# Part2Part

Iacob Sergiu B1

#### Introducere

Part2Part este o aplicatie de tip peer-to-peer ce satisface nevoia utilizatorilor de a partaja documente cu colegii de lucru. Facilitatile acestui proiect sunt adresate oricui are nevoia de a trimite fisiere unui alt utilizator al acestei aplicatii intr-un mod rapid si sigur. Orice utilizator conectat la platforma va pune la dispozitie mai multe fisiere de pe sistemul de lucru local al acestuia, aceste fisiere putand fi descarcate de ceilalti utilizatori.

## Tehnologiile Utilizate

Acest proiect va implementa functionalitatii protocolului TCP (Transmission Control Protocol)

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# Hamiltonian Mechanics unter besonderer Berücksichtigung der höhreren Lehranstalten

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WWW home page: http://users/~iekeland/web/welcome.html <sup>2</sup> Université de Paris-Sud, Laboratoire d'Analyse Numérique, Bâtiment 425, F-91405 Orsay Cedex, France

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Keywords: computational geometry, graph theory, Hamilton cycles

#### 1 Fixed-Period Problems: The Sublinear Case

With this chapter, the preliminaries are over, and we begin the search for periodic solutions to Hamiltonian systems. All this will be done in the convex case; that is, we shall study the boundary-value problem

$$\dot{x} = JH'(t, x)$$
$$x(0) = x(T)$$

with  $H(t,\cdot)$  a convex function of x, going to  $+\infty$  when  $||x|| \to \infty$ .

#### 1.1 Autonomous Systems

In this section, we will consider the case when the Hamiltonian H(x) is autonomous. For the sake of simplicity, we shall also assume that it is  $C^1$ .

We shall first consider the question of nontriviality, within the general framework of  $(A_{\infty}, B_{\infty})$ -subquadratic Hamiltonians. In the second subsection, we shall look into the special case when H is  $(0, b_{\infty})$ -subquadratic, and we shall try to derive additional information.

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$$\lambda := \text{largest negative eigenvalue of } J \frac{d}{dt} + A_{\infty} .$$
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Theorem 1 tells us that if  $\lambda + \gamma < 0$ , the boundary-value problem:

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**Proposition 1.** Assume H'(0) = 0 and H(0) = 0. Set:

$$\delta := \liminf_{x \to 0} 2N(x) \|x\|^{-2} . \tag{7}$$

If  $\gamma < -\lambda < \delta$ , the solution  $\overline{u}$  is non-zero:

$$\overline{x}(t) \neq 0 \quad \forall t \ .$$
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*Proof.* Condition (7) means that, for every  $\delta' > \delta$ , there is some  $\varepsilon > 0$  such that

$$||x|| \le \varepsilon \Rightarrow N(x) \le \frac{\delta'}{2} ||x||^2$$
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It is an exercise in convex analysis, into which we shall not go, to show that this implies that there is an  $\eta > 0$  such that

$$f \|x\| \le \eta \Rightarrow N^*(y) \le \frac{1}{2\delta'} \|y\|^2$$
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Since  $u_1$  is a smooth function, we will have  $||hu_1||_{\infty} \leq \eta$  for h small enough, and inequality (10) will hold, yielding thereby:

$$\psi(hu_1) \le \frac{h^2}{2} \frac{1}{\lambda} \|u_1\|_2^2 + \frac{h^2}{2} \frac{1}{\delta'} \|u_1\|^2 . \tag{11}$$

Fig. 1. This is the caption of the figure displaying a white eagle and a white horse on a snow field

If we choose  $\delta'$  close enough to  $\delta$ , the quantity  $\left(\frac{1}{\lambda} + \frac{1}{\delta'}\right)$  will be negative, and we end up with

$$\psi(hu_1) < 0 \quad \text{for } h \neq 0 \text{ small }. \tag{12}$$

On the other hand, we check directly that  $\psi(0) = 0$ . This shows that 0 cannot be a minimizer of  $\psi$ , not even a local one. So  $\overline{u} \neq 0$  and  $\overline{u} \neq \Lambda_o^{-1}(0) = 0$ .

