FYS4560; ELEMENTARY PARTICLE PHYSICS

FINAL PROJECT

Higher Dimensions; Theoretical and Experimental Aspects

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Abstract

The Randall-Sundrum (RS) model of compacted dimensions will be studied, together with the postulate of the graviton particle. Quark-anti-quark production of the simplest possible massive graviton (1st order tower of Kaluza-Klein excitations) will be calculated.

1 Introduction

Higher dimensions, also called extra dimensions, are physical models for the dimensionality of our universe, mostly suggested with the goal of explaining the hierarchy problem of the standard model. There are many theories for higher dimensions. The most recognised are:

- Large extra dimensions: The often-heard theory that gravity acts through several dimensions, therefore becoming weaker. It originates from the ADD model as an attempt to solve the hierarchy problem¹.
- Warped extra dimensions: Describing our universe as a five-dimensional anti-de Sitter space, and claiming the SM particles are localized on a (3 + 1)-dimensional brane(s).
- Universal extra dimensions: All particles move universally through the extra dimensions, unlike to two other models where only gravity propagates through them.

Obviously, a thorough description of any of these models is near impossible for such a small paper, let alone all the models together. Therefore, a brief outline of the theory behind the two currently most promising² models will be given.

The first is the large extra dimension model by Arkani-Hamed, Dimopoulos, and Dvali (ADD). Originally, it was proposed as a model to explain the hierarchy problem (why the weak force is 10^{32} times stronger than gravity, among other problems). The extra dimensions³ are then suggested as planes into which gravity, assumed just as strong as the other forces, spreads. Therefore gravity becomes "diluted", while the known SM particles stay in (1,3)-spacetime.

The second model is the warped extra dimension model by Randall and Sundrum (RS), made due to disliking the current universal extra dimensions models. They assumed that, rather than having universal extra dimensions in which all particles propagate, there is a small extra dimension. This means they model our world as a 5-dimensional anti-de Sitter space⁴. By small, it means the extra dimension has a large curvature, or is *warped*. From general relativity, gravity and curvature are very much the same thing, and therefore the extra dimension, called the Planckbrane, can easily host gravitons.

A question that then springs to mind is why exactly gravitons and extra dimensions are connected (other than gravitons "carrying" gravity). If the standard model is expanded, but without inclusion of extra dimensions, to include a graviton field, then measuring it would be, at best, very optimistic⁵. Should any extra dimensional model be true, it would certainly be desirable to prove it by measurement. Finding a particle with the properties of the graviton would mean that at least *some* extra dimensions model is true.

2 Kaluza-Klein theory

The reason Kaluza-Klein theory is discussed is because one requires knowledge of the so-called "Kaluza-Klein towers"; massive excitations of an expansion $model^6$ of the spacetime metric.

As already mentioned, one, or several, extra spatial dimension(s) is assumed, but only one extra will be considered for the sake of simplicity. This means that the spacetime metric goes "up in size", i.e.

$$\tilde{g}_{ab}(x,\theta)) = \begin{pmatrix} g_{\mu\nu}(x,\theta) - B_{\mu}(x,\theta)B_{\nu}(x,\theta)\Phi(x,\theta) & B_{\mu}(x,\theta)\Phi(x,\theta) \\ B_{\nu}(x,\theta)\Phi(x,\theta) & -\Phi(x,\theta) \end{pmatrix}$$
(1)

where

¹The "hierarchy problem" is the problem in explaining why gravity and the weak force are so weak compared to QED and QCD.

²"Promising" in the sense that they provide measurable outcomes, and fit very well with what we already know from the standard model. The problem is that so far they have predicted nothing. Finding a graviton would possibly confirm either theory.

³Today, 6 dimensions is the most common expansion.

⁴This will be explained later on.

 $^{^5}$ Refer to "Can Gravitons be Detected?" by Rothman and Boughn (https://arxiv.org/pdf/gr-qc/0601043v3.pdf) for an impression of the problem.

⁶This expansion depends on the model, and is where RS and ADD differ.

- Latin indices represent all dimensions, greek indices represent 4-dimensional spacetime.
- $g_{\mu\nu}$ is the 4-dimensional spacetime metric from before.
- B_{μ} is a gauge field, defined by $B_{\mu} \equiv \frac{g_{\mu 5}}{g_{55}}$
- Φ is a scalar field, defined by $\Phi = g_{55}$.
- x are the normal spacetime coordinates, while θ is the extra coordinate. The notation is due to the extra dimension(s) often being "circular" (such as the S^1 space).

As can be seen, a complete description depends on $g_{\mu 5}$ and g_{55} . However, this is precisely where the complication of extra dimensions arises. The metric depends on the space, and the shape of the space is what we do not know⁷.

