

# Numerical Algorithms for HPC

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Solution of Time-Dependent PDEs



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## Overview

- Boundary Value Problem: pollution model
- Solution using Jacobi
- Discretisation of time-dependent problem
- Euler equations
- Stability
- Implicit methods
- Error Analysis
- Other equations



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## Pollution Model

- Previously posed as a boundary value problem
  - i.e. static in time
- What is the solution to:  $-\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right)u(x,y) = 0$ 
  - with  $u = 0$  on north, south and west boundaries
  - and  $u =$  hump function on east boundary (location of chimney)
- Discretised equations using standard recipes
  - e.g. 5-point stencil for derivative
- Can be solved using standard methods
  - e.g. Jacobi, Over-Relaxed Gauss-Seidel, Conjugate Gradient (later), ...



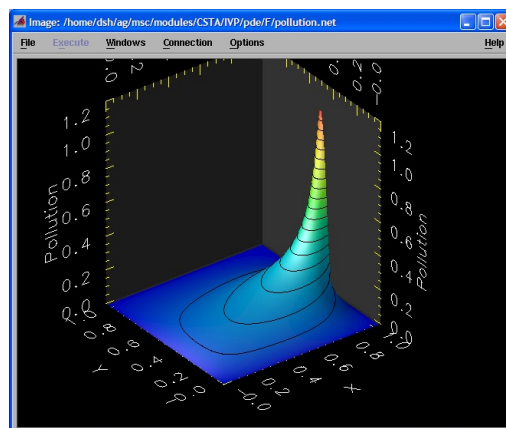
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## Example Solution

- Use a 72×72 grid
  - included a north-easterly wind of strength (10.0, 4.0)



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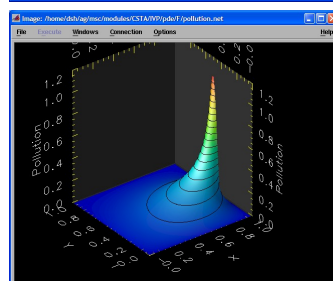
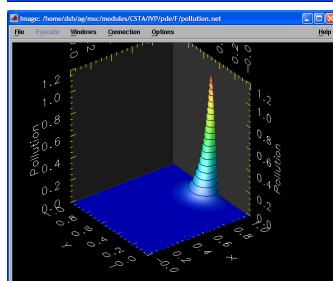
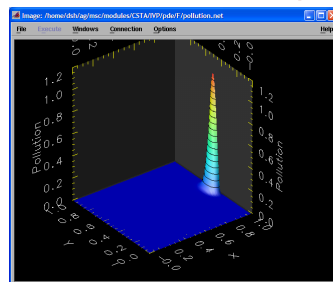
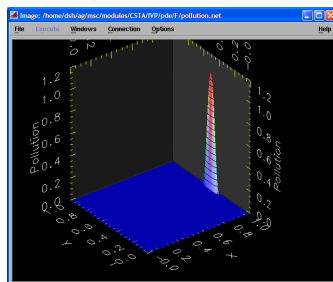


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## Evolution of the Solution

- Initial guess to solution is the hump function
  - i.e. the situation when chimney is just switched on
- Final output is the static solution
  - i.e. the situation when the chimney has been on for a long time
- What about the intermediate solutions
  - are they related in any way to the actual evolution in time?
- With Jacobi, yes ...

## Jacobi (after 0, 10, 100, 1000 iterations)



## How do we get real time evolution?

- This is an Initial Value Problem (IVP)
  - specify the solution at time  $t = 0$
  - need to know the solution at some later time  $t$
- Similar issues to orbits practical
  - initialise position
  - compute force
  - update position
- Except we have many variables to update
  - here, pollution values at more than 5000 points (on  $72 \times 72$  grid)
  - force term is calculated from the grid



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## Diffusion Equation

- We have considered :  $\nabla^2 u(x, y) = \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) u(x, y) = 0$ 
  - this is an example of an *elliptic equation*
- Full problem is actually:  $\frac{\partial u}{\partial t} = \nabla^2 u(x, y)$ 
  - this is an example of a *parabolic equation*
- We have already solved the static solution
  - i.e. have set time derivative to zero
  - but what if we want to track the evolution in time?



