

Kane-Mele model

In 2005, Kane and Mele [PRL 95, 226801 (2005)] generalized the Haldane model to the graphene lattice model with the time reversal invariant spin-orbit coupling, which is known as the Kane-Mele model.

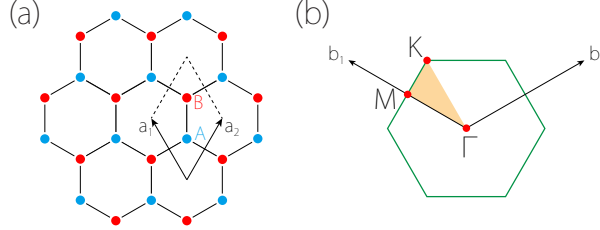


Figure 1. (a) Honeycomb structure. (b) Hexagonal Brillouin zone.

The graphene lattice has the following translation vectors:

$$\mathbf{a}_1 = a\left(-\frac{\sqrt{3}}{2}, \frac{3}{2}\right), \quad \mathbf{a}_2 = a\left(\frac{\sqrt{3}}{2}, \frac{3}{2}\right)$$

where a is the bond length. The fractional coordinates of A and B sites are written as

$$\mathbf{r}_A = \left(\frac{1}{3}, \frac{1}{3}\right), \quad \mathbf{r}_B = \left(\frac{2}{3}, \frac{2}{3}\right)$$

The Kane-Mele model is expressed as,

$$\mathcal{H} = t \sum_{\langle i,j \rangle} c_i^\dagger c_j + i\lambda_{SO} \sum_{\langle\langle i,j \rangle\rangle} v_{ij} c_i^\dagger s_z c_j + i\lambda_R \sum_{\langle i,j \rangle} c_i^\dagger (\mathbf{s} \times \mathbf{d}_{ij})_z c_j + \lambda_v \sum_i \xi_i c_i^\dagger c_i,$$

where $c_i^\dagger = (c_{i,\uparrow}^\dagger, c_{i,\downarrow}^\dagger)$ and the Pauli matrices s_i describe the electron spin. The first term is the nearest neighbor hopping term on a graphene lattice. The second term is a mirror symmetric spin-orbit interaction which connects second neighbors with a spin dependent amplitude. Here, $v_{ij} = \frac{2}{\sqrt{3}}(\mathbf{d}_i \times \mathbf{d}_j)_z = \pm 1$, where \mathbf{d}_i and \mathbf{d}_j are two unit vectors along the two bonds the electron traverses going from site j to i . The third term is the nearest neighbor Rashba term, which violates the M_z mirror symmetry, breaks the conservation of s_z and couples the electrons with spin-up and spin-down. The last term is a staggered sublattice potential with $\xi_i = \pm 1$.

A magnetic field \mathbf{B} breaks time-reversal symmetry, which allows the terms in the Hamiltonian as

$$H_B = \sum_i c_i^\dagger (\mathbf{B} \cdot \mathbf{s}) c_i$$