Firstly, consider a simple model with a massive scalar field $\hat{\phi}(x,\theta)$,

$$S_5 = \int dx^4 d\theta \sqrt{-\tilde{G}} \left(-\frac{1}{2} (\partial_a \hat{\phi}) (\partial^a \hat{\phi})^{\dagger} - \frac{m^2}{2} \hat{\phi}^2 - \frac{c_5}{4!} \hat{\phi}^4 \right) \tag{2}$$

where $\tilde{G} \equiv \det(\tilde{g}_{ab})$ and c_5 is some coupling constant we do not yet know. The scalar field can be Fourier expanded as

$$\hat{\phi}(x,\theta) = \sum_{\substack{n = \{n_1, n_2, \dots, n_5\};\\ n_1 \in \mathbb{Z} \ \forall \ i}} \phi(n)(x) Y_n(\theta) \tag{3}$$

where $Y_n(\theta)$ are orthogonal, normalized eigenfuntions of the Laplace operator on the internal space⁸:

$$\Delta_{K_d} Y_n(\theta) = \frac{\lambda_n}{R^2} Y_n(\theta) \tag{4}$$

where R is the "characteristic size" of the space K_d (obviously a radius in the case of circular objects). Inserting this into equation 2 and performing the integral over K_d gives

$$S_{5} = \int dx^{4} \sqrt{-\tilde{G}} \left(-\frac{1}{2} (\partial_{\mu} \phi^{(0)}) (\partial^{\mu} \phi^{(0)})^{\dagger} - \frac{m^{2}}{2} (\phi^{(0)})^{2} - \sum_{n \neq 0} \left[(\partial_{\mu} \phi^{(n)}) (\partial^{\mu} \phi^{(n)})^{\dagger} + m_{n}^{2} \phi^{(n)} \phi^{(n)\dagger} \right] - \frac{c_{5}}{4!} (\phi^{(0)})^{4} - \frac{c_{5}}{4} (\phi^{(0)})^{2} \sum_{n \neq 0} \phi^{(n)} \phi^{(n)\dagger} - \dots \right),$$

$$(5)$$

and by comparison one sees $m_n^2 = m^2 + \frac{\lambda_n}{R^2}$. The dots represent all terms without any $\phi^{(0)}$ -coefficients. As one sees, there is an infinite number of masses, since there is an infinite number of eigenvalues; a mathematical problem without any current answer. The hope is that these modes⁹ can be tied to the graviton, and will be explained better in the following sections.

A relation that will be of importance later is the reduction formula:

$$M_{Pl}^2 = V_d M^{d+2} (6)$$

where $M_{Pl}=G_{N(4)}^{-1/2}$ is the 4-dimensional Planck mass and $M^{d+2}=G_{N(4+d)}^{-1/2}$ is the fundamental mass scale in the new model. The relation is derived by demanding the Einstein-Hilbert action to be the same

⁷This is why the field of extra dimensions can be so large and complex, as there are many ways to guess at a shape of the new dimension(s), and each model gives a characteristic outcome which must fit experiments at, for example, the LHC.

⁸The extra dimension/space, denoted K_d .

⁹The infinite set of which is called the Kaluza-Klein tower of modes. (Often one hears of KK towers.)

with and without the new dimension(s) 10 , and performing the integrals by using KK mode expansion on the integrand. I.e one demands:

$$S_{E(4)} = S_{E(d+4)} \tag{7}$$

where

$$S_{E(D)} = \int d^{D}x \sqrt{-G_{D}} \frac{1}{16\pi G_{N(D)}} \mathcal{R}^{D}(G_{ab})$$
(8)

3 The Arkani-Hamed-Dimopoulos-Dvali model

The ADD model is mainly based on 3 features:

- There exists n new spatial compact dimensions, with compactification volume V_n .
- The Planck scale is very low, at the order of one TeV,
- The SM degrees of freedom are localized on a 3D-brane, stretching along 3 non-compact spatial dimensions (i.e. the SM particles move in normal spacetime, not in the new dimension(s)).

In the ADD model, the standard metric is expanded upon by introducing a field $\hat{h}_{ab}(x,\theta)$ such that:

$$\tilde{g}_{ab}(x,\theta) = \eta_{ab} + \frac{2}{M^{1+d/2}}\hat{h}_{ab}(x,\theta) \tag{9}$$

4 The Randall-Sundrum model

4.1 The hierarchy problem in RS1

The RS model assumes that there are two points on the S^1/\mathbb{Z}^2 ordibfold in which 3-branes are compactified. What this means is that there are two 3-dimensional, non-compact branes that are connected at every point by a "circle", see figure SOME FIG.

