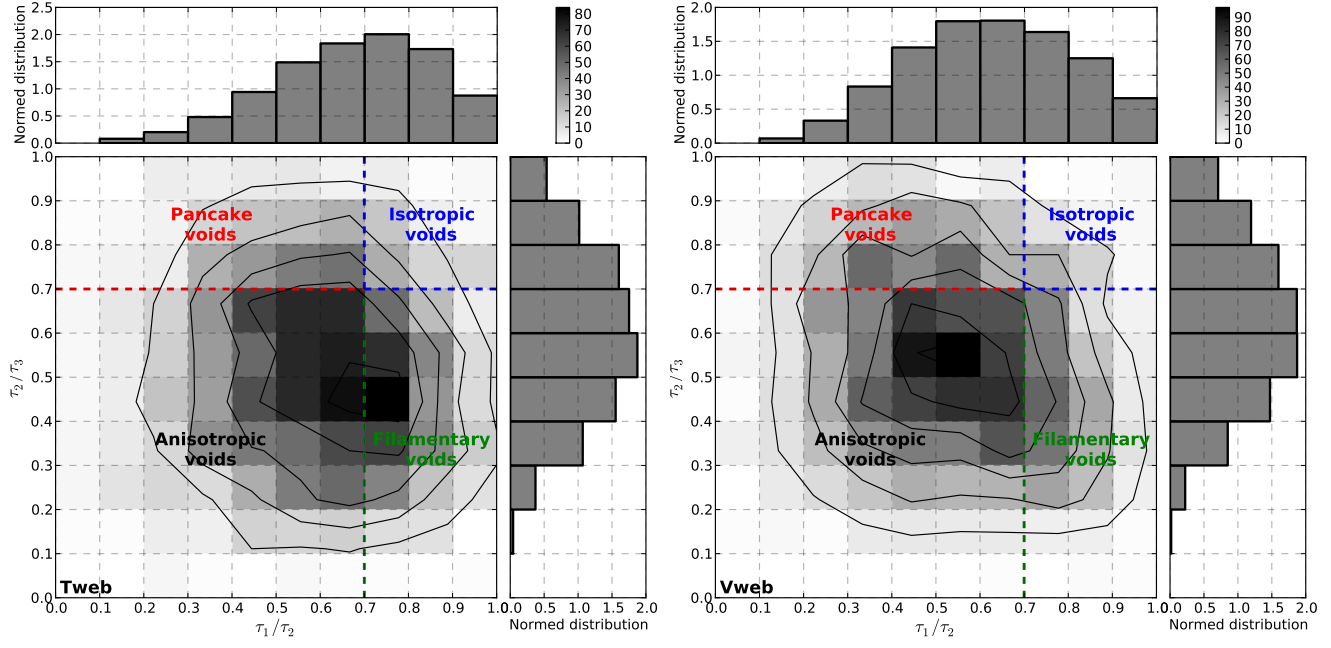


Figure 4. Histogram of eigenvalue ratio τ_1/τ_2 vs τ_2/τ_3 for the inertia tensor of void regions. T-web scheme (left). V-web scheme (right).