

## Wave Nature of Particle &

## Schrodinger Wave Equation

Wave Nature of Particles →

Quantum Mechanics →

Quantum Mechanics is the branch of physics which deals with the microscopic particles.

De-Broglie proposed that a material particle such as electrons, protons, etc is associated with the wave if  $m$  is the mass of the material particle moving with velocity  $v$  then the wavelength of the wave associated with it is given by →

$$\lambda = \frac{h}{mv} = \frac{\lambda}{p}$$

where,  $p$  = momentum of particle

This is called De-Broglie wavelength. The wave associated with the material particle is called Matter wave OR De-Broglie wave. Hence, the material particle has dual nature.

Free Particle →

A particle which is not acted upon by any force is called Free particle.

According to Newton 2<sup>nd</sup> Law of motion →

$$F = \frac{dp}{dt}$$

For Free particle,  $F = 0$

$$\frac{dp}{dt} = 0$$

$$P = \text{constant}$$

This means that the free particle has a constant momentum. The total energy of free particle is only kinetic energy and is given by

$$E = \frac{1}{2}mv^2$$

$$E = \frac{1}{2}mv^2 \times m$$

$$E = \frac{1}{2}m^2v^2$$

$$E = \frac{p^2}{2m}$$

since,  $P$  is a constant. Therefore, the energy of the free particle is also constant.

Hence, a free particle is characterised by definite momentum & energy.

Wave Function for a Free Particle  $\rightarrow$

Wave function is a function which describes the matter wave associated with the material particle and is represented by  $\rightarrow$

$$\psi(x, t) = Ae^{i(kx - \omega t)}$$

$$k = \text{propagation constant} = \frac{2\pi}{\lambda}$$

$$t = \text{Time}$$

$$A = \text{Amplitude of wave}$$

$$\omega = \text{Angular frequency}$$

$$x = \text{location of particle}$$

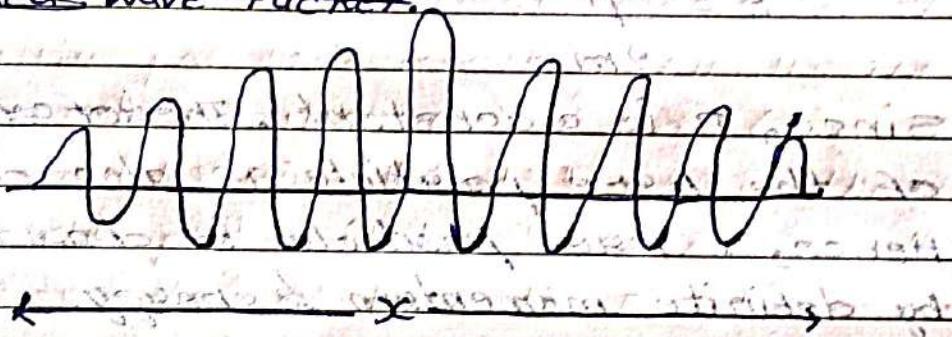
This wave function describes the matter wave associated with free particle and moving along the  $x$ -axis.

Wave Packet →

The material particle is represented by wave function  $\Psi(x, t) = A e^{i(kx - \omega t)}$

This wave spreads out over the infinite extent in one dimensional space. In order to describe the particle localised in a certain region ( $\Delta x$ ). It is necessary that the wave function should be zero everywhere except in the region ( $\Delta x$ ) where the particle is located.

A wave of finite extent is shown in the fig. ie ~~wave packet~~ wave packet.



In order to obtain a wave packet. It is necessary to modulate the amplitude of the wave in such a way that it is non-zero in the region ( $\Delta x$ ) and zero outside. This can be done by adding the no. of waves of slightly different frequencies & amplitude. If we suppose ( $N$ ) no. of waves then the function can be written as

$$\Psi(x, t) = \sum_{i=1}^N A_i e^{i(k_i x - \omega_i t)}$$

OR

$$\Psi(x, t) = \int_{-\infty}^{+\infty} A(k) e^{i(kx - \omega t)} dk$$

This wave function represents a wave packet which describes a localised particle. Hence, we conclude

that a localised particle is always associated with the wave packet.

