

# Introduction to Quantum Computation and Quantum Algorithms

January 8, 2024

## Part I : Foundations, Protocols and Algorithms

1. Axioms of Quantum Mechanics
2. Quantum Gates
3. Quantum Protocols
4. Quantum Algorithms

# Focus on Ideas. Contrast with Conventional (Classical) bits

## 1. Simplicity and Ideas at the cost of Generality

- ▶ Ex.  $\mathbb{R}^2, \mathbb{R}^3$  or Finite Dim. Inner product spaces instead of Hilbert spaces.

## 2. Comparison with classical bits, notions - A Running Thread.

## 3. Pictorial. Dont get bogged down by the math.

- ▶ Fine to not grasp text on a slide.

The **Power** of Quantum Algorithms, Quantum Cryptography

crucially relies on

**Unique Behaviour** of Quantum Systems - Superposition, Entanglement, etc.

To understand, design, leverage this power,

An Understanding of the Behaviour of Quantum Systems is Necessary.

Behaviour of Quantum Systems described through

Axioms of Quantum Mechanics ← Our First Topic

# Axioms of Quantum Mechanics

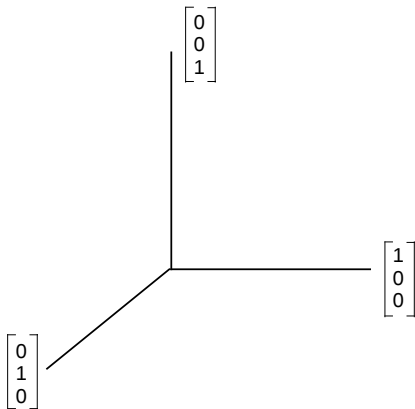
A bit lives in  $\{0, 1\}$  ( it's **state space**). It is 0 or 1.

Where does a **Qubit** live?

Axiom 1

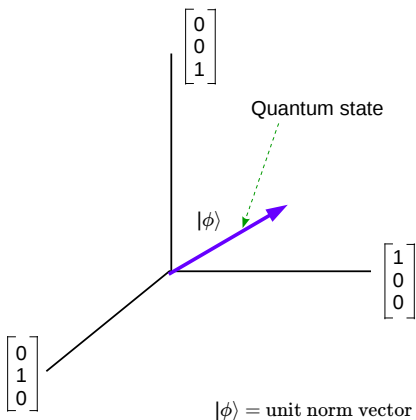


## Axiom 1 : How is a Quantum system described?



State Space of a quantum system is an Inner Product Space (IPS).

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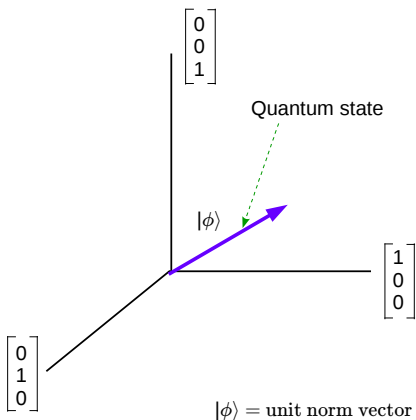


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The state of a quantum system is described through a Unit vector in an IPS  $\mathcal{H}$ .

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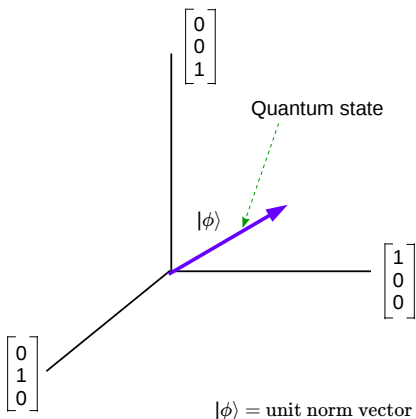
$|\phi\rangle$  : **unit vector** in  $\mathcal{H}$ .

$$\text{Ex. : } |\phi\rangle = \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \\ 0 \end{bmatrix} \in \mathcal{H} = \mathbb{R}^3.$$

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Polarization of photon, spin of electron.

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The state of a quantum system is described through a **Unit vector** in an **IPS**  $\mathcal{H}$ .

# Why $\mathcal{H}$ ? What is the General Theory?

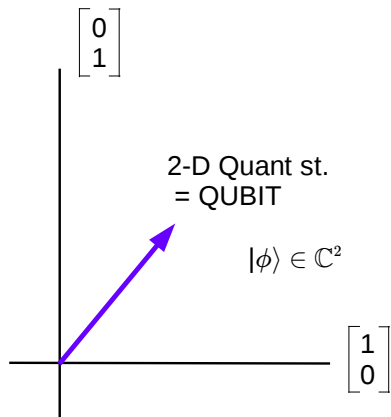
General Quantum Theory is based on a **Hilbert space**. Hence  $\mathcal{H}$ .

Mathematician : **Hilbert space** is a complete  $\infty$ -dimensional inner product space.

This tutorial : Euclidean space with std. inner product suffices  $\leftarrow$  **our Hilbert space**.

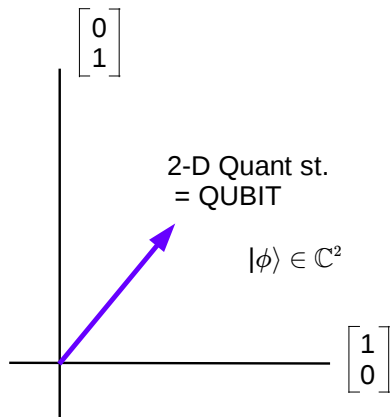
$\mathbb{R}^d$  suffices. But we denote it as  $\mathbb{C}^d$ . **Pretend  $\mathbb{C} = \mathbb{R}$ .**

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A 2- dimensional quantum state is a QUBIT.

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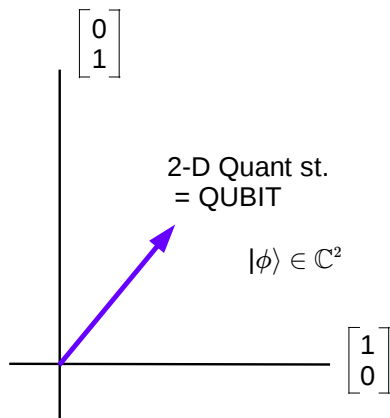


A 2- dimensional quantum state is a QUBIT.

Two Special Qubits

$$|0\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \quad \text{and} \quad |1\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

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### Two Special Qubits

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### CAUTION

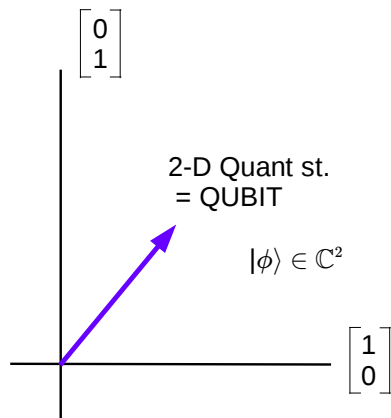
For **any**  $\alpha, \beta \in \mathbb{C}$  s.t  $|\alpha|^2 + |\beta|^2 = 1$

$\alpha|0\rangle + \beta|1\rangle$  is **valid qubit**

Valid state of a quantum system.



## Axiom 1 : How is a Quantum system described?



A 2- dimensional quantum state is a QUBIT.

**Incorrect illustration** : Scalars are Complex numbers.

**Correct illustration** via 3-dimensional **Bloch** sphere.

Two Special Qubits

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## Axiom 1 : Superposition and Inner Products.

Suppose System is in state  $|\phi\rangle = \alpha|0\rangle + \beta|1\rangle$ .  $|\phi\rangle$  is a **Superposition state**.

**INCORRECT:** System is in state  $|0\rangle$  with prob.  $|\alpha|^2$  and in state  $|1\rangle$  with prob.  $|\beta|^2$ .

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The inner product (IP) between  $|x\rangle \in \mathcal{H}$  and  $|y\rangle \in \mathcal{H}$  is denoted  $\langle y|x\rangle$ .

Example :  $|x\rangle = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} \in \mathbb{C}^2$ ,  $|y\rangle = \begin{bmatrix} y_1 \\ y_2 \end{bmatrix} \in \mathbb{C}^2$ ,

$\langle y|x\rangle = y_1^* x_1 + y_2^* x_2$ . Note : First argument is  $\mathbb{C}$ -conjugated. Physics Notation.

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$$\langle y|x\rangle = y_1x_1 + y_2x_2.$$

Qubits are our Information Carriers. Analogous to Bits.

# Axiom 1 : Contrasting Quantum and Classical Worlds

## Quantum World

Qubit : Unit vector in a Inner product space.

$\mathcal{H} \equiv$  Inner product space.

$|\phi\rangle$  : where we encode our information.

$|\phi\rangle \in \mathbb{R}^2$  is a qubit.

## Classical World

Bit : Element in a Finite Set

$\mathcal{X}$  - Our Finite set

$x \equiv$  the information we wish to encode.

$x$  in  $\mathcal{X} = \{0, 1\}$  is a bit.

## Points to Keep in Mind

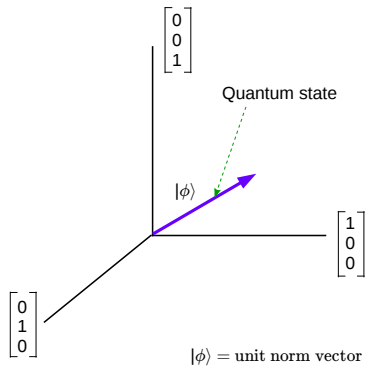
Unit norm.

# Acronyms, Abbreviations and Short Forms

IP FDIPS dim.



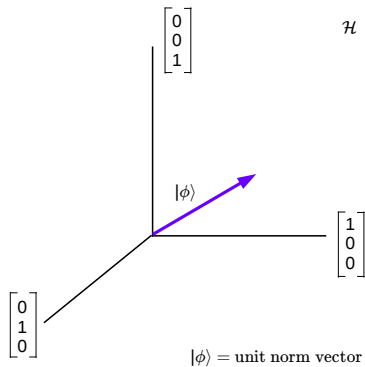
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Linear Transformation (LT) :  $T : \mathcal{H} \rightarrow \mathcal{H}$

$$T|\phi\rangle$$



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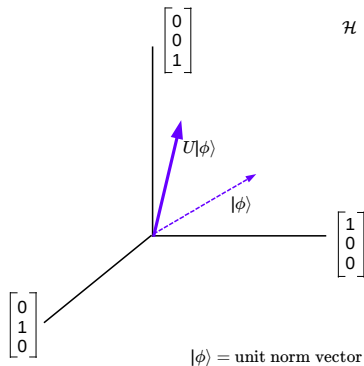
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$\mathcal{H}$  Unitary Transf. : LT that preserves length.

Just a rotation

$$U : \mathcal{H} \rightarrow \mathcal{H}$$



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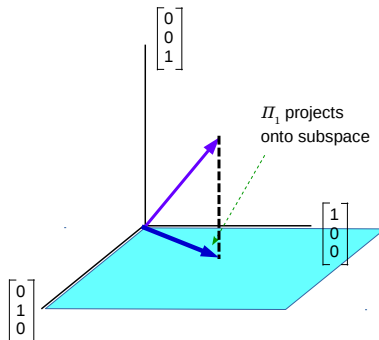
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$$U : \mathcal{H} \rightarrow \mathcal{H}$$

Projection : LT that projects.