Corollary 1. Assume H is  $C^2$  and  $(a_{\infty}, b_{\infty})$ -subquadratic at infinity. Let  $\xi_1, \ldots, \xi_N$  be the equilibria, that is, the solutions of  $H'(\xi) = 0$ . Denote by  $\omega_k$  the smallest eigenvalue of  $H''(\xi_k)$ , and set:

$$\omega := \operatorname{Min} \left\{ \omega_1, \dots, \omega_k \right\} . \tag{13}$$

*If:* 

$$\frac{T}{2\pi}b_{\infty} < -E\left[-\frac{T}{2\pi}a_{\infty}\right] < \frac{T}{2\pi}\omega\tag{14}$$

then minimization of  $\psi$  yields a non-constant T-periodic solution  $\overline{x}$ .

We recall once more that by the integer part  $E[\alpha]$  of  $\alpha \in \mathbb{R}$ , we mean the  $a \in \mathbb{Z}$  such that  $a < \alpha \le a+1$ . For instance, if we take  $a_{\infty} = 0$ , Corollary 2 tells us that  $\overline{x}$  exists and is non-constant provided that:

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Hence:

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The condition  $\gamma < -\lambda < \delta$  now becomes:

$$b_{\infty} - a_{\infty} < -\frac{2\pi}{T}k_o - a_{\infty} < \omega - a_{\infty} \tag{19}$$

which is precisely condition (14).

**Lemma 1.** Assume that H is  $C^2$  on  $\mathbb{R}^{2n} \setminus \{0\}$  and that H''(x) is non-degenerate for any  $x \neq 0$ . Then any local minimizer  $\widetilde{x}$  of  $\psi$  has minimal period T.

*Proof.* We know that  $\widetilde{x}$ , or  $\widetilde{x} + \xi$  for some constant  $\xi \in \mathbb{R}^{2n}$ , is a T-periodic solution of the Hamiltonian system:

$$\dot{x} = JH'(x) \ . \tag{20}$$

There is no loss of generality in taking  $\xi = 0$ . So  $\psi(x) \ge \psi(\widetilde{x})$  for all  $\widetilde{x}$  in some neighbourhood of x in  $W^{1,2}\left(\mathbb{R}/T\mathbb{Z};\mathbb{R}^{2n}\right)$ .

But this index is precisely the index  $i_T(\widetilde{x})$  of the *T*-periodic solution  $\widetilde{x}$  over the interval (0,T), as defined in Sect. 2.6. So

$$i_T(\widetilde{x}) = 0. (21)$$

Now if  $\tilde{x}$  has a lower period, T/k say, we would have, by Corollary 31:

$$i_T(\widetilde{x}) = i_{kT/k}(\widetilde{x}) \ge ki_{T/k}(\widetilde{x}) + k - 1 \ge k - 1 \ge 1.$$
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This would contradict (21), and thus cannot happen.

*Notes and Comments.* The results in this section are a refined version of [1]; the minimality result of Proposition 14 was the first of its kind.

To understand the nontriviality conditions, such as the one in formula (16), one may think of a one-parameter family  $x_T$ ,  $T \in (2\pi\omega^{-1}, 2\pi b_{\infty}^{-1})$  of periodic solutions,  $x_T(0) = x_T(T)$ , with  $x_T$  going away to infinity when  $T \to 2\pi\omega^{-1}$ , which is the period of the linearized system at 0.