Uncertainty Principle  $\rightarrow$

According to uncertainty principle, it is impossible to measure simultaneously and precisely both the position and momentum of particle. If  $(\Delta x)$  is the uncertainty in position &  $(\Delta p_x)$  is the uncertainty in momentum. Then,

$$\Delta x \cdot \Delta p_x \geq \hbar$$

It states that, if we try to measure the position & momentum of particle simultaneously then the product of uncertainties  $(\Delta x)$  and  $(\Delta p_x)$  cannot be less than  $\hbar$ .

$$\Delta p_x \geq \frac{\hbar}{\Delta x}$$

$$\Delta x = 0$$

$$\Delta p_x \geq \frac{\hbar}{0} \Rightarrow \Delta p_x \geq \infty$$

This means that the momentum of the particle is completely unknown.

$$\text{IF } \Delta p_x = 0 \Rightarrow \Delta x \geq \infty$$

Hence, the particle is completely unlocalised & spreads out all over the space.

In 3D, the uncertainty principle becomes

$$\Delta y \cdot \Delta p_y \geq \hbar$$

$$\Delta z \cdot \Delta p_z \geq \hbar$$

The uncertainty principle indicates that we don't know the exact position and momentum of particle of atomic dimensions.

Born Interpretation of Wave Functions →

Max Born →

Extend the concept of wave packet associated with a material particle the motion of the electron travelling ( $x$ ) direction with constant momentum ( $P$ ) is controlled by the matter waves associated with it. The wave function of such a wave is given by →

$$\Psi(x, t) = A e^{i(kx - \omega t)}$$

The quantity  $(\Psi)^2$  play the important role for the matter waves. The square of the absolute value of wave function  $\Psi(x, t)$  measures the probability per unit length of finding the electron at position ( $x$ ) & time ( $t$ ).

The product  $(P) = \Psi(x, t) \Psi^*(x, t)$  is the probability per unit length of finding the particle at position ( $x$ ) at time ( $t$ ).  $P(x, t)$  is the probability density. The probability of finding the particle is large whenever  $(\Psi)$  is large & vice-versa.

Hence,  $\Psi(x, t)$  is called probability amplitude.

→ Born Interpretation of wave function mean the location of the particle cannot be defined precisely but the wave function can predict only the probability density for finding the particle.

## Time Dependent Schrödinger Wave Equation →

Newton's law cannot describe the motion of the particles of atomic dimensions whose behaviour is explained by De-Broglie waves associated with them OR by the wave function. The knowledge of  $\Psi(x, t)$  at one point in space & time is not sufficient. We need to know  $\Psi(x, t)$  at all the times to describe the motion of the particles in space and time. For this it is necessary to find a differential equation which controls the space-time behaviour of the wave function.

→ For free particle :-

The wave function for a free particle is given by

$$\Psi(x, t) = A e^{i(kx - \omega t)}$$

The total energy of a free particle of mass 'm' is only kinetic & is given by :-

$$E = \frac{1}{2} m v^2 = \frac{p^2}{2m} \rightarrow 1st$$

The momentum & the energy is related to the propagation constant & angular frequency by the relation

$$E = \hbar \omega$$

$$p = \hbar k$$

$$\hbar \omega = \frac{\hbar^2 k^2}{2m} \rightarrow 2nd$$

$$\lambda = \frac{h}{p}$$

$$k = \frac{2\pi}{\lambda}, \lambda = \frac{2\pi}{k}$$

$$\frac{2\pi}{\lambda} = \frac{h}{p}, \Rightarrow p = \frac{h}{2\pi} \times k$$

Differentiate the wave function  $\Psi(x, t)$  with respect to  $x$  & time. We get →

$$\frac{\partial \Psi}{\partial x} = A e^{i(kx - \omega t)} \cdot ik$$

$$\frac{\partial \Psi}{\partial t}$$

$$\frac{\partial \Psi}{\partial t} = A e^{i(kx - \omega t)} \cdot i\hbar \omega$$

