

Homework - 3

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Question 1

The following are the point group operations which leaves the cluster unchanged:

Identity (E) : $\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$, Inverse (I) : $\begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix}$, Mirror-x (M_x) : $\begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$, Mirror-y (M_y) : $\begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$.

The force constant matrix can be expressed as a sum over tensor basis:

$$\begin{bmatrix} \phi_{xx} & \phi_{xy} \\ \phi_{yx} & \phi_{yy} \end{bmatrix} = \phi_{xx} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} + \phi_{xy} \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} + \phi_{yx} \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix} + \phi_{yy} \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \quad (1)$$

To find the symmetry invariant basis, apply Reynold's operator :

$$\tilde{\Lambda}_i = \frac{1}{|S|} \sum_S S^T \Lambda_i S \quad (2)$$

where S can be any of the symmetry operator which leaves the cluster unchanged (E, I, M_x , M_y).

Applying Reynold's operator on $\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$ gives:

$$\begin{aligned} \frac{1}{4} \left(\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} + \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} + \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \right) \\ = \frac{1}{4} \begin{bmatrix} 4 & 0 \\ 0 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \end{aligned}$$

Applying Reynold's operator on $\begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$ gives:

$$\begin{aligned} \frac{1}{4} \left(\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} + \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} + \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \right) \\ = \frac{1}{4} \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} \end{aligned}$$

Applying Reynold's operator on $\begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}$ gives:

$$\begin{aligned} \frac{1}{4} \left(\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} + \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} + \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \right) \\ = \frac{1}{4} \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} \end{aligned}$$

Applying Reynold's operator on $\begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$ gives:

$$\begin{aligned} \frac{1}{4} \left(\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix} + \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} + \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \right) \\ = \frac{1}{4} \begin{bmatrix} 0 & 0 \\ 0 & 4 \end{bmatrix} = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \end{aligned}$$

The force constant matrix (Φ) expressed in the symmetry invariant tensor basis is,

$$\Phi = \alpha_1 \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} + \alpha_2 \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} = \boxed{\begin{bmatrix} \alpha_1 & 0 \\ 0 & \alpha_2 \end{bmatrix}} \quad (3)$$

Question 2

Dynamical matrix (D) is given by,

$$D = \frac{1}{m} \sum_{\vec{r}_l} \Phi(0, \vec{r}_l) e^{-i\vec{k} \cdot \vec{r}_l} \quad (4)$$

where \vec{r}_l is position vector of the neighboring atoms, \vec{k} is the crystal momentum and Φ is the force constant matrix. To obtain D, we need Φ of all the nearest neighbors. It can be obtained using the relation: $\Phi(0, \vec{r}_l) = S^T \Phi(0, \vec{r}_1) S$, where S is the symmetry operation which transforms the cluster $(0, \vec{r}_1)$ to $(0, \vec{r}_l)$. The following are the force constant matrices for all the 6 pair clusters.

$$\begin{aligned} \Phi(0, r_1) &= \begin{bmatrix} \alpha_1 & 0 \\ 0 & \alpha_2 \end{bmatrix} \\ \Phi(0, r_2) &= \begin{bmatrix} \frac{1}{4}\alpha_1 + \frac{3}{4}\alpha_2 & \frac{\sqrt{3}}{4}(\alpha_1 - \alpha_2) \\ \frac{\sqrt{3}}{4}(\alpha_1 - \alpha_2) & \frac{3}{4}\alpha_1 + \frac{1}{4}\alpha_2 \end{bmatrix} \\ \Phi(0, r_3) &= \begin{bmatrix} \frac{1}{4}\alpha_1 + \frac{3}{4}\alpha_2 & \frac{\sqrt{3}}{4}(\alpha_2 - \alpha_1) \\ \frac{\sqrt{3}}{4}(\alpha_2 - \alpha_1) & \frac{3}{4}\alpha_1 + \frac{1}{4}\alpha_2 \end{bmatrix} \\ \Phi(0, r_4) &= \begin{bmatrix} \alpha_1 & 0 \\ 0 & \alpha_2 \end{bmatrix} \\ \Phi(0, r_5) &= \begin{bmatrix} \frac{1}{4}\alpha_1 + \frac{3}{4}\alpha_2 & \frac{\sqrt{3}}{4}(\alpha_1 - \alpha_2) \\ \frac{\sqrt{3}}{4}(\alpha_1 - \alpha_2) & \frac{3}{4}\alpha_1 + \frac{1}{4}\alpha_2 \end{bmatrix} \\ \Phi(0, r_6) &= \begin{bmatrix} \frac{1}{4}\alpha_1 + \frac{3}{4}\alpha_2 & \frac{\sqrt{3}}{4}(\alpha_2 - \alpha_1) \\ \frac{\sqrt{3}}{4}(\alpha_2 - \alpha_1) & \frac{3}{4}\alpha_1 + \frac{1}{4}\alpha_2 \end{bmatrix} \end{aligned}$$

