

ORBITAL MECHANICS

We go on expressing a, e, p as a function r_p, r_a

$$\begin{aligned} r_p \text{ (1.45)} & \rightarrow \frac{r_p}{r_a} = \frac{1-e}{1+e} \rightarrow r_p(1+e) = r_a(1-e) \\ r_a \text{ (1.46)} & \end{aligned}$$

$$(r_p + r_a)e = r_a - r_p \Rightarrow e = \frac{r_a - r_p}{r_a + r_p} \quad (1.47)$$

↑
very useful

$$\begin{aligned} r_a - r_p &= \overline{FF'} \\ r_a + r_p &= 2a \end{aligned} \quad \text{Distance of foci} \quad \left\{ \Rightarrow e = \frac{\overline{FF'}}{2a} \quad (1.48) \right.$$

SEMI-MAJOR AXIS

$$(1.40) \rightarrow$$

$$a = \frac{r_p + r_a}{2} \quad (1.49)$$

$$r_p = a(1-e)$$

$$r_a = a(1+e)$$

$$p = a(1-e^2)$$

SEMI LATUS RECTUM

$$\left\{ \begin{aligned} r_p &= \frac{p}{1+e} \\ e &= \frac{r_a - r_p}{r_a + r_p} \end{aligned} \right. \rightarrow r_p(1+e) = p \rightarrow p = r_p \left(1 + \frac{r_a - r_p}{r_a + r_p} \right)$$

$$p = \frac{2r_a r_p}{r_a + r_p} \quad (1.50)$$

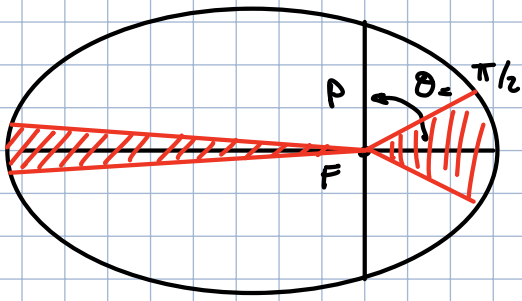
VELOCITY

$$(1.33) \quad (1.34) \quad v_\theta, v_r$$

$$\begin{aligned} v &= \sqrt{v_\theta^2 + v_r^2} \quad \leftarrow \text{magnitude of the velocity} \\ &= \sqrt{\frac{\mu^2}{h^2} (e^2 \sin^2 \theta + (1 + e \cos \theta)^2)} \quad \text{but } p = \frac{h^2}{\mu} \quad (1.38) \end{aligned}$$

$$\Rightarrow v = \sqrt{\frac{\mu}{p} (e^2 \sin^2 \theta + (1 + e \cos \theta)^2)}^{1/2}$$

$$r = \left[\frac{\mu}{p} (e^2 + 1 + 2e \cos \theta) \right]^{1/2}$$



$$p = r(\theta = \frac{\pi}{2})$$

- ① $r_p \rightarrow r_p = \sqrt{\frac{\mu}{p}} (1+e)$ periaapsis maximum velocity (1.52)
- ② $r_a \rightarrow r_a = \sqrt{\frac{\mu}{p}} (1-e)$ apoapsis minimum velocity (1.53)

↳ we do not have any velocity along the radial direction

ORBIT ENERGY AND VELOCITY

$E = \text{CONSTANT}$ TOTAL ENERGY \forall point in orbit.

$$\textcircled{1} \quad E_p = \left(\frac{v_p^2}{2} - \frac{\mu}{r_p} \right) = \frac{\mu}{2p} (1+e)^2 - \frac{\mu(1+e)}{p}$$

(1.52) (1.45)

$$E_p = \frac{\mu}{2p} [(1+e)^2 - 2(1+e)] = \frac{\mu}{2p} (1 + 2e + e^2 - 2 - 2e) \\ = \frac{\mu}{2p} (e^2 - 1)$$

$$E_p = \frac{\mu}{2p} (e^2 - 1) = - \frac{\mu}{2p} (1 - e^2) = - \frac{\mu^2}{2h^2} (1 - e^2)$$

$$p = \frac{a^2}{\mu}$$

$$E_p = - \frac{1}{2} \frac{\mu^2}{h^2} (1 - e^2) \quad (1.54)$$

$$p = a(1-e^2) \quad (1.41)$$

$$\Rightarrow \boxed{e_p = -\frac{\mu}{2\varepsilon}} \quad (1.55)$$

It's very important because it gives us e in function of the semi major-axis

↳ As a function of the h and $e \rightarrow$ clearly the orbital energy ε is not an independent quantity.