$$\frac{\partial^2 \Psi}{\partial x^2} = -A k^2 e^{i(kx - \omega t)}$$

$$\frac{\partial^2 \Psi}{\partial x^2} = -k^2 \Psi$$

$$k^2 = -1 \cdot \frac{\partial^2 \Psi}{\partial x^2} \rightarrow \text{3rd}$$

$$\Psi = A i \omega e^{i(kx - \omega t)}$$

$$\frac{\partial^2 \Psi}{\partial t^2} = A i^2 \omega^2 e^{i(kx - \omega t)}$$

$$\frac{\partial^2 \Psi}{\partial t^2} = -A \omega^2 e^{i(kx - \omega t)}$$

$$\frac{\partial^2 \Psi}{\partial t^2} = -\omega^2 \Psi$$

$$\frac{\partial^2}{\partial t^2}$$

$$\omega = -1 \cdot \frac{\partial \Psi}{\partial t} \quad (kx - \omega t)$$

$$\omega = -i^2 \frac{\partial \Psi}{\partial t}$$

$$\omega = i \frac{\partial \Psi}{\partial t} \rightarrow 4\pi h = 2\pi m \lambda = 7$$

$$\Psi \frac{\partial}{\partial t}$$

From 2nd, 3rd & 4th

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{m} \cdot \frac{\partial^2 \Psi}{\partial x^2}$$

$$\Psi \frac{\partial}{\partial t} = 2m \Psi \frac{\partial^2}{\partial x^2}$$

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{m} \frac{\partial^2 \Psi}{\partial x^2} \rightarrow 5th$$

$$\frac{\partial}{\partial t} = 2m \frac{\partial^2}{\partial x^2}$$

This is the time dependent wave equation for a free particle.

→ For a particle subjected to a force:

When the particle is acted upon by a force then the total energy is the sum of the K.E & P.E. Hence,

$$E = \frac{1}{2} m v^2 + V(x, t)$$

Then, the wave equation becomes →

$$i\frac{\hbar}{m}\frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m}\frac{\partial^2 \Psi}{\partial x^2} + V(x, t) \quad [\text{From 5th}]$$

$$i\frac{\hbar}{m}\frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m}\nabla^2 \Psi + V(x, t)\Psi$$

This is one dimensional time dependent wave equation for a particle subjected to a force.

For 3D →

$$i\frac{\hbar}{m}\frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m}\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\right]\Psi + V(x, y, z)\Psi(x, y, z)$$

$$i\frac{\hbar}{m}\frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m}\nabla^2 \Psi + V(x, y, z)\Psi(x, y, z)$$

Note →  $E = h\nu \times 2\pi$

$$2\pi$$

$$E = \hbar\omega$$

• ~~REMARK~~  $\lambda = h$

$$p$$

$$p = h \times \frac{1}{2\pi}$$

$$\lambda = \frac{h}{2\pi}$$

$$p = \hbar k$$

Time Independent Schrödinger's Wave Equation →

The one dimensional time independent Schrödinger's Wave equation is

$$i\frac{\hbar}{m}\frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m}\frac{\partial^2 \Psi}{\partial x^2} + V\Psi$$

$$i\frac{\hbar}{m}\frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m}\frac{\partial^2 \Psi}{\partial x^2}$$

In quantum mechanics, the potential energy is not a function of time. In such case, it is possible

to separate the space & time variable in the Schrodinger wave equation. Let us assume that the wave function  $\Psi(x, t)$  is a product of two functions  $\Psi(x) T(t)$ . Then, we can write  $\rightarrow$

$$\Psi(x, t) = \Psi(x) T(t)$$

$$i\hbar \frac{\partial \Psi(x) T(t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi(x) T(t)}{\partial x^2} + V(x) \Psi(x) T(t)$$

$$i\hbar \Psi(x) \frac{\partial T(t)}{\partial t} = -\frac{\hbar^2}{2m} T(t) \frac{\partial^2 \Psi(x)}{\partial x^2} + V(x) \Psi(x) T(t)$$

divide the above equation by  $\Psi(x) T(t) \rightarrow$

$$i\hbar \frac{\partial T(t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi(x)}{\partial x^2} + V(x)$$

The Right Hand Side of this equation is a function of  $x$  only. and the LHS is a function of  $T(t)$  only. The only case in which both sides can be equal is that each expression should be equal to a constant.