Just a projection  $\Pi_1 : \mathcal{H} \rightarrow \mathcal{H}$



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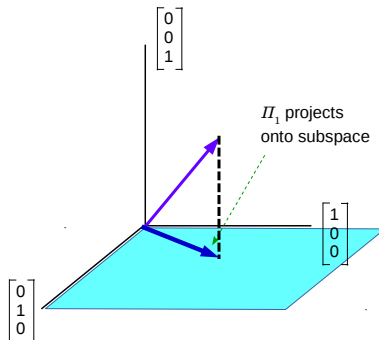
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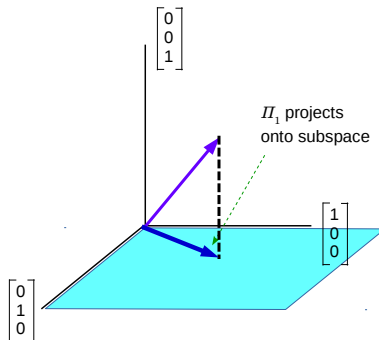
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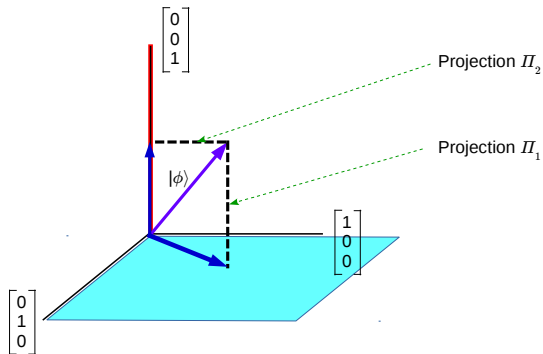
$$\Pi_1 = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Important

Any projector  $\Pi$  satisfies  $\Pi^2 = \Pi^\dagger = \Pi$ .

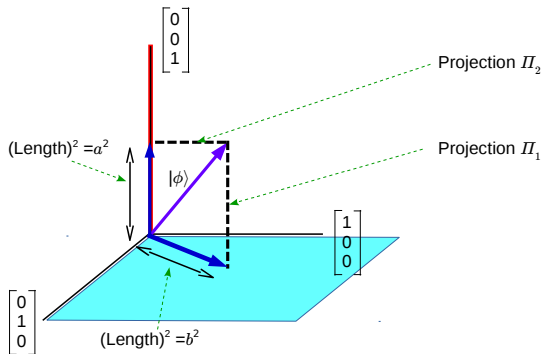


## More on Projections



Projections  $\Pi_1, \Pi_2 : \mathcal{H} \rightarrow \mathcal{H}$ .

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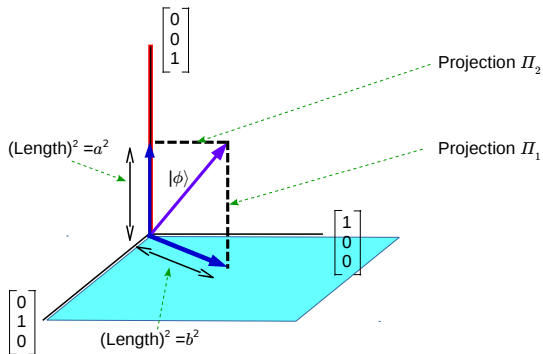


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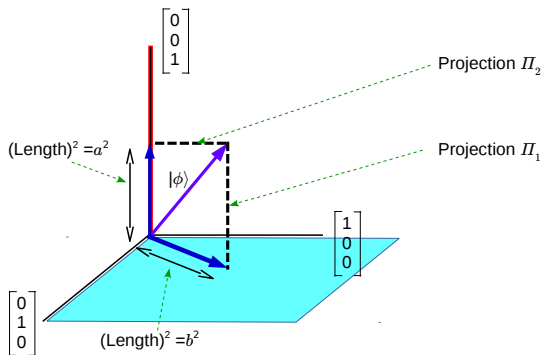


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$$\begin{aligned} a^2 + b^2 &= (\text{length of } |\phi\rangle)^2 \\ &= |\langle \phi | \phi \rangle|^2 = 1 \end{aligned}$$

## Axiom 2 : How does a Closed system evolve?

The evolution of a closed (isolated) quantum system evolves through a Unitary Transformation.

$|x\rangle_{t_1} \equiv$  State of System at time  $t_1$ ,  $|x\rangle_{t_2} \equiv$  State of System at time  $t_2$

$|x\rangle_{t_2}$  is related to  $|x\rangle_{t_1}$  through a Unitary transformation  $U$ .

$$|x\rangle_{t_2} = U|x\rangle_{t_1}$$

Axiom 3 :

Our Interaction with a Quantum System and the  
Rules that Govern this Interaction

Axiom 3 is the **Measurement Axiom**

## Axiom 3 - The Measurement Axiom - A Very Important Axiom

Can eye-ball/read-out a bit. Cannot eye-ball/stare at qubit.

Your **interaction** is via a **Measurement**.

Axiom 3 describes this interaction and the rules governing this interaction.

## Axiom 3 : The Measurement Axiom

A measurement is described through

a collection  $\{\Pi_{\alpha_1}, \Pi_{\alpha_2}, \dots, \Pi_{\alpha_K}\}$  of projectors acting on inner product Space  $\mathcal{H}$

that satisfy the Completeness Relation

$$\sum_{k=1}^K \Pi_{\alpha_k} = \Pi_{\alpha_1} + \dots + \Pi_{\alpha_K} = I \quad (I \equiv \text{the Identity on } \mathcal{H}).$$

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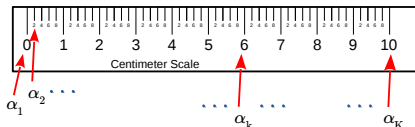
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Indices  $\alpha_1, \dots, \alpha_K$  : possible outcomes.

Each Projector  $\Pi_{\alpha_k}$  corresponds to its outcome  $\alpha_k$ .

Completeness Relation "You must get atleast one of the possible outcomes."



Simplify, Simplify, Simplify, ...

Just call  $\alpha_1, \alpha_2, \alpha_K$  as  $1, 2, \dots, K$

Outcomes are  $1, 2, \dots, K$ .

Reduce notation.

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$$P(\text{Outcome} = k) = (\text{Length of proj. } \Pi_k|\phi\rangle)^2 = \text{Inn. prod. between } \Pi_k|\phi\rangle \text{ and } \Pi_k|\phi\rangle$$

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Note :

$$\sum_{k=1}^K P(\text{Outcome} = k) = \sum_{k=1}^K \langle \phi | \Pi_k | \phi \rangle = \langle \phi | \sum_{k=1}^K \Pi_k | \phi \rangle = \langle \phi | I | \phi \rangle = 1 \quad \begin{array}{l} \text{Completeness} \\ + \text{unit-norm} \end{array}$$

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3. Moreover, if you observe outcome  $j$ , then the state collapses to

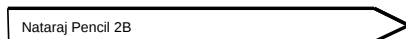
$$\frac{\Pi_j|\phi\rangle}{\sqrt{\langle \phi | \Pi_j | \phi \rangle}}$$



# Understanding the Measurement Axiom : Classical World

Classical world

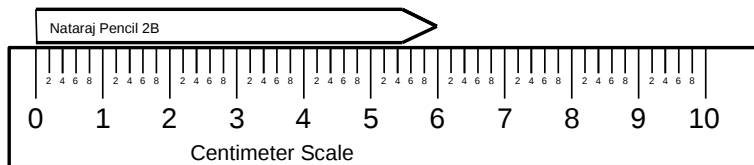
Wish to measure pencil's length



# Understanding the Measurement Axiom : Classical World

Classical world

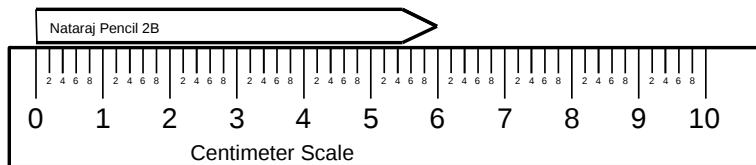
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# Understanding the Measurement Axiom : Classical World

## Classical world

1. Length is accurately read- 6cm.  
No uncertainty.



# Understanding the Measurement Axiom : Classical World

## Classical world

1. Length is accurately read- 6cm.  
No uncertainty.

2. Pencil's length does **NOT**  
**change post-measurement**



Nataraj Pencil 2B

# Understanding the Measurement Axiom : Quantum World

Quantum World

Wish to measure pencil's  
(quantum state) length

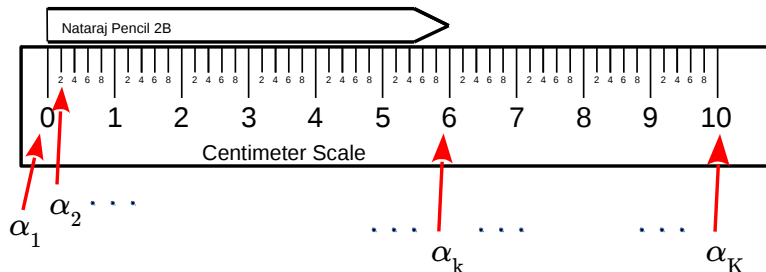


Nataraj Pencil 2B

# Understanding the Measurement Axiom : Quantum World

Quantum World

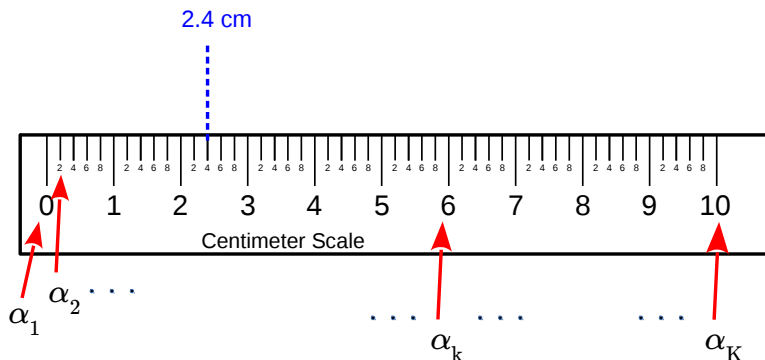
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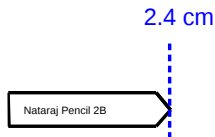
## Quantum World

1. Outcome is **RANDOM**.



# Understanding the Measurement Axiom : Quantum World

## Quantum World

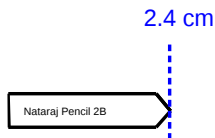


1. Outcome is **RANDOM**.
2. Pencil's length **CHANGES** post-measurement



# Understanding the Measurement Axiom : Quantum World

## Quantum World



1. Outcome is **RANDOM**.

2. Pencil's length **CHANGES**  
post-measurement

Welcome to the QUANTUM WORLD.

## Measurement Axiom : An Example

### Example

Quantum system in state  $|\phi\rangle = \frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}|1\rangle = \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{bmatrix} \in \mathcal{H} = \mathbb{C}^2$ .

Perform measurement with two outcome  $\{-0.5, +0.5\}$ .

Two meas. operators  $\Pi_{-0.5} = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$ ,  $\Pi_{+0.5} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$ .

