# QED corrections to the Bethe-Heitler process in the $\gamma p \to l^+ l^- p$ reaction

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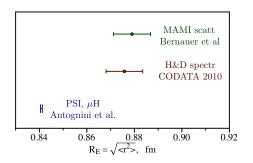
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## Proton radius puzzle

- discrepancy between measurements from muonic spectrocospy and electron data
- at least 4 standard derivations difference
- physics beyond the Standard model?



## Lepton universality

- Assume universal proton form factor for all leptons
   → results in same proton radius for muons and electrons
- Broken universality could be an explanation for proton radius puzzle
- Test this with upcoming experiments:
  - MUSE @ PSI (electron vs muon scattering)
  - MAMI  $\gamma p 
    ightarrow e^- e^+ p$  vs.  $\gamma p 
    ightarrow \mu^- \mu^+ p$

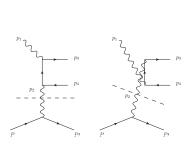
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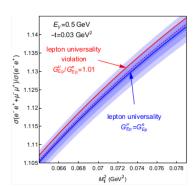
## Bethe-Heitler process

phtotproduction of lepton pair on a proton target

$$\gamma p \rightarrow I^+ I^- p$$

 $\bullet$  comparsion of cross sections for  $\mathit{I}=\mathit{e}$  and  $\mathit{I}=\mu$  can be used to test lepton universality

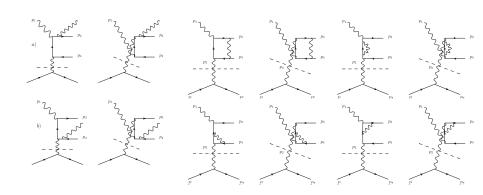




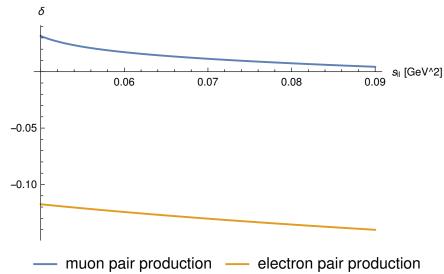
## **QED** corrections

#### Real corrections

#### Virtual corrections

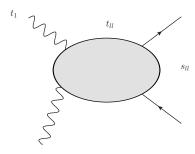


## First result in Soft-photon approximation

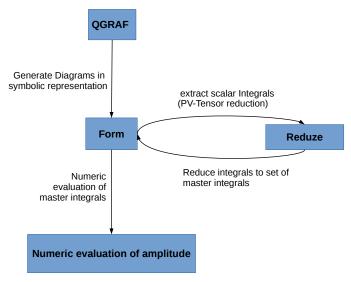


## Difficulties in the full calculation

- many scales
  - 4 external scales
  - 1 internal mass (lepton)
- due to massive Dirac chains huge expressions
- Passarino-Veltman tensor reduction up to rank 3



## Overview of full calculation



#### **IBP** relations

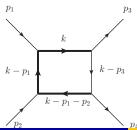
- relations between scalar integrals based on Integration-By-Part
- in dimensional regularisation, any scalar integra obeys

$$\int d^d k_i \frac{\partial}{\partial k_i^{\mu}} \left[ q^{\mu} I(p_1, ..., p_n, k_1, ... k_m) \right] = 0$$
 (1)

 leads to rational system of equations, that can be solved for a set of so called master integrals

## Integral families

- set of scalar propagators  $P_1, ... P_n$ , such that any scalar product between external momenta and loop momenta or between two loop momenta can be re-expressed by these
- member of family can be labeled by integer exponents of the  $P_i$
- for Bethe-Heitler process we have



Example for labeling:

$$\begin{split} &\int \frac{d^dk}{(2\pi)^d} k^2 \frac{1}{(k-p_1)^2 - m^2} \frac{1}{(k-p_1-p_2)^2 - m^2} \frac{1}{(k-p_3)^2} \\ &= \int \frac{d^dk}{(2\pi)^d} (P_1^{-1} + m^2) P_2 P_3 P_4 \\ &= I(-1,1,1,1) + m^2 I(0,1,1,1) \end{split}$$

## Solving master integrals

- master integrals can be solved numerically or analytically
- for numerical solution, one can use Sector Decomposition and MC-Integration
- for analytic solution, use method of differential equations

$$\partial_{x}\vec{I} = R(x,\epsilon)\vec{I} \tag{2}$$

 by particular change of master integrals, one can get canonical form, where solution gets significant simpler

$$\partial_{x}\vec{I} = \epsilon R(x)\vec{I} \tag{3}$$



## Conclusion and Outlook

- As a first step, we presented the first order QED calculation for the Bethe Heitler process in the Soft-photon limit
- Full calculation by method of master integrals is work in progress
- Upcoming experiment at MAMI (Mainz) aims to test lepton universality
- Complementary  $\gamma p \to l^+ l^- p$  experiment could shed light on the proton radius puzzle

# Example: Bubble-Integral

Define the integral family:

$$I(a,b) = \int \frac{d^d k}{(2\pi)^d} \left(\frac{1}{k^2}\right)^a \left(\frac{1}{(k-p)^2 - m^2}\right)^b$$
 (4)

IBP identities:

$$\frac{\partial}{\partial k^{\mu}} \left( \frac{1}{k^2} \right)^{s} \left( \frac{1}{(k-\rho)^2 - m^2} \right)^{b} \tag{5}$$

$$= -2a \left(\frac{1}{k^2}\right)^{(a+1)} \left(\frac{1}{(k-\rho)^2 - m^2}\right)^b k_{\mu} - 2b \left(\frac{1}{k^2}\right)^a \left(\frac{1}{(k-\rho)^2 - m^2}\right)^{(b+1)} (k_{\mu} - p_{\mu}) \tag{6}$$

This can be contracted with  $k_{\mu}$  or  $p_{\mu}$  (2 identities).

Any integral of the family can be related with these identities to two master integrals, e.g. I(0,1) and I(1,1).

# Example: Bubble-Integral

Choose

$$\vec{I}_1 = \begin{pmatrix} I(0,1) \\ I(1,1) \end{pmatrix}, \tag{7}$$

then

$$\partial_{m}\vec{I}_{1} = \frac{2}{m(m^{2} - s)} \begin{pmatrix} -(m^{2} - s)(-1 + \epsilon) & 0\\ (-1 + \epsilon) & -m^{2}(-1 + 2\epsilon) \end{pmatrix} \vec{I}_{1}, \tag{8}$$

$$\partial_{s}\vec{l}_{1} = \frac{1}{s(m^{2} - s)} \begin{pmatrix} 0 & 0 \\ -(-1 + \epsilon) & m^{2} + (-2 + \epsilon)s + m^{2}(-2 + \epsilon) + 2s \end{pmatrix} \vec{l}_{1}. \tag{9}$$

"Better" choice:

$$\vec{l}_2 = \begin{pmatrix} \epsilon m^{2\epsilon} \ l(0,2) \\ \epsilon(-s) \ m^{2\epsilon} \ l(1,2) \end{pmatrix}, \quad x = \frac{(-s)}{m^2}.$$
 (10)

Then:

$$\partial_{x}\vec{l}_{2} = \epsilon \begin{bmatrix} \frac{1}{x} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} + \frac{1}{x+1} \begin{pmatrix} 0 & 0 \\ -1 & -2 \end{pmatrix} \end{bmatrix} \vec{l}_{2}$$
 (11)