

Baroclinic Instabilities in the Wave-Forced Ocean Mixed Layer

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¹ Keck Center, Room 254 (Come say hello!)

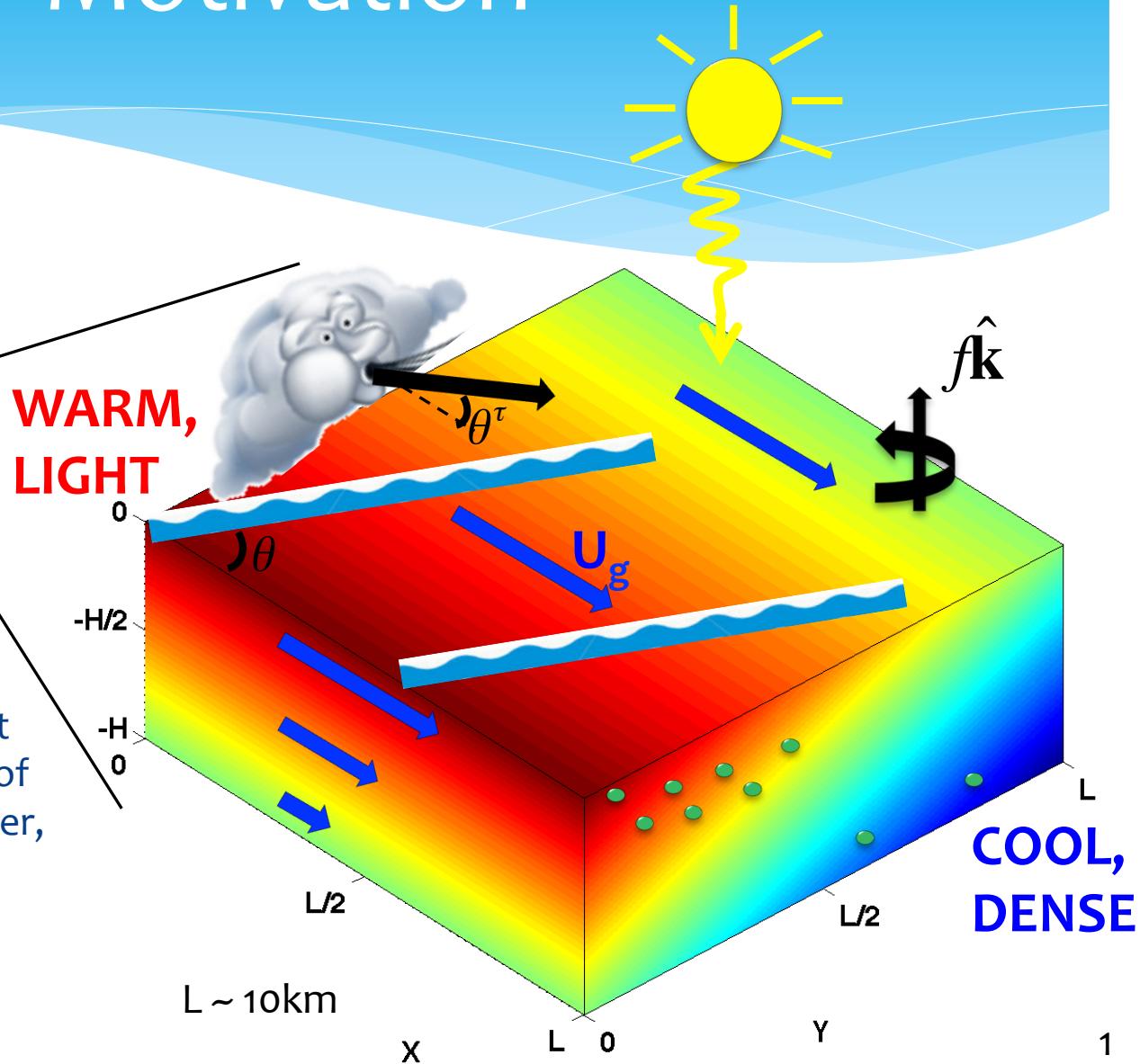
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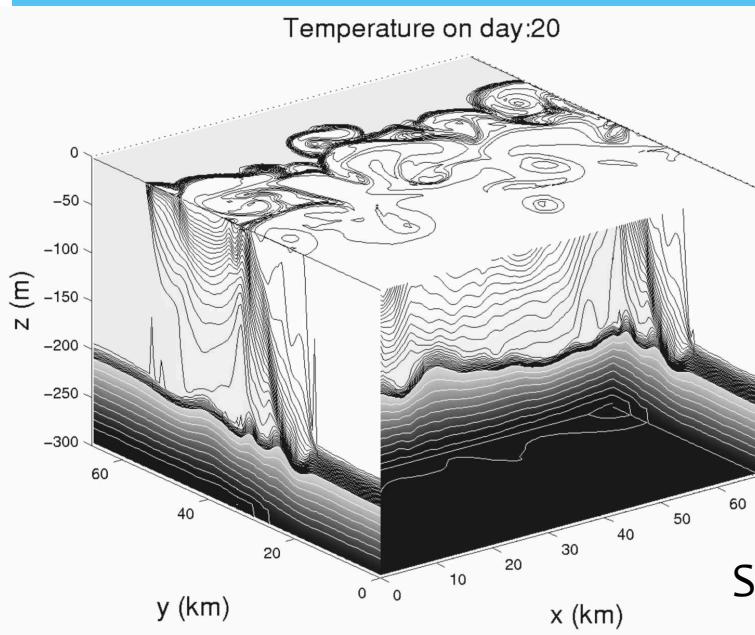
Motivation



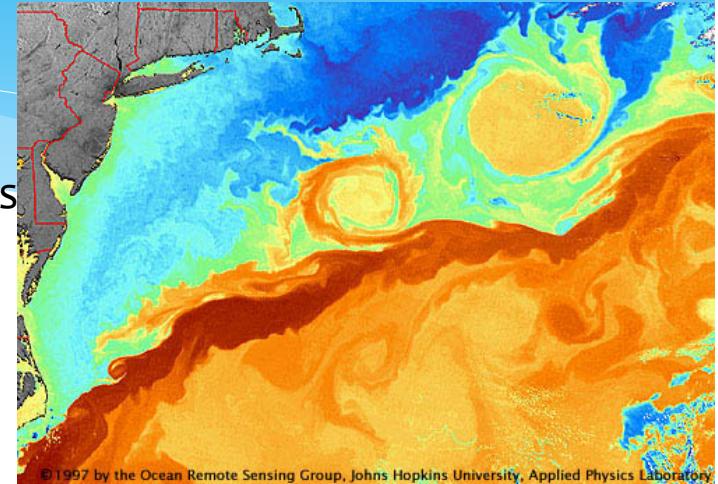
- * Coarse global models do not accurately represent fluxes of heat, momentum, fresh water, and gases.
- * Phytoplankton blooms are influenced by mixed layer dynamics

What's Out There?

Restratifying:



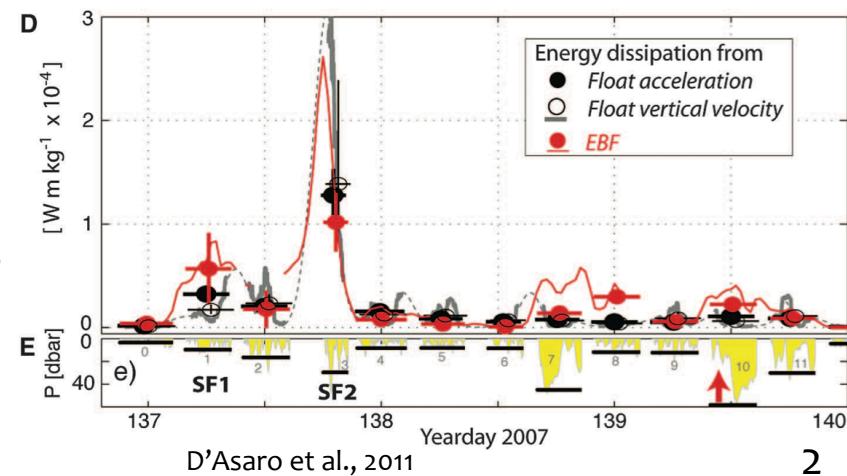
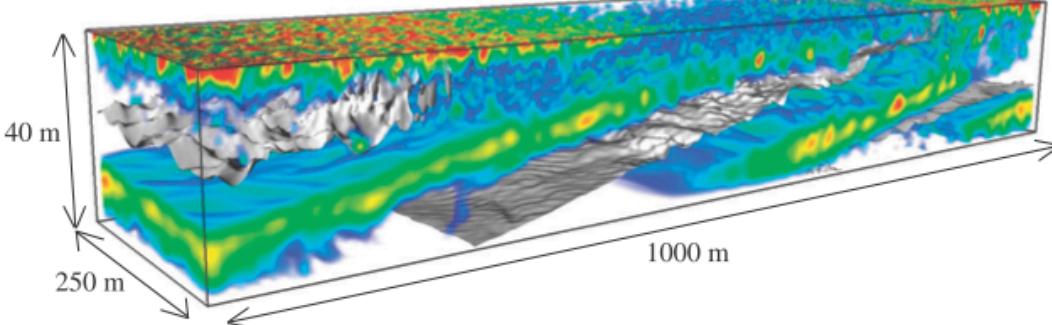
Mixed layer eddies
(Geostrophic
Instabilities; GI)



Symmetric instabilities (SI)

Fox-Kemper and Ferarri 2008

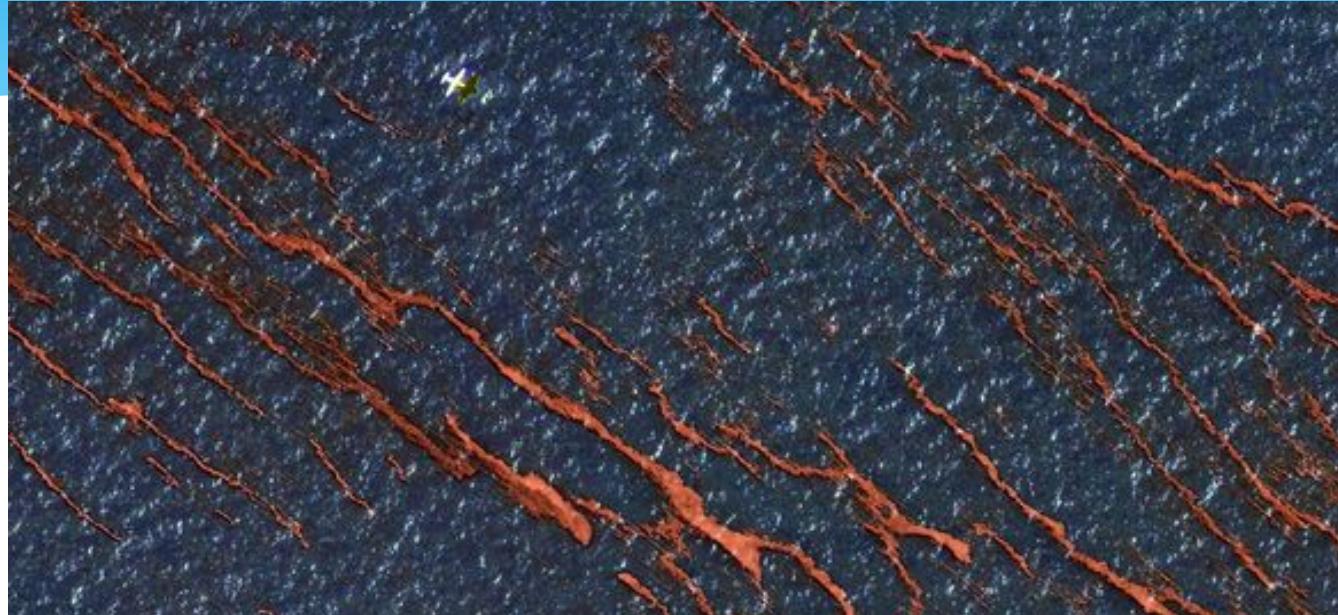
Taylor and Ferarri 2010



What's Out There?