Plugging in the force constant matrices in Equation. 4, gives:

$$D = \frac{1}{m} \begin{bmatrix} 2\alpha_1 \cos(ak_x) + (\alpha_1 + 3\alpha_2) \cos(\frac{a}{2}k_x) \cos(\frac{a\sqrt{3}}{2}k_y) & -\sqrt{3}(\alpha_1 - \alpha_2) \sin(\frac{a}{2}k_x) \sin(\frac{a\sqrt{3}}{2}k_y) \\ -\sqrt{3}(\alpha_1 - \alpha_2) \sin(\frac{a}{2}k_x) \sin(\frac{a\sqrt{3}}{2}k_y) & 2\alpha_2 \cos(ak_x) + (3\alpha_1 + \alpha_2) \cos(\frac{a}{2}k_x) \cos(\frac{a\sqrt{3}}{2}k_y) \end{bmatrix} \quad (5)$$

Question 3

Considering $(a, 0)$, $(\frac{a}{2}, \frac{a\sqrt{3}}{2})$ to be the real space lattice vectors of the triangular lattice, the reciprocal lattice vectors are given by: $\frac{2\pi}{a}(0, \frac{2}{\sqrt{3}})$ and $\frac{2\pi}{a}(1, \frac{-1}{\sqrt{3}})$.

Figure.1 shows the first Brillouin zone (shaded portion) and some of the high symmetry points (black dots) of the triangular lattice.

Question 4

The analytical expression for the eigen values of the Dynamical matrix obtained in Question 2 is given by,

$$\begin{aligned} e &= \frac{1}{m} \left((\alpha_1 + \alpha_2) \left(\cos(ak_x) + 2 \cos(\frac{a}{2}k_x) \cos(\frac{a\sqrt{3}}{2}k_y) \right) \right. \\ &\quad \left. \pm (\alpha_1 - \alpha_2) \sqrt{\left(\cos(ak_x) - \cos(\frac{a}{2}k_x) \cos(\frac{a\sqrt{3}}{2}k_y) \right)^2 + 3 \sin^2(\frac{a}{2}k_x) \sin^2(\frac{a\sqrt{3}}{2}k_y)} \right) \end{aligned} \quad (6)$$

Figure.2 shows the dispersion curves along some high symmetry directions (Γ, K, M) for $m = 1$, $\alpha_1 = 0.1$, and $\alpha_2 = 0.2$.

Question 5

The displacement field corresponding to an eigen vector (e_α) at a k point (\vec{k}) is given by:

$$\vec{u}(\vec{k}) = c \vec{e}_\alpha(\vec{k}) \quad (7)$$

Figure. 3 shows the displacement field corresponding to the two eigen vectors at M point. The black and blue arrows represent displacement fields corresponding to 2 different eigen vectors.

