

1) Time to collision

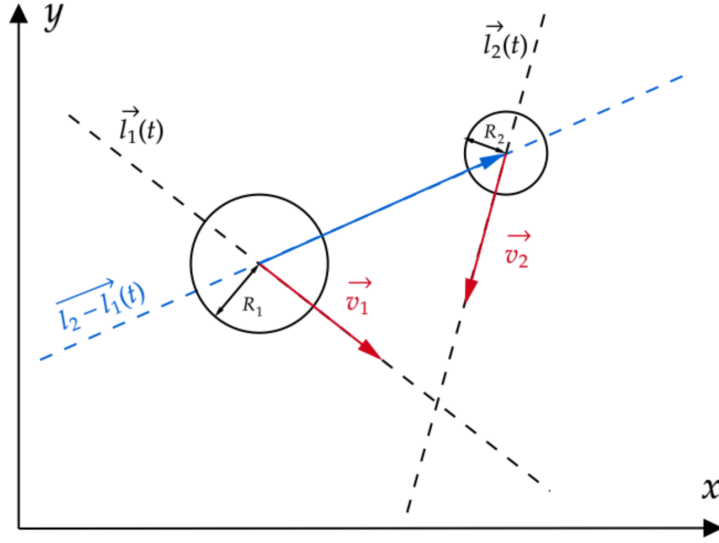


Figure 1: Two particles moving with constant speed on a collision trajectory

Parameterization of each trajectory as:

$$\vec{l}_1(t) = \vec{v}_1 \cdot t + \vec{r}_1 \quad (1)$$

$$\vec{l}_2(t) = \vec{v}_2 \cdot t + \vec{r}_2 \quad (2)$$

yields the distance between the two particles as:

$$d(t) = \|\vec{l}_2(t) - \vec{l}_1(t)\| = \|(\vec{v}_2 - \vec{v}_1)t + (\vec{r}_2 - \vec{r}_1)\| \quad (3)$$

Therefore the time of a collision can be found as the solution to:

$$d(t) = R_1 + R_2 \leftrightarrow d(t) - (R_1 + R_2) = 0 \quad (4)$$

which, in turn, corresponds to solving the 2 degree polynomial:

$$d(t)^2 - (R_1 + R_2)^2 = 0. \quad (5)$$

with:

$$d(t)^2 = \|\vec{l}_2(t) - \vec{l}_1(t)\|^2 \quad (6)$$

$$= \|\vec{v}_2 - \vec{v}_1\|^2 t^2 + 2\langle \vec{v}_2 - \vec{v}_1 | \vec{r}_2 - \vec{r}_1 \rangle t + \|\vec{r}_2 - \vec{r}_1\|^2 \quad (7)$$

the solution to (5), therefore is given by the standard formula, and yields the following:

$$t_1, t_2 = \frac{-2\langle \vec{v}_2 - \vec{v}_1 | \vec{r}_2 - \vec{r}_1 \rangle \pm \sqrt{4\langle \vec{v}_2 - \vec{v}_1 | \vec{r}_2 - \vec{r}_1 \rangle^2 - 4\|\vec{v}_2 - \vec{v}_1\|^2(\|\vec{r}_2 - \vec{r}_1\|^2 - (R_1 + R_2)^2)}}{2\|\vec{v}_2 - \vec{v}_1\|^2}$$

2) Updating velocity

Neglecting gravity, conservation of energy and momentum in 1D:

$$\frac{1}{2}m_1v_{1,i}^2 + \frac{1}{2}m_2v_{2,i}^2 = \frac{1}{2}m_1v_{1,f}^2 + \frac{1}{2}m_2v_{2,f}^2$$

$$m_1v_{1,i} + m_2v_{2,i} = m_1v_{1,f} + m_2v_{2,f}$$

yields:

$$v_{1,f} = \frac{m_1 v_{1,i} - m_2(v_{1,i} - 2v_{2,i})}{m_1 + m_2} \quad \text{and} \quad v_{2,f} = \frac{m_1(2v_{1,i} - v_{2,i}) + m_2 v_{2,i}}{m_1 + m_2} \quad (8)$$