Langmuir Circulation/mixing/turbulence (LC)

Deep Water
Horizon oil
slick. Digital
Globe (2010)



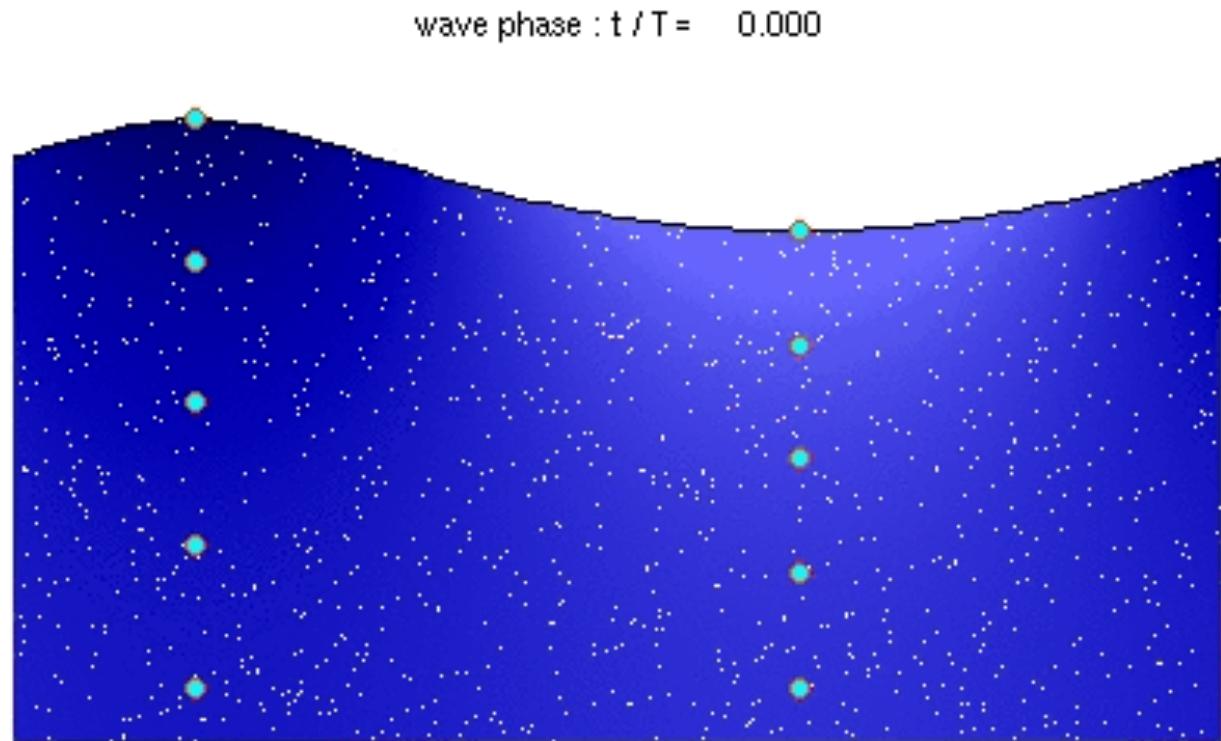
A lake in
Chilean
Patagonia



Wave-Induced Currents

$$U^s = \frac{1}{T} \int_0^T [(\text{displacements}) \cdot \nabla] (\text{wave velocity}) dt$$

- Stokes drift is the wave-averaged velocity following a particle.
- Lagrangian, not Eulerian.
- Decays steeply with depth.



The Wave-Averaged Boussinesq Equations

$$\partial_t \mathbf{u} + (\mathbf{u}^L \cdot \nabla) \mathbf{u} + \hat{f} \mathbf{k} \times \mathbf{u}^L + \frac{\nabla p}{\rho_0} + u^{L,j} \nabla U^{S,j} = b \mathbf{k} + \nu \nabla^2 \mathbf{u}$$

Acceleration Lagrangian Advection Coriolis and Stokes Coriolis Pressure Stokes Shear Buoyancy Dissipation

$$\partial_t \mathbf{u} + (\mathbf{u}^L \cdot \nabla) b = 0$$

$$\nabla \cdot \mathbf{u} = 0$$

Where, $\mathbf{u}^L = \mathbf{u} + \mathbf{U}^S$, and $\mathbf{U}^S(z)$ is prescribed

The Stokes Shear Force

- * The Stokes shear force knocks the vertical momentum out of hydrostatic equilibrium.
- * Horizontal variations in the Eulerian velocity induce LC.

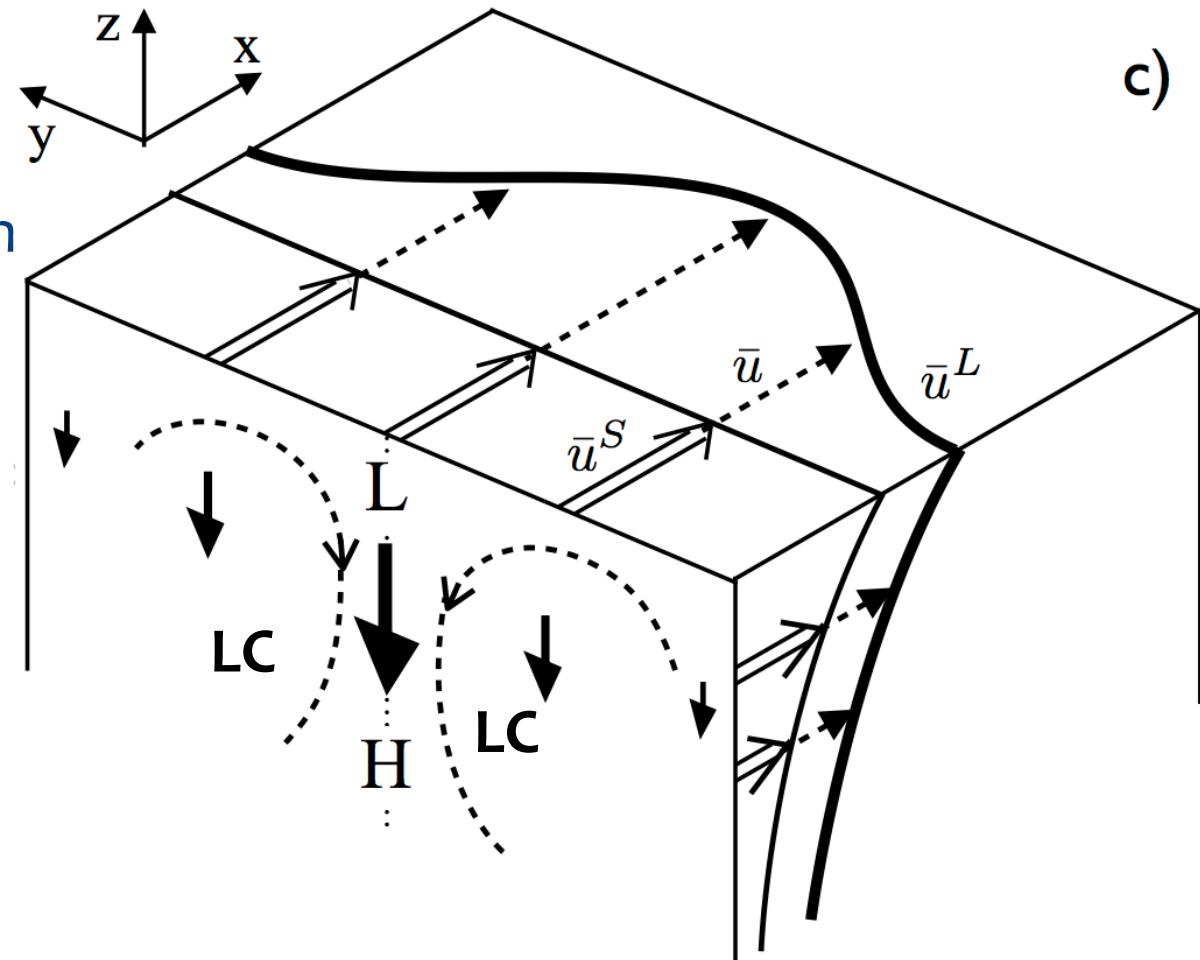
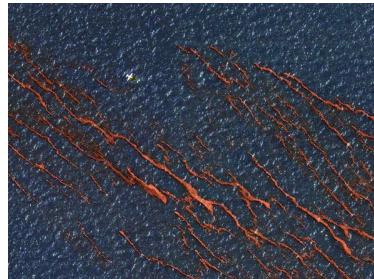


Figure 1c from Suzuki and Fox-Kemper 2015

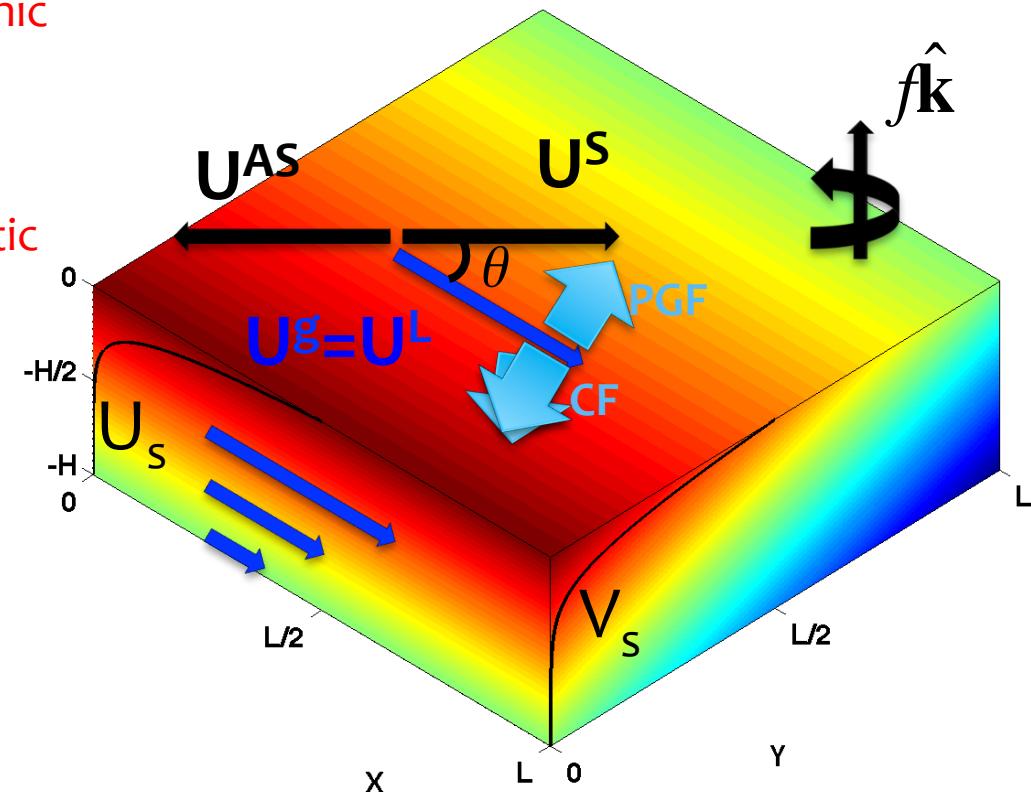
Lagrangian Thermal Wind Balance and Anti-Stokes Flow