NOTE

$e_p(1.55)$ and $e_p(1.54)$ valid for \forall orbit (elliptical, parabolic, hyperbolic)

$$\text{From (1.10) and (1.55)} \rightarrow \boxed{\frac{v^2}{2} - \frac{\mu}{r} = -\frac{\mu}{2a}} \quad (1.56)$$

a semi major-axis tells us the energy level of the orbit.

↳ depends on orbit
a defines energy level

$$(1.56) \rightarrow v$$

Knowing ε and $r \rightarrow$ velocity at any point of the orbit

$$\boxed{v = \sqrt{\mu \left(\frac{2}{r} - \frac{1}{a} \right)}}$$

$$(1.57)$$

ORBITAL PERIOD OF A CLOSED ORBIT \rightarrow (valid only for circular and elliptical orbit)

We start from Area of an ellipse $\frac{1}{2} \pi a b$

$$A = \pi a b$$

(1.58)

From here we use Kepler's second law

$$\frac{dA}{dt} = \frac{1}{2} h \quad \text{one revolution} \quad \pi a b = \frac{1}{2} h T$$

\uparrow

Kepler's second law

$$T = \frac{2\pi a b}{h}$$

$$b = a \sqrt{1-e^2} \quad (1.44)$$

$$h = \sqrt{\mu p} \quad (1.38)$$

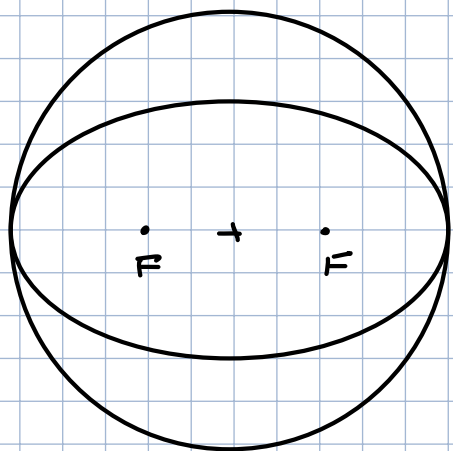
$$p = a(1-e^2) \quad (1.41)$$

$$\Rightarrow T = \frac{2\pi a^2 \sqrt{1-e^2}}{\sqrt{\mu a(1-e^2)}}$$

Period of a closed orbit

$$T = 2\pi \sqrt{\frac{a^3}{\mu}} \quad (1.59)$$

it depends only on the semi-major axis (like the total energy)



NOTE

These two orbits have the same energy and orbital period.

CIRCULAR ORBIT

same characteristic of the elliptical orbit but it is necessary to set $e = 0$.

$e = 0$ in eq (1.26)

$$r = \frac{h^2}{\mu} \quad \text{or} \quad r = a \quad (1.6)$$

$r = \text{constant}$

$$\text{Eq (1.33)} \quad \dot{r} = \frac{\mu}{h} e \sin \theta \rightarrow r_r = \dot{r} = 0$$

it is present only the azimuthal component of r for circular orbit $r = r_0 \rightarrow h$ reduces

$$h_{\text{circular}} = r_0 v = r \cdot v \quad (1.61)$$

$$h = \frac{h^2}{\mu} = \frac{r v^2}{\mu}$$

$$v_{\text{circular}} = \sqrt{\frac{\mu}{r}} \quad (1.62)$$

$$T = 2\pi \sqrt{\frac{a^3}{\mu}} \rightarrow T = \frac{\text{circumference}}{\text{speed}} = \frac{2\pi r}{\sqrt{\mu/r}} = 2\pi \sqrt{\frac{r^3}{\mu}}$$

$$T_c = 2\pi \sqrt{\frac{r^3}{\mu}} \quad (1.63)$$

$$a = r$$

$$\text{Eq (1.54)} \quad E = -\frac{1}{2} \frac{\mu^2}{h^2} (1 - e^2)$$

$$E_{\text{circular}} = -\frac{1}{2} \frac{\mu^2}{h^2}$$

$$E_T \text{ (1.55)}$$

$$E = -\frac{\mu}{2a}$$

$$E_{\text{circular}} = -\frac{\mu}{2r} \quad (1.64)$$

as r gets bigger E_c gets smaller.