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## Discretisation in space and time

- Use a simple five-point stencil for the RHS with spacing  $h$
- Use a simple forward difference for LHS with timestep  $dt$

$$\frac{\partial u}{\partial t} \approx \frac{u^{(t+dt)} - u^{(t)}}{dt}$$

– superscript refers to real time  $t$  and not “computer time”  $n$

- Full equations are:  $\frac{\partial u}{\partial t} = \nabla^2 u(x, y)$

– discretised:  $\frac{u_{i,j}^{(t+dt)} - u_{i,j}^{(t)}}{dt} = u_{i,j-1}^{(t)} + u_{i-1,j}^{(t)} - 4u_{i,j}^{(t)} + u_{i+1,j}^{(t)} + u_{i,j+1}^{(t)}$

–  $u_{i,j}^{(t+dt)} = u_{i,j}^{(t)} + dt \left( u_{i,j-1}^{(t)} + u_{i-1,j}^{(t)} - 4u_{i,j}^{(t)} + u_{i+1,j}^{(t)} + u_{i,j+1}^{(t)} \right)$

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## Euler update

- A simple *explicit* scheme

• Jacobi:  $u_{i,j}^{(n+1)} = \frac{1}{4} \left( u_{i,j-1}^{(n)} + u_{i-1,j}^{(n)} + u_{i+1,j}^{(n)} + u_{i,j+1}^{(n)} \right)$

•  $\Rightarrow u_{i,j}^{(n+1)} = u_{i,j}^{(n)} + \frac{1}{4} \left( u_{i,j-1}^{(n)} + u_{i-1,j}^{(n)} - 4u_{i,j}^{(n)} + u_{i+1,j}^{(n)} + u_{i,j+1}^{(n)} \right)$

- Time dependent:

$$u_{i,j}^{(t+dt)} = u_{i,j}^{(t)} + dt \left( u_{i,j-1}^{(t)} + u_{i-1,j}^{(t)} - 4u_{i,j}^{(t)} + u_{i+1,j}^{(t)} + u_{i,j+1}^{(t)} \right)$$

- Notes

– Jacobi update corresponds to *time integration* with  $dt = 0.25!$

– smaller values of  $dt$  will give more accurate intermediate solutions

• but will require more timesteps and hence more work

– should always arrive at the same static solution (eventually) if it exists

– in real situations there might be no static solution (e.g. turbulent flow)

– Update of the form  $u^{(t+dt)} = u^{(t)} + dt f(u^{(t)})$  known as *Euler update*

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## Checking the accuracy

- Initial Value Problems are a leap in the dark
  - how do we know if the solution is correct?
  - Can no longer simply check convergence
- Need to be careful
  - can sometimes monitor conserved quantities, for example the total amount of pollution
  - also look at solution visually
- Start with the hump in the middle (Following week's exercise)
  - easier to understand: pollution stays away from boundaries for longer
  - pollution does not disappear for the initial timesteps
- Note that the height actually decreases as the width widens
  - visualisation software used in next slide (unfortunately) rescales the z-axis

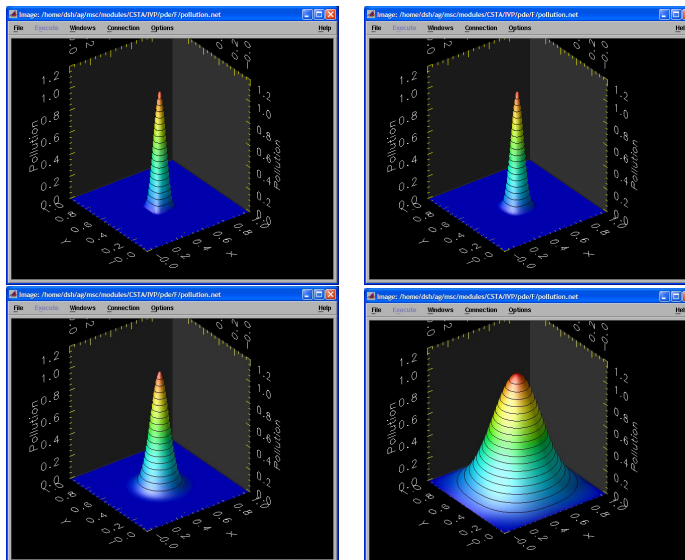
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$$dt = 0.1; t = 0.0, 1.0, 10.0, 100.0$$



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## Stability

- Can perform a formal analysis of stability
  - following Von Neumann
- This shows that, for stability in the 2D diffusion problem we require  $dt \leq 0.25$ 
  - Jacobi algorithm therefore uses the maximum timestep
- But what happens if we go beyond this?
  - see exercise!