Background Flow

$$f(\hat{\mathbf{k}} \times \bar{\mathbf{U}}^L) = -\frac{\nabla \bar{P}}{\rho_0} \quad \} \text{ Geostrophic}$$

$$\bar{P}_z = \bar{B} \quad \} \text{ Hydrostatic}$$

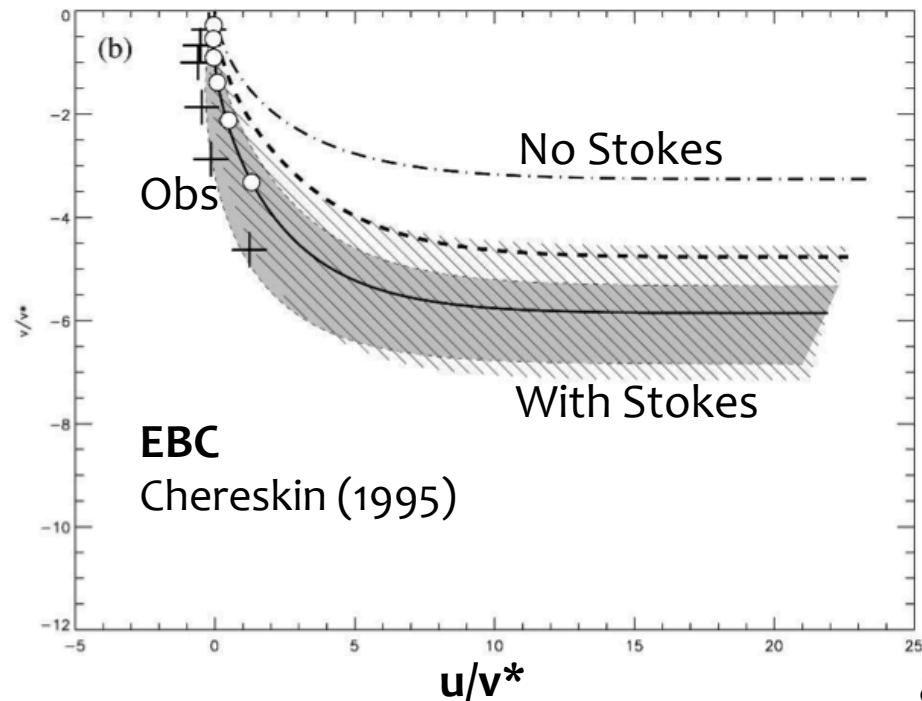
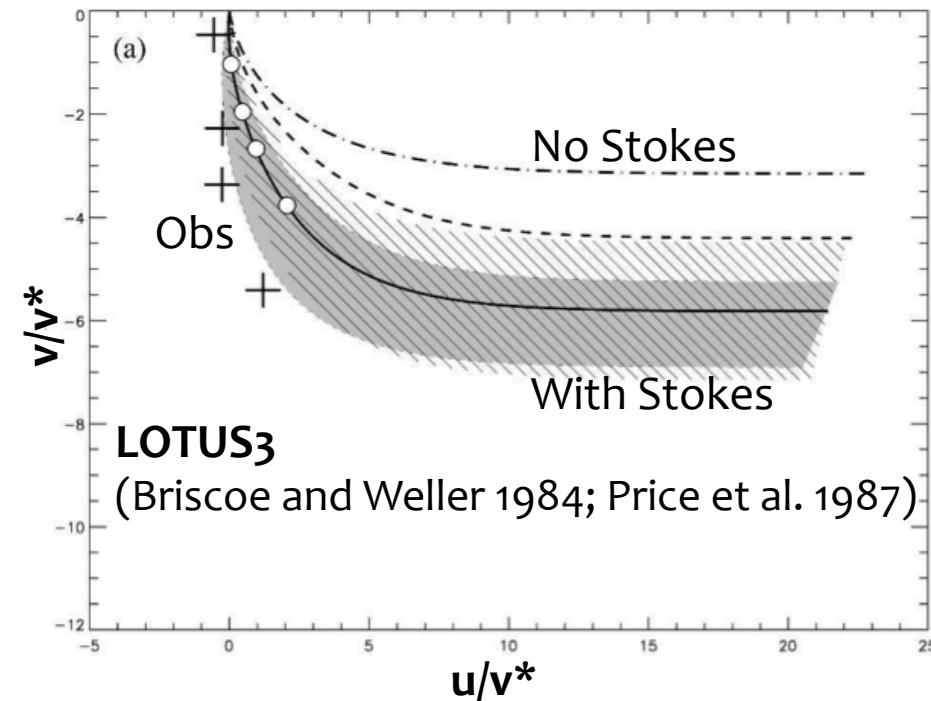
$$f \bar{U}_z^L = -\bar{B}_Y \quad } \text{ Thermal Wind}$$



Observational Evidence of “Anti-Stokes Flow”

- * Accounting for Stokes drift in the Ekman spiral may account for differences between analytic solutions and observations

Polton et al. 2005



Geostrophic Instabilities

- * Switching A and C is stable.
- * Switching A and B is unstable.
- * Once the flow gets moving, the Coriolis force turns it to the right

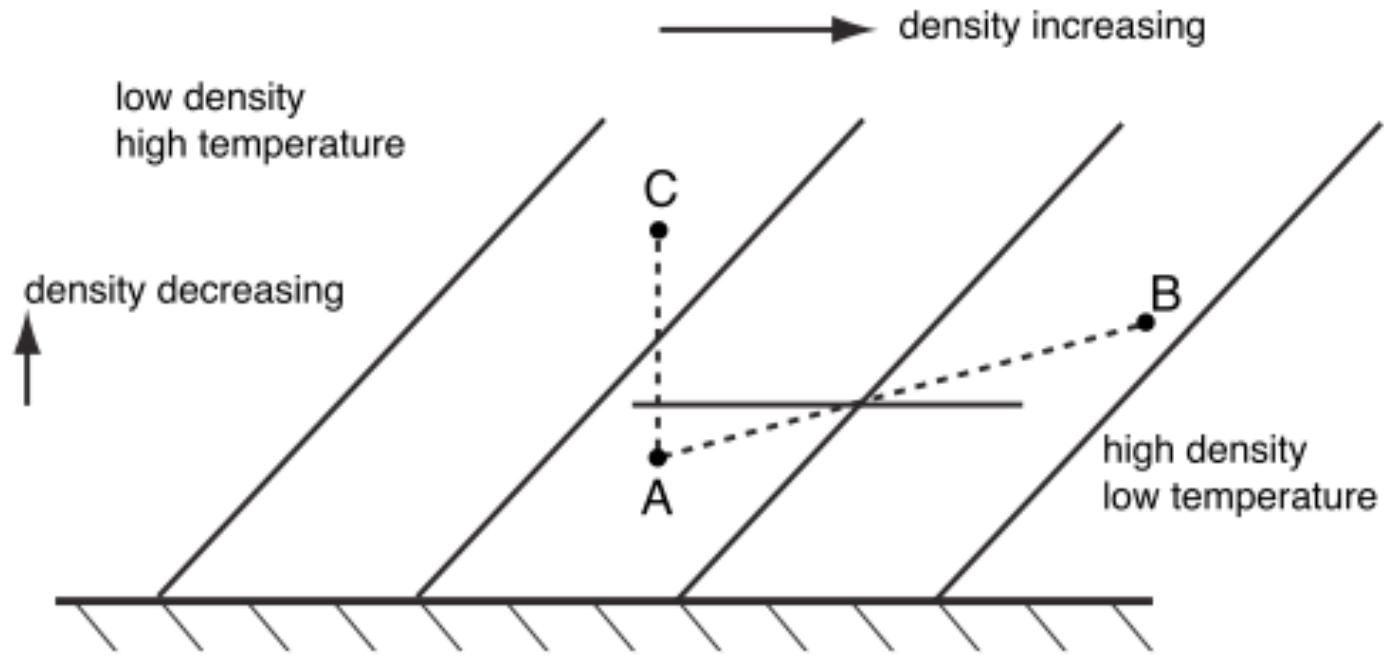
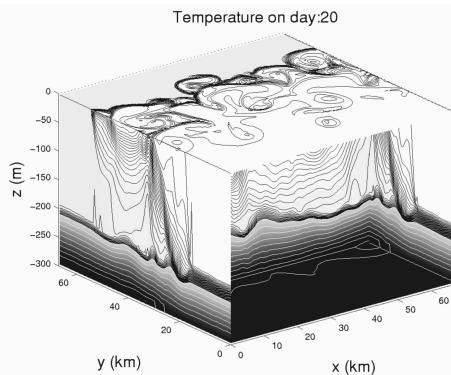


Figure 6.9 from Vallis, 2006

Analytic Stability Criteria: Geostrophic Instabilities

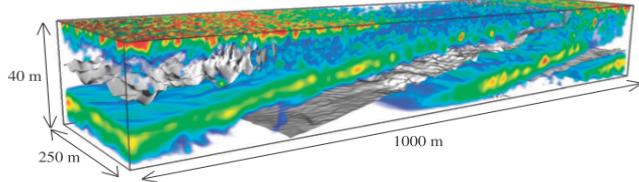
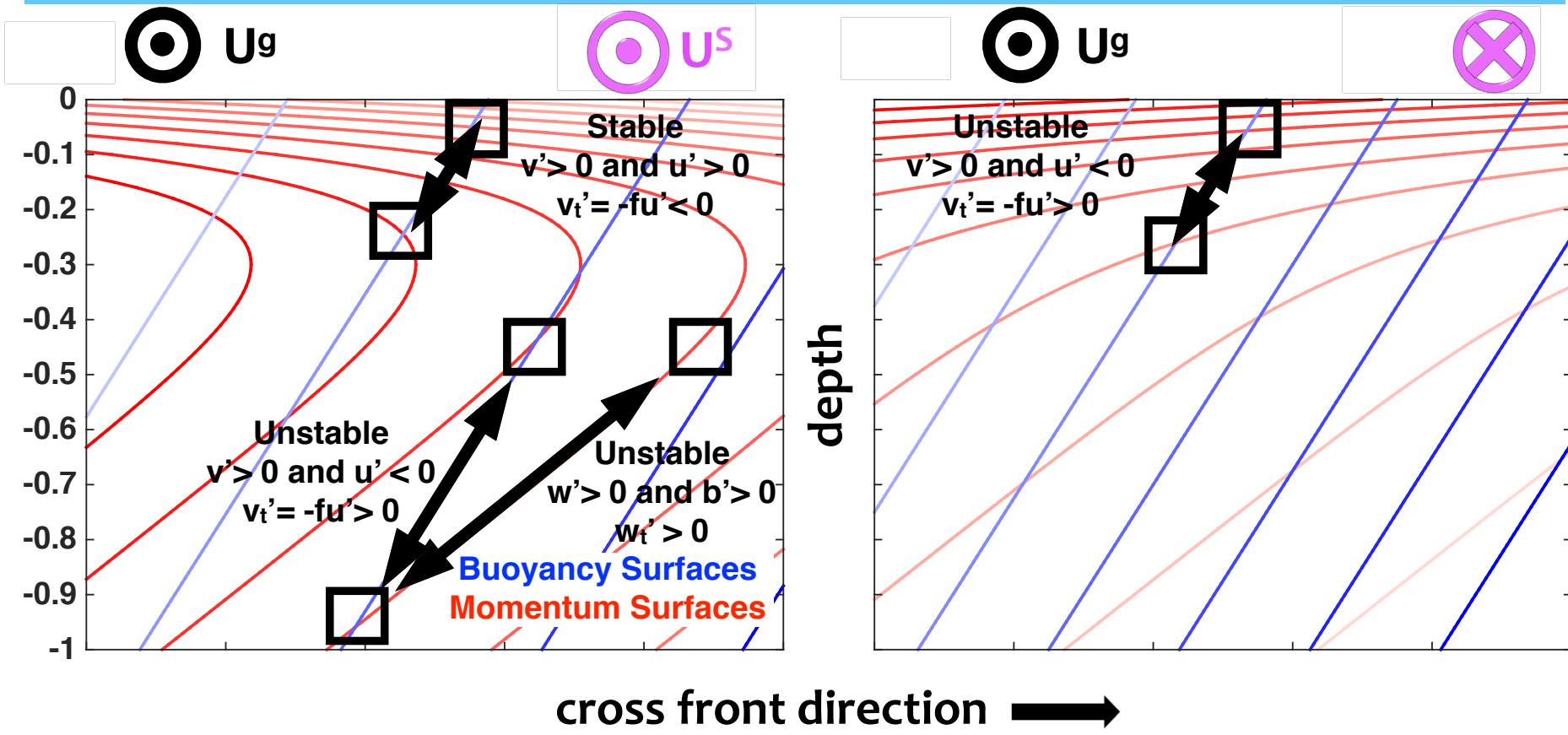
* Charney, Stern, and Pedlosky showed, that geostrophic instability exists only if any of the following is true:

1. Q_y changes sign in the interior of the domain.
2. Q_y is the opposite sign to U_z^L at the surface.
3. Q_y is the same sign to U_z^L at the bottom.
4. U_z^L has the same sign at the surface and bottom.

