

# Notes on Fourier Analysis

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# 1 The Laplace Transform

$$\begin{aligned} L[1] &= \int_0^{\infty} e^{-st} dt \\ &= -\frac{1}{s} e^{-st} \Big|_0^{\infty} \\ &= \frac{1}{s} \end{aligned}$$

$$\begin{aligned} L[t] &= \int_0^{\infty} t e^{-st} dt \rightarrow \begin{bmatrix} u = t & v = -\frac{1}{s} e^{-st} \\ du = dt & dv = e^{-st} dt \end{bmatrix} \\ &= -\frac{t}{s} e^{-st} \Big|_0^{\infty} + \int_0^{\infty} \frac{1}{s} e^{-st} dt \\ &= \frac{1}{s^2} \end{aligned}$$

$$\begin{aligned} L[t^2] &= \int_0^{\infty} t^2 e^{-st} dt \rightarrow \begin{bmatrix} u = t^2 & v = -\frac{1}{s} e^{-st} \\ du = 2t dt & dv = e^{-st} dt \end{bmatrix} \\ &= -\frac{t^2}{s} e^{-st} \Big|_0^{\infty} + \int_0^{\infty} \frac{2}{s} t e^{-st} dt \\ &= \frac{2}{s} \int_0^{\infty} t e^{-st} dt \\ &= \frac{2}{s^3} \end{aligned}$$

There is an interesting pattern showing here, let's try a more generic case now.

$$\begin{aligned} L[t^n] &= \int_0^{\infty} t^n e^{-st} dt \rightarrow \begin{bmatrix} u = t^n & v = -\frac{1}{s} e^{-st} \\ du = n t^{n-1} dt & dv = e^{-st} dt \end{bmatrix} \\ &= -\cancel{\frac{t^n}{s} e^{-st}} \Big|_0^{\infty} + \frac{n}{s} \int_0^{\infty} t^{n-1} e^{-st} dt \end{aligned}$$

Note that we can tell that the "boundary term" will be zero for any  $n$  as an exponential grows faster than  $t^n$ .

However, we get an interesting recurrence relation:

$$L[t^n] = \int_0^{\infty} t^n e^{-st} dt = \frac{n}{s} \int_0^{\infty} t^{n-1} e^{-st} dt$$

By doing a couple terms we can even see how

$$L[t^n] = \frac{n!}{s^{n+1}}$$

Now, let's look into some other functions.

$$\begin{aligned} L[e^{kt}] &= \int_0^\infty e^{kt} e^{-st} dt \\ &= \int_0^\infty e^{-(s-k)t} dt \end{aligned}$$

Note that if  $s < k$ , the the integral blows up. So let's assume  $s > k$ ,

$$\begin{aligned} L[e^{kt}] &= \int_0^\infty e^{-(s-k)t} dt \\ &= \frac{1}{s-k} \end{aligned}$$

$$\begin{aligned} L[\sin(kt)] &= \int_0^\infty \sin(kt) e^{-st} dt \rightarrow \begin{bmatrix} u = \sin kt & v = -\frac{1}{s} e^{-st} \\ du = k \cos(kt) dt & dv = e^{-st} dt \end{bmatrix} \\ &= \cancel{-\frac{1}{s} \sin(kt) e^{-st}} \Big|_0^\infty + \frac{k}{s} \int_0^\infty \cos(kt) e^{-st} dt \end{aligned}$$

We now need to calculate  $L[\cos(kt)]$ ,

$$\begin{aligned} L[\cos(kt)] &= \int_0^\infty \cos(kt) e^{-st} dt \rightarrow \begin{bmatrix} u = \cos kt & v = -\frac{1}{s} e^{-st} \\ du = -k \sin(kt) dt & dv = e^{-st} dt \end{bmatrix} \\ &= -\frac{1}{s} \cos(kt) e^{-st} \Big|_0^\infty - \frac{k}{s} \int_0^\infty \sin(kt) e^{-st} dt \\ &= \frac{1}{s} - \frac{k}{s} \int_0^\infty \sin(kt) e^{-st} dt \end{aligned}$$

We get two interesting results here together. First,

$$\begin{aligned} L[\sin(kt)] &= \int_0^\infty \sin(kt) e^{-st} dt \\ &= \frac{k}{s} \int_0^\infty \cos(kt) e^{-st} dt \\ &= \frac{k}{s} \left( \frac{1}{s} - \frac{k}{s} \int_0^\infty \sin(kt) e^{-st} dt \right) \\ &= \frac{k}{s^2} - \frac{k^2}{s^2} \int_0^\infty \sin(kt) e^{-st} dt \end{aligned}$$

Which can be re-arranged into,

$$\begin{aligned}\left(1 + \frac{k^2}{s^2}\right) \int_0^\infty \sin(kt)e^{-st} dt &= \frac{s^2 + k^2}{s^2} \int_0^\infty \sin(kt)e^{-st} dt \\ &= \frac{k}{s^2}\end{aligned}$$