$\Pi_{-0.5} + \Pi_{+0.5} = I$ . Completeness Relation satisfied.

$$\begin{aligned} P(\text{Outcome} = -0.5) &= (\text{Length of } \Pi_{-0.5}|\phi\rangle)^2 = \left\| \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{bmatrix} \right\|^2 \\ &= \left\| \begin{bmatrix} 0 \\ \frac{1}{\sqrt{2}} \end{bmatrix} \right\|^2 = \frac{1}{2} \\ P(\text{Outcome} = +0.5) &= \left\| \begin{bmatrix} \frac{1}{\sqrt{2}} \\ 0 \end{bmatrix} \right\|^2 = \frac{1}{2} \end{aligned}$$

If Outcome = -0.5, state collapses to  $|1\rangle$ . If Outcome = +0.5, state collapses to  $|0\rangle$ .

## Points to Keep in Mind

- ▶ Non-orthogonal states cannot be distinguished with certainty.
- ▶ Computation/Communication results need to be projected to orthogonal states.

## Axiom 4 : Description of a Joint/Composite Quantum System

### Quantum World

Suppose Quantum System 1 is in state  $|\phi_1\rangle \in \mathcal{H}_1$

Quantum System 2 is in state  $|\phi_2\rangle \in \mathcal{H}_2$

$\vdots$

Quantum System  $n$  is in state  $|\phi_n\rangle \in \mathcal{H}_n$

State space of composite Quant Sys. is the **tensor product**

$\mathcal{H}_1 \otimes \mathcal{H}_2 \otimes \cdots \otimes \mathcal{H}_n$  of constituent state spaces.

Composite System is described by State

$$|\phi_1\rangle \otimes |\phi_2\rangle \otimes \cdots \otimes |\phi_n\rangle \in \mathcal{H}_1 \otimes \mathcal{H}_2 \otimes \cdots \otimes \mathcal{H}_n.$$

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System 2 in state  $x_2 \in \mathcal{X}_2$

$\vdots$

System  $n$  in state  $x_n \in \mathcal{X}_n$

**Cartesian product**

$$\mathcal{X}_1 \times \mathcal{X}_2 \times \cdots \times \mathcal{X}_n.$$

$n$ -tuple

$$(x_1, \dots, x_n) \in \mathcal{X}_1 \times \cdots \times \mathcal{X}_n.$$

# What is a Tensor Product and what are the rules governing it?

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Suppose  $V$  is a  $m$ -dimensional IPS,

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All possible linear combinations of tensor product  $|v\rangle \otimes |w\rangle$  of elements  $|v\rangle \in V$  and  $|w\rangle \in W$ .

Just an (ordered) pair of vectors from respective spaces

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# Rules Governing Linear Combinations in Tensor Product Spaces

## Rules governing Linear combination

$$\underbrace{|v_1\rangle \otimes |w\rangle + |v_2\rangle \otimes |w\rangle}_{\uparrow} = \underbrace{(|v_1\rangle + |v_2\rangle)}_{\uparrow} \otimes \underbrace{|w\rangle} \quad \text{State Distributv Law (SDL) 1}$$

$$\underbrace{|v\rangle \otimes |w_1\rangle + |v\rangle \otimes |w_2\rangle}_{\uparrow} = \underbrace{|v\rangle}_{\uparrow} \otimes \underbrace{(|w_1\rangle + |w_2\rangle)}_{\uparrow} \quad \text{State Distributv Law (SDL) 2}$$

$$\alpha \cdot (|v\rangle \otimes |w\rangle) = (\alpha \cdot |v\rangle) \otimes |w\rangle = |v\rangle \otimes (\alpha \cdot |w\rangle) \quad \text{State Distributv Law (SDL) 3}$$

The above rules tell you how and when to combine terms.

In general, if the above rules do not apply, the sum

$$|v_1\rangle \otimes |w_1\rangle + |v_2\rangle \otimes |w_2\rangle = |v_1\rangle \otimes |w_1\rangle + |v_2\rangle \otimes |w_2\rangle$$

is a distinct element of  $V \otimes W$ .

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$$(A \otimes B)\left(\sum_i |v_i\rangle \otimes |w_i\rangle\right) = \sum_i A|v_i\rangle \otimes B|w_i\rangle \quad \text{Operator Dist. Law (ODL) 2}$$

$$A \otimes (B_1 + B_2) = A \otimes B_1 + A \otimes B_2 \quad \text{Operator Dist. Law (ODL) 3}$$

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What about the inner product on  $V \otimes W$

Ans : **Product of inner products.**

$$\text{IP between } |v_1\rangle \otimes |w_1\rangle \text{ and } |v_2\rangle \otimes |w_2\rangle = \langle v_1|v_2\rangle \langle w_1|w_2\rangle.$$

## Tensor Product : A Concrete Example

$$V = \mathbb{C}^2, W = \mathbb{C}^2, \quad |v\rangle = \begin{bmatrix} 1 \\ 2 \end{bmatrix}, \quad |w\rangle = \begin{bmatrix} 3 \\ 4 \end{bmatrix}, \quad |v\rangle \otimes |w\rangle = \begin{bmatrix} 1 \times 3 \\ 1 \times 4 \\ 2 \times 3 \\ 2 \times 4 \end{bmatrix}$$

$$A = \begin{bmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{bmatrix}, \quad B = \begin{bmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{bmatrix} \quad A \otimes B = \begin{bmatrix} a_{11}B & a_{12}B \\ a_{21}B & a_{22}B \end{bmatrix}$$

### Simple Consequences

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1.  $\dim(V \otimes W) = \dim(V) \times \dim(W)$ .
2. If  $\{|\alpha_1\rangle, \dots, |\alpha_m\rangle\}$  is **orthonormal** basis for  $V$ ,  
 $\{|\beta_1\rangle, \dots, |\beta_n\rangle\}$  is **orthonormal** basis for  $W$ ,

then  $\{|\alpha_i\rangle \otimes |\beta_j\rangle : 1 \leq i \leq m, 1 \leq j \leq n\}$  is **orthonormal** basis for  $V \otimes W$ .



## Our Basis in $\mathbb{C}^2 \otimes \mathbb{C}^2$

**Example**  $|0\rangle, |1\rangle$  forms an orthonormal basis for  $\mathcal{H}_A = \mathcal{H}_B = \mathbb{C}^2$

$|0\rangle \otimes |0\rangle, |0\rangle \otimes |1\rangle, |1\rangle \otimes |0\rangle, |1\rangle \otimes |1\rangle$  forms an orthonormal basis for  $\mathcal{H}_A \otimes \mathcal{H}_B$

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If  $|v_1\rangle, |v_2\rangle, |v_3\rangle, |v_4\rangle$  are orthonormal  $\Rightarrow \left\{ \begin{aligned} &\frac{1}{\sqrt{2}} |v_1\rangle + \frac{1}{\sqrt{2}} |v_2\rangle, \frac{1}{\sqrt{2}} |v_1\rangle - \frac{1}{\sqrt{2}} |v_2\rangle \\ &\frac{1}{\sqrt{2}} |v_3\rangle + \frac{1}{\sqrt{2}} |v_4\rangle, \frac{1}{\sqrt{2}} |v_3\rangle - \frac{1}{\sqrt{2}} |v_4\rangle \end{aligned} \right\}$

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$\frac{1}{\sqrt{2}} |00\rangle \pm \frac{1}{\sqrt{2}} |11\rangle, \quad \frac{1}{\sqrt{2}} |01\rangle \pm \frac{1}{\sqrt{2}} |10\rangle,$  are orthonormal

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# Separable and Entangled states

## Definition

Consider a joint quantum system consisting of 2 constituent quantum systems. The joint state vector  $|\phi\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B$  is **separable** if  $|\phi\rangle$  can be expressed as a tensor product of constituent state vectors  $|\phi_1\rangle \in \mathcal{H}_A$ ,  $|\phi_2\rangle \in \mathcal{H}_B$ , i.e.,

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individual **constituent components** have **no definite description**.

What is state of the first component  $|\Phi^+\rangle$  : **Invalid Qn..**

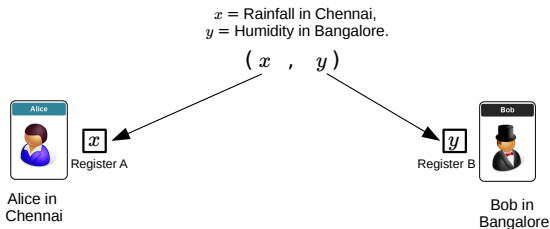
Only **state of a joint system**.

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Analogous to a pair of registers storing the values of two quantities.

### Joint System in our Classical World



In spite of (potentially) correlated/ or related, each element of the pair  $(x, y)$  has **its identity, description**.

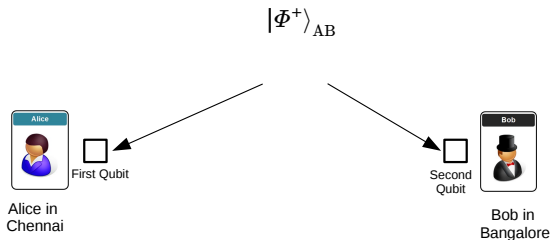


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**Joint System in our Quantum World**



Alice and Bob can share a pair of qubits describing the joint system.

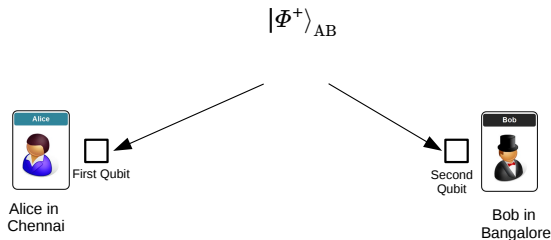
However, **each qubit has no definite description, identity.**

## Entanglement has NO Classical Analogue

The **entangled** state  $|\Phi^+\rangle = \left( \frac{1}{\sqrt{2}} |00\rangle + \frac{1}{\sqrt{2}} |11\rangle \right)$  represents the state of a **joint** system.

Analogous to a pair of registers storing the values of two quantities.

**Joint System in our Quantum World**



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The joint system is in a superposition of states  $|00\rangle$  and  $|11\rangle$ .

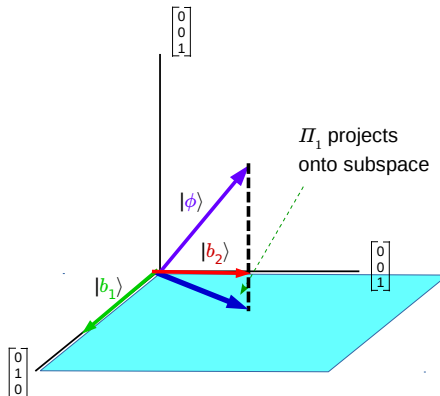
Entanglement

+

Randomness in measurement outcomes  
yield new information processing resources.