The two branes, 1 and 2, are located on points $\theta = 0$ and $\theta = \pi$, which means the branes are separated by a distance 2L = 2R, where R is the circle radius.

As mentioned, the model set out to explain the hierarchy problem, and will be briefly explained how below. Firstly, the action of the model is given by the sum of the Hilbert-Einstein action and the matter "part":

$$S_5 = S_E + S_M = \int d^4x \int_{-L}^{L} dy \sqrt{-\tilde{G}} (M^3 \mathcal{R}_5 - \Lambda_5)$$
 (10)

where Λ_5 is the five dimensional cosmological constant. In order to match real world observations, the new metric must uphold Poinecaré invariance, which can be shown to lead to:

$$ds^{2} = e^{-2A(y)} \eta_{\mu\nu} dx^{\mu} dx^{\nu} + dy^{2}$$
(11)

where A(y) is called the warp factor. Solving the Einstein equations lets one find that $A(y)=\pm ky$, where $k\equiv\sqrt{\frac{-\Lambda}{12M^3}}$ is a constant. Since the orbifold abides \mathbb{Z}^2 symmetry, $A(y)=A(-y)\to A(y)=k|y|$. There is a problem with action above, which is that it does not include the energy densities of the two branes, that are:

$$S_1 = \int_{B_1} \int_{S^1/\mathbb{Z}^2} d^4x dy \sqrt{-\tilde{G}(x,y)} \lambda_1 \delta(y)$$
(12)

$$S_2 = \int_{B_2} \int_{S^1/\mathbb{Z}^2} d^4x dy \sqrt{-\tilde{G}(x,y)} \lambda_2 \delta(y-L)$$
(13)

¹⁰This is because the new dimension(s) must reproduce what we observe, and 4D spacetime fits very well with observations.

where $\tilde{G}(x,0) = G_1$ and $\tilde{G}(x,L) = G_2$ have been used. So, the total action for the space is $S = S_E + S_M + S_1 + S_2$. To fulfil the Einstein equations, one needs to impose $\lambda_1 = -\lambda_2 = 12kM^3$. THINK ON THIS

4.2 KK mode expansion/Finding gravitons

Much in the same manner as for QED and other gauge theories, one can create gravitational bosons. This means doing a metric transform, given by:

$$ds^{2} = e^{-2k|y|} (\eta_{\mu\nu} + \tilde{h}_{\mu\nu}(x,y)) dx^{\mu} dx^{\nu} + (1+\phi(x)) dy^{2}$$
(14)

However, this means there will be cross-terms ($\phi \tilde{h}_{\mu\nu}$ -terms) in the Lagrangian that prevent mass eigenstates. It can be shown¹¹ that it can be diagonalized. One will end up with two fields, $h_{\mu\nu}$ and φ , which will be used hereafter.

Then, a KK mode expansion is done on $h_{\mu\nu}(x,y)$, giving

$$h_{\mu\nu}(x,y) = \sum_{n=0}^{\infty} h_{\mu\nu}^{(n)}(x) \frac{\chi_n(y)}{R},$$
(15)

where $\chi_0(y) = 2\sqrt{kR}e^{-2k|y|}$ and

$$\chi_n(y) = N_n \left[C_1 Y_2 \left(\frac{m_n}{k} e^{k|y|} \right) + C_2 J_2 \left(\frac{m_n}{k} e^{k|y|} \right) \right] \quad , \quad n \neq 0$$
 (16)

While this seems out of the blue, this is the result of diagonalizing the Lagrangian, which is a rather intricate procedure. The functions χ_n are eigenfunction of an equation met during the diagonalising of the Lagrangian. J_2 and Y_2 are the Bessel functions of the first and second kind, respectively. The constants C_1 and C_2 can be determined by the boundary conditions on the vacuum energy terms of the energy-momentum tensor (the delta functions),

$$T_{ab} = \lambda_1 \sqrt{G_1} g_{\mu\nu}^{(1)} \delta_a^{\mu} \delta_b^{\nu} \delta(y) + \lambda_2 \sqrt{G_2} g_{\mu\nu}^{(2)} \delta_a^{\mu} \delta_b^{\nu} \delta(y - \pi R), \tag{17}$$

which give $C_1 = Y_1\left(\frac{m_n}{k}\right)$ and $C_2 = -J_1\left(\frac{m_n}{k}\right)$ SOMETHING MORE HERE

The $h_{\mu\nu}^{(0)}(x)$ field describes the massless graviton, while the $n \ge 1$ states are the massless KK modes, and finally φ describes the massless radion.

5 Graviton production at the LHC

6 Conclusions

 $^{^{11}\}mathrm{See}$ http://arxiv.org/abs/hep-th/0105304v3, section 2, for a rigorous derivation.