

RINGDOWN TESTS OF GRAVITY

COSMO24 KYOTO • OCTOBER 2024

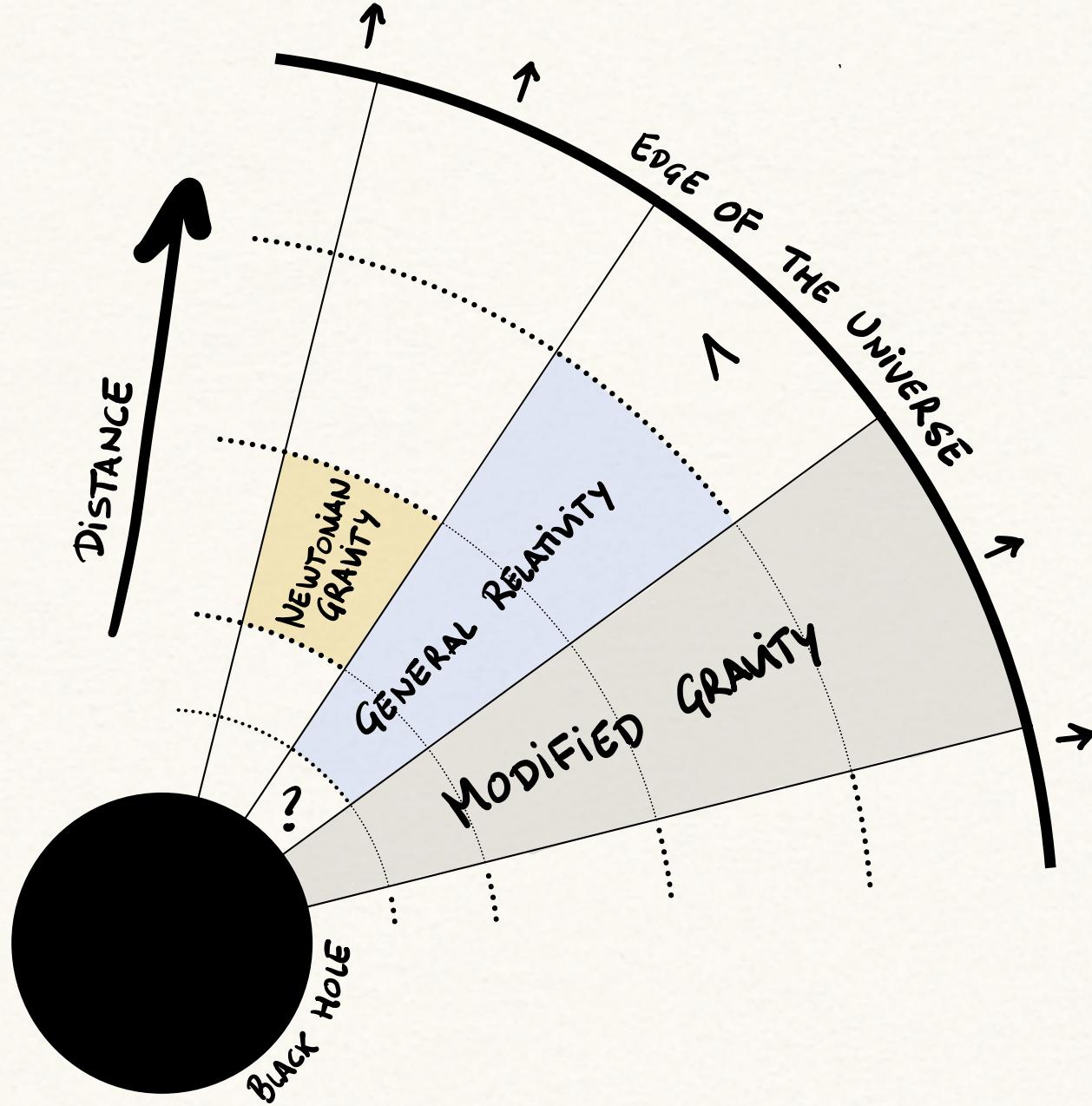
SERGI SIRERA

[2301.10272 , 2408.01720] SS + JOHANNES NOLLER



JAPAN SOCIETY FOR THE PROMOTION OF SCIENCE
日本学術振興会

TESTING GRAVITY



TESTING GRAVITY

$$\left. \begin{array}{l} 4D \\ g_{\mu\nu} \\ \text{LOCAL} \\ 2^{\text{nd}} \text{order EOM} \end{array} \right\} \text{GR} \rightarrow S = \int d^4x \sqrt{-g} R[g_{\mu\nu}] \quad (\text{LOVELOCK'S THEOREM})$$

$$\left. \begin{array}{l} 4D \\ g_{\mu\nu} + \phi \\ \text{LOCAL} \\ 2^{\text{nd}} \text{order EOM} \end{array} \right\} \text{HORNDESKI} \rightarrow S = \int d^4x \sqrt{-g} H[g_{\mu\nu}, \phi]$$

TESTING GRAVITY

$$S = \int d^4x \sqrt{-g} (\mathcal{L}_2 + \mathcal{L}_3 + \mathcal{L}_4 + \mathcal{L}_5)$$

HORNDESKI GRAVITY

$$\mathcal{L}_2 = G_2(\phi, X)$$

$$\mathcal{L}_3 = G_3(\phi, X) \square \phi$$

$$\mathcal{L}_4 = G_4(\phi, X) R + G_{4X}(\phi, X) [(\square \phi)^2 - (\phi_{\mu\nu})^2]$$

$$\mathcal{L}_5 = G_5(\phi, X) G_{\mu\nu} \phi^{\mu\nu} - \frac{1}{6} G_{5X}(\phi, X) [(\square \phi)^3 - 3(\phi_{\mu\nu})^2 \square \phi + 2(\phi_{\mu\nu})^3]$$

WHERE $X = -\frac{1}{2} \nabla_\mu \phi \nabla^\mu \phi$, $\phi_\mu := \nabla_\mu \phi$, $\phi_{\mu\nu} := \nabla_\nu \nabla_\mu \phi$, ...

$$G_{4X} := \partial_X G_4 \quad (\phi_{\mu\nu})^2 := \phi_{\mu\nu} \phi^{\mu\nu}$$