# Getting used to the Ket-Bra notation

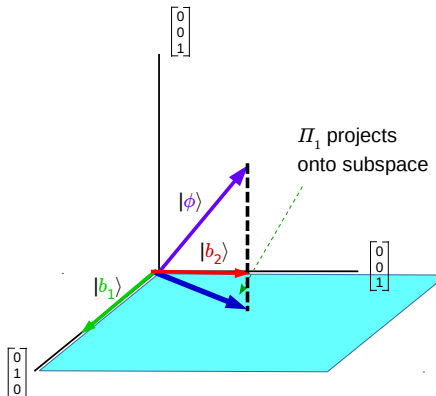
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## Getting used to the Ket-Bra notation

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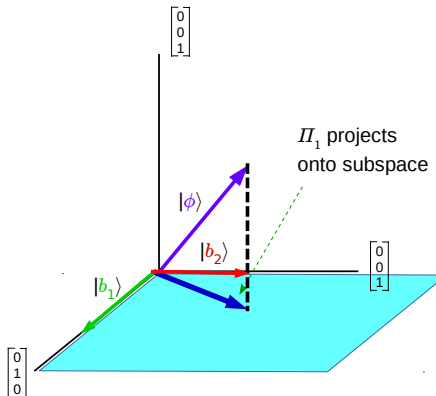
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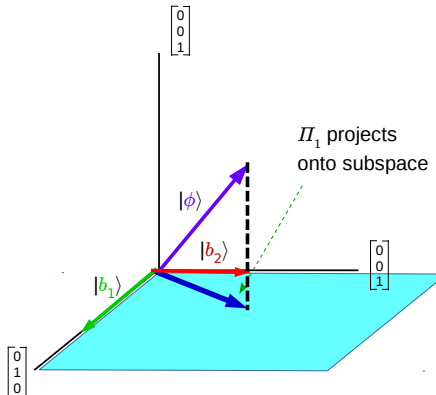


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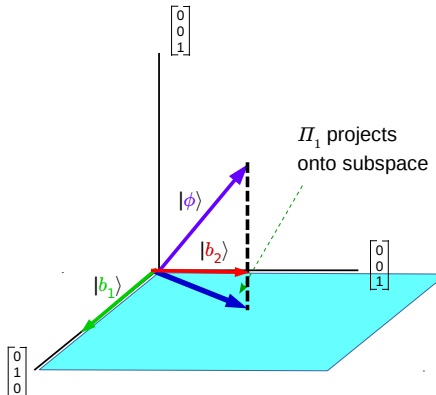
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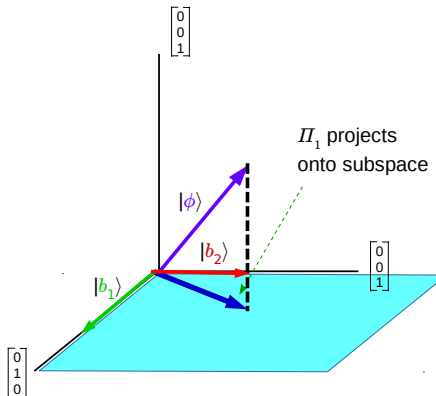
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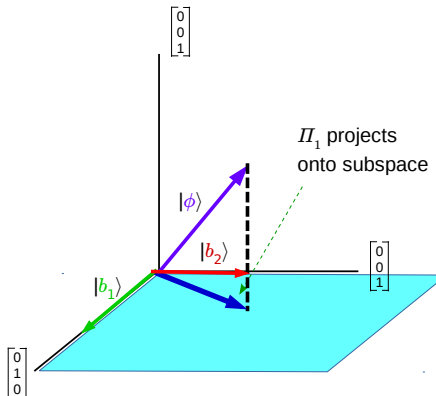
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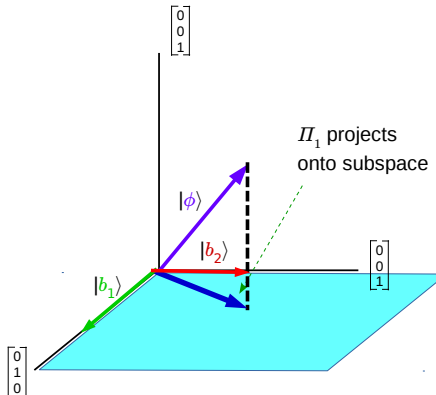
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$$\underbrace{\langle a | b \rangle}_{\text{bra-ket}} \quad \underbrace{\langle a |}_{\text{bra}} \quad \underbrace{|b \rangle}_{\text{ket}} \quad \underbrace{|b \rangle \langle b |}_{\text{ket-bra}}$$

## Ket-Bra Notation

The ket-bra notation is very useful in simplifying computation.

Suppose  $\Pi$  is a projection onto subspace  $\mathcal{W} \subseteq \mathcal{H}$ .

Suppose  $\{|v_1\rangle, \dots, |v_r\rangle\} \in \mathcal{W}$  is an orthonormal basis (ONB).

$$\Pi = |v_1\rangle\langle v_1| + |v_2\rangle\langle v_2| + \dots + |v_r\rangle\langle v_r| = \sum_{i=1}^r |v_i\rangle\langle v_i|$$

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$$\text{Bra-ket Notation } |a\rangle\langle b|c\rangle = \langle b|c\rangle |a\rangle$$

Suppose  $\mathcal{H}_A$  has ONB  $\{|v_1\rangle, \dots, |v_d\rangle\}$ , then any linear transformation  $T: \mathcal{H} \rightarrow \mathcal{H}$  can be expressed as

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## Ket-Bra Notation

$$T = |a\rangle \langle b|$$

$$T|c\rangle = |a\rangle \langle b|c\rangle = \langle b|c\rangle |a\rangle$$

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$|\cdot\rangle \langle \cdot|$  is an Operator       $\langle \cdot | |\cdot\rangle = \langle \cdot | \cdot \rangle$  is a scalar

$|\cdot\rangle$  is an vector       $\langle \cdot |$  is a linear functional



# Examples of Ket-Bra notation with Tensor Products

Recall Operative Distributive Law

$$(A \otimes B)(|v\rangle \otimes |w\rangle) = A|v\rangle \otimes B|w\rangle \quad \text{Operator Dist. Law (ODL) 1}$$

## Examples of Ket-Bra notation with Tensor Products

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$$\text{Recall ODL 3} \quad A \otimes (B_1 + B_2) = A \otimes B_1 + A \otimes B_2$$

$$\begin{aligned} |0\rangle\langle 0| \otimes (|0\rangle\langle 0| + |1\rangle\langle 1|) &= |0\rangle\langle 0| \otimes |0\rangle\langle 0| + |0\rangle\langle 0| \otimes |1\rangle\langle 1| \\ &= |00\rangle\langle 00| + |01\rangle\langle 01| \end{aligned}$$

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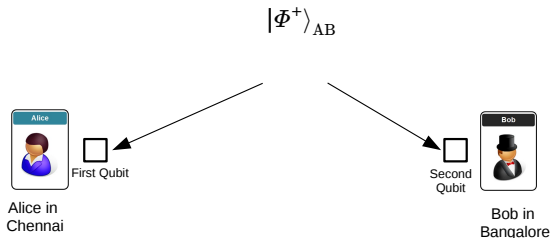
## Points to keep in mind

$|0\rangle\langle 0| \otimes I_B = |00\rangle\langle 00| + |01\rangle\langle 01| = \text{Projection on subspace spanned by } |00\rangle, |01\rangle.$

$|1\rangle\langle 1| \otimes I_B = |10\rangle\langle 10| + |11\rangle\langle 11| = \text{Projection on subspace spanned by } |10\rangle, |11\rangle.$

## Entangled pair can be separated, Acted upon Individually

Components of the joint system can be separated, Acted upon Individually



Suppose Alice performs measurement  $\{\Pi_1, \dots, \Pi_K\}$ . Bob remains silent.

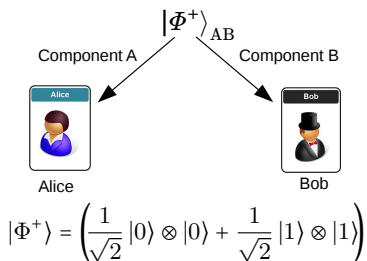
?? Effect on Joint system ??

Equivalent to measurement  $\{\Pi_1 \otimes I_B, \dots, \Pi_K \otimes I_B\}$  on joint system.

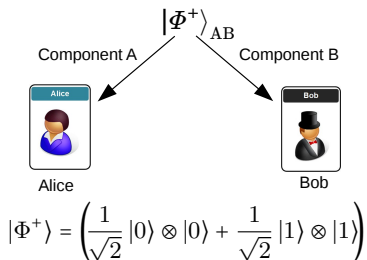
Only Alice sees outcome. Joint state collapses.



# Measurement on a Distributed Entangled State



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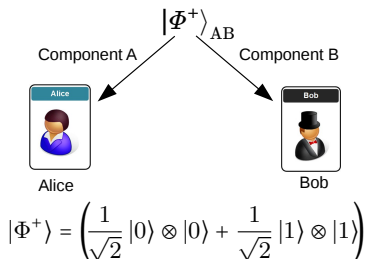


Alice performs measurement  
 $\{\Pi_0 = |0\rangle\langle 0|, \Pi_1 = |1\rangle\langle 1|\}$

Bob does nothing.

Measurement on joint system  
 $\{|0\rangle\langle 0| \otimes I_B, |1\rangle\langle 1| \otimes I_B\}$

# Measurement on a Distributed Entangled State



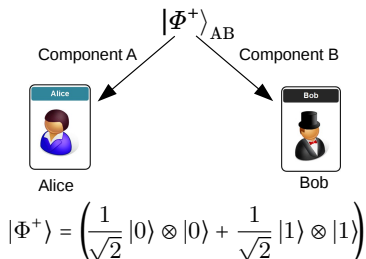
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equivalent to

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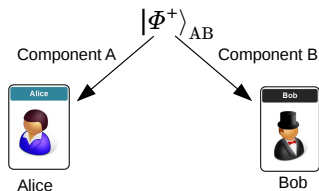


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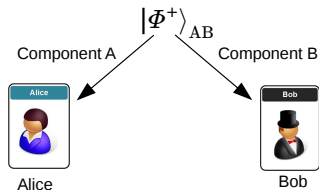
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Bob does nothing.

Measurement on joint system

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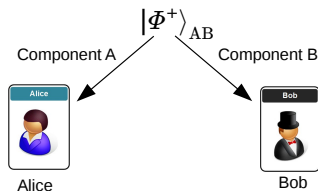
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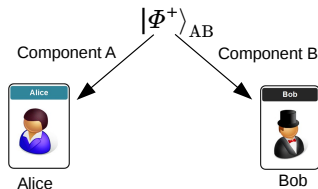
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Post Measurement on joint system

Outcome 0 with prob.  $\frac{1}{2}$ .

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States are **UNENTANGLED**

Outcome 0 and state  $|0\rangle$

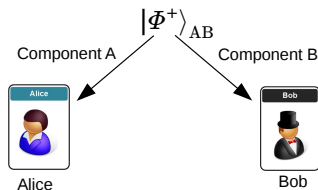
No Outcome. State  $|0\rangle$

Meas.  $\{|0\rangle\langle 0|, |1\rangle\langle 1|\}$

**Sure shot outcome 0.**



## Distributed Generation of common randomness



$$|\Phi^+\rangle = \left( \frac{1}{\sqrt{2}} |00\rangle + \frac{1}{\sqrt{2}} |11\rangle \right)$$

Alice performs measurement  
 $\{\Pi_0 = |0\rangle\langle 0|, \Pi_1 = |1\rangle\langle 1|\}$

Bob does nothing.

Post Measurement on joint system

Outcome 1 with prob.  $\frac{1}{2}$ .