On the surface, this doesn't seem to provide any viable solution to expressing the post-collision velocity in 2D. However, expressing their velocities in the basis of a line parallel to their contact-point, and one orthogonal to their contact-point:

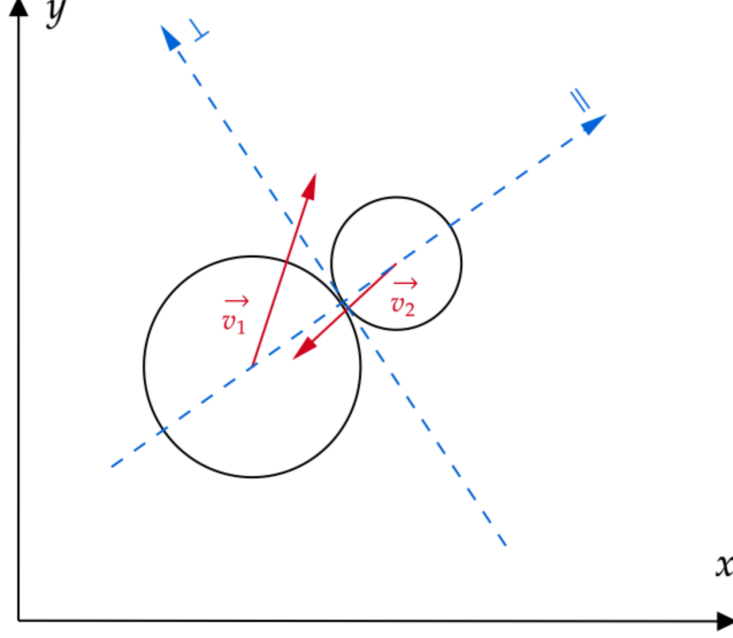


Figure 2: Two particles moving with constant speed on a collision trajectory

s.t.:

$$\vec{v}_i = v_{i,\parallel} \hat{e}_{\parallel} + v_{i,\perp} \hat{e}_{\perp} \quad (9)$$

gives the advantage that there is only an exchange of forces in direction of the contact, i.e. $\frac{d\vec{p}_{\perp}}{dt} = 0$, and therefore that only the component of the velocity in this direction changes.

In this case, a unit-vector in the \parallel direction is given as:

$$\hat{e}_{\parallel} = \frac{\vec{r}_2 - \vec{r}_1}{\|\vec{r}_2 - \vec{r}_1\|} \quad (10)$$

and, the unit-vector in the orthogonal direction is found by simple rotation:

$$\hat{e}_{\perp} = \theta\left(\frac{\pi}{2}\right) \hat{e}_{\parallel} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \hat{e}_{\parallel} \quad (11)$$

finally the velocity component in each direction can be found by means of projection:

$$v_{i,\parallel} = \langle \vec{v}_i | \hat{e}_{\parallel} \rangle \hat{e}_{\parallel} \quad \text{and} \quad v_{i,\perp} = \langle \vec{v}_i | \hat{e}_{\perp} \rangle \hat{e}_{\perp} \quad (12)$$

Applying (8) on the parallel component finally gives:

$$\vec{v}_{1,f} = v_{1,i,\parallel} \hat{e}_{\parallel} + v_{1,f,\perp} \hat{e}_{\perp} = \vec{v}_{1,i} - \frac{2m_2}{m_1 + m_2} \frac{\langle \vec{v}_{1,i} - \vec{v}_{2,i} | \vec{r}_{1,i} - \vec{r}_{2,i} \rangle}{\|\vec{r}_{1,i} - \vec{r}_{2,i}\|^2} (\vec{r}_{1,i} - \vec{r}_{2,i}) \quad (13)$$

$$\vec{v}_{2,f} = v_{2,i,\parallel} \hat{e}_{\parallel} + v_{2,f,\perp} \hat{e}_{\perp} = \vec{v}_{2,i} - \frac{2m_1}{m_1 + m_2} \frac{\langle \vec{v}_{2,i} - \vec{v}_{1,i} | \vec{r}_{2,i} - \vec{r}_{1,i} \rangle}{\|\vec{r}_{2,i} - \vec{r}_{1,i}\|^2} (\vec{r}_{2,i} - \vec{r}_{1,i}) \quad (14)$$