Let us assume the constant be  $E$ .

$$i\hbar \frac{\partial T}{\partial t} = E$$

$$T(t) \frac{d}{dt}$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V(x) = E$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2}$$

$$2m \Psi(x) \frac{\partial^2}{\partial x^2}$$

$$-\frac{\partial^2 \Psi}{\partial x^2} = 2m(E - V) \Psi(x)$$

$$\frac{\partial^2 \Psi}{\partial x^2}$$

$$\frac{\partial^2 \Psi}{\partial x^2} + 2m(E - V) \Psi(x) = 0$$

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{\hbar^2}{2m} E \Psi(x) = 0$$

This equation is known as Time independent Schrodinger wave equation in one dimension.

$$\nabla^2 \psi + \frac{2m}{\hbar^2} [ (E - V) - \frac{\hbar^2}{8m} ] \psi(x) = 0$$

This equation is in 3-D. Let  $\vec{r} = x\hat{i} + y\hat{j} + z\hat{k}$ ,

Expectation Value  $\rightarrow$

The average value of a function in quantum mechanics is called Expectation value. The expectation value of a function  $F(x)$  is represented as  $\langle F(x) \rangle$ .

$$\langle F(x) \rangle = \int \psi^*(x, t) F(x) \psi(x, t) dx$$

Expectation Value of Momentum  $\rightarrow$

The expectation value of  $x$  component of momentum is written as  $\rightarrow$

$$\langle P_x \rangle = \int \psi^*(x, t) P_x \psi(x, t) dx \rightarrow 1st$$

Consider a wave function of a free particle moving in  $x$  direction is given by

$$\psi(x, t) = A e^{i(kx - \omega t)}$$

$$\frac{\partial \psi}{\partial x} = iAk e^{i(kx - \omega t)}$$

$\partial x$

$$P_x = \hbar k \Rightarrow k = \frac{P_x}{\hbar} \quad \left[ \hbar = \frac{h}{2\pi} \right]$$

$$\frac{\partial \psi}{\partial x} = i P_x A e^{i(kx - \omega t)}$$

$\frac{\partial x}{\partial x} = \hbar$

$$\hbar \frac{\partial \psi}{\partial x} = P_x \psi(x, t)$$

$$\int \frac{\partial \psi}{\partial x} dx = \int P_x \psi(x, t) dx$$

Put in 1st

$$\langle P_x \rangle = \int \psi^*(x, t) \cdot \hbar \frac{\partial \psi}{\partial x} dx$$

$$\langle P_x \rangle = \int \psi^*(x, t) \hbar \frac{\partial \psi}{\partial x} dx$$

$$\langle P_x \rangle = - \int \Psi(x, t) i\hbar \frac{\partial \Psi^*(x, t)}{\partial x} dx$$

$$P_x \rightarrow -i\hbar \frac{\partial}{\partial x}, \quad P_y \rightarrow -i\hbar \frac{\partial}{\partial y}, \quad P_z \rightarrow -i\hbar \frac{\partial}{\partial z}$$

Expectation value of Energy  $E \rightarrow \epsilon i\hbar \frac{\partial}{\partial t}$

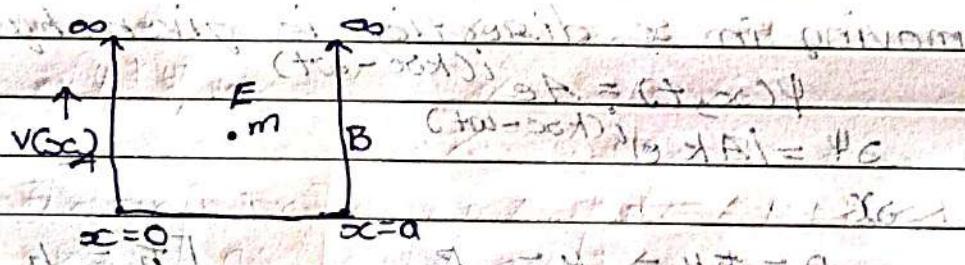
Solution of stationary state Schrödinger's Equation

For one dimensional Problem Particle in 1-D box  $\rightarrow$   
consider a one-dimensional closed box of width  $a$ .  
This box is represented by infinite square well.