Where Q is the quasi-geostrophic potential vorticity:

$$\bar{Q} = \nabla_H^2 \bar{\psi} + \beta Y + \partial_z \left(\frac{f_0^2}{N^2} \bar{\psi}_z^L \right)$$

Symmetric Instabilities



Analytic Stability Criterion: Symmetric Instability

- * The Ertel potential vorticity (PV) depends on the alignment of the vorticity and buoyancy gradient.

$$\bar{Q} = (\nabla \times \bar{\mathbf{U}} + f\hat{\mathbf{k}}) \cdot \nabla \bar{B}$$

- * The Stokes drift (Lagrangian mean of the leading order, *irrotational* wave velocity) produces no vorticity. However, if the flow is in Lagrangian thermal wind balance, the vorticity is modified by the anti-Stokes Eulerian flow.
- * Hoskins (1974) showed that if a front in thermal wind balance is symmetrically unstable, the PV must be negative.
- * Extends to flows in Lagrangian thermal wind balance in the special case:

$$U^S = \mu z, \quad V^S = 0$$

$$SI \Rightarrow f\bar{Q} = \underbrace{f^2 N^2 - M^4}_{\text{Geostrophic}} - \underbrace{fM^2 U_z^S}_{\text{Anti-Stokes}} < 0$$

Linear Stability Method

- * Rescale equation
- * Multiple scales of horizontal variation: x, X, y, Y, t, T .
- * Decompose into mean and perturbation:

$$\mathbf{u} = \bar{\mathbf{U}}(X, Y, z, T) + \mathbf{u}'(x, X, y, Y, z, t, T)$$

- * Find a solution to the mean (averaged over x,y) equations
- * Force the perturbation equations with the mean flow solution
- * Assume

$$u' = \tilde{u}(z) e^{i(kx+ly+\sigma t)}$$

Vertical structure of instabilities

Growth rate of instabilities

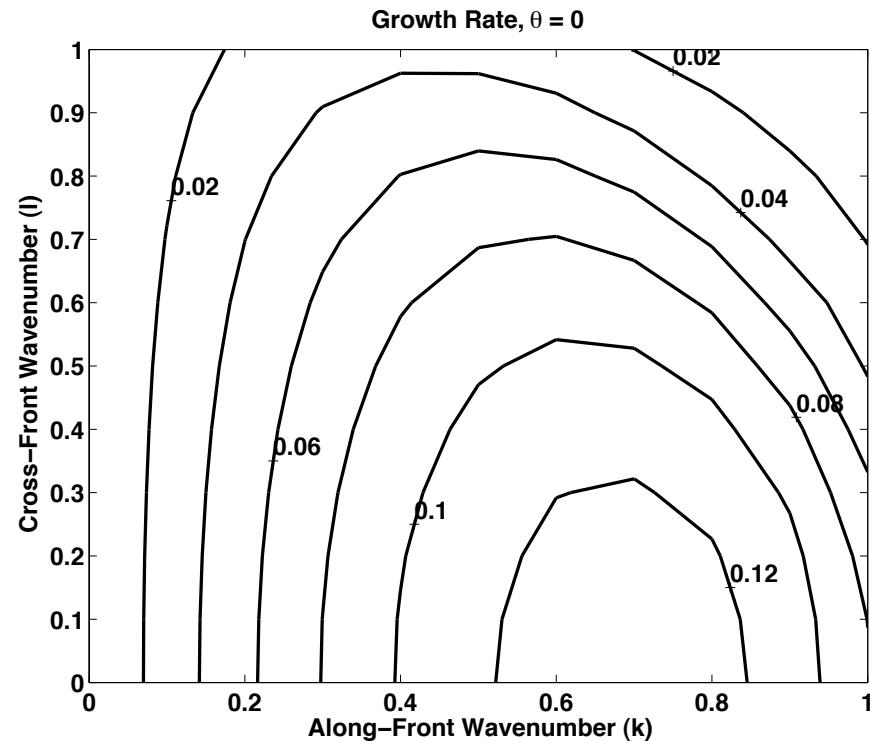
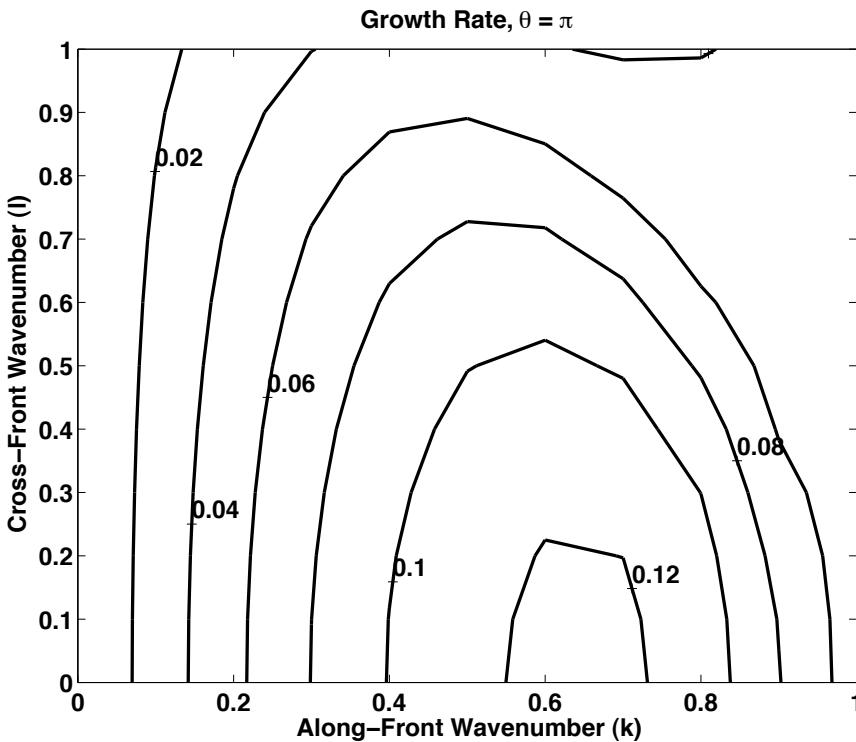
$$\tilde{u}, \sigma = F(Ri, \mu\lambda, \gamma, \lambda, \theta, Ek, Ro, \alpha)$$

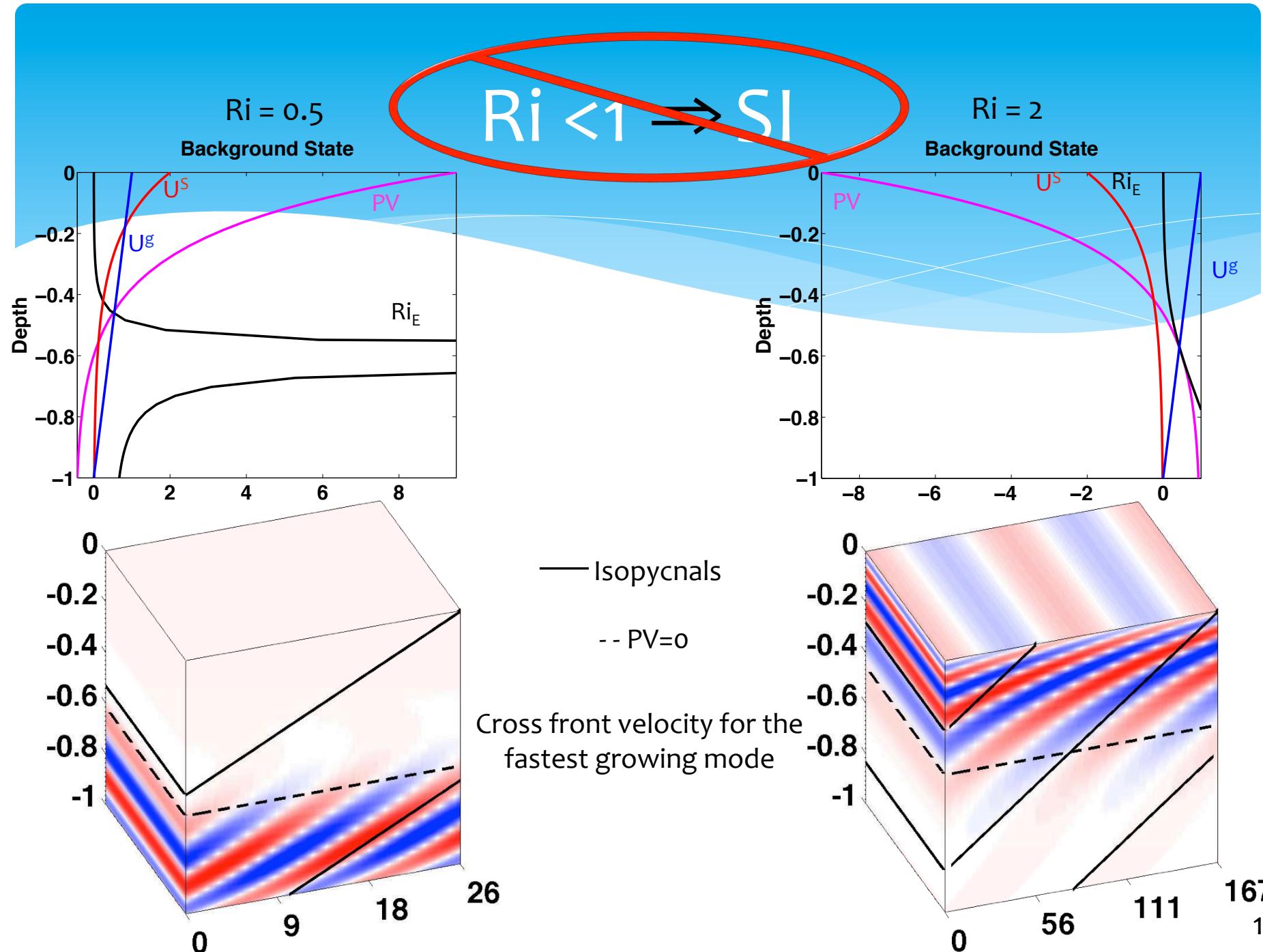
$$Ri = \frac{N^2}{U_z^{L,2}} = \frac{\text{Vertical Stratification}}{\text{Lagrangian Shear Squared}}$$

$$\mu\lambda = \frac{U_z^S}{U_z^L} = \frac{\text{Stokes Shear}}{\text{Lagrangian Shear}}$$

Geostrophic Instabilities

- * When the Stokes drift and geostrophic flow are aligned, the anti-Stokes flow yields reduced Eulerian shear.
- * Less Eulerian shear near the surface results in higher growth rates and wavenumbers for GI.





Energetics

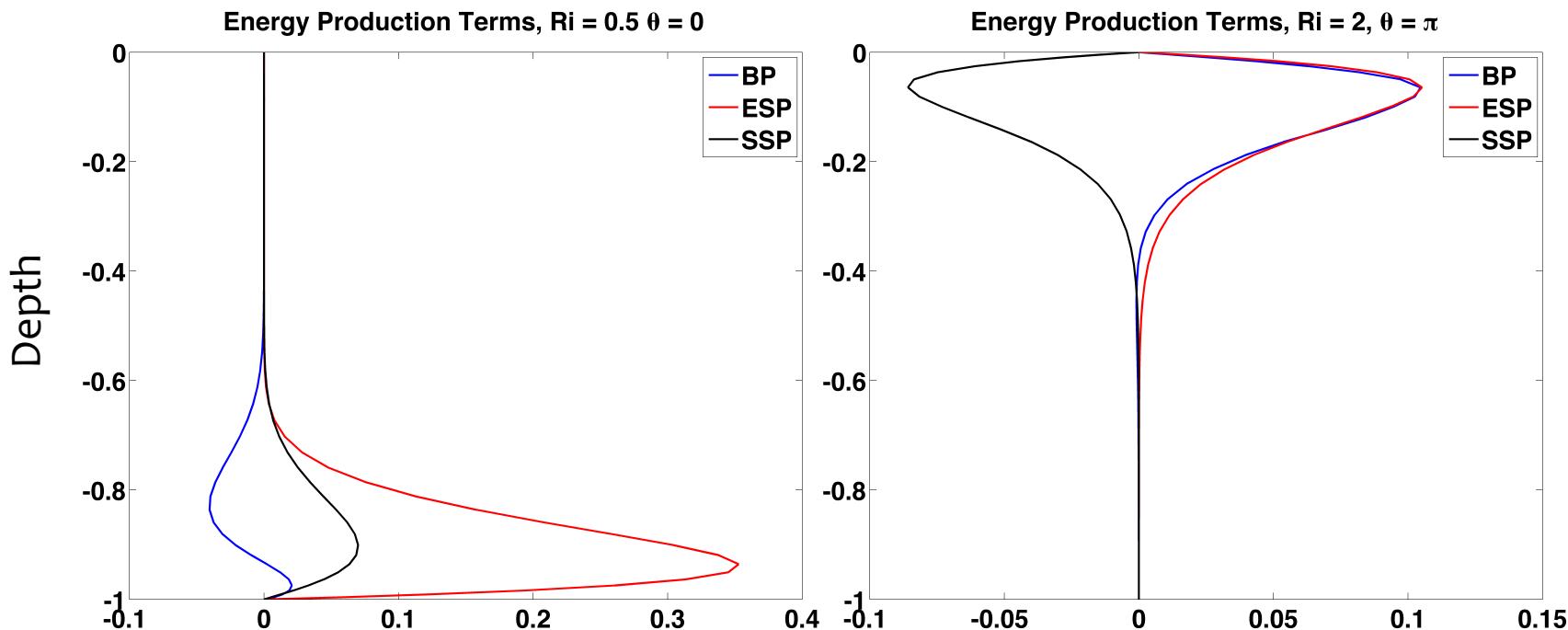
- * Energetics are a useful tool to distinguish modes.