Meaning that,

$$\begin{aligned}L[\sin(kt)] &= \int_0^\infty \sin(kt)e^{-st} dt \\ &= \frac{k}{s^2 + k^2}\end{aligned}$$

And now that we have this result, we can go back to,

$$\begin{aligned}L[\cos(kt)] &= \int_0^\infty \cos(kt)e^{-st} dt \\ &= \frac{1}{s} - \frac{k}{s} \int_0^\infty \sin(kt)e^{-st} dt \\ &= \frac{1}{s} - \frac{k}{s} \frac{k}{s^2 + k^2} \\ &= \frac{1}{s} - \frac{k^2}{s(s^2 + k^2)} \\ &= \frac{s^2 + k^2 - k^2}{s(s^2 + k^2)} \\ &= \frac{s}{s^2 + k^2}\end{aligned}$$

## 2 Fourier Series

### 2.1 It all Adds Up

$$\frac{a_0}{2} + \sum_{n=1}^N (a_n \cos(2\pi nt) + b_n \sin(2\pi nt))$$

Given that

$$e^{2\pi int} = \cos(2\pi nt) + i \sin(2\pi nt)$$

and

$$e^{-2\pi int} = \cos(2\pi nt) - i \sin(2\pi nt)$$

Since  $\cos(-x) = \cos(x)$  while  $\sin(-x) = -\sin(x)$ .

The above also mean that

$$\cos(2\pi nt) = \frac{1}{2} (e^{2\pi int} + e^{-2\pi int})$$

and

$$\sin(2\pi nt) = \frac{1}{2i} (e^{2\pi int} - e^{-2\pi int}) = -\frac{i}{2} (e^{2\pi int} - e^{-2\pi int})$$

In the last step we used the fact that  $i^{-1} = -i$ . ( $1 = (-i)i = i \cdot i^{-1}$ ).

Using these expressions to rewrite our series of sines and cosines we get

$$\begin{aligned} \frac{a_0}{2} + \sum_{n=1}^N (a_n \cos(2\pi nt) + b_n \sin(2\pi nt)) &= \frac{a_0}{2} + \sum_{n=1}^N \left[ \frac{a_n}{2} (e^{2\pi int} + e^{-2\pi int}) + \frac{b_n}{2i} (e^{2\pi int} - e^{-2\pi int}) \right] \\ &= \frac{a_0}{2} + \sum_{n=1}^N \left[ \frac{a_n}{2} (e^{2\pi int} + e^{-2\pi int}) - \frac{ib_n}{2} (e^{2\pi int} - e^{-2\pi int}) \right] \\ &= \frac{a_0}{2} + \sum_{n=1}^N \left[ \frac{a_n}{2} e^{2\pi int} + \frac{a_n}{2} e^{-2\pi int} - \frac{ib_n}{2} e^{2\pi int} + \frac{ib_n}{2} e^{-2\pi int} \right] \\ &= \frac{a_0}{2} + \sum_{n=1}^N \left[ \frac{1}{2} (a_n - ib_n) e^{2\pi int} + \frac{1}{2} (a_n + ib_n) e^{-2\pi int} \right] \end{aligned}$$

If we look at the terms within the square braces we see that we have a complex number times and exponent and its complex conjugate,

$$c_n := \frac{1}{2} (a_n - ib_n)$$

and

$$\bar{c}_n = \frac{1}{2} (a_n + ib_n)$$

So

$$\frac{a_0}{2} + \sum_{n=1}^N (a_n \cos(2\pi nt) + b_n \sin(2\pi nt)) = \frac{a_0}{2} + \sum_{n=1}^N [c_n e^{2\pi int} + \bar{c}_n e^{-2\pi int}]$$

Here is where the additional requirement comes in. If we have  $\bar{c}_n = c_{-n}$ , then  $a_{-n} = a_n$  and  $b_{-n} = -b_n$ . We can then reindex our series,

$$\begin{aligned} \frac{a_0}{2} + \sum_{n=1}^N (a_n \cos(2\pi nt) + b_n \sin(2\pi nt)) &= \frac{a_0}{2} + \sum_{n=1}^N \left[ c_n e^{2\pi i n t} + \bar{c}_n e^{\overline{2\pi i n t}} \right] \\ &= \frac{a_0}{2} + \sum_{n=1}^N c_n e^{2\pi i n t} + \sum_{n=-1}^{-N} c_{-n} e^{-2\pi i n t} \end{aligned}$$

and if we have  $c_0 = \frac{1}{2}(a_0 - ib_0) = \frac{a_0}{2}$ , ( $b_0 = 0$ ), then we can see where

$$\sum_{n=-N}^N c_n e^{2\pi i n t}$$

comes from.

### 3 Fourier Series