## Other Schemes

- Imagine evaluating the derivative at time  $t + dt$

$$\frac{u_{i,j}^{(t+dt)} - u_{i,j}^{(t)}}{dt} = \nabla^2 u^{(t+dt)}$$

- One arrangement gives an *implicit* scheme:

$$(1 - dt \nabla^2) u^{(t+dt)} = u^{(t)}$$

- must solve a full boundary value problem at every timestep!
  - however, it is always stable so can use a much larger  $dt$

- Many other integration schemes exist
  - more stable than explicit scheme
  - allow larger timesteps

## Error analysis: second derivative

$$\nabla^2 u(x) \approx u_{i-1} - 2u_i + u_{i+1}$$

$$u_{i-1} = u(x-h) = u(x) - h \frac{du}{dx} + \frac{h^2}{2} \frac{d^2u}{dx^2} - \frac{h^3}{6} \frac{d^3u}{dx^3} + O(h^4)$$

$$u_i = u(x)$$

$$u_{i+1} = u(x+h) = u(x) + h \frac{du}{dx} + \frac{h^2}{2} \frac{d^2u}{dx^2} + \frac{h^3}{6} \frac{d^3u}{dx^3} + O(h^4)$$

$$\frac{1}{h^2} (u_{i-1} - 2u_i + u_{i+1}) = \frac{d^2u}{dx^2} + O(h^2)$$

- Expression is therefore accurate to *second order* in  $h$ 
  - straightforward to extend to 2D
  - note we've previously been rather lax about units, i.e. factors of  $h^2$  etc.

## Error analysis: first derivative

$$\frac{\partial u}{\partial t} \approx \frac{u^{(t+dt)} - u^{(t)}}{dt}$$

$$\frac{u^{(t+dt)} - u^{(t)}}{dt} = \frac{\partial u}{\partial t} + O(dt)$$

- Euler integration is only accurate to first order
  - but very simple!



## Units for diffusion equation

- Actual diffusion equation is:  $\frac{\partial u}{\partial t} = D \nabla^2 u$

– diffusion constant  $D$  is large for a gas, small for treacle, ...

- Full update equations are:

$$u_{i,j}^{(t+dt)} = u_{i,j}^{(t)} + \frac{Ddt}{h^2} (u_{i,j-1}^{(t)} + u_{i-1,j}^{(t)} - 4u_{i,j}^{(t)} + u_{i+1,j}^{(t)} + u_{i,j+1}^{(t)})$$

– previously we set  $D = h^2$  for simplicity

- Stability condition is  $dt \leq h^2/(4D)$ 
  - this is the Courant–Friedrichs–Lewy (CFL) condition
- Euler is impractical for the diffusion equation
  - halving  $h$  means reducing  $dt$  by a factor of 4 ...



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## Other equations

- An important equation is the wave equation (1D and 2D)

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

$$\frac{\partial^2 u}{\partial t^2} = c^2 \nabla^2 u$$

- Describes waves travelling with velocity  $c$ 
  - sound
  - water
  - electromagnetic
  - ...
- This is a *hyperbolic equation*



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## Boundary conditions

- We have used very simple fixed boundary conditions
  - these are Dirichlet conditions
  - e.g. pollution is zero on the boundary
  - e.g. simply set  $u_{i,0} = 0$  to force this on the southern boundary
- May want conditions on the derivative
  - e.g. pollution is constant across the southern boundary
  - i.e.  $\frac{\partial u}{\partial y} = 0$  for all values of  $x$  at  $y = 0$
- Compute discrete form using Taylor expansion
  - here it is simply  $u_{i,0} = u_{i,1}$  on the southern boundary
  - re-impose this condition every time we update  $u$



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## Exercise (next week)

- Look at stability of Euler Integration for 2D diffusion
  - vary  $dt$
  - what happens when the CFL condition is violated?
  - is the total amount of pollution conserved?
  - what happens when boundary conditions are changed?



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