$$(\phi_{\mu\nu})^3 := \phi_{\mu\nu} \phi^{\mu\sigma} \phi_{\sigma}{}^{\nu}$$

TESTING GRAVITY

THEORY



OBSERVABLE

$$S = \int d^4x \sqrt{-g} R$$



$$\alpha = 0$$

$$S = \int d^4x \sqrt{-g} H$$

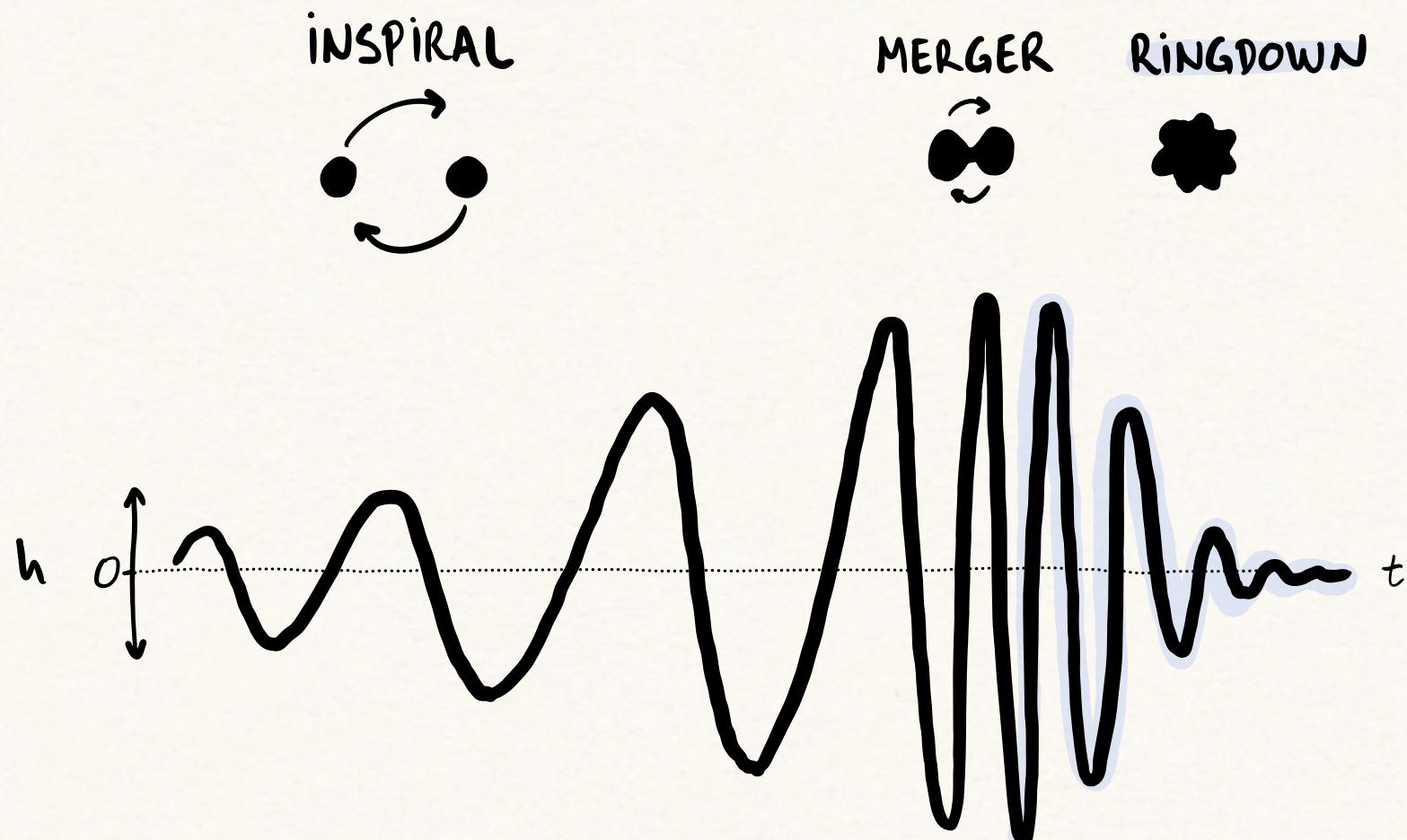


$$\alpha \neq 0$$

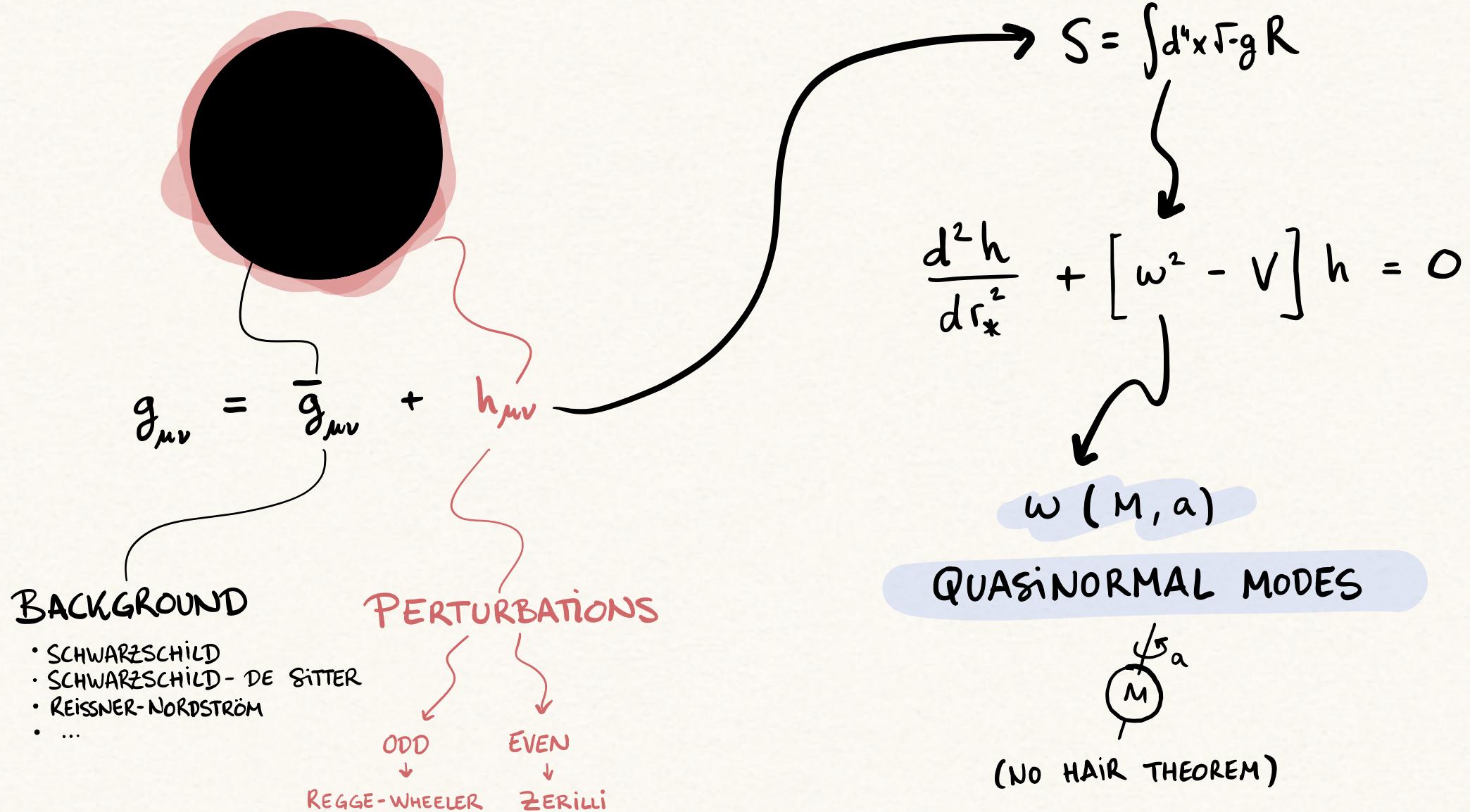
"SMOKING
GUN SIGNAL"

GRAVITATIONAL WAVES : BH PERTURBATION THEORY

ASTROPHYSICAL SOURCES : MERGERS (BLACK HOLES / NEUTRON STARS)