State collapses to  $|11\rangle$

States are **UNENTANGLED**

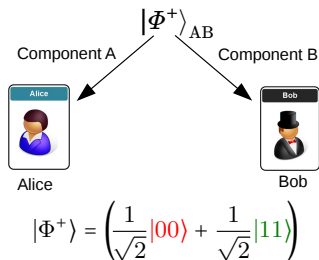
Outcome 1 and state  $|1\rangle$

No Outcome. State  $|1\rangle$

Meas.  $\{|0\rangle\langle 0|, |1\rangle\langle 1|\}$

**Sure shot outcome 1.**

## Distributed Generation of common randomness



Alice in Nice, Bob in Paris can generate common randomness.

Experimentally, components of entangled pair are separated by 1100 kms!!!!

## Idea Points to take Home

Entangled particles evolve simultaneously.

If you perturb one, the other gets perturbed.

If you wish to perturb the other, you can perturb your qubit!!!

# A Quantum system cannot be Cloned - The No-Cloning Theorem

The contents of a (classical) register can be copied onto another register.

However, the state of a quantum system cannot be duplicated or cloned.

Given an arbitrary state  $|\phi\rangle$ , there exists no unitary transformation that can duplicate this state.

## Theorem

*There exists no unitary transformation  $U : \mathcal{H} \otimes \mathcal{H} \rightarrow \mathcal{H} \otimes \mathcal{H}$  and a state  $|\omega\rangle \in \mathcal{H}$  such that*

$$U(|\phi\rangle \otimes |\omega\rangle) = |\phi\rangle \otimes |\phi\rangle$$

*holds for every  $|\phi\rangle \in \mathcal{H}$ .*

## 2. Quantum Gates

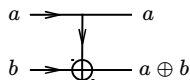
# Boolean Gates

A classical computation  $\equiv$  a map  $f : \{0, 1\}^n \rightarrow \{0, 1\}^n$ .

Computation is **reversible** if the input bits can be determined from the output bits, i.e.,  $f$  is invertible (1 : 1 and ONTO).

**Example** : NAND is **NOT** reversible.

**Example** : Controlled NOT (C-NOT)



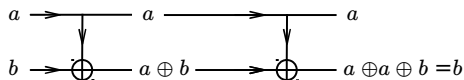
## Boolean Gates

A classical computation  $\equiv$  a map  $f : \{0, 1\}^n \rightarrow \{0, 1\}^n$ .

Computation is **reversible** if the input bits can be determined from the output bits, i.e.,  $f$  is invertible (1 : 1 and ONTO).

**Example** : NAND is **NOT** reversible.

**Example** : Controlled NOT (C-NOT)



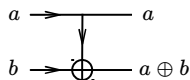
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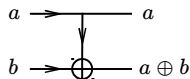
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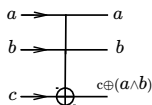
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**Example** : CC-NOT (C-NOT) is reversible.



# Quantum Gates and Operations are Unitary Transformations

Quantum circuits map superposition of  $n$  qubits into a superposition of  $n$  qubits.

$$\text{Quantum Gate} : |\phi\rangle \mapsto |\omega\rangle.$$

Valid Transformations : 1) Norm Preservation  $\langle\phi|\phi\rangle = \langle\omega|\omega\rangle$ . 2) Linearity.

Non-Linearity results in physical unrealizability.

Quantum Gate is a Unitary Transformation.

# Quantum Gates and Operations are Unitary Transformations

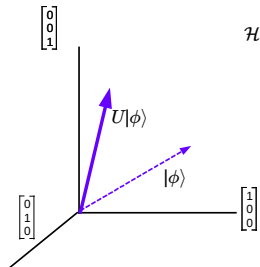
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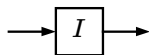
Quantum Operations : Unitary Transformations Mapping  $n$  qubits to  $n$  qubits.

Operation of a Quantum Gate : Completely specified by action on its bases.

Only need  $|0\rangle \mapsto ?$  and  $|1\rangle \mapsto ?$

## Single Qubit Gates - Pauli gates

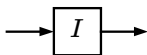
Identity Gate



$$I: \begin{array}{l} |0\rangle \mapsto |0\rangle \\ |1\rangle \mapsto |1\rangle \end{array}$$

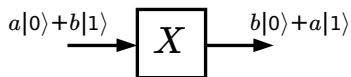
## Single Qubit Gates - Pauli gates

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### Pauli $X$ - Gate



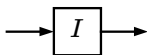
$$X: \begin{array}{l} |0\rangle \mapsto |1\rangle \\ |1\rangle \mapsto |0\rangle \end{array}$$

### Matrix Representation

$$X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

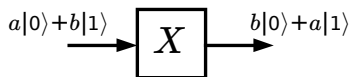
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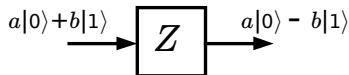


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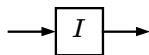


$$Z: \begin{array}{l} |0\rangle \mapsto |0\rangle \\ |1\rangle \mapsto -|1\rangle \end{array}$$

$$\text{Matrix Representation } Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

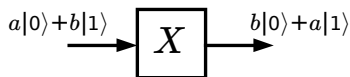
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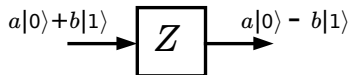


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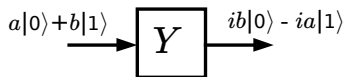
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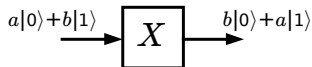
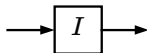
## Pauli Y-Gate



$$Y: \begin{array}{l} |0\rangle \mapsto i|1\rangle \\ |1\rangle \mapsto -i|0\rangle \end{array}$$

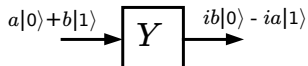
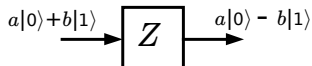
Matrix Representation  $Y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$

## Playing with the Pauli $I, X, Y, Z$ Gates



Your task is to **recover** the qubit  $a|0\rangle + b|1\rangle$ .

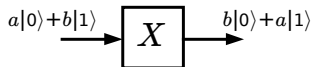
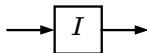
**Which operator** will you use if you are given



State you are given	Operator to use
$a 0\rangle - b 1\rangle$	
$b 0\rangle + a 1\rangle$	
$b 0\rangle - a 1\rangle$	

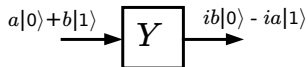
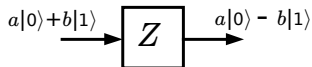


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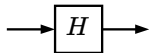
**Which operator** will you use if you are given



State you are given	Operator to use
$a 0\rangle - b 1\rangle$	$Z$
$b 0\rangle + a 1\rangle$	$X$
$b 0\rangle - a 1\rangle$	First $Z$ , then $X$

# The Hadamard Gate and Some Interesting Properties

## Hadamard Gate



$$H: \begin{array}{l} |0\rangle \mapsto \frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |1\rangle =: |+\rangle \\ |1\rangle \mapsto \frac{1}{\sqrt{2}} |0\rangle - \frac{1}{\sqrt{2}} |1\rangle =: |-\rangle \end{array}$$

## Matrix Representation

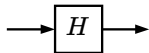
$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

# The Hadamard Gate and Some Interesting Properties

## Property 1

$$\begin{aligned}(H \otimes H)(|0\rangle \otimes |0\rangle) &= \frac{1}{2}(|0\rangle + |1\rangle) \otimes (|0\rangle + |1\rangle) \\ &= \frac{1}{2}(|00\rangle + |01\rangle + |10\rangle + |11\rangle)\end{aligned}$$

## Hadamard Gate



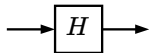
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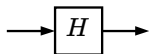
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$$H^{\otimes n} |0\rangle^{\otimes n} = \sum_{\substack{(b_1, \dots, b_n) \\ \in \{0,1\}^n}} \frac{1}{\sqrt{2^n}} |b_1 \dots b_n\rangle$$

All possible bit combinations are stored in  $n$ -qubits.

# The Hadamard Gate and Some Interesting Properties

## Hadamard Gate



$$H : \begin{array}{l} |0\rangle \mapsto \frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |1\rangle =: |+\rangle \\ |1\rangle \mapsto \frac{1}{\sqrt{2}} |0\rangle - \frac{1}{\sqrt{2}} |1\rangle =: |-\rangle \end{array}$$

## Matrix Representation

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

## Property 2

$$|0\rangle \xrightarrow{H} |0\rangle + |1\rangle, \quad |1\rangle \xrightarrow{H} |0\rangle - |1\rangle$$

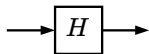
$$|0\rangle \xrightarrow{H} (-1)^{0 \cdot 0} |0\rangle + (-1)^{0 \cdot 1} |1\rangle,$$

$$|1\rangle \xrightarrow{H} (-1)^{1 \cdot 0} |0\rangle + (-1)^{1 \cdot 1} |1\rangle,$$

$$|b\rangle \xrightarrow{H} \sum_{z=0}^1 (-1)^{b \cdot z} |z\rangle,$$

# The Hadamard Gate and Some Interesting Properties

## Hadamard Gate



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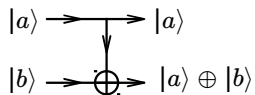
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$$|b\rangle \xrightarrow{H} \sum_{z=0}^1 (-1)^{b \cdot z} |z\rangle,$$

$$|b_1 \dots b_n\rangle \xrightarrow{H^{\otimes n}} \sum_{z \in \{0,1\}^n} (-1)^{b_1 z_1 + \dots + b_n z_n} |z_1 \dots z_n\rangle$$

# Our Two Qubit C-NOT Gate

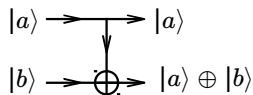
## C-NOT Gate



$$\overline{C}: \begin{array}{l} |00\rangle \mapsto |00\rangle \\ |01\rangle \mapsto |01\rangle \\ |10\rangle \mapsto |11\rangle \\ |11\rangle \mapsto |10\rangle \end{array}$$

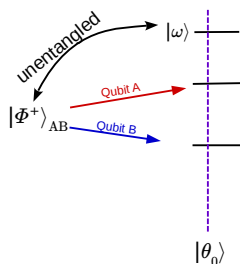
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An Application : Entangle two unentangled systems.