It is defined as  $V(x) = \infty, 0 < x < a$

where  $V(x) = \infty, x < 0, x > a$

$V(x)$  = potential difference



Suppose a particle of mass 'm' & energy 'E' is confined between the two rigid walls A & B at  $x=0$  &  $x=a$  respectively. The one-dimensional Schrödinger's Wave equation is given by  $\rightarrow$

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{2m(E-V)}{\hbar^2} \Psi = 0$$

Inside the well  $V(x) = 0$ , so, equation becomes  $\rightarrow$

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{2mE}{\hbar^2} \Psi = 0$$

$$\frac{\partial^2 \Psi}{\partial x^2} + k^2 \Psi = 0 \quad ; \quad k = \sqrt{\frac{2mE}{\hbar^2}}$$

$$\frac{\partial^2 \Psi}{\partial x^2} + k^2 \Psi = 0 \quad ; \quad k = \sqrt{\frac{2mE}{\hbar^2}}$$

$$\frac{\partial^2 \Psi}{\partial x^2} + k^2 \Psi = 0 \quad ; \quad k = \sqrt{\frac{2mE}{\hbar^2}}$$

where  $\kappa^2 = \frac{2mE}{\hbar^2} \rightarrow$  2nd

The general solution of equation 1st is  $\rightarrow$

$$\psi(x) = A\sin(kx) + B\cos(kx) \rightarrow$$
 3rd

where  $A$  &  $B$  are the arbitrary constants.

$$\text{At } x=0, \psi(x)=0$$

Put in 3rd

$$0 = A\sin kx_0 + B\cos kx_0$$

$$0 = 0 + B$$

$$B = 0 \quad \text{Now, } \psi(x) = A\sin(kx)$$

$$\text{At } x=a, \psi(x)=0$$

$$0 = A\sin(ka) \quad \cancel{\text{sinka}} = 0$$

$$\sin ka = 0$$

$ka = n\pi$  where,  $n$  is a integer ( $n=1, 2, 3, \dots$ )

$$K = n\pi$$

From eqn 2nd  $\rightarrow$

$$n^2\pi^2 = 2mE$$

$$a^2 = \hbar^2$$

$$E_n = \frac{n^2\pi^2\hbar^2}{2ma^2}$$

$$E_1 = \pi^2\hbar^2$$

$$2ma^2$$

It is called the ground state energy of the particle. This minimum energy of particle is called zero point energy.

For  $n=2$

$$E_2 = \frac{4\pi^2\hbar^2}{2ma^2}$$

$$2ma^2$$

The states of the particle corresponding to  $n=2, 3, \dots$  are called Excited states.

$$E_3 = \frac{9\pi^2\hbar^2}{2ma^2}$$

Hence, the particle inside the potential well cannot have any arbitrary energy but only discrete energy values given by  $E_1, E_2, E_3, \dots$   
So, the energy of the particle in the box is quantized.

Eigen values  $\rightarrow E_1, E_2, E_3$  are the eigen values.

Eigen Function  $\rightarrow$

Put the value of  $k$  in equation  $\psi(x) = A \sin(kx)$

$$\psi(x) = A \sin[n\pi x]$$

To find the value of  $k$  apply the normalization condition.

$$\int \psi_n^* \psi_n dx = 1$$

$$\int_a^a A \sin[n\pi x] A \sin[n\pi x] dx = 1$$

$$\int_0^a \sin^2[n\pi x] dx = 1$$

Using the standard cond. integral

$$\int_0^a \sin^2[n\pi x] dx = a$$

Now,

$$a^2/a = 1$$

$$2$$

$$A = \sqrt{\frac{2}{a}}$$

$$\psi_n(x) = \frac{1}{\sqrt{a}} \sin \frac{n\pi}{a} x$$

This function is known as Eigen Function.

Linear Harmonic Oscillator →

A system on which the force is directly proportional to the displacement from equilibrium position & ex. in a direction opposite to the displacement is called a Harmonic Oscillator.