GRAVITATIONAL WAVES : BH PERTURBATION THEORY



GRAVITATIONAL WAVES : BH PERTURBATION THEORY

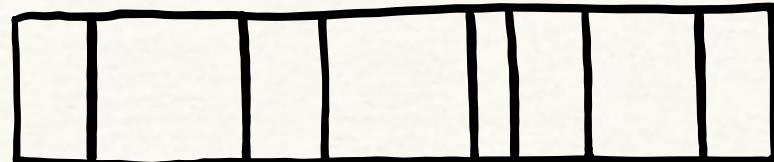
BLACK HOLE SPECTROSCOPY

$\omega(M, a)$

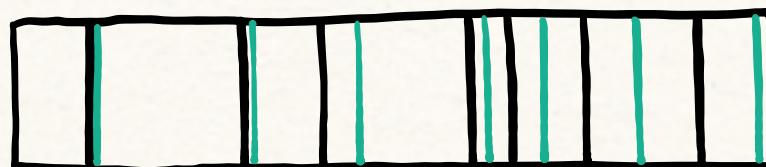
- 1st QNM sets (M, a)



- All other QNMs are fixed in GR



MEASURING QNMs PROVIDES CLEAN TESTS OF
BACKGROUND GEOMETRY AND UNDERLYING THEORY



GR
MG

$\omega(M, a)$
 $\omega(M, a, \alpha)$

RINGDOWN OF HAIRY BHs

IN SHIFT-SYMMETRIC SCALAR-TENSOR THEORIES, BLACK HOLES CANNOT SUPPORT SCALAR HAIR IF:

- ① SPACETIME IS SPHERICALLY SYMMETRIC, STATIC AND ASSYMPTOTICALLY FLAT
- ② SCALAR FIELD IS STATIC $\phi(r)$ AND HAS A VANISHING DERIVATIVE ϕ' AT INFINITY
- ③ NORM OF THE CURRENT ASSOCIATED W/ SHIFT SYMMETRY IS FINITE DOWN TO THE HORIZON
- ④ ACTION CONTAINS A CANONICAL KINETIC TERM $X \in G_2$
- ⑤ ALL $G_i(X)$ FUNCTIONS ARE ANALYTICAL AT $X=0$

[1312.5742] HU, RAVERI, FRUSCINANTE, SILVESTRI

RINGDOWN OF HAIRY BHs

STATIC $\Phi(r)$

- E.G. SCALAR GAUSS-BONET

$$G_2 = \gamma X, G_4 = \zeta, G_5 = \alpha \ln |X|$$

[1312.3622, 1408.1698] SOTIRIOU, ZHOU

- MORE GENERAL

$$G_2 \geq \sqrt{X}, G_3 \geq \ln |X|,$$

$$G_4 \geq \sqrt{X}, G_5 \geq \ln |X|$$

[1702.01938] BABICHEV, CHARMOUSIS, LEHÉBEL

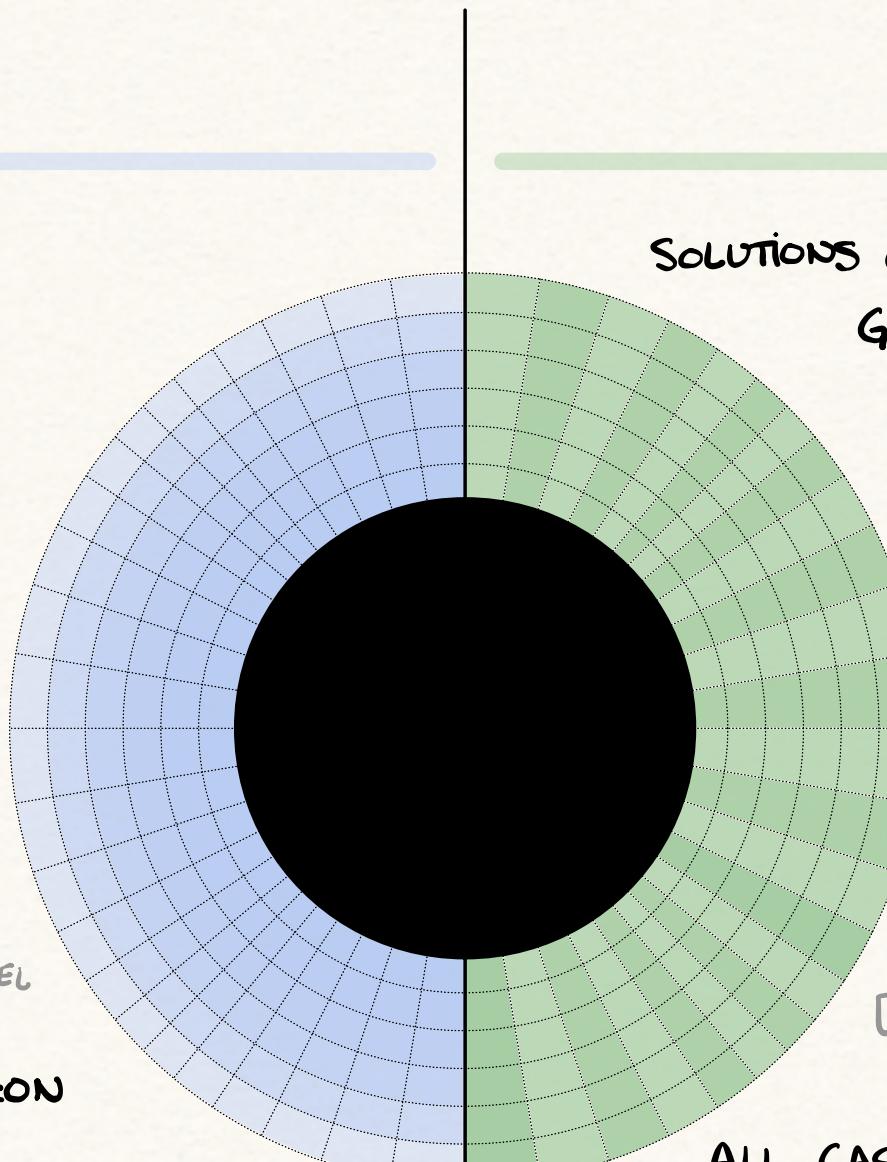
→ FINITE-NORM CURRENT AT HORIZON

→ LACK LORENTZ INVARIANT ($X=0$) SOLUTION

IN MINKOWSKI SPACETIME

[1903.02055] SARAVANI, SOTIRIOU

REAL PROBLEM OR NOT?



