

## Problem Sheet

### Mathematical models in nonlinear acoustics

Westervelt equation is a classical model used for describing the propagation of nonlinear sound waves. In thermoviscous fluid media, it is given by

$$(1) \quad u_{tt} - c^2 \Delta u - b \Delta u_t = k(u^2)_{tt} + f.$$

Here  $u$  denotes the acoustic pressure,  $c > 0$  is the speed of sound in the medium,  $b > 0$  the sound diffusivity, and  $k \in \mathbb{R}$  the nonlinearity coefficient. The function  $f = f(x, t)$  acts as the source of sound.

Consider a linearized version of (1) with  $k = 0$ :

$$(2) \quad u_{tt} - c^2 \Delta u - b \Delta u_t = f,$$

on  $\Omega \times (0, T)$ , and supplement it with homogeneous initial data  $(u, u_t)|_{t=0} = (0, 0)$  and homogeneous Dirichlet boundary conditions. Let  $f \in L^2(0, T; L^2(\Omega))$ . Assuming the solution of this problem exists and is sufficiently smooth, by using suitable test functions, show that the following (higher-order) energy inequality holds:

$$\begin{aligned} & \int_0^t \|u_{tt}(s)\|_{L^2(\Omega)}^2 ds + \|\Delta u(t)\|_{L^2(\Omega)}^2 + \|u_t(t)\|_{H^1(\Omega)}^2 + \int_0^t \|\Delta u(s)\|_{L^2(\Omega)}^2 ds \\ & \leq C(T) \int_0^t \|f(s)\|_{L^2(\Omega)}^2 ds \end{aligned}$$

for all  $t \in (0, T)$ .

These energy arguments can be transferred to the study of Westervelt equation (1) by first carrying out the estimates on a linearization and then tying them to the original equation via Banach's fixed-point theorem on a ball of a sufficiently small radius. If you have time, you can try to work out the details.

You might need to rely on Grönwall's inequality: *Let  $w, v \in L^\infty(0, T)$  be almost everywhere non-negative functions that satisfy*

$$w(t) + v(t) \leq a_1 + \int_0^t a_2(s)w(s) ds \quad \text{for a.e. } t \in [0, T],$$

*where  $a_1 \geq 0$  and  $a_2 \in L^1(0, T)$  is an almost everywhere non-negative function. Then the following Grönwall inequality holds:*

$$w(t) + v(t) \leq a_1 e^{\int_0^t a_2(s) ds} \quad \text{for a.e. } t \in [0, T].$$

# 1 Uniform energy bounds

The estimate can be obtained through two testing steps.

- Testing with  $u_{tt}$

We take  $\phi = u_{tt}$  as the test function and integrate over  $(0, t)$ :

$$\int_0^t (u_{tt} - c^2 \Delta u - b \Delta u_t, u_{tt})_{L^2(\Omega)} ds = \int_0^t (f, u_{tt})_{L^2(\Omega)} ds.$$

We can use integration by parts in time and space to treat the  $b$  term on the left-hand side:

$$\begin{aligned} -b \int_0^t (\Delta u_t, u_{tt})_{L^2(\Omega)} ds &= b \int_0^t (\nabla u_t, \nabla u_{tt})_{L^2(\Omega)} ds \\ &= b(\nabla u_t, \nabla u_t)_{L^2(\Omega)} \Big|_0^t - b \int_0^t (\nabla u_{tt}, \nabla u_{tt})_{L^2(\Omega)} ds. \end{aligned}$$

From here we have

$$-b \int_0^t (\Delta u_t, u_{tt})_{L^2(\Omega)} ds = b \int_0^t (\nabla u_t, \nabla u_{tt})_{L^2(\Omega)} ds = \frac{b}{2} \|\nabla u_t(t)\|_{L^2(\Omega)}^2$$

The  $f$  term on the right-hand side can be estimated by using Cauchy–Schwarz inequality combined with Young’s  $\varepsilon$  inequality:

$$(3) \quad xy \leq \varepsilon x^2 + \frac{1}{4\varepsilon} y^2, \quad \varepsilon > 0.$$

In this manner, we obtain

$$\int_0^t (f, u_{tt})_{L^2(\Omega)} ds \leq \varepsilon \int_0^t \|u_{tt}\|_{L^2(\Omega)}^2 ds + \frac{1}{4\varepsilon} \int_0^t \|f\|_{L^2(\Omega)}^2 ds.$$

Similarly,

$$c^2 \int_0^t (\Delta u, u_{tt})_{L^2(\Omega)} ds \leq \varepsilon \int_0^t \|u_{tt}\|_{L^2(\Omega)}^2 ds + \frac{1}{4\varepsilon} c^4 \int_0^t \|\Delta u\|_{L^2(\Omega)}^2 ds$$

for any  $\varepsilon > 0$ . By picking  $\varepsilon$  to be small enough, we arrive at

$$(4) \quad \begin{aligned} &\int_0^t \|u_{tt}\|_{L^2(\Omega)}^2 ds + \frac{b}{2} \|\nabla u_t(t)\|_{L^2(\Omega)}^2 \\ &\leq \frac{1}{4\varepsilon} \int_0^t \|f\|_{L^2(\Omega)}^2 ds + \frac{1}{4\varepsilon} c^4 \int_0^t \|\Delta u\|_{L^2(\Omega)}^2 ds. \end{aligned}$$

- Testing with  $-\Delta u$

With this choice of the test function, we have, after integration in time,

$$\int_0^t (u_{tt} - c^2 \Delta u - b \Delta u_t, -\Delta u)_{L^2(\Omega)} ds = \int_0^t (f, -\Delta u)_{L^2(\Omega)} ds.$$

Similarly to before, we can estimate

$$\int_0^t (f, -\Delta u)_{L^2(\Omega)} ds \leq \frac{1}{2} \int_0^t \|\Delta u\|_{L^2(\Omega)}^2 ds + \frac{1}{2} \int_0^t \|f\|_{L^2(\Omega)}^2 ds$$

and (after transferring it to the right-hand side of the equation)

$$\int_0^t (u_{tt}, \Delta u)_{L^2(\Omega)} ds \leq \varepsilon \int_0^t \|u_{tt}\|_{L^2(\Omega)}^2 ds + \frac{1}{4\varepsilon} \int_0^t \|\Delta u\|_{L^2(\Omega)}^2 ds.$$

From here we have

$$\begin{aligned} & \frac{c^2}{2} \|\Delta u(t)\|_{L^2(\Omega)}^2 + b \int_0^t \|\Delta u_t\|_{L^2(\Omega)}^2 ds \\ & \leq \varepsilon \int_0^t \|u_{tt}\|_{L^2(\Omega)}^2 ds + \int_0^t \left(1 + \frac{1}{4\varepsilon}\right) \|\Delta u\|_{L^2(\Omega)}^2 ds + \frac{1}{2} \int_0^t \|f\|_{L^2(\Omega)}^2 ds. \end{aligned}$$

for any  $\varepsilon > 0$  (different than the one we had before).

We can now add this bound to (4), choose  $\varepsilon > 0$  small so that we can absorb the  $u_{tt}$  term on the right and then employ Grönwall's inequality. This leads to the desired estimate.