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**Theorem 1 (Ghoussoub-Preiss).** Assume H(t,x) is  $(0,\varepsilon)$ -subquadratic at infinity for all  $\varepsilon > 0$ , and T-periodic in t

$$H(t, \cdot)$$
 is convex  $\forall t$  (23)

$$H(\cdot, x)$$
 is  $T$ -periodic  $\forall x$  (24)

$$H(t,x) \ge n(\|x\|)$$
 with  $n(s)s^{-1} \to \infty$  as  $s \to \infty$  (25)

$$\forall \varepsilon > 0 , \quad \exists c : H(t, x) \le \frac{\varepsilon}{2} \|x\|^2 + c .$$
 (26)

Assume also that H is  $C^2$ , and H''(t,x) is positive definite everywhere. Then there is a sequence  $x_k$ ,  $k \in \mathbb{N}$ , of kT-periodic solutions of the system

$$\dot{x} = JH'(t, x) \tag{27}$$

such that, for every  $k \in \mathbb{N}$ , there is some  $p_o \in \mathbb{N}$  with:

$$p \ge p_o \Rightarrow x_{pk} \ne x_k \ . \tag{28}$$

Example 1 (External forcing). Consider the system:

$$\dot{x} = JH'(x) + f(t) \tag{29}$$

where the Hamiltonian H is  $(0, b_{\infty})$ -subquadratic, and the forcing term is a distribution on the circle:

$$f = \frac{d}{dt}F + f_o \quad \text{with } F \in L^2\left(\mathbb{R}/T\mathbb{Z}; \mathbb{R}^{2n}\right) ,$$
 (30)

where  $f_o := T^{-1} \int_o^T f(t) dt$ . For instance,

$$f(t) = \sum_{k \in \mathbb{N}} \delta_k \xi , \qquad (31)$$

where  $\delta_k$  is the Dirac mass at t = k and  $\xi \in \mathbb{R}^{2n}$  is a constant, fits the prescription. This means that the system  $\dot{x} = JH'(x)$  is being excited by a series of identical shocks at interval T.

**Definition 1.** Let  $A_{\infty}(t)$  and  $B_{\infty}(t)$  be symmetric operators in  $\mathbb{R}^{2n}$ , depending continuously on  $t \in [0,T]$ , such that  $A_{\infty}(t) \leq B_{\infty}(t)$  for all t.

continuously on  $t \in [0,T]$ , such that  $A_{\infty}(t) \leq B_{\infty}(t)$  for all t. A Borelian function  $H: [0,T] \times \mathbb{R}^{2n} \to \mathbb{R}$  is called  $(A_{\infty}, B_{\infty})$ -subquadratic at infinity if there exists a function N(t,x) such that:

$$H(t,x) = \frac{1}{2} (A_{\infty}(t)x, x) + N(t,x)$$
 (32)

$$\forall t$$
,  $N(t,x)$  is convex with respect to  $x$  (33)

$$N(t,x) \ge n(\|x\|)$$
 with  $n(s)s^{-1} \to +\infty$  as  $s \to +\infty$  (34)

$$\exists c \in \mathbb{R} : H(t,x) \le \frac{1}{2} (B_{\infty}(t)x, x) + c \quad \forall x .$$
 (35)

If  $A_{\infty}(t) = a_{\infty}I$  and  $B_{\infty}(t) = b_{\infty}I$ , with  $a_{\infty} \leq b_{\infty} \in \mathbb{R}$ , we shall say that H is  $(a_{\infty}, b_{\infty})$ -subquadratic at infinity. As an example, the function  $||x||^{\alpha}$ , with  $1 \leq \alpha < 2$ , is  $(0, \varepsilon)$ -subquadratic at infinity for every  $\varepsilon > 0$ . Similarly, the Hamiltonian

$$H(t,x) = \frac{1}{2}k \|k\|^2 + \|x\|^{\alpha}$$
(36)

is  $(k, k + \varepsilon)$ -subquadratic for every  $\varepsilon > 0$ . Note that, if k < 0, it is not convex.

Notes and Comments. The first results on subharmonics were obtained by Rabinowitz in [5], who showed the existence of infinitely many subharmonics both in the subquadratic and superquadratic case, with suitable growth conditions on H'. Again the duality approach enabled Clarke and Ekeland in [2] to treat the same problem in the convex-subquadratic case, with growth conditions on H only.