TIME-DEPENDENT $\Phi(t,r)$

SOLUTIONS OF SHIFT + REFLECTION SYMM HORNDESKI
 $G_2(x), G_4(x)$

[1312.3204] BABICHEV, CHARMOUSIS

[1403.4364] KOBAYASHI, TANAHASHI

HOWEVER, PRONE TO INSTABILITIES

[1510.07400] OGAWA, KOBAYASHI, SUYAMA

[1803.11444] BABICHEV, CHARMOUSIS,
ESPOSITO-FARÈSE, LEHÉBEL

[1610.00432] TAKAHASHI, SUYAMA

[1904.03554] TAKAHASHI, MOTOHASHI, MINAMITSUJI

[1907.00699] DE RHAM, ZHANG

ALL CASES STUDIED ASSUMED $X = \text{constant}$

SOLUTION WITH $X \neq \text{constant}$

[2310.11919] BAKOPOULOS, CHARMOUSIS, KANTI, LECOEUR, NAKAS

STATIC HAIR AND SPEED OF GRAVITY

[2301.10272] SS + JOHANNES NOUER

SET-UP

THEORY : HORNDESKI WITH $G_{4\phi} = 0$

BACKGROUND : METRIC : SCHWARZSCHILD

SCALAR : STATIC HAIRY DEVIATION $\bar{\Phi} = \hat{\Phi} + \varepsilon \delta\phi(r)$, $\delta\phi = q_c \frac{2M}{r}$

MODIFIED REGGE - WHEELER EQUATION

$$\frac{d^2 Q}{dr_*^2} + [w^2(1 + \alpha_T) + V + \alpha_T \delta V] Q = 0$$

$$\delta V = -\frac{1}{f} \omega_0^2 + \frac{M(2r-5M)}{r^3(r-2M)} + \frac{(l+2)(l-1)}{r^2} - \frac{r-2M}{2r} \left(\left(\frac{\delta\phi''}{\delta\phi'} \right)^2 - \frac{\delta\phi'''}{\delta\phi'} \right) + \frac{r-5M}{r^2} \frac{\delta\phi''}{\delta\phi'}$$

GRAVITATIONAL WAVE
SPEED EXCESS

$$\alpha_T = -f(2M)^2 G_T \delta\phi'^2$$

$$G_T = \frac{G_{4x} - G_{5\phi}}{G_4}$$

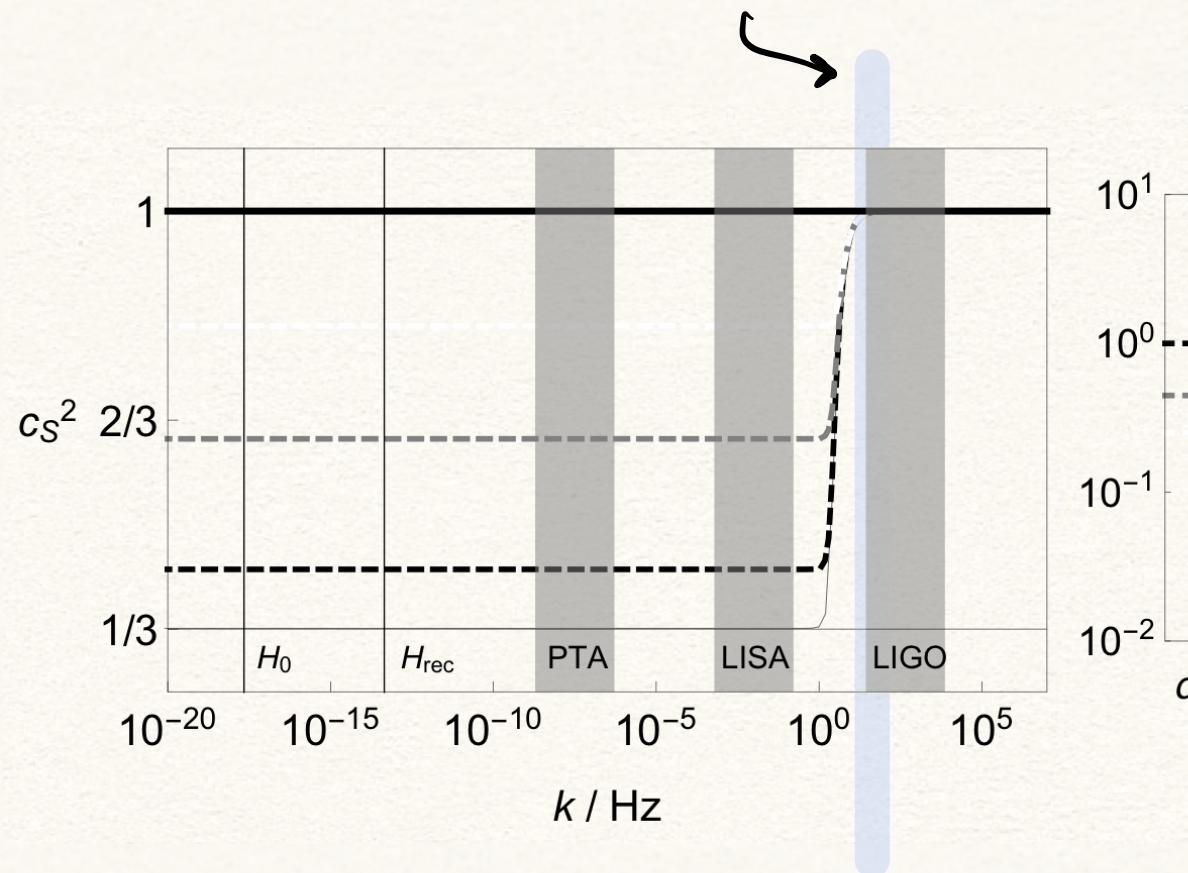
STATIC HAIR AND SPEED OF GRAVITY

[2301.10272] SS + JOHANNES NOUER

WHAT DO WE KNOW ABOUT α_T ?

- LIGO : $\alpha_T \lesssim 10^{-15}$ (GW170817)