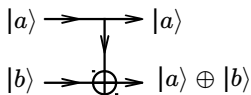


$$\text{Let } |\omega\rangle = a|0\rangle + b|1\rangle. \quad |\theta_0\rangle = |\omega\rangle \otimes |\Phi^+\rangle_{AB}$$



# Our Two Qubit C-NOT Gate

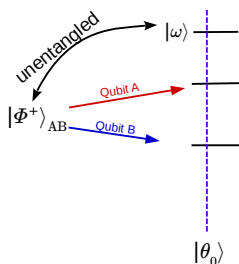
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$$\overline{C} \otimes I: \begin{array}{l} |000\rangle \mapsto |000\rangle \\ |001\rangle \mapsto |001\rangle \\ |010\rangle \mapsto |010\rangle \\ |011\rangle \mapsto |011\rangle \\ |100\rangle \mapsto |110\rangle \\ |111\rangle \mapsto |101\rangle \\ |100\rangle \mapsto |110\rangle \\ |111\rangle \mapsto |101\rangle \end{array}$$

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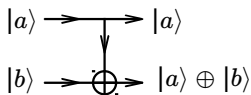


$$\text{Let } |\omega\rangle = a|0\rangle + b|1\rangle. \quad |\theta_0\rangle = |\omega\rangle \otimes |\Phi^+\rangle_{AB}$$

Use gate  $\overline{C} \otimes I$

# Our Two Qubit C-NOT Gate

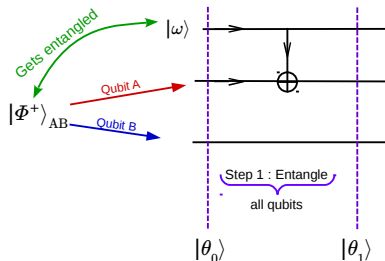
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$$\text{Let } |\omega\rangle = a|0\rangle + b|1\rangle. \quad |\theta_0\rangle = |\omega\rangle \otimes |\Phi^+\rangle_{AB}$$

Use gate  $\overline{C} \otimes I$

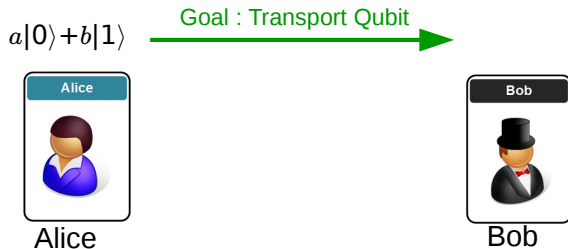
$$|\theta_1\rangle = (\overline{C} \otimes I) |\theta_0\rangle = (\overline{C} \otimes I)(|\omega\rangle \otimes |\Phi^+\rangle_{AB})$$

### 3. Quantum Protocols

# Quantum Teleportation

No Cloning Theorem : No Replication of qubits.

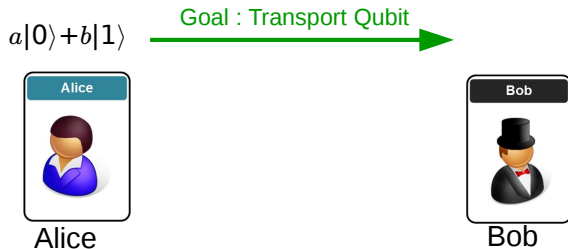
Can we transport them?



# Quantum Teleportation

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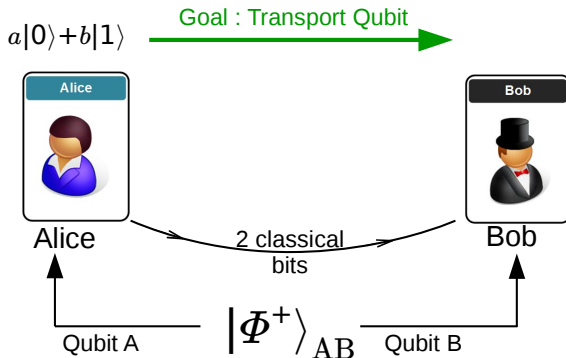
Can we transport them? If YES, what **resources** do we need?



# Quantum Teleportation

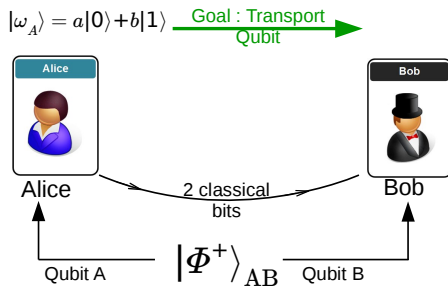
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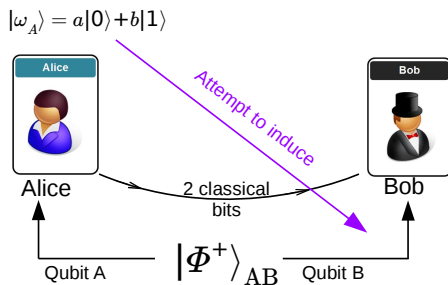


**Resource** : Shared entangled state  $|\Phi^+\rangle = \frac{1}{\sqrt{2}}|00\rangle + \frac{1}{\sqrt{2}}|11\rangle$  and 2 classical bits suffice.

# The Technique behind Teleportation



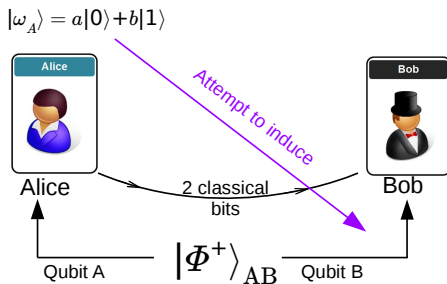
# The Technique behind Teleportation



High-Level Technique : 'Induce'  $|\omega_A\rangle$  on to Qubit B of  $|\Phi^+\rangle_{AB}$ .

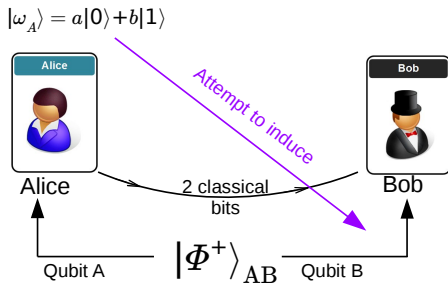


# The Technique behind Teleportation



High-Level Technique : 'Induce'  $|\omega_A\rangle$  on to Qubit B of  $|\Phi^+\rangle_{AB}$ . **HOW?**

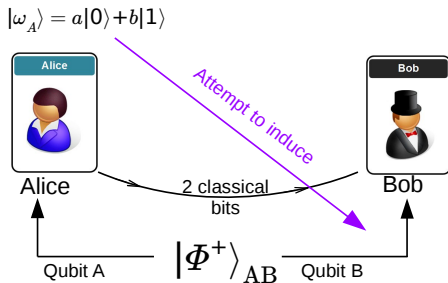
# The Technique behind Teleportation



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Entangled qubits evolve simultaneously.

# The Technique behind Teleportation

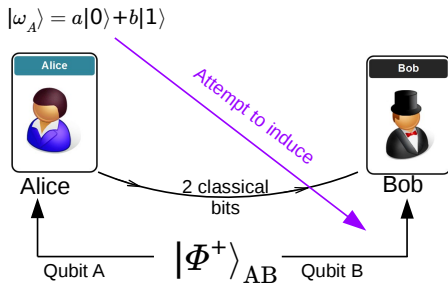


High-Level Technique : 'Induce'  $|\omega_A\rangle$  on to Qubit B of  $|\Phi^+\rangle_{AB}$ . **HOW?**

Entangled qubits evolve simultaneously.

Alice has  $|\omega_A\rangle$  AND first qubit of  $|\Phi^+\rangle_{AB}$ .

# The Technique behind Teleportation



High-Level Technique : 'Induce'  $|\omega_A\rangle$  on to Qubit B of  $|\Phi^+\rangle_{AB}$ . **HOW?**

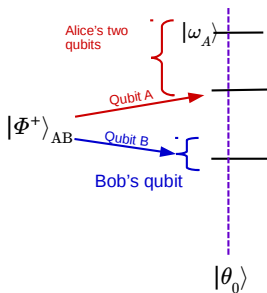
Entangled qubits evolve simultaneously.

Alice has  $|\omega_A\rangle$  AND first qubit of  $|\Phi^+\rangle_{AB}$ .

Step 1 : Entangle  $|\Phi^+\rangle_{AB}$  with  $|\omega_A\rangle$  by Alice entangling her two qubits.

Step 2 : Alice **cleverly evolves** her two qubits. Bob's entangled qubit evolves!!!

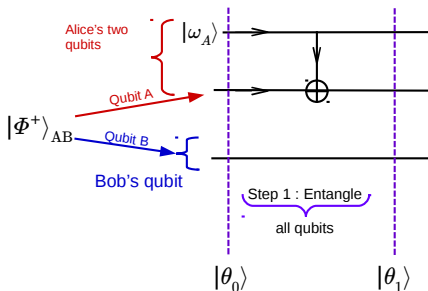
# Quantum Teleportation



$$|\theta_0\rangle = |\omega_A\rangle \otimes |\Phi^+\rangle_{AB} = (a|0\rangle + b|1\rangle) \otimes \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

$$|\theta_0\rangle = \frac{1}{\sqrt{2}}(a|000\rangle + a|011\rangle + b|100\rangle + b|111\rangle)$$

# Quantum Teleportation

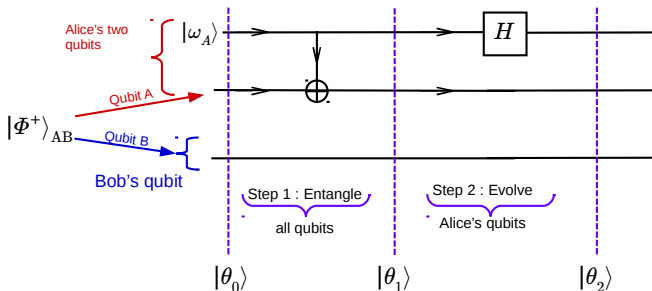


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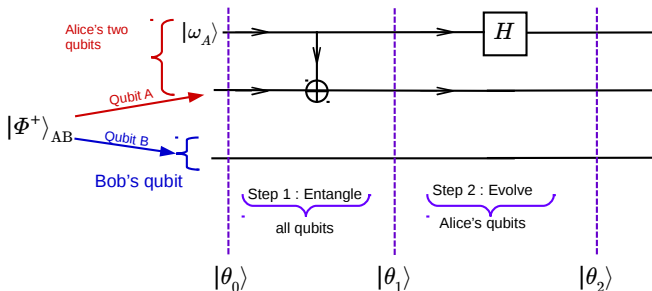
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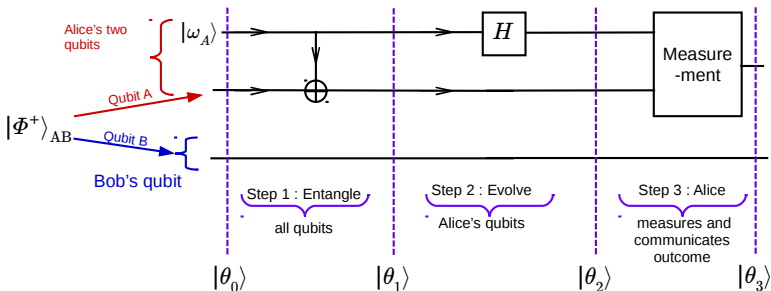
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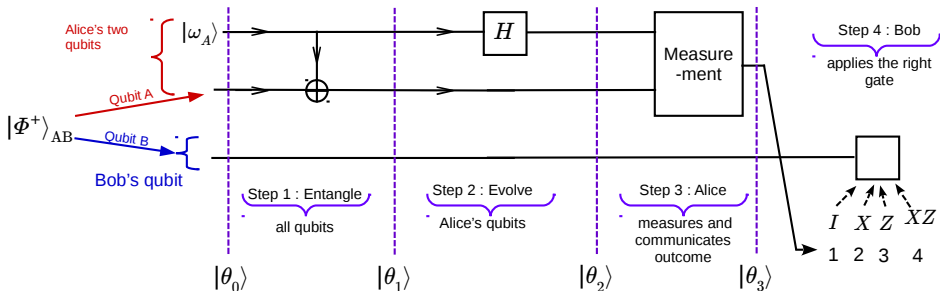
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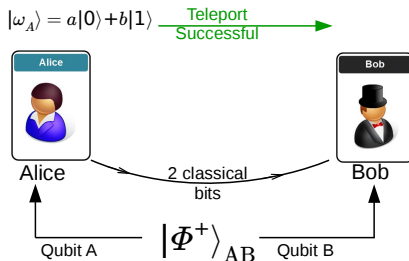
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# Super Dense Coding

How many classical bits of information can you pack in one qubit?