$$\frac{\overline{D^L e'}}{Dt} = \underbrace{-\overline{\mathbf{u}' w'} \cdot \overline{\mathbf{U}_z}}_{\text{ESP}} - \underbrace{-\overline{\mathbf{u}' w'} \cdot \overline{\mathbf{U}_z^S}}_{\text{SSP}} - \underbrace{\overline{w' b'}}_{\text{BP}} - PW + D$$

- * BP dominant: instability extracts potential energy to RE-stratify the mixed layer (typical of GI).
- * SSP, ESP dominant: instability extracts kinetic energy (typical of SI, LC, KH)
- * Hybrid modes with various mixed of energy production terms exist.

Stokes Drift Induces more Restratiification by SI

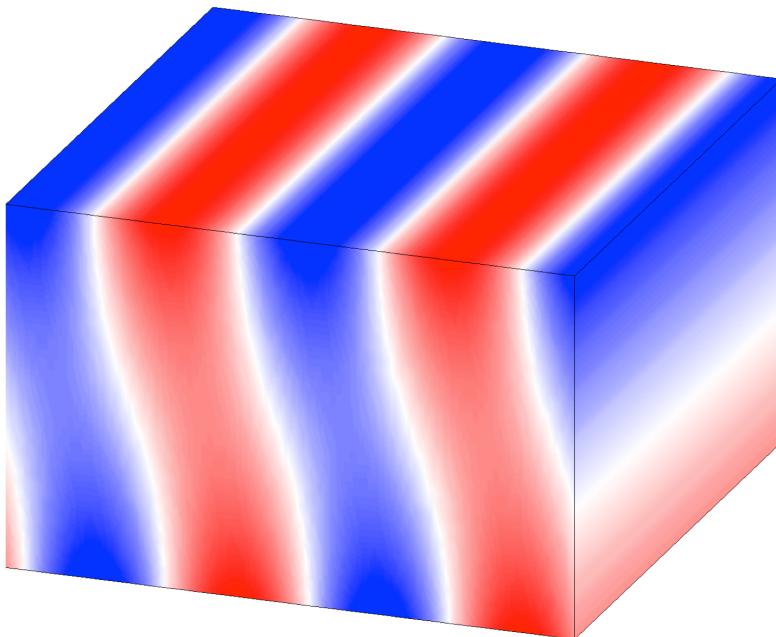
- * Stokes drift changes the path along which SI move, favoring more cross isopycnal motion near the surface.
- * This increases BP (restratification).
- * Anti-aligned Stokes drift \Rightarrow SSP < 0 (the work done by the Stokes shear force).



Stokes-Modified GI

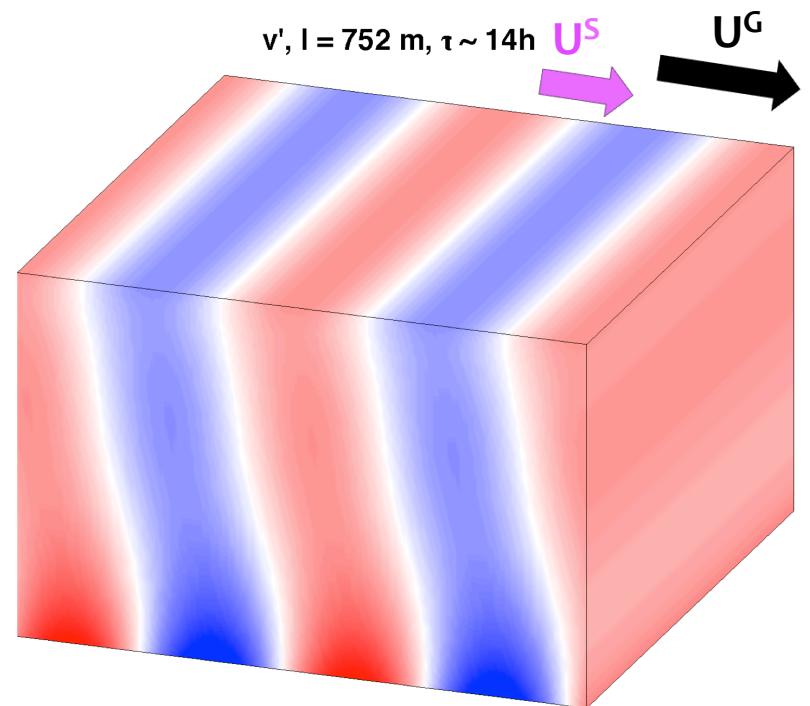
No Stokes

$v', l = 1233 \text{ m}, \tau \sim 16\text{h}$

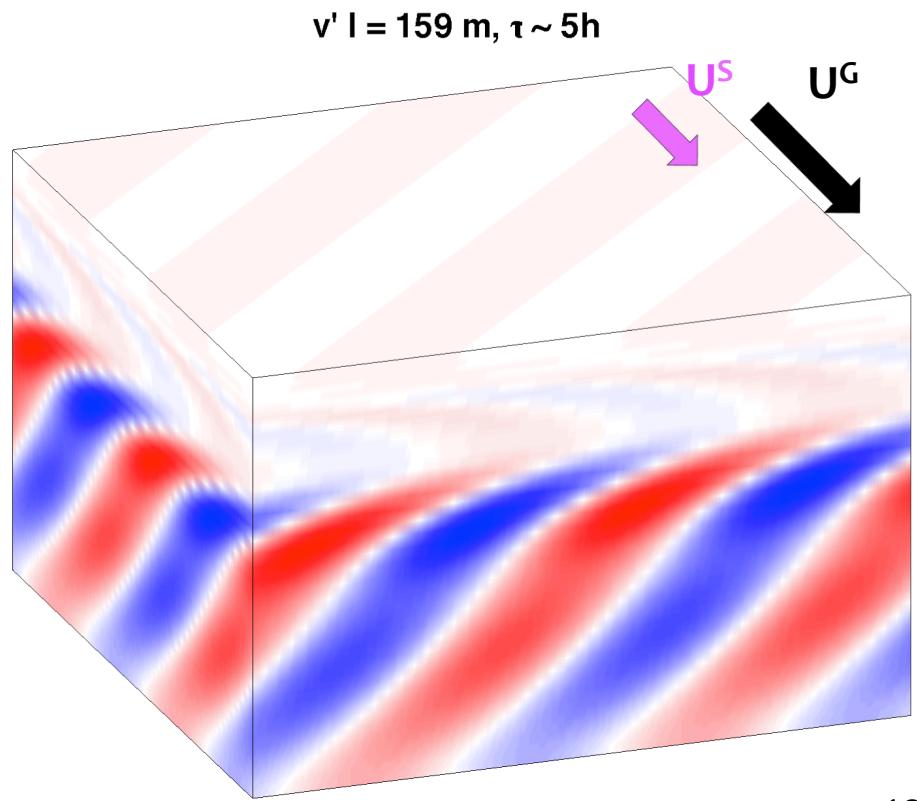
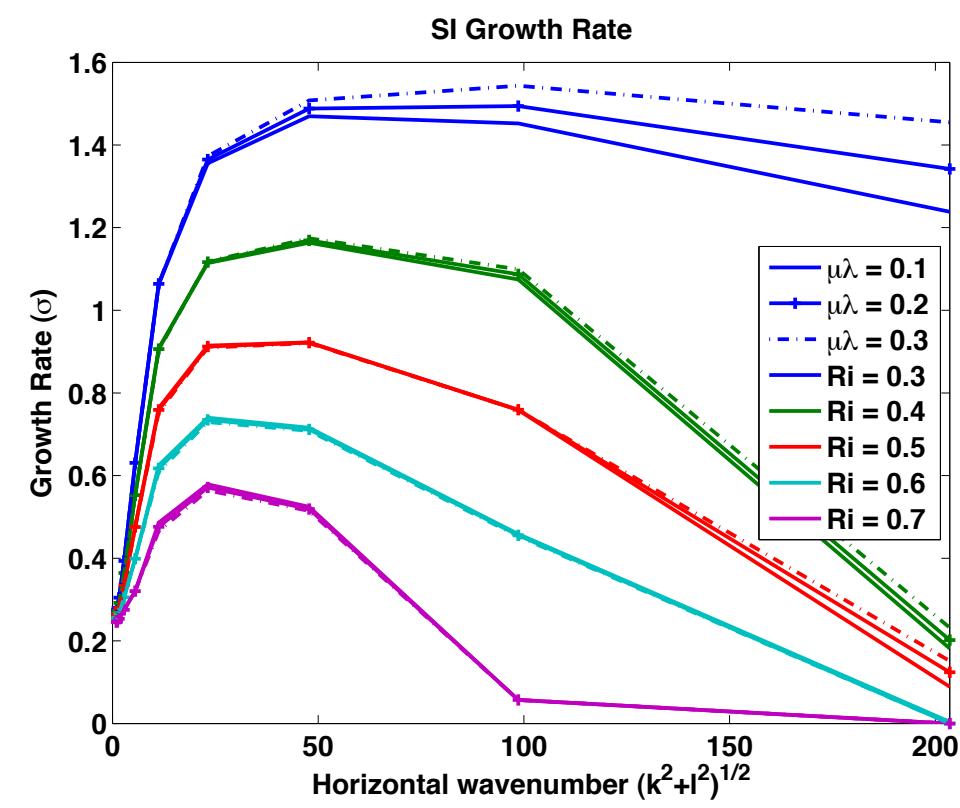