- DARK ENERGY EFTs : CUTOFF AT $\sim 10^2$ Hz [1806.09417] MELVILLE, DE RHAM



STATIC HAIR AND SPEED OF GRAVITY

[2301.10272] SS + JOHANNES NOUER

OBTAIN QNM CORRECTIONS WITH PARAMETRISED RINGDOWN FORMALISM

[1901.01265] CARDOSO, KIMURA, MASELLI, BERTI, MACEDO, McMANUS

$$\omega = \omega_0 + \delta\omega(\alpha_T)$$

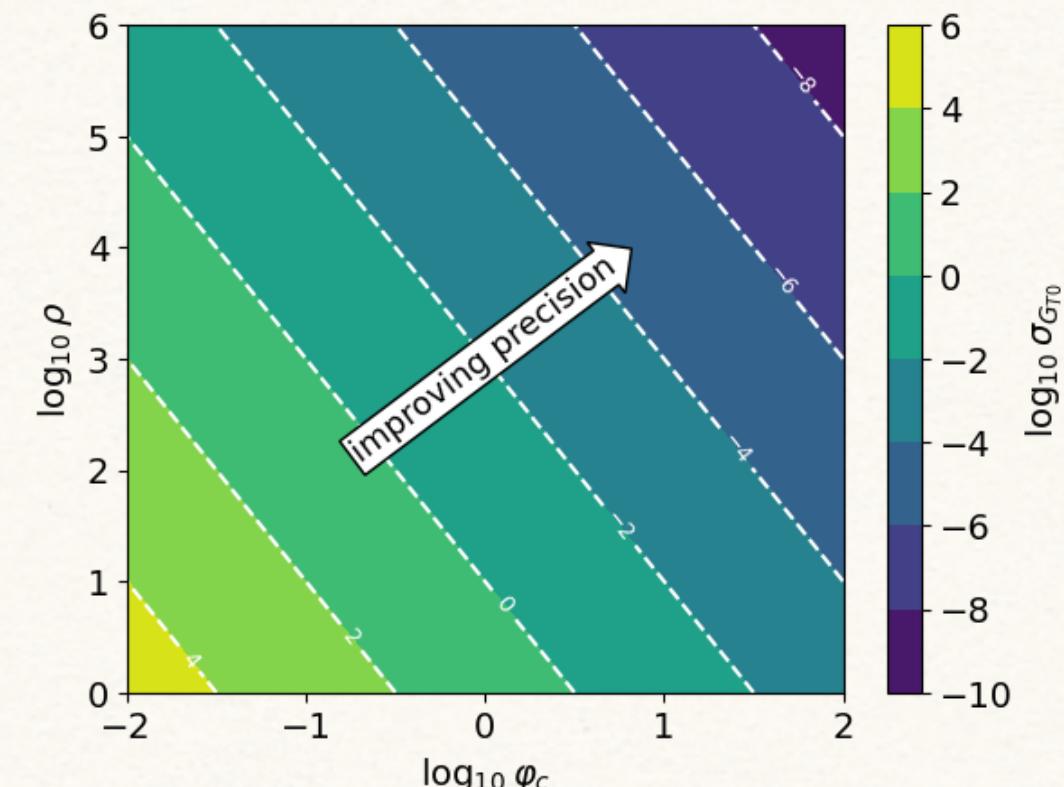
GR ↪

$$\delta\omega = \sum_{i=0}^{\infty} A_{Ti} \left[(2M\omega_0)^2 e_{4+i} - (l(l+1) - 9) e_{6+i} + (l(l+1) - 20) e_{7+i} + \frac{45}{4} e_{8+i} \right]$$

$$A_{Ti} = G_{Ti} \Phi_c^2 , \quad G_T = \frac{1}{(2M)^2} \frac{G_{4x} - G_{5\phi}}{G_4} = \sum_i G_{Ti} \left(\frac{2M}{r} \right)^i$$

FISHER FORECAST ANALYSIS

Detector(s)	Ringdown SNR (ρ)	Error on α_T
LVK	10 [136–138]	1
ET / CE	10^2 [138–141]	10^{-1}
pre-DECIGO	10^2 [142]	10^{-1}
DECIGO / AEDGE	10^3 [143, 144]*	10^{-2}
LISA	10^5 [137, 145]	10^{-4}
TianQin	10^5 [145]	10^{-4}
AMIGO	10^5 [130]	10^{-4}



TIME-DEPENDENT HAIR

[2408.01720] SS + JOHANNES NOUER

$$\phi = q^t + \Psi(r)$$

\mathcal{L}	Background solution		Stability	
	$g_{\mu\nu}$	X	odd	even
(Shift + refl)-sym Horndeski [29, 30]	$S(dS)^*$	$\frac{q^2}{2} = \text{const}$	✓	✗
Cubic Galileon [31, 32]	$S(dS)^*$ (non-exact)	non-const (non-exact)	?	?
Shift-sym breaking Horndeski [65]	✗ (for large subclasses)	✗ (for large subclasses)	-	-
$G_2 = \eta X, G_4 = \zeta + \beta \sqrt{X}$ [33]	$S(dS)^* + RN(dS)^*$ (non-exact)	non-const	?	?
Shift-sym beyond Horndeski [34]	SdS^*	$\frac{q^2}{2} = \text{const}$	✓	✗
Shift-sym breaking quadratic DHOST [35]	$S(dS)^*$	$\frac{q^2}{2} = \text{const}$	✓	✗
Shift-sym quadratic DHOST [36, 66]	$S(dS)^* + K^*$	const	✓	✗
Quadratic DHOST [37]	$S(dS)^* + (K)RN(dS)^*$	const	✓	✗
$G_2 = -2\Lambda + 2\eta \sqrt{X}, G_4 = 1 + \lambda \sqrt{X}$ [1]	$S(dS)$	$\frac{1}{2} \frac{q^2 \lambda}{\lambda + \eta r^2}$	✓(this work)	?

TIME-DEPENDENT HAIR

[2408.01720] SS + JOHANNES NOUER

SET-UP

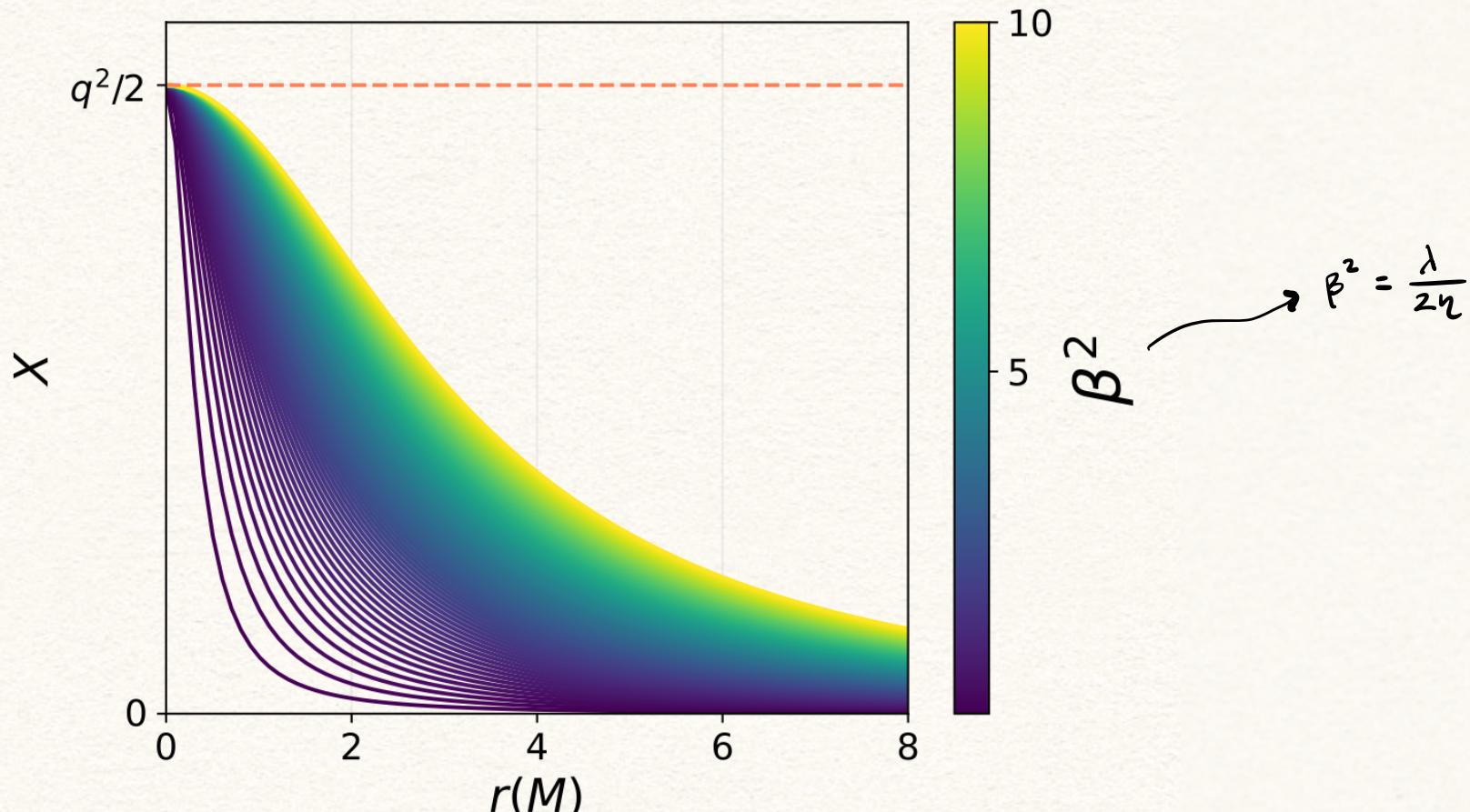
THEORY :

$$\mathcal{L} = -2\Lambda + 2\eta\sqrt{X} + (1 + \lambda\sqrt{X})R + \frac{\lambda}{2\sqrt{X}}(\square\phi^2 + \phi_{\mu\nu}^2)$$