NOTE

With the same energy a launcher can put a heavy satellite in a low earth-orbit or a smaller and lighter satellite in a elliptical or ionion orbit with a bigger semi-major axis.

example LEO altitude [150 km 2000 km]

- remote sensing
- imaging
- navigation satellites
- telecommunication (internet broadband / large constellation).

example GEO geostationary Earth orbit

s/c remains stationary above same point on Earth.

rotational velocity of Earth

$$\omega_E = \frac{2\pi}{\text{sidereal day}} \quad (1.65)$$

sidereal day: time it takes to Earth to complete 1 full revolution wrt inertial direction (fixed star)
23h and 56 min.

synodic day (solar day): 24h time Sun takes full apparent rotation around Earth (noon to noon).

$$\omega_E = \frac{2\pi}{23 \text{ h } 56 \text{ min}}$$

angular velocity of the Earth.

GEO s/c are used for → -telecoms
-television broadcast

$$r_{GEO} = 42164 \text{ km}$$

$$\left. \begin{aligned} v_{GEO} &= \omega_E r_{GEO} \\ v_{GEO} &= \sqrt{\frac{\mu}{r_{GEO}}} \end{aligned} \right\} \omega_E^2 r_{GEO}^2 = \frac{\mu}{r_{GEO}}$$

$$r_{GEO}^3 = \frac{\mu}{\omega_E^2}$$

$$r_{GEO} = \sqrt[3]{\frac{\mu_E}{\omega_E^2}}$$

$$\mu_E = G(M_E + m_{s/c}) = G m_E = 398600 \frac{\text{km}^3}{\text{s}^2}$$

$$h_{GEO} = r_{GEO} - R_E = 35786 \text{ km}$$

$$v_{GEO} = \sqrt{\frac{\mu_E}{r_{GEO}}} = 3.075 \frac{\text{km}}{\text{s}}$$

PARABOLIC TRAJECTORY

$$e=1$$

$$\text{From eq (1.26)} \quad r = \frac{a^2}{\mu} \frac{1}{1+e \cos \theta}$$

$$\text{as } \theta \rightarrow \pi \quad 1 + e \cos \theta \rightarrow 0 \quad \text{and} \quad r \rightarrow \infty \quad \text{OPEN ORBIT}$$

$$\text{From eq (1.54)} \quad E = -\frac{1}{2} \frac{\mu^2}{h^2} (1-e^2)$$

$$E_{PAR} = 0$$

$$(1.66)$$

$$G_p (1.56) \quad \frac{v^2}{2} - \frac{\mu}{r} = -\frac{\mu}{2a} \rightarrow \text{Parabolic trajectory} \quad \frac{v^2}{2} - \frac{\mu}{r} = 0$$

\Rightarrow

$$v_{par} = \sqrt{\frac{2M}{r}}$$

(1.67)

If part see the parabola,

$r \rightarrow \infty$

$$\lim_{\substack{r \rightarrow \infty \\ \theta \rightarrow \pi}} v_{par} = 0$$

m_2 escapes from m_1 at ϕ relative velocity.

\Rightarrow MINIMUM ESCAPE ENERGY (PATH)

At r from m_1 the minimum escape velocity is the one of a parabola

FIRST COSMIC SPEED

$$v_{cir} = \sqrt{\frac{M}{r}}$$

(1.68)

SECOND COSMIC SPEED

$$v_{esc} = \sqrt{\frac{2M}{r}}$$

(1.69)

$$v_{es} = \sqrt{2} v_{cir}$$

① r on circular orbit (v_{cir}) to escape boost velocity of 41.4%

After escaping from the Earth $v_{par} \xrightarrow{r \rightarrow \infty} 0$ we are at ϕ but with a ϕ relative velocity \rightarrow at this point the ϕ is in orbit around the sun with the same velocity of the Earth. Parabolic orbit is the limit between the Earth and the Sun.