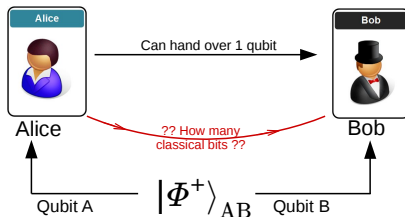
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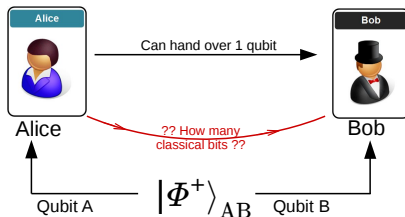


Shared entangled state  $|\Phi^+\rangle + \text{Hand over 1 qubit} = \text{?? number of classical bits ??}$

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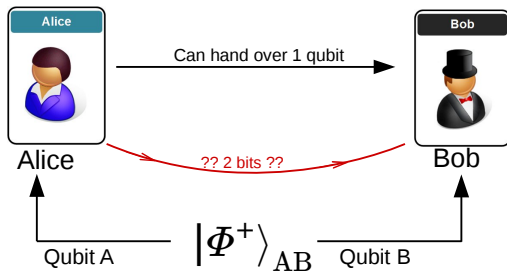
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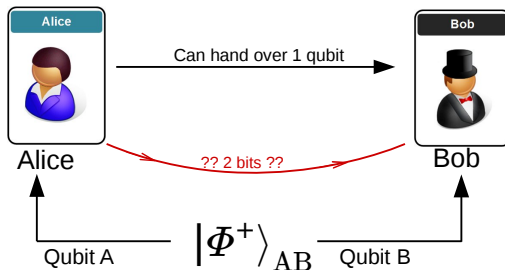
Shared entangled state  $|\Phi^+\rangle + \text{Hand over 1 qubit} = \quad !!! \text{ Answer is 2 } !!!$

Super Dense Coding

## Super Dense Coding : What is the idea? How do we do it?

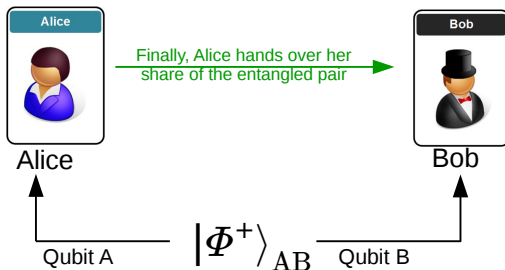


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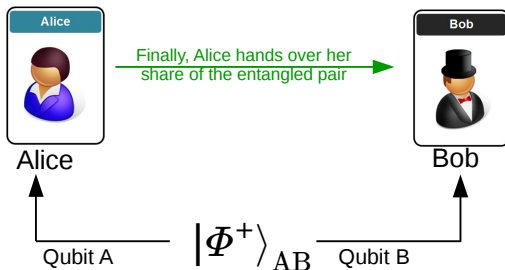
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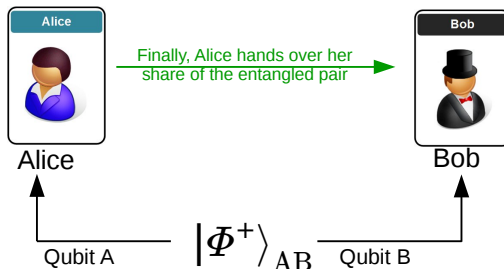
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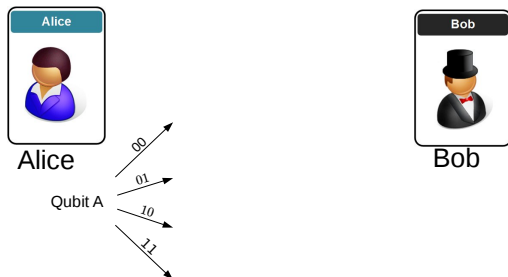
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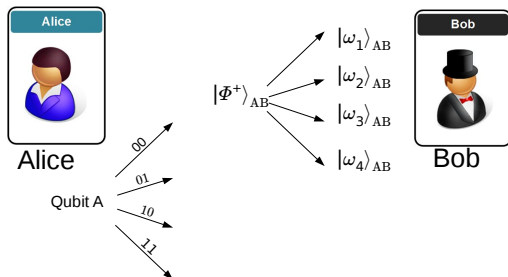


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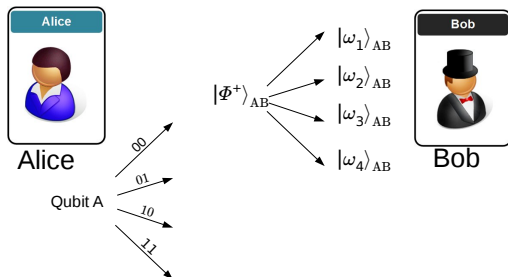
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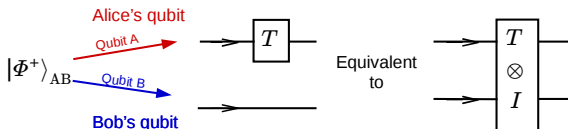
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The entangled pair evolves.

If  $|\omega_1\rangle, |\omega_2\rangle, |\omega_3\rangle, |\omega_4\rangle$  are perfectly distinguishable, Bob can recover the two bits.

Need  $|\omega_1\rangle, |\omega_2\rangle, |\omega_3\rangle, |\omega_4\rangle$  mutually orthonormal in  $\mathbb{R}^4$ .

## The Pauli Gates to our rescue



Alice applying gate  $T$  is equivalent to transformation  $T \otimes I$  on composite system.

Information bits	Gate	Resulting State
00	$I \otimes I$	$\frac{1}{\sqrt{2}}  00\rangle + \frac{1}{\sqrt{2}}  11\rangle$
01	$Z \otimes I$	$\frac{1}{\sqrt{2}}  00\rangle - \frac{1}{\sqrt{2}}  11\rangle$
10	$X \otimes I$	$\frac{1}{\sqrt{2}}  01\rangle + \frac{1}{\sqrt{2}}  10\rangle$
11	$iY \otimes I$	$\frac{1}{\sqrt{2}}  01\rangle - \frac{1}{\sqrt{2}}  10\rangle$

On receiving the Qubit  $A$  from Alice, Bob performs the measurement

$$\{\Pi_{00} = |00\rangle\langle 00|, \Pi_{01} = |01\rangle\langle 01|, \Pi_{10} = |10\rangle\langle 10|, \Pi_{11} = |11\rangle\langle 11|\}$$

## 4. Quantum Algorithms

# Comparing Classical and Quantum Computational Powers

- ▶ Side-Step a formal definition of a Quantum Turing Machine and Quantum complexity classes.
- ▶ Single-Qubit Unitary operator  $\equiv$  single-input Boolean gate.
- ▶ Proxy for run-time  $\sim$  No. of quantum gates and No. of unitary operations

**BPP** : Problem  $\Pi$  is in BPP if  $\exists$  a poly-time algo on a **probabilistic Classical** Turing Machine that returns correct answer with prob. atleast  $\frac{3}{4}$ .

**BQP** : Problem  $\Pi$  is in BPP if  $\exists$  a poly-time algo on a probabilistic **Quantum** Turing Machine that returns correct answer with probability atleast  $\frac{3}{4}$ .

**Informal Analysis. Techniques to exploit Superposition.**



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We have an oracle who, given  $x^n$ , will compute  $f(x^n)$ .

One usage : Binary oracle will provide us  $f(x^n)$ .

One usage : Quantum oracle will provide us  $|f(x^n)\rangle$ .

How many times should we poll our oracles?

# Known algorithms on Classical Computers

## Worst Case Analysis with guaranteed correctness

Must poll  $\geq 2^{n-1} + 1$  sequences in  $\{0, 1\}^n$

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## Performance of Probabilistic (Randomized) Algorithms

Algorithm : Pick  $k$  boolean inputs uniformly and randomly.

Poll  $f$ -values for chosen random inputs.

If all  $f$ -values for chosen random inputs are same, declare  $f$  is **constant**, (Category 1).

Otherwise, declare  $f$  is **balanced**, i.e., **Category 2**.

Performance : If you declare  $f$  is **balanced**,  $f$  is definitely **balanced**.

$\Rightarrow P(f \text{ is } \text{constant} \mid \text{you declare } \text{balanced}) = 0.$

$$P(f \text{ is } \text{balanced} \mid \text{you declare } \text{constant}) = \frac{2^{-k} P(f \text{ is } \text{balanced})}{1 - P(f \text{ is } \text{balanced})} \xrightarrow{k \rightarrow \infty} 0.$$

Problem in BPP.

$2^{n-1} + 1$  computations for certain answer.



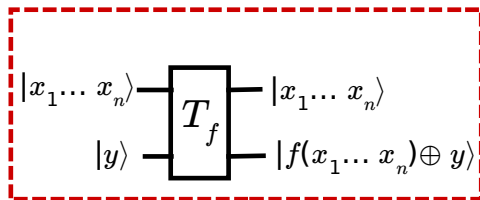
# Deutsch Jozsa discovered an efficient quantum algorithm

What is the idea?

Prepare a  $(n + 1)$ -qubit state  $|\phi\rangle$  based on the function  $f$  such that is

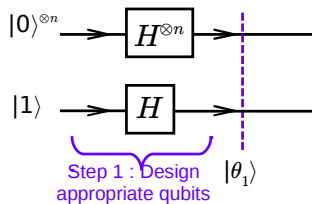
- (a) if  $f$  is **constant** state  $|\phi\rangle$  lies in subspace  $W$  and
- (b) if  $f$  is **balanced**, then  $|\phi\rangle$  lies in subspace  $W^\perp$ .
- (c) Preparation of  $|\phi\rangle$  has low **quantum complexity**.

## Our Quantum Oracle



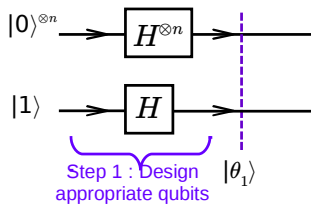
How many times will we need poll this quantum oracle?