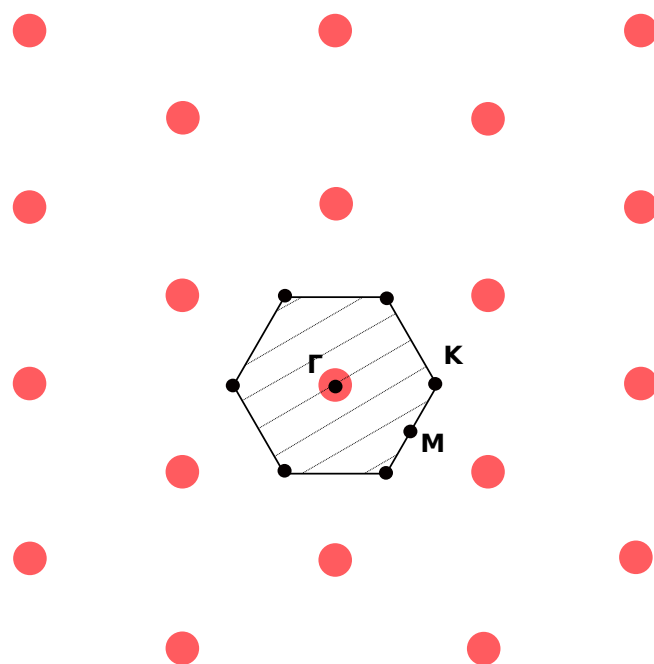


Figure 1: Reciprocal lattice

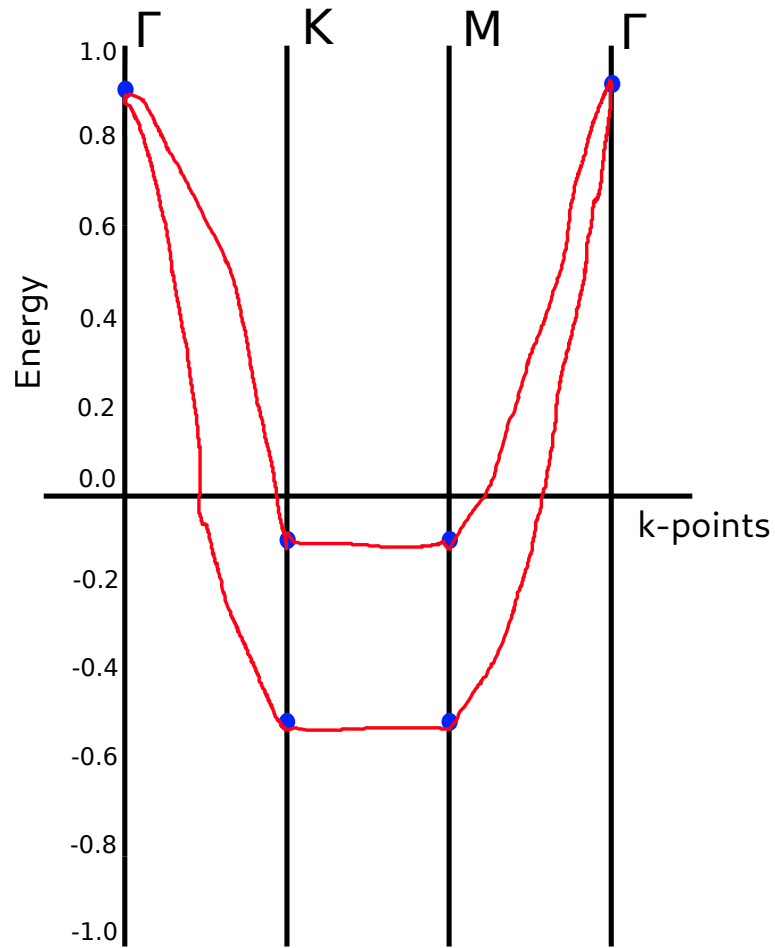


Figure 2: Dispersion curves

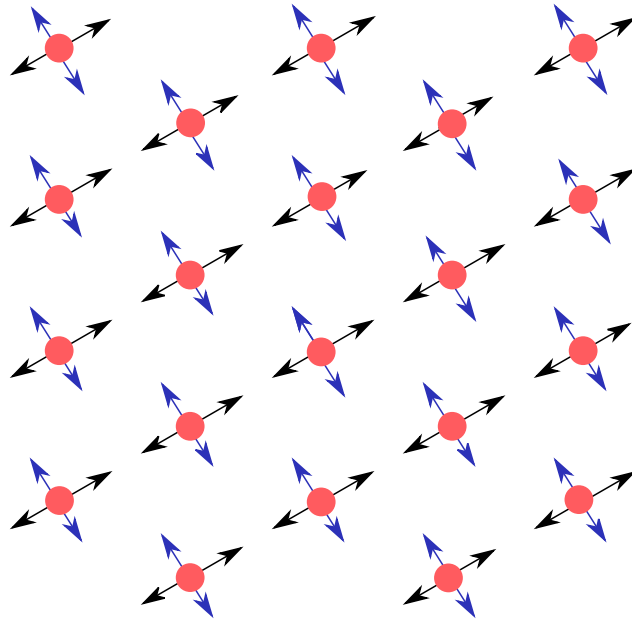


Figure 3: Displacement field