with Stokes

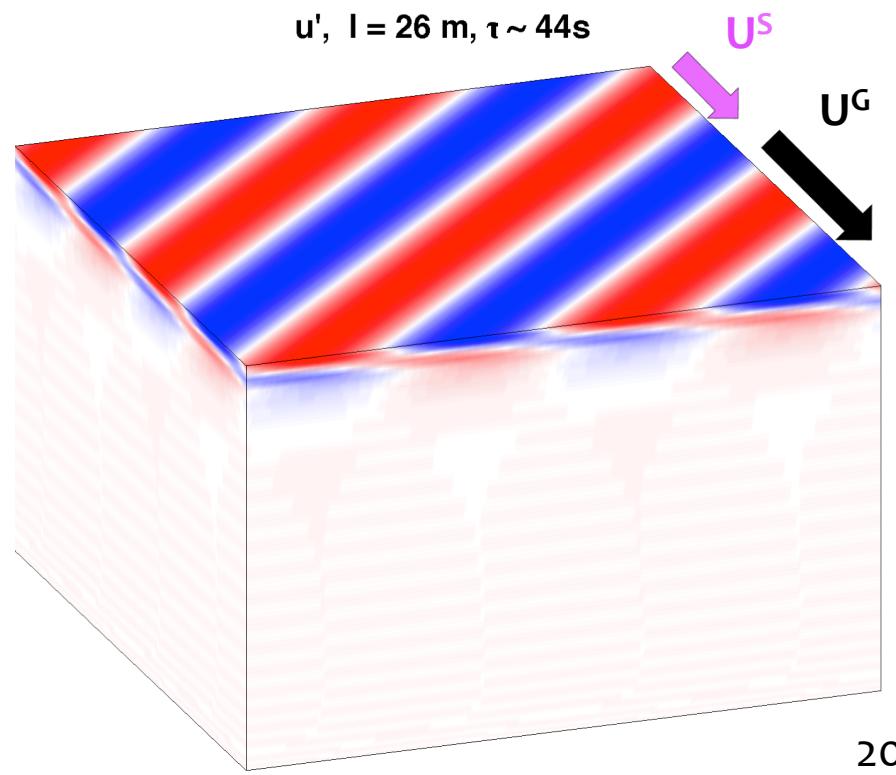
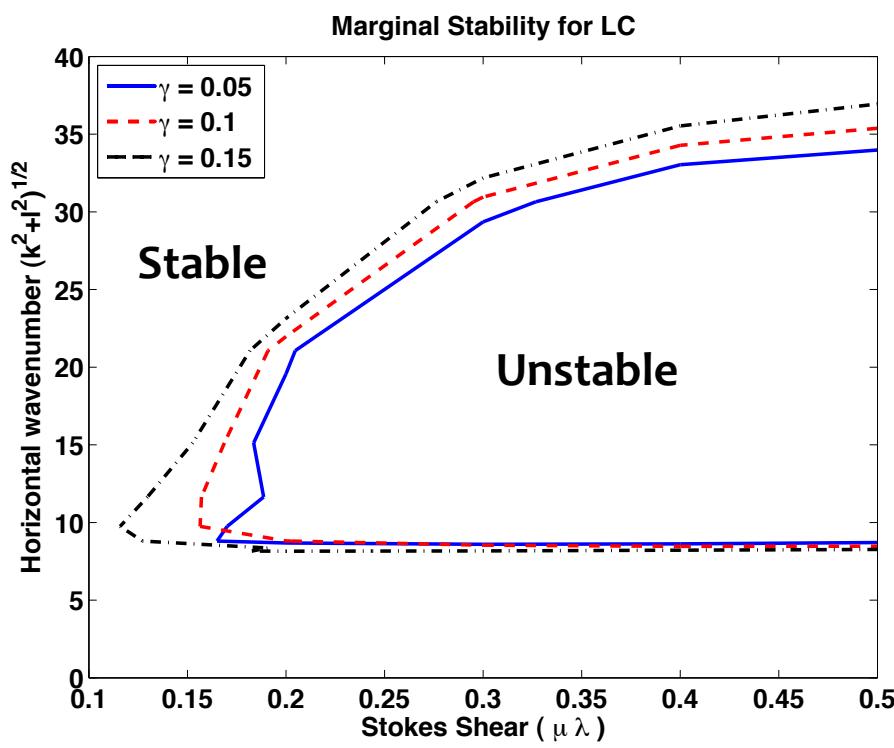
$v', l = 752 \text{ m}, \tau \sim 14\text{h}$ U^S



Stokes-Modified SI



LC within a Front

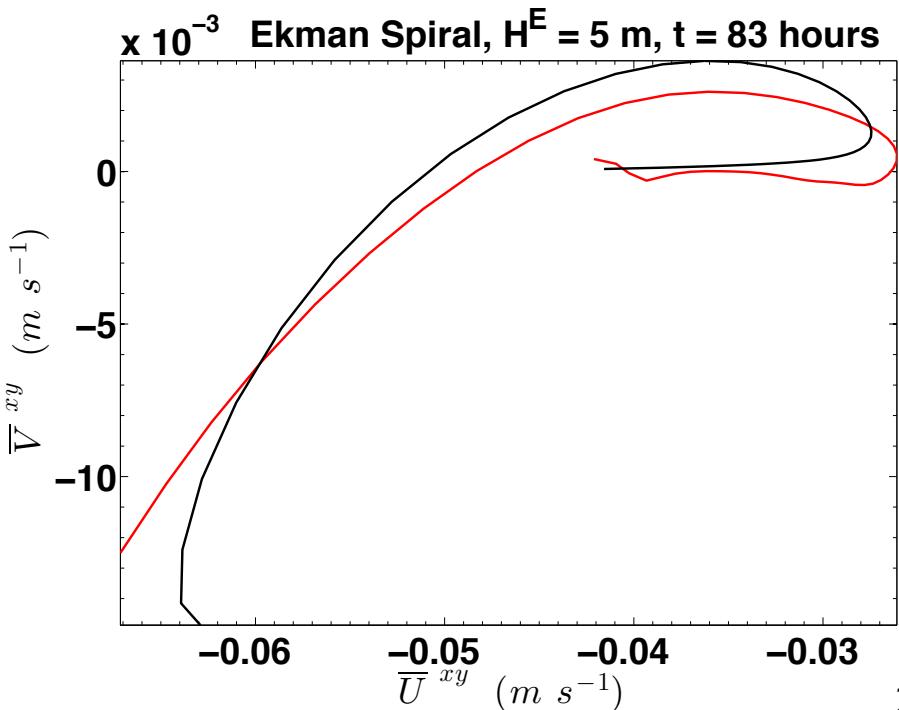
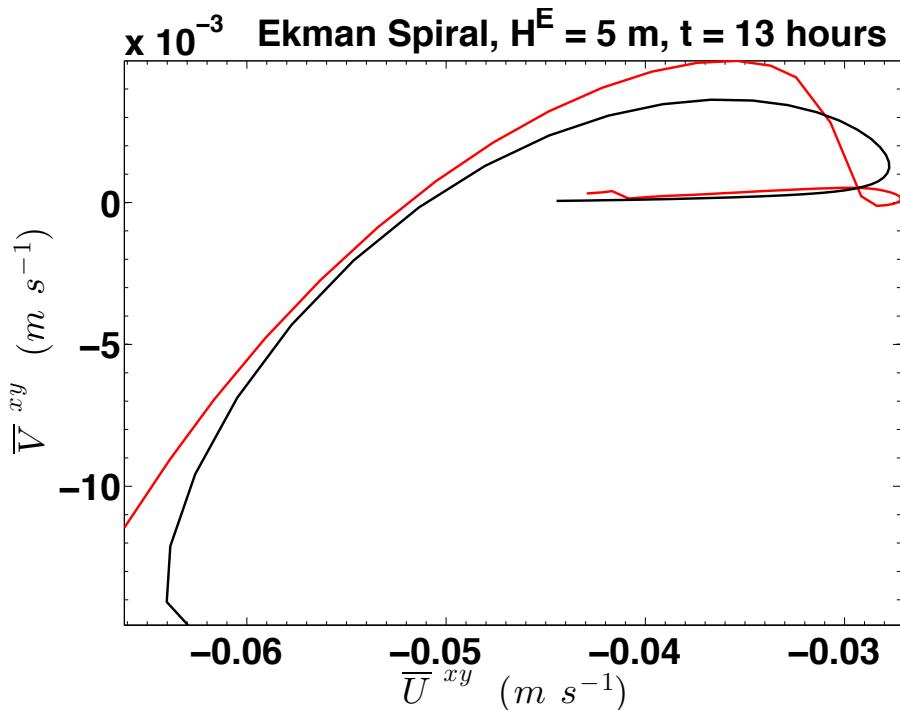


Nonlinear Simulations

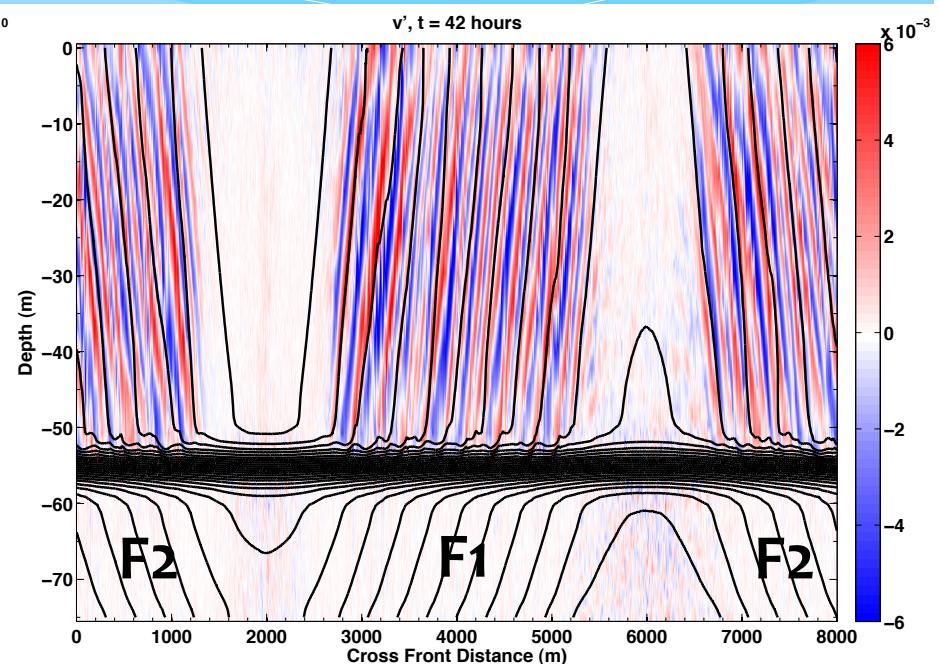
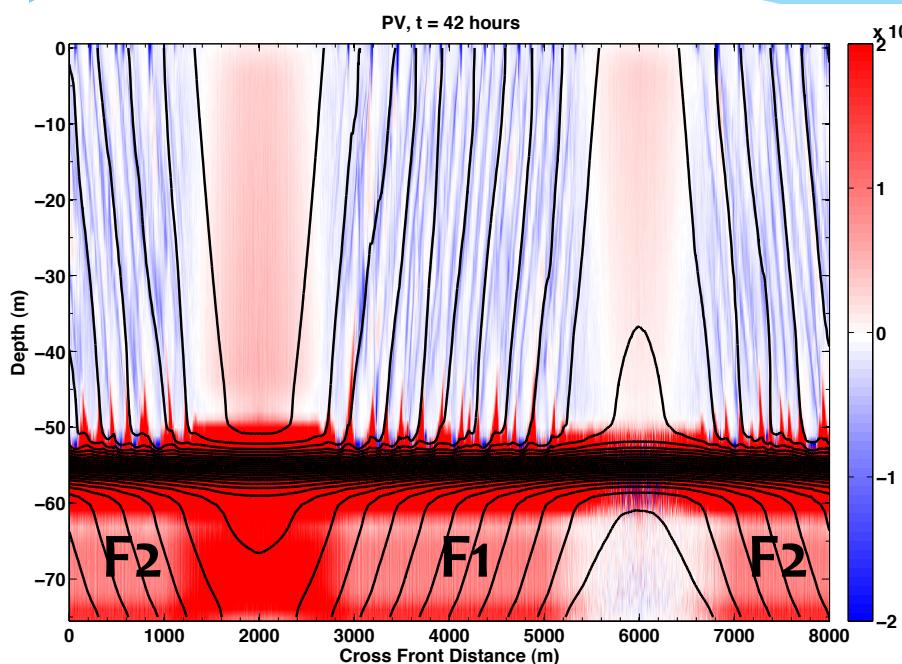
- * National Center for Atmospheric Research (NCAR) LES model.
- * Configuration
 - 500m (along-front) x 8km (cross-front) x 75m deep
 - Intentionally along-front limited to prohibit GI. Simulating GI and LC would require $O(10^6 \text{ cpuh})$ vs the $O(10^4 \text{ cpuh})$ required for the simulations performed.
 - $\sim 4\text{m}$ horizontal x $\sim 1\text{m}$ vertical resolution.
 - Periodic BC's in the horizontal (requires simulating 2 fronts)
 - No flux on top and bottom
 - No wind stress on top
- * Cases
 1. $\text{PV} < 0$, no Stokes ($\text{Ri} = 0.5, \mu = 0$): control case
 2. $\text{PV} < 0$ at depth, with Stokes ($\text{Ri} = 0.5, \mu = 2$)
 3. $\text{PV} > 0$ at depth, with Stokes ($\text{Ri} = 2, \mu = 1$)