If  $x=0$  is the equilibrium position then the force is  $F \propto -x$ .

where,  $k$  is the constant of proportionality.  
the relation b/w force & the potential energy is given by

$$F = -kx$$

$$+kx = dv$$

$$(integ) dx$$

$$kx = dv$$

$$(integ) dx$$

$$kx dx = dv$$

Integrating both sides →

$$v = \frac{1}{2} kx^2$$

This gives the potential energy of harmonic oscillator. We know that the time independent Schrodinger wave equation is

$$\frac{d^2\psi}{dx^2} + \frac{2m(E-V)}{\hbar^2} \psi = 0$$

For harmonic oscillator  $V = \frac{1}{2} kx^2$

$$\frac{d^2\psi}{dx^2} + \frac{2m}{\hbar^2} [E - \frac{1}{2}kx^2] \psi = 0$$

$$\frac{d^2\psi}{dx^2} + \frac{2mE\psi}{\hbar^2} - \frac{kx^2\psi}{\hbar^2} = 0$$

We know that  $k = \omega^2 m$ . So, eqn becomes

$$\frac{d^2\psi}{dx^2} + \left[ \frac{2mE}{\hbar^2} - \frac{\omega^2 m^2 x^2}{\hbar^2} \right] \psi = 0$$

$$\frac{d^2\psi}{dx^2} + (\beta^2 - \alpha^2 x^2) \psi = 0$$

where,  $\beta^2 = \frac{2mE}{\hbar^2}$  &  $\alpha^2 = \frac{\omega^2 m^2}{\hbar^2}$

On solving the above equation we obtain

$$\beta = 2n+1$$

Now,

$$\sqrt{\frac{2mE}{\hbar^2}} = (2n+1)$$

$$\frac{\sqrt{2mE}}{\hbar} = 2n+1$$

com

$\hbar$

$$2mE = (2n+1)^2$$

$$2mE = 4(2n+1)^2$$

$$\omega^2 m^2$$

$$F = 1/4 \hbar (2n+1)$$

$$E = 1/2 \omega^2 m (2n+1)^2$$

9

$$E = \hbar \omega f n + 1$$

2

where  $n = 0, 1, 2, 3, 4, \dots$

since, the energy depends on  $n$ .

$$E_n = [n+1] \hbar \omega$$

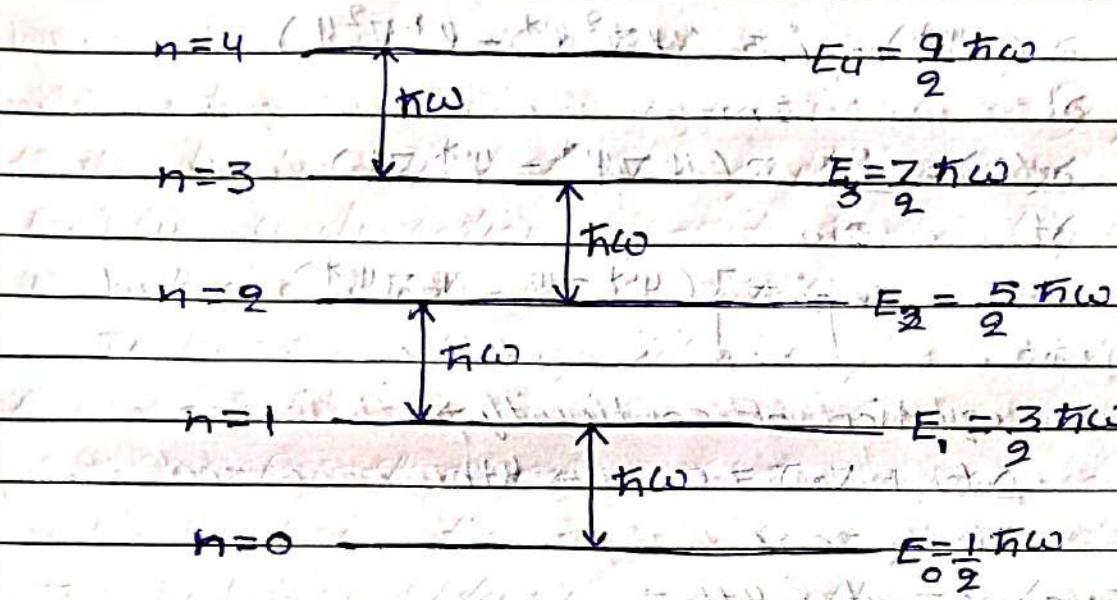
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when  $n=0$