BACKGROUND :

METRIC : SCHWARZSCHILD - DE SITTER

$$\text{SCALAR : } \bar{\phi} = q t + \Psi(r), \quad \Psi'(r)^2 = \frac{q^2}{B} \left(1 - \frac{\lambda B}{\lambda + \eta r^2}\right), \quad X = \frac{1}{2} \frac{q^2 \lambda}{\lambda + \eta r^2}$$



QUADRATIC ACTION

$$S^{(2)} = \frac{1}{2} \int d^4x \sqrt{-g} \left[\mathcal{L}_{GR}^{(2)} + \mathcal{L}_\eta^{(2)} + \mathcal{L}_\lambda^{(2)} \right],$$

$$\begin{aligned} \mathcal{L}_{GR}^{(2)} &= \frac{1}{2} (\nabla^\sigma h^{\mu\nu} (2\nabla_\nu h_{\mu\sigma} - \nabla_\sigma h_{\mu\nu}) + 2\Lambda h_{\mu\nu} h^{\mu\nu}), \\ \mathcal{L}_\eta^{(2)} &= \frac{-\eta}{\sqrt{X}} (X h_{\mu\nu} h^{\mu\nu} + \phi^\mu \phi^\nu h_\mu^\sigma h_{\nu\sigma}), \\ \mathcal{L}_\lambda^{(2)} &= \frac{-\lambda}{2\sqrt{X}} \left[\nabla_\sigma h_{\mu\nu} \left(X (\nabla^\sigma h^{\mu\nu} - 2\nabla^\nu h^{\mu\sigma}) + \phi^\sigma \phi^\rho \left(\frac{1}{2} \nabla_\rho h^{\mu\nu} - 2\nabla^\nu h_\rho^\mu \right) + \phi^\mu \phi^\nu \nabla_\rho h^{\sigma\rho} + 2\phi^\mu \phi^\rho \nabla^\nu h_\rho^\sigma \right) \right. \\ &\quad + h_{\mu\nu} \left(\frac{1}{2} h^{\mu\nu} ((\square\phi)^2 - \phi_{\rho\gamma} \phi^{\rho\gamma}) + 4h_\sigma^\nu (\Lambda \phi^\mu \phi^\sigma - \phi^{\sigma\mu} \square\phi + \phi_\rho^\mu \phi^{\rho\sigma}) + 4h_{\sigma\rho} \phi^{\mu[\sigma} \phi^{\nu]\rho} \right. \\ &\quad \quad \left. + \frac{1}{2X} \left(h_\sigma^\nu \phi^\mu \phi^\sigma (\phi_{\rho\gamma} \phi^{\rho\gamma} - (\square\phi)^2) + 2\phi^{\sigma\rho} (\phi^\mu \phi^\nu \phi_{\{\sigma}^\lambda h_{\rho\}\lambda} - h_{\sigma\rho} \phi^\mu \phi^\nu \square\phi) \right) \right. \\ &\quad \quad \left. + 4\phi^\rho \phi^{\mu\sigma} (2\nabla_{[\sigma} h_{\rho]}^\nu + \nabla^\nu h_{\rho\sigma}) + 2\phi^\mu \phi^{\sigma\rho} (2\nabla_\sigma h_\rho^\nu - \nabla^\nu h_{\sigma\rho}) + 2\phi^\mu \phi^{\nu\sigma} \nabla_\rho h_\sigma^\rho \right. \\ &\quad \quad \left. + 2\phi^\sigma \phi_\sigma^\mu \nabla_\rho h^{\nu\rho} + 2\phi^\sigma \phi_{\sigma\rho} \nabla^{[\nu} h^{\rho]\mu} + 2\square\phi (\phi^\sigma \nabla_\sigma h^{\mu\nu} - 2\phi^\mu \nabla_\sigma h^{\nu\sigma} - 2\phi^\sigma \nabla^\nu h_\sigma^\mu) \right. \\ &\quad \quad \left. + \frac{1}{2X} \phi^\mu \phi^\nu \phi^\sigma (\phi_\sigma^\rho \nabla_\gamma h_\rho^\gamma + \phi^{\rho\gamma} (2\nabla_{[\rho} h_{\gamma]\sigma} - \nabla_\sigma h_{\rho\gamma})) \right), \end{aligned}$$