Stokes-Ekman-Front Layer

- Analytic Solution
- Horizontal average from LES



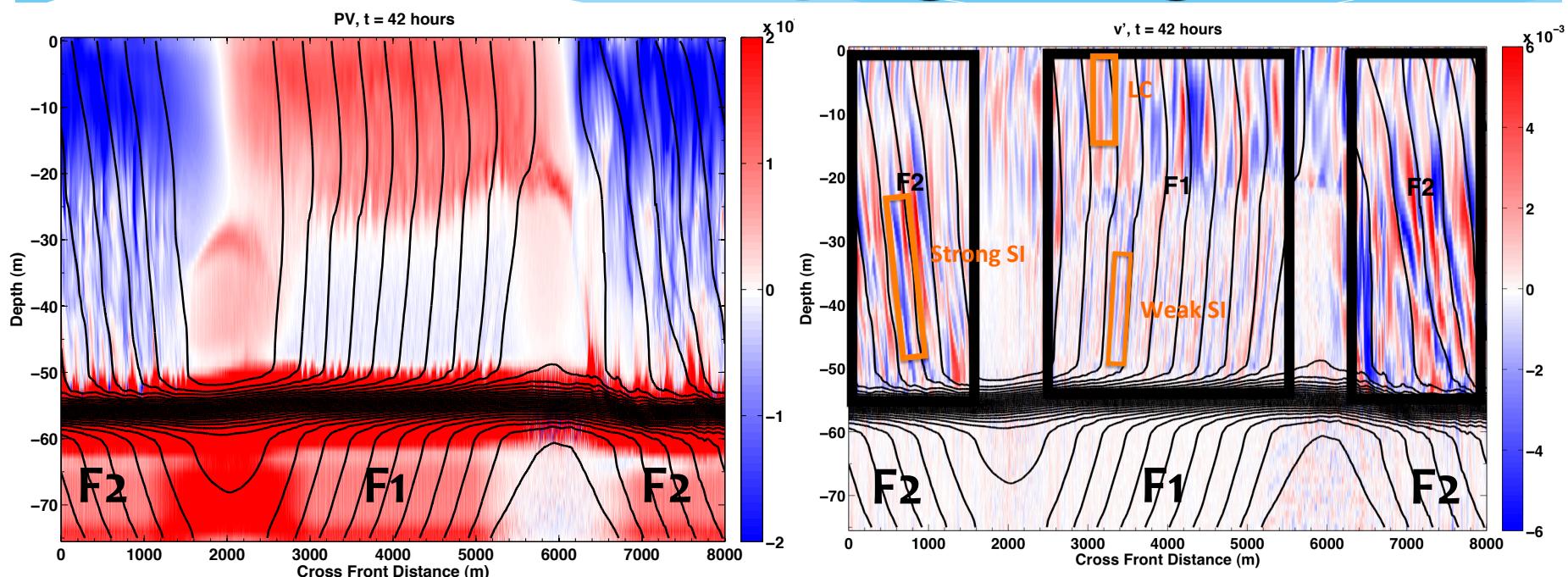
LES Case 1 (Control Case)



- * PV<0, no Sokes ($Ri = 0.5, \mu = 0$)
- * We expect SI in the within the fronts

LES Case 2

$Ri = 0.5, \mu = 2$



- * We expect SI in the regions where $PV < 0$
- * Stokes-Ekman-Front layer yields an Ekman transport to the left, destabilizing F1 while stabilizing F2.

LES Case 3

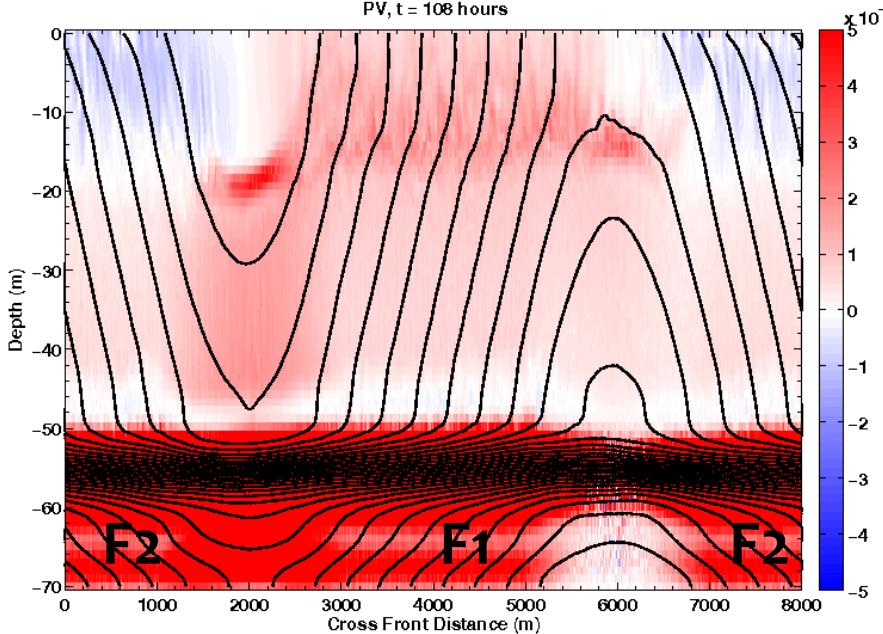
$Ri = 2, \mu = 1$



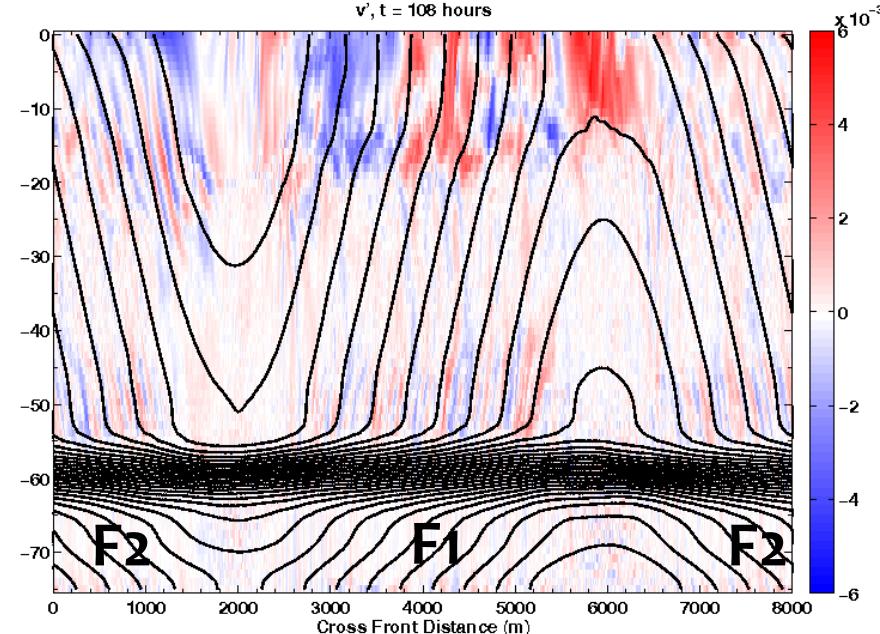
U_g



PV, t = 108 hours



v' , t = 108 hours



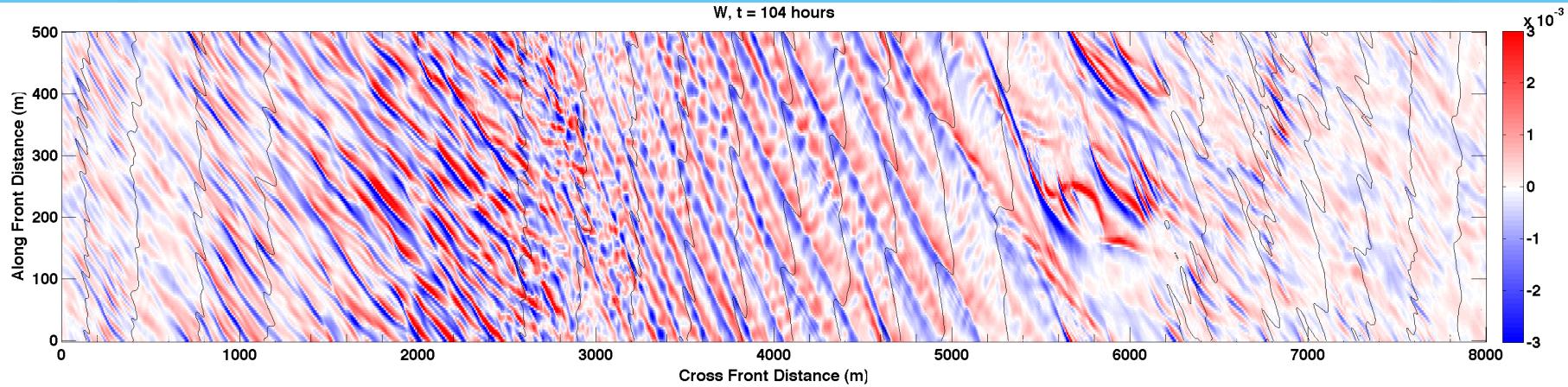
- * We expect SI in the regions where $PV < 0$
- * Stokes-Ekman-Front layer yields an Ekman transport to the left, destabilizing F1 while stabilizing F2.

Vertical Stratification Slows LC

Horizontal slice of vertical velocity at $\sim 5\text{m}$ deep.