Recently, Michalek and Tarantello (see [3] and [4]) have obtained lower bound on the number of subharmonics of period kT, based on symmetry considerations and on pinching estimates, as in Sect. 5.2 of this article.

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- Clarke, F., Ekeland, I.: Nonlinear oscillations and boundary-value problems for Hamiltonian systems. Arch. Rat. Mech. Anal. 78, 315–333 (1982)
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## Hamiltonian Mechanics Two

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$$\dot{x} = JH'(t, x) \tag{27}$$

such that, for every  $k \in \mathbb{N}$ , there is some  $p_o \in \mathbb{N}$  with:

$$p \ge p_o \Rightarrow x_{pk} \ne x_k \ . \tag{28}$$

Example 1 (External forcing). Consider the system:

$$\dot{x} = JH'(x) + f(t) \tag{29}$$

where the Hamiltonian H is  $(0, b_{\infty})$ -subquadratic, and the forcing term is a distribution on the circle:

$$f = \frac{d}{dt}F + f_o \quad \text{with } F \in L^2\left(\mathbb{R}/T\mathbb{Z}; \mathbb{R}^{2n}\right) ,$$
 (30)

where  $f_o := T^{-1} \int_0^T f(t) dt$ . For instance,

$$f(t) = \sum_{k \in \mathbb{N}} \delta_k \xi , \qquad (31)$$

where  $\delta_k$  is the Dirac mass at t = k and  $\xi \in \mathbb{R}^{2n}$  is a constant, fits the prescription. This means that the system  $\dot{x} = JH'(x)$  is being excited by a series of identical shocks at interval T.

**Definition 1.** Let  $A_{\infty}(t)$  and  $B_{\infty}(t)$  be symmetric operators in  $\mathbb{R}^{2n}$ , depending continuously on  $t \in [0,T]$ , such that  $A_{\infty}(t) \leq B_{\infty}(t)$  for all t.

A Borelian function  $H:[0,T]\times\mathbb{R}^{2n}\to\mathbb{R}$  is called  $(A_{\infty},B_{\infty})$ -subquadratic at infinity if there exists a function N(t,x) such that:

$$H(t,x) = \frac{1}{2} (A_{\infty}(t)x, x) + N(t,x)$$
 (32)

$$\forall t$$
,  $N(t,x)$  is convex with respect to  $x$  (33)

$$N(t,x) \ge n(\|x\|)$$
 with  $n(s)s^{-1} \to +\infty$  as  $s \to +\infty$  (34)

$$\exists c \in \mathbb{R} : H(t,x) \le \frac{1}{2} (B_{\infty}(t)x, x) + c \quad \forall x .$$
 (35)

If  $A_{\infty}(t) = a_{\infty}I$  and  $B_{\infty}(t) = b_{\infty}I$ , with  $a_{\infty} \leq b_{\infty} \in \mathbb{R}$ , we shall say that H is  $(a_{\infty}, b_{\infty})$ -subquadratic at infinity. As an example, the function  $||x||^{\alpha}$ , with  $1 \leq \alpha < 2$ , is  $(0, \varepsilon)$ -subquadratic at infinity for every  $\varepsilon > 0$ . Similarly, the Hamiltonian

$$H(t,x) = \frac{1}{2}k \|k\|^2 + \|x\|^{\alpha}$$
(36)

is  $(k, k + \varepsilon)$ -subquadratic for every  $\varepsilon > 0$ . Note that, if k < 0, it is not convex.

Notes and Comments. The first results on subharmonics were obtained by Rabinowitz in 1985, who showed the existence of infinitely many subharmonics both in the subquadratic and superquadratic case, with suitable growth conditions on H'. Again the duality approach enabled Clarke and Ekeland in 1981 to treat the same problem in the convex-subquadratic case, with growth conditions on H only.

Recently, Michalek and Tarantello (see Michalek, R., Tarantello, G. 1982 and Tarantello, G. 1983) have obtained lower bound on the number of subharmonics of period kT, based on symmetry considerations and on pinching estimates, as in Sect. 5.2 of this article.

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