

Problem 1.

- (a) Derive Equation 8.1, from the Feynman rules for QED.
- (b) Obtain Equation 8.2 from Equation 8.1
- (c) Derive Equation 8.3 from Equation 8.2
- (d) Derive Equation 8.4 from Equation 8.3

### Feynman Rules

1. Notation :

external lines with momentum  $p_1, p_2 \dots$ , arrow direction forwards in time  
internal lines with momentum  $q_1, q_2 \dots$ , arrow direction forwards in time

2. External lines :

$$\begin{aligned} \text{Electrons} : & \begin{cases} \text{Incoming} : v \\ \text{Outgoing} : \bar{v} \end{cases} \\ \text{Quarks} : & \begin{cases} \text{Incoming} : \bar{\mu} \\ \text{Outgoing} : \mu \end{cases} \\ \text{Photons} : & \begin{cases} \text{Incoming} : \epsilon_e \\ \text{Outgoing} : \epsilon_{e^*} \end{cases} \end{aligned}$$

3. Vertex factors (where lines meet) :

$$\text{QED vertex, gamma} : ig_e \gamma^u$$

4. Propagator : each internal lines' factors

$$\text{Electrons and positrons} : \frac{i(\gamma^u q_u + mc)}{q^2 - m^2 c^2}$$

$$\text{Quarks and antiquarks} : \frac{i(q' + mc)}{q^2 - m^2 c^2}$$

(The problem said use only QED Feynman rules, so this was not used.)

$$\text{Photons} : \frac{-ig_v \mu}{q^2}$$

5. Conservation of E and P for each vertex :

$$(2\pi)^4 \delta^4(k_1+k_2+k_3)$$

where k's are the three four-momenta

6. Integrate over internal momenta : For internal moment q, the factors are

$$\frac{d^4 q}{(2\pi)^4}$$

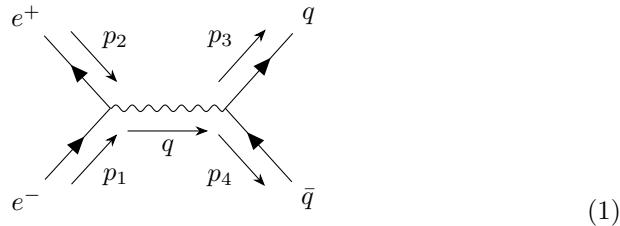
And then integrate

7. Cancel the delta function :

$$(2\pi)^4 \delta^4(p_1 + p_2 + \dots - p_n)$$

This will simplify the integral  $\rightarrow$  only one of q left by the conservation and return 1 by the delta function.

Solution



A transition amplitude for the first order of perturbation quantum-mechanics, peskin eqn(1.2),

In QM perturbation theory, to first order, the amplitude is,

$$\langle \text{final state} | H_1 | \text{initial state} \rangle$$

This is the first order, but the hamiltonian can not mediate the two state, but gamma does it.

So, expand this equation to the next order with  $\gamma$ .

For  $(e^- + e^+ \rightarrow \bar{q} + q)$ ,

$$M \sim \langle \bar{q} q | H_1 | \gamma \rangle^u \langle \gamma | H_1 | e^+ e^- \rangle_u$$

1. External electron lines :  $|e^+e^- \rangle$
2. External quark lines :  $\langle q \bar{q} |$
3. The vertices :  $H_1 = ig_e \gamma^u$  always for QED. The object  $\gamma^u$  are  $4 \times 4$  matrices.
4. Internal photon line :  $|\gamma\rangle\langle\gamma| = \frac{-ig_{v\mu}}{q^2}$

$$(2\pi)^4 \int [\bar{u}(3)(iQg_e \gamma^v)v(4)] \frac{-ig_{v\mu}}{q^2} [\bar{v}(2)(ig_e \gamma^u)\mu(1)] \times \delta^4(q-p_3-p_4) \delta^4(p_1+p_2-q) dq^4$$

And, by the momentum conservation,  $q = -p_1 - p_2 = p_3 + p_4$ .

$$M = \frac{Qg_e^2}{(p_1 + p_2)^2} [\bar{v}(p_2)\gamma^u\mu(p_1)][\bar{u}(p_3)\gamma_u v(p_4)]$$

This is the equation (8.1)

(small notes for myself :

Q. why do we only use spinors ?

A. chapter 7.2 The solution for dirac equation

Q. General solution of QM

A. Get the E.S and E.V from assuming  $p = 0 \rightarrow$  Get the solutions for the general  $p$  with the states when  $p = 0$ , and here, when  $p = 0$ , the spinors are the eigenstates.)

(b) Obtain Equation (8.2) from (8.1)

(8.2)

$$\langle|M|^2\rangle = \frac{1}{4} \left[ \frac{Qg_e^2}{(p_1 + p_2)^2} \right]^2 Tr[\gamma^\mu(p_1+mc)\gamma^\nu(p_2-mc)] \times Tr[\gamma_\mu(p_1-Mc)\gamma_\nu(p_2+Mc)]$$

(8.1)

$$M = \frac{Qg_e^2}{(p_1 + p_2)^2} [\bar{v}(p_2)\gamma^u\mu(p_1)][\bar{u}(p_3)\gamma_u v(p_4)]$$

Where the slash notation means,  $\not{a} \equiv a^\mu \gamma_\mu$ .

(Following of chapter 7.7)

$\langle|M|^2\rangle \equiv$  average over initial spins, sum over final spins of  $|M(s_i \rightarrow s_f)|^2$ .

For electron-muon scattering amplitude,

$$|M|^2 = \frac{g_e^4}{(p_1 - p_3)^4} [\bar{u}(3)\gamma^\mu u(1)][\bar{u}(4)\gamma_\mu u(2)][\bar{u}(3)\gamma^\nu u(1)]^*[\bar{u}(4)\gamma_\nu u(2)]^*$$

$$G \equiv [\bar{u}(a)\Gamma_1 u(b)][\bar{u}(a)\Gamma_2 u(b)]^*$$

We can evaluate the complex conjugate because the quantity in the square brackets is a  $1 \times 1$  matrix.

$$[\bar{u}(a)\Gamma_2 u(b)]^* = [u(a)^\dagger \gamma^0 \Gamma_2 u(b)]^\dagger = u(b)^\dagger \Gamma_2^\dagger \gamma^{0\dagger} u(a)$$

Now,  $\gamma^{0\dagger} = \gamma^0$ , and  $(\gamma^0)^2 = 1$ , so

$$u(b)^\dagger \Gamma_2^\dagger \gamma^{0\dagger} u(a) = u(b)^\dagger \gamma^0 \gamma^0 \Gamma_2^\dagger \gamma^0 u(a) = \bar{u}(a) \bar{\Gamma}_2 u(a)$$

Where

$$\bar{\Gamma}_2 \equiv \gamma^0 \Gamma_2^\dagger \gamma^0$$

Thus

$$G = [\bar{u}(a)\Gamma_1 u(b)][\bar{u}(b)\bar{\Gamma}_2 u(a)]$$