$$Ri = 0.5, \mu = 2$$



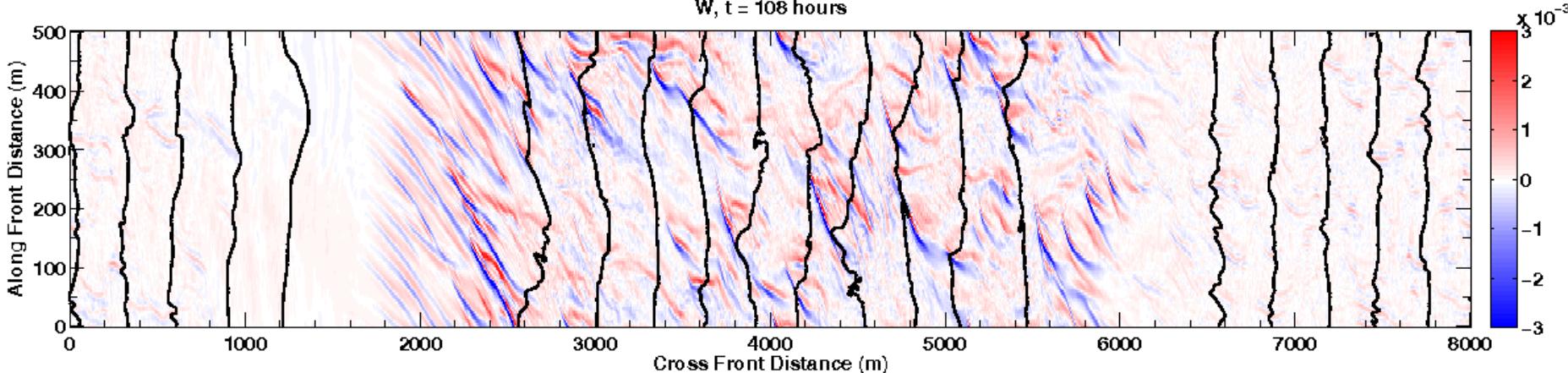
HOT, LIGHT

$$Ri = 2, \mu = 1$$

COLD, DENSE

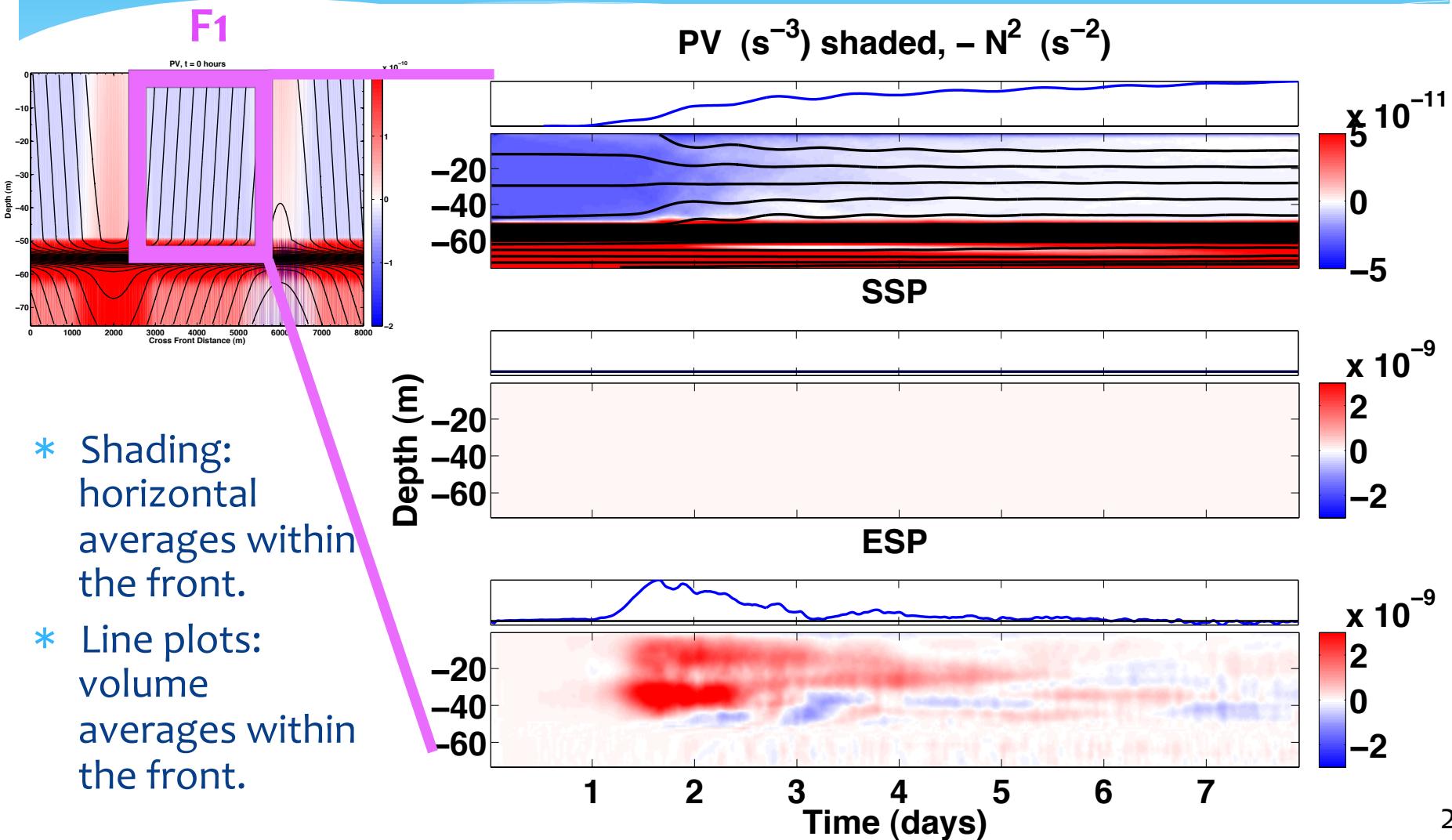
W, t = 108 hours

This figure shows a horizontal slice of the vertical velocity field for a different set of parameters. The plot displays a more organized, periodic pattern of red and blue lines. The x-axis is labeled 'Cross Front Distance (m)' and ranges from 0 to 8000. The y-axis is labeled 'Along Front Distance (m)' and ranges from 0 to 500. A color bar on the right indicates velocity values ranging from -3 to 3, with a multiplier of 10^{-3} . The plot is titled 'W, t = 108 hours'. The flow is primarily downward, indicated by arrows and labeled U^S , with upward components indicated by arrows and labeled U^G .



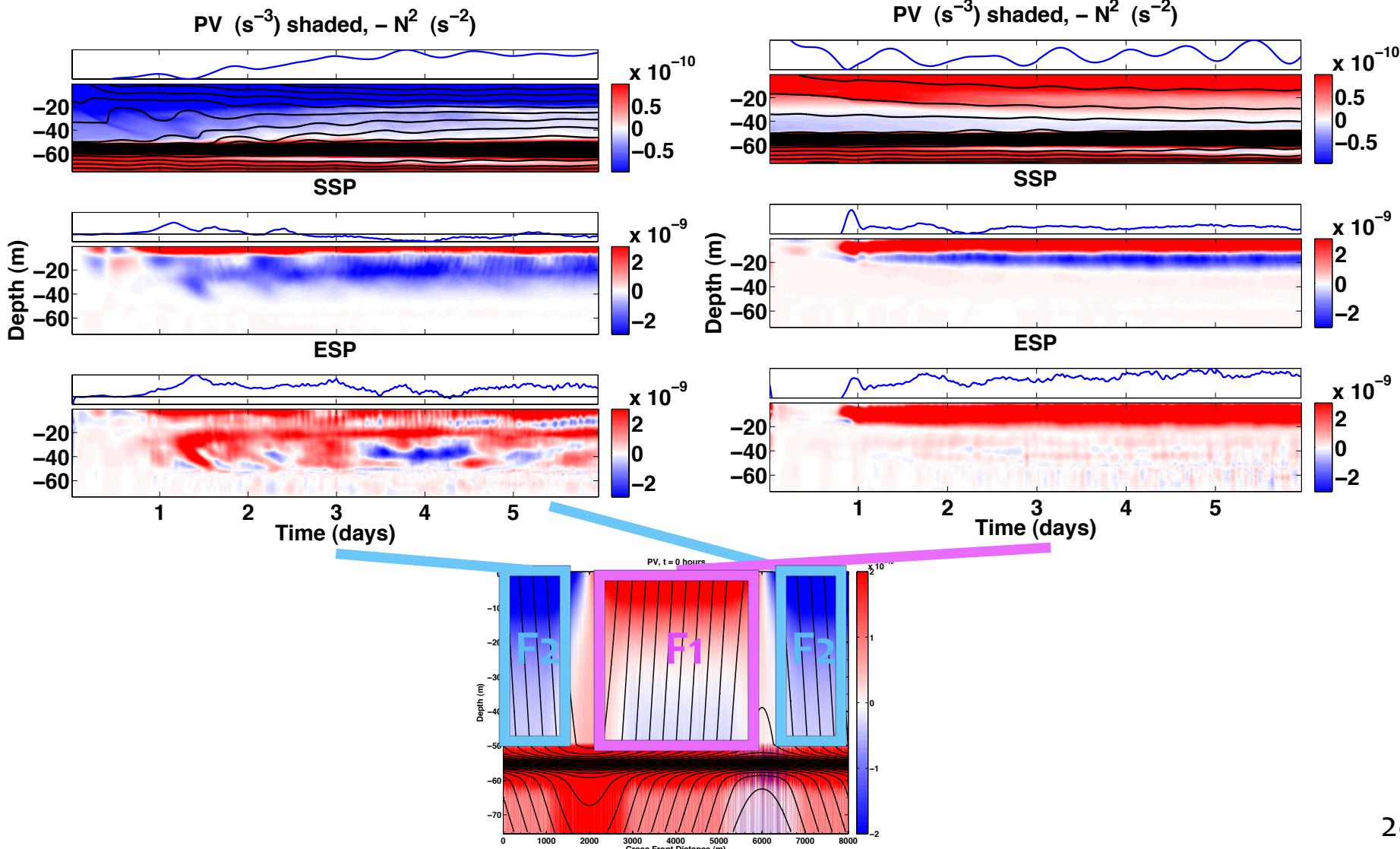
Energetics and PV: Control Case

$$Ri = 0.5, \mu = 0$$



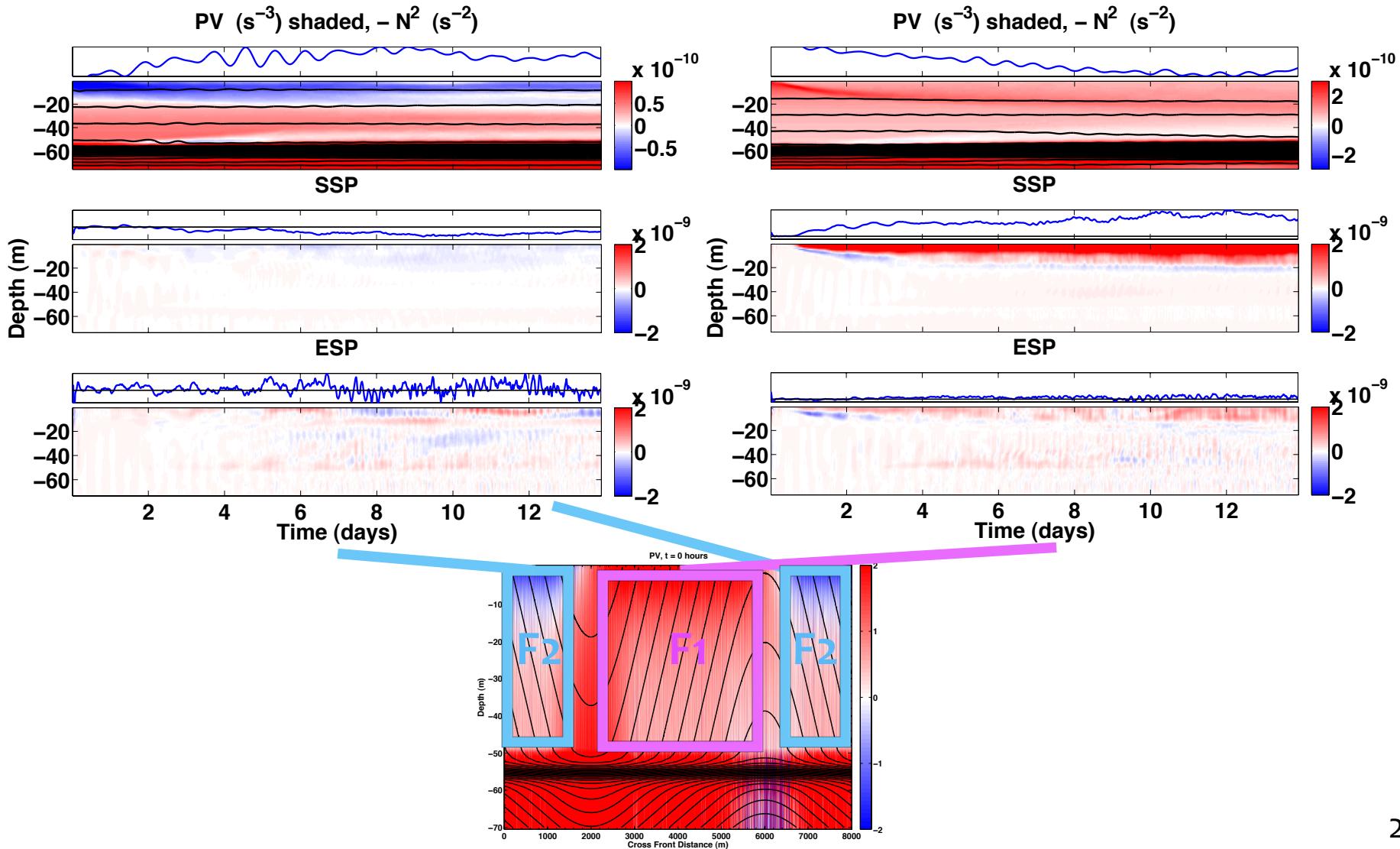
Energetics and PV: Case 2

$Ri = 0.5, \mu = 2$



Energetics and PV: Case 3

$$Ri = 2, \mu = 1$$



Conclusions

- * GI are only weakly affected by Stokes drift, but their instability depends on the Lagrangian, not just Eulerian shear.
- * If the flow is indeed unstable to GI, increased anti-Stokes Eulerian shear reduces the growth rate and wavenumber.
- * If the flow is unstable to SI, then $PV < 0$ (and the implication appears to go the other way as well), and anti-Stokes Eulerian flow modifies the PV.
 - Observational estimates of PV must be based on Eulerian shear if SI are of interest.
- * Stokes forced SI do more BP than their no Stokes counterparts.
- * SSP does work against SI when the Eulerian and Stokes shears oppose each other.
- * Stokes drift can indirectly induce restratification (rather than mixing with LC) by modifying the PV and shear, causing SI to do more BP.
- * LC are suppressed by the Ekman induced restratification of the front.



Thanks!