TIME-DEPENDENT HAIR

[2408.01720] SS + JOHANNES NOUER

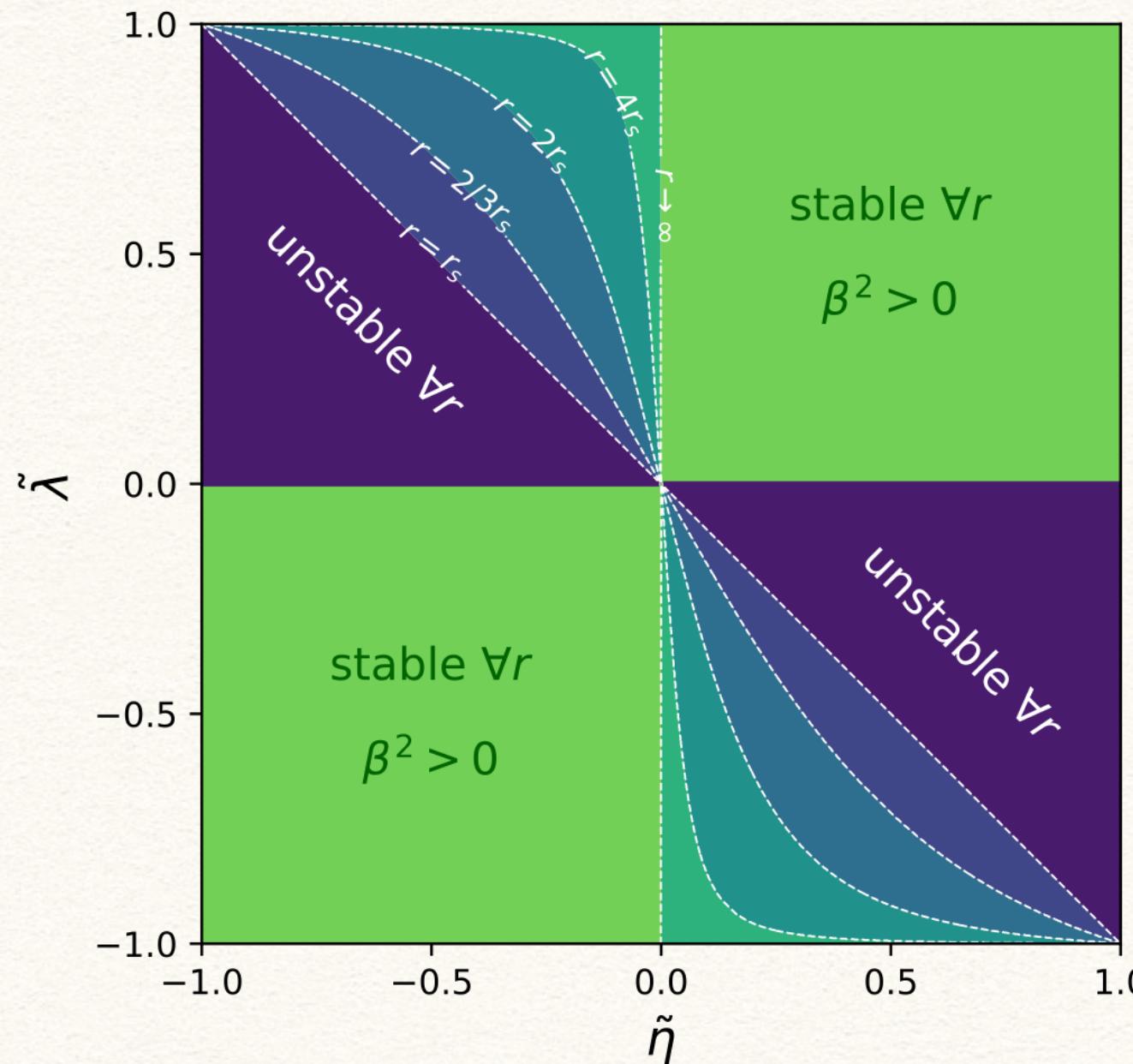
STABILITY CONDITIONS

$$S^{(2)} = \frac{\ell(\ell+1)}{4(\ell-1)(\ell+2)} \int dt dr \left[\tilde{b}_1 (\partial_t Q)^2 - b_2 Q'^2 - (\ell(\ell+1)b_4 + V_{\text{eff}}) Q^2 \right]$$

$$\tilde{b}_1 > 0$$

$$b_2 \geq 0$$

$$b_4 \geq 0$$



MODIFIED REGGE-WHEELER EQUATION

$$\frac{d^2 Q}{dr_*^2} + [w^2 - V] Q = 0$$

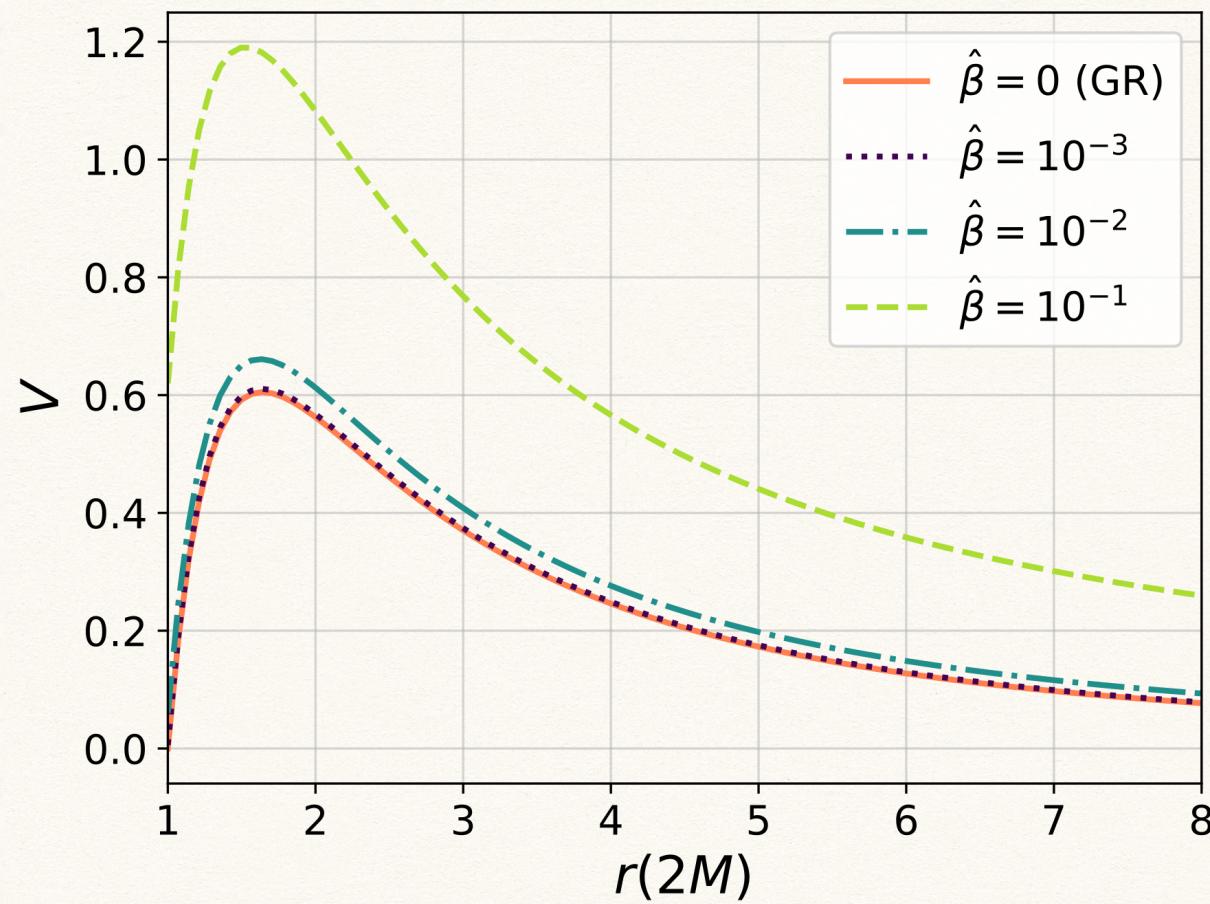
MODIFIED REGGE-WHEELER POTENTIAL

$$\begin{aligned}
 V = & V_{RW} \left(1 + \frac{q\beta\sqrt{2\beta^2+r^2}}{B} \right) + \frac{q\beta(B+q\beta\sqrt{2\beta^2+r^2})}{4r^2(2\beta^2+r^2)(2q\beta^3+\sqrt{2\beta^2+r^2})^3} \times \\
 & \times \left[q^2\beta^6(192\beta^4+104\beta^2r^2+3r^4) + 2(2\beta^2+r^2)(24\beta^4+17\beta^2r^2+2r^4) + 6q\beta^3\sqrt{2\beta^2+r^2}(384\beta^4+416\beta^2r^2+117r^4) \right. \\
 & \left. - \beta^2B(2(48q^2\beta^6(2\beta^2+r^2)+48\beta^4+56\beta^2r^2+19r^4)) + \frac{q\beta^3}{\sqrt{2\beta^2+r^2}}(384\beta^4+416\beta^2r^2+117r^4) \right]. \tag{48}
 \end{aligned}$$