The oscillator have minimum  $E_0 = 1/2 \hbar \omega$

9

This energy is called zero point energy. Hence, the particle executing simple harmonic motion can only have discrete energies, i.e.,  $1\frac{1}{2}\hbar\omega, 3\frac{3}{2}\hbar\omega, 5\frac{5}{2}\hbar\omega$



Probability current  $\psi^* \nabla \psi - \psi \nabla \psi^*$

consider a plane wave is represented by a

wave function  $\psi(x, t) = A e^{i(kx - \omega t)}$ . The time dependent Schrödinger wave equation is →

$$i\hbar \frac{\partial \psi}{\partial t} = \frac{\hbar^2}{m} \nabla^2 \psi + V \psi$$

Multiply by  $\psi^*$  on both sides →

$$i\hbar \psi^* \frac{\partial \psi}{\partial t} = \frac{\hbar^2}{m} \psi^* \nabla^2 \psi + V \psi \psi^* \rightarrow 1st$$

Taking the complex conjugate of equation 1st

$$-i\hbar \psi \frac{\partial \psi^*}{\partial t} = \frac{\hbar^2}{m} \psi \nabla^2 \psi^* + V \psi \psi^* \rightarrow 2nd$$

$$\frac{\partial t}{dt} = 2m$$

Subtract eqn 2nd From 1st →

$$i\hbar \psi^* \frac{\partial \psi}{\partial t} + i\hbar \psi \frac{\partial \psi^*}{\partial t} = \frac{\hbar^2}{m} \psi^* \nabla^2 \psi + V \psi \psi^* - \frac{\hbar^2}{m} \psi \nabla^2 \psi^* - V \psi \psi^*$$

$$\frac{i\hbar}{dt} [\psi^* \frac{\partial \psi}{\partial t} + \psi \frac{\partial \psi^*}{\partial t}] = \frac{\hbar^2}{m} [\psi^* \nabla^2 \psi - \psi \nabla^2 \psi^*]$$

$$\text{if } \frac{\partial}{\partial t} (\Psi \Psi^*) = -\frac{1}{\hbar^2} (\Psi \nabla^2 \Psi^* - \Psi^* \nabla^2 \Psi)$$

at  $x = 2m$

$$\text{if } \frac{\partial}{\partial t} (\Psi \Psi^*) = \frac{i}{\hbar} (\Psi \nabla^2 \Psi^* - \Psi^* \nabla^2 \Psi)$$

at  $x = 2m$

$$\frac{\partial}{\partial t} (\Psi \Psi^*) = i \frac{1}{\hbar} (\Psi \nabla^2 \Psi^* - \Psi^* \nabla^2 \Psi)$$

at  $x = 2m$

$$\frac{\partial}{\partial t} (\Psi \Psi^*) = i \frac{1}{\hbar} \nabla (\Psi \nabla \Psi^* - \Psi^* \nabla \Psi)$$

at  $x = 2m$

$$\frac{\partial P}{\partial t} = \nabla \cdot \left[ \frac{1}{\hbar} (\Psi^* \nabla \Psi - \Psi \nabla \Psi^*) \right] \rightarrow \text{3rd}$$

at  $x = 2m$

From equation of continuity  $\rightarrow$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot J = 0 \rightarrow \text{4th}$$

at  $x = 2m$

From 3rd & 4th  $\rightarrow$

$$J = i \frac{1}{\hbar} (\Psi^* \nabla \Psi - \Psi \nabla \Psi^*)$$

at  $x = 2m$

This gives you the probability current for a plane wave.

### Properties of the Wave Function $\rightarrow$

- The wave function  $\Psi(x)$  &  $\partial \Psi / \partial x$  must satisfy the following conditions  $\rightarrow$ 
  - The wave function  $\Psi(x, t)$  &  $\partial \Psi / \partial x$  must be finite everywhere.
  - The wave function  $\Psi(x, t)$  &  $\partial \Psi / \partial x$  must be a single value.
  - The wave function  $\Psi(x)$  must be continuous everywhere & its  $\partial \Psi / \partial x$  must be continuous everywhere.
  - The  $\Psi(x)$  must be zero vanish at infinity.

$$\Psi(x) = 0 \text{ at } x \rightarrow \infty$$