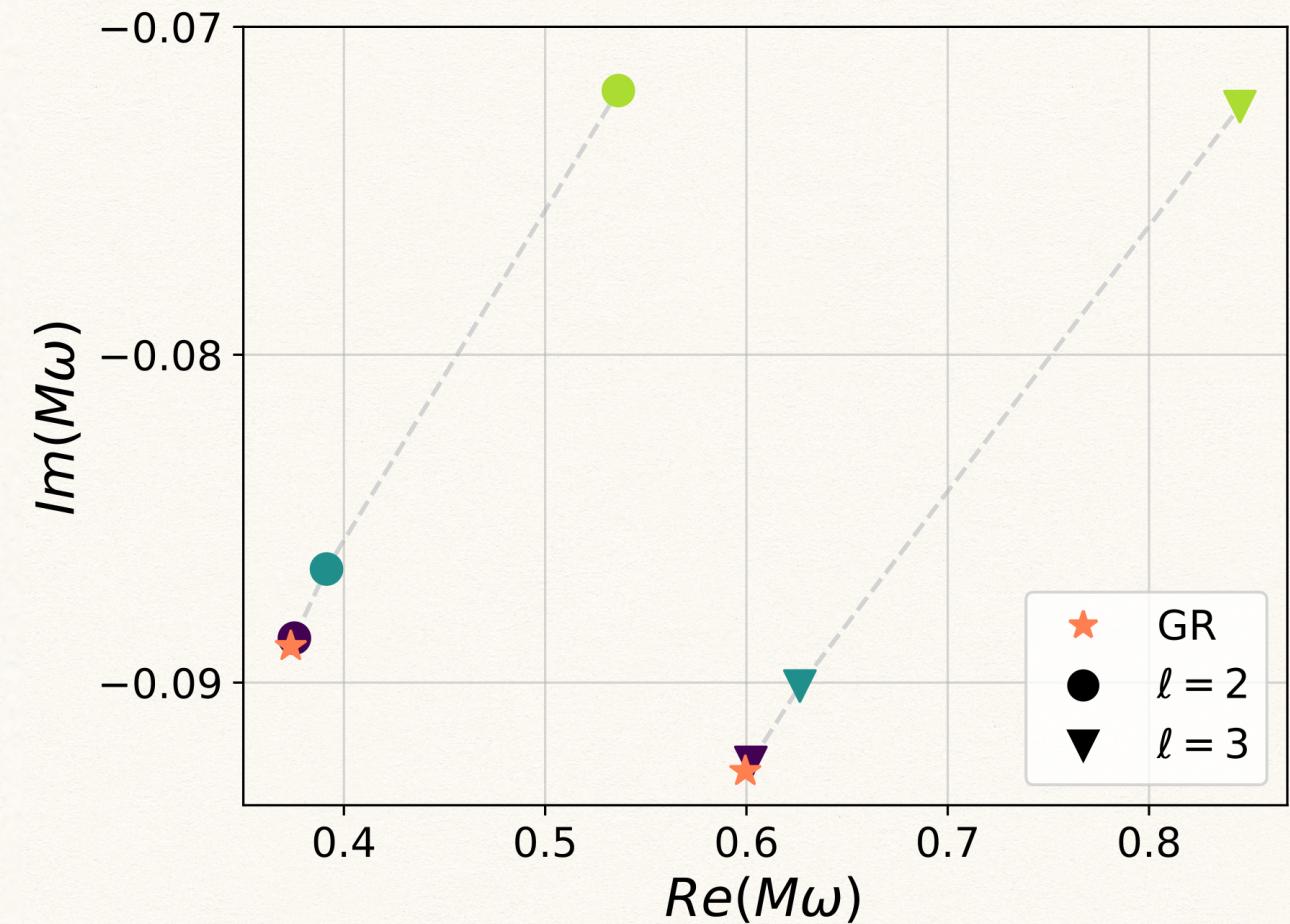
TIME-DEPENDENT HAIR

[2408.01720] SS + JOHANNES NOUER

MODIFIED REGGE-WHEELER POTENTIAL



QNM DEVIATIONS



TIME-DEPENDENT HAIR

[2408.01720] SS + JOHANNES NOUER

QNMs CALCULATED WITH WKB APPROXIMATION

[1904.1033] KONOPLYA, ZHIDENKO, ZINHAILO

$$\omega^2 = V_0 - i \left(n + \frac{1}{2} \right) \sqrt{2V_0^{(2)}} - i \sqrt{2V_0^{(2)}} \sum_{i=2}^N \Lambda_j.$$

LIGHT RING EXPANSION

[1810.07706] FRANCIOULINI, Hui, PENCO, SANTONI, TRINCHERINI

POTENTIAL MAXIMUM

$$\begin{aligned} \frac{M_{\text{Pl}}^2}{M} r_*^{\max} &= 3.2808 - 3.0306 \cdot \hat{\beta} + 0.8316 \cdot \hat{\beta}^2 \\ &\quad - \left(0.22819 + 1.8502 \frac{M_{\text{Pl}}^8}{M^4 q^2} \right) \cdot \hat{\beta}^3 + \mathcal{O}(\hat{\beta}^4) \end{aligned}$$

SEMI-ANALYTIC QNMs

$$\begin{aligned} \frac{M\omega}{M_{\text{Pl}}^2} &= \frac{M\omega_0}{M_{\text{Pl}}^2} + (1.80 + 0.25i)\hat{\beta} \\ &\quad - (2.24 + 1.57i)\hat{\beta}^2 + \mathcal{O}(\hat{\beta}^3), \end{aligned}$$

	$\text{Re}(M\omega)$	$\text{Im}(M\omega)$
$\hat{\beta} = 0$ (GR)	0.3736	-0.0889
$\hat{\beta} = 10^{-3}$	0.3754	-0.0886
$\hat{\beta} = 10^{-2}$	0.3914	-0.0865
$\hat{\beta} = 10^{-1}$	0.5364	-0.0719

FISHER FORECAST ANALYSIS

Detector(s)	Ringdown SNR (ρ)	Error on $\hat{\beta}$
LVK	10 [132–134]	10^{-2}
ET / CE	10^2 [134–137]	10^{-3}
pre-DECIGO	10^2 [138]	10^{-3}
DECIGO / AEDGE	10^3 [139, 140]*	10^{-4}
LISA	10^5 [133, 141]	10^{-6}
TianQin	10^5 [141]	10^{-6}
AMIGO	10^5 [142]	10^{-6}

SUMMARY

- QNMS CAN BE USED TO TEST NEW GRAVITATIONAL PHYSICS

STATIC

- SET UP: HORNDESKI W/ $G_{4\phi} = 0$
PARAMETRISED STATIC HAIRY DEVIATIONS
- QNMS ARE α_T -DEPENDENT IFF \exists SCALAR HAIR
- CALCULATED QNM CORRECTIONS $S_W(\alpha_T)$
W/ PARAMETRISED RINGDOWN FORMALISM
- LISA CAN CONSTRAIN $|\alpha_T| \leq 10^{-4}$
W/ RINGDOWN OF ONE SMBH MERGER

TIME-DEPENDENT

- SET UP: SPECIFIC THEORY W/ $G_2, G_4 \supseteq \sqrt{x}$
LINEAR T-DEP SCALAR W/ $x \neq \text{const}$
- IDENTIFIED STABLE PARAMETER SPACE
- CALCULATED QNM CORRECTIONS $S_W(\hat{\beta})$
W/ WKB APPROXIMATION
- LISA CAN CONSTRAIN $\hat{\beta} \leq 10^{-6}$
W/ RINGDOWN OF ONE SMBH MERGER

THANK YOU!



[ringdown-calculations](#)

Public

A collection of notebooks for black hole perturbation theory calculations in GR and modified gravity.

• Mathematica ★ 7 ⚡ 4