# Deutsch Jozsa Algorithm



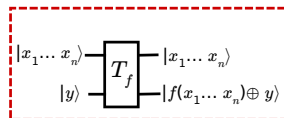
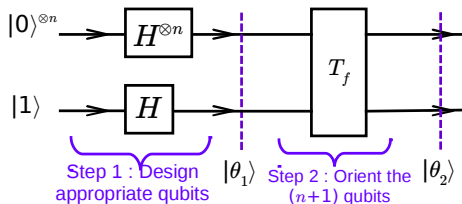
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$$\begin{aligned}
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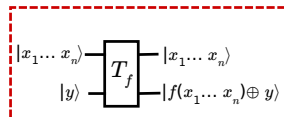
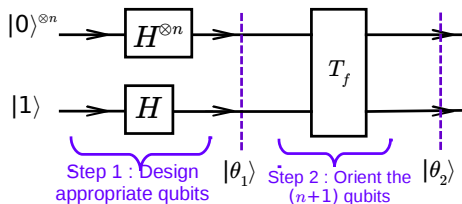
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Suppose  $f$  were a **constant**  $f(b^n) = 0$  for all  $b^n$

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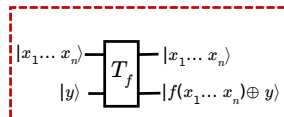
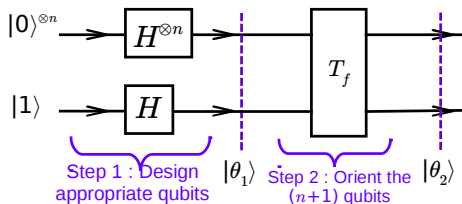
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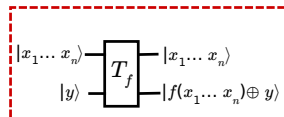
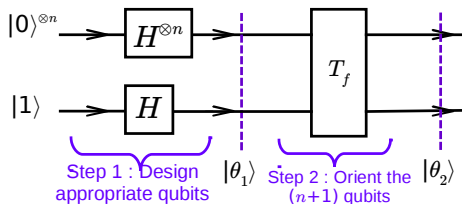
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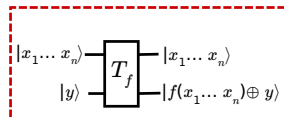
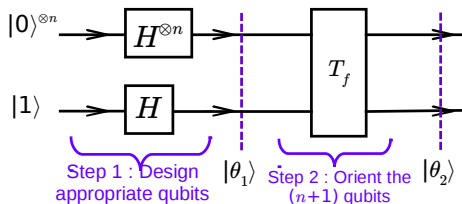
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Suppose  $f$  were a **constant**

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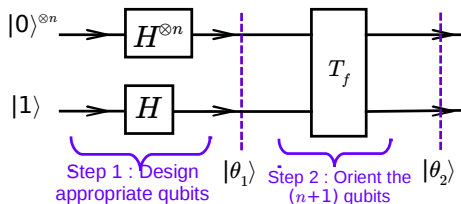
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Suppose  $f$  were a **balanced**

$$|\theta_2\rangle_{\text{blncd}} = \sum_{b^n: f(b^n)=0} |b_1 \dots b_n\rangle \otimes (|0\rangle - |1\rangle) - \sum_{b^n: f(b^n)=1} |b_1 \dots b_n\rangle \otimes (|0\rangle - |1\rangle)$$

# Deutsch Jozsa Algorithm



$|\theta_2\rangle_{\text{cnst}}$  and  $|\theta_2\rangle_{\text{blncd}}$   
are orthogonal!!!

$$\begin{aligned}
 |\theta_1\rangle &= H^{\otimes(n+1)} |0\rangle^{\otimes n} |1\rangle = H^{\otimes n}(|0\rangle^{\otimes n}) \otimes H(|1\rangle) \quad \text{Ignoring } \frac{1}{\sqrt{2}} \text{ factors} \\
 &= \sum_{b^n \in \{0,1\}^n} |b_1 \cdots b_n 0\rangle - \sum_{b^n \in \{0,1\}^n} |b_1 \cdots b_n 1\rangle \\
 |\theta_2\rangle &= \sum_{b^n \in \{0,1\}^n} |b_1 \cdots b_n f(b_1 \cdots b_n)\rangle - \sum_{b^n \in \{0,1\}^n} |b_1 \cdots b_n 1 \oplus f(b_1 \cdots b_n)\rangle
 \end{aligned}$$

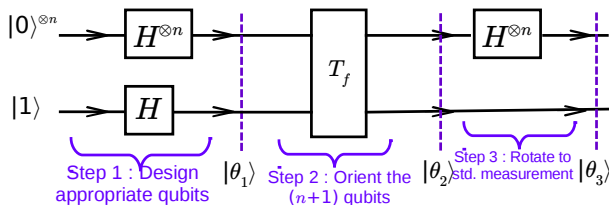
Suppose  $f$  were a **constant**

$$|\theta_2\rangle_{\text{cnst}} = \pm \sum_{b^n \in \{0,1\}^n} |b_1 \cdots b_n\rangle \otimes (|0\rangle - |1\rangle)$$

Suppose  $f$  were a **balanced**

$$|\theta_2\rangle_{\text{blncd}} = \sum_{b^n: f(b^n)=0} |b_1 \cdots b_n\rangle \otimes (|0\rangle - |1\rangle) - \sum_{b^n: f(b^n)=1} |b_1 \cdots b_n\rangle \otimes (|0\rangle - |1\rangle)$$

# Deutsch Jozsa Algorithm



$$|\theta_1\rangle = H^{\otimes(n+1)} |0\rangle^{\otimes n} |1\rangle = H^{\otimes n}(|0\rangle^{\otimes n}) \otimes H(|1\rangle) \quad \text{Ignoring } \frac{1}{\sqrt{2}} \text{ factors}$$

$$= \sum_{b^n \in \{0,1\}^n} |b_1 \dots b_n 0\rangle - \sum_{b^n \in \{0,1\}^n} |b_1 \dots b_n 1\rangle$$

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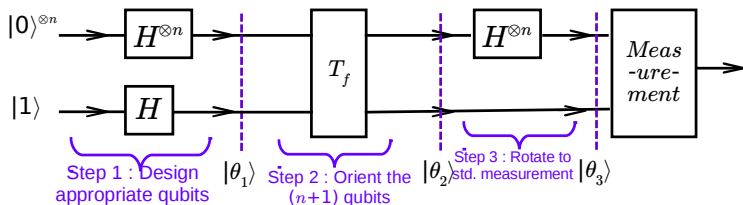
Suppose  $f$  were a **constant**

$$|\theta_2\rangle_{\text{cns}} = \pm \sum_{b^n \in \{0,1\}^n} |b_1 \dots b_n\rangle \otimes (|0\rangle - |1\rangle)$$

Suppose  $f$  were a **balanced**

$$|\theta_2\rangle_{\text{blncd}} = \sum_{b^n: f(b^n)=0} |b_1 \dots b_n\rangle \otimes (|0\rangle - |1\rangle) - \sum_{b^n: f(b^n)=1} |b_1 \dots b_n\rangle \otimes (|0\rangle - |1\rangle)$$

# Deutsch Jozsa Algorithm



$$|\theta_1\rangle = H^{\otimes(n+1)} |0\rangle^{\otimes n} |1\rangle = H^{\otimes n} (|0\rangle^{\otimes n}) \otimes H(|1\rangle) \quad \text{Ignoring } \frac{1}{\sqrt{2}} \text{ factors}$$

$$= \sum_{b^n \in \{0,1\}^n} |b_1 \dots b_n 0\rangle - \sum_{b^n \in \{0,1\}^n} |b_1 \dots b_n 1\rangle$$

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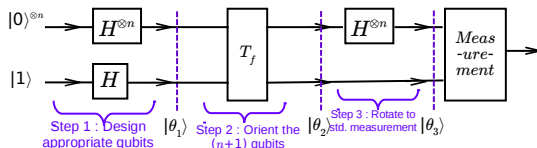
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# Analyzing Quantum Complexity

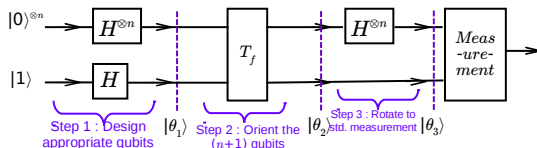


## Quantum Algorithm

Computes **Correct** Answer with **CERTAINTY**.

No. of Unitary Operations =  $O(n)!!!$

# Analyzing Quantum Complexity



Quantum Algorithm

Classical Computer

Computes **Correct** Answer with **CERTAINTY**.

$2^{n-1} + 1$  computations for certain answer.

No. of Unitary Operations =  $O(n)$ !!!

Problem in **BPP**.

Problem in  $\text{BPP} \cap \text{BQP}$ . No insights on  $\text{BQP} \setminus \text{BPP}$ .

## Finding the Unknown Period in $(\mathbb{Z}_2)^n$

$f : \{0, 1\}^n \rightarrow \{0, 1\}^n$  is **2-to-1** and **periodic** with **unknown period**  $(a_1, \dots, a_n)$ .

Exactly two  $n$ -bit sequences yield same output and  $f(x_1, \dots, x_n) = f(x_1 \oplus a_1, \dots, x_n \oplus a_n)$ .

On how many  $n$ -bit inputs must you poll  $f(\cdot)$ -values to figure out period  $(a_1, \dots, a_n)$ ?

Classically, if you poll for  $2^{\alpha n}$   $n$ -bit sequences, you have  $f(\cdot)$ -values for at most  $\binom{2^{\alpha n}}{2} \leq 2^{2\alpha n}$  input pairs.

$$P(\text{Finding } a^n) = \frac{2^{2\alpha n}}{2^n} = 2^{-n(1-2\alpha)} \xrightarrow{n \rightarrow \infty} 0 \quad \text{if } \alpha < \frac{1}{2}$$

Need to poll  $f(\cdot)$ -values for  $2^{\frac{n}{2}}$  inputs to obtain reasonable success.

## Recall Property 2 of the Hadamard Gate

For  $x \in \{0, 1\}$  or  $x^n \in \{0, 1\}^n$

$$|x\rangle \xrightarrow{H} \sum_{z=0}^1 (-1)^{x \cdot z} |z\rangle$$

$$|x_1 \cdots x_n\rangle \xrightarrow{H^{\otimes n}} \sum_{z^n \in \{0,1\}^n} (-1)^{x_1 \cdot z_1 + \cdots + x_n \cdot z_n} |z_1 \cdots z_n\rangle = \sum_{z^n \in \{0,1\}^n} (-1)^{x \cdot z} |z_1 \cdots z_n\rangle$$



# Exploring the Versatility of the The Hadamard Gate

**Problem :** Prepared State :

$$|\phi\rangle = \frac{1}{\sqrt{2}}(|x_1 \cdots x_n \ b_1 \cdots b_n\rangle + |(x_1 \oplus a_1) \cdots (x_n \oplus a_n) \ b_1 \cdots b_n\rangle)$$

$x_1, \dots, x_n, a_1, \dots, a_n$  unknown. Find  $a_1, \dots, a_n$ . !!! Cannot eye-ball A State!!!

# Exploring the Versatility of the The Hadamard Gate

**Problem :** Prepared State :  $x_1 \oplus a_1 = y_1, \dots, x_n \oplus a_n = y_n$

$$|\phi\rangle = \frac{1}{\sqrt{2}}(|x_1 \dots x_n \ b_1 \dots b_n\rangle + |y_1 \dots y_n \ b_1 \dots b_n\rangle)$$

$x_1, \dots, x_n, a_1, \dots, a_n$  unknown. Find  $a_1, \dots, a_n$ . !!! Cannot eye-ball A State!!!

## Exploring the Versatility of the The Hadamard Gate

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$x_1, \dots, x_n, a_1, \dots, a_n$  unknown. Find  $a_1, \dots, a_n$ .

Step 1: Apply  $H^{\otimes n} \otimes I_2^{\otimes n}$

$$\begin{aligned} |\phi\rangle &\mapsto \sum_{z_1, \dots, z_n \in \{0,1\}^n} \left[ (-1)^{\underline{x} \cdot \underline{z}} + (-1)^{\underline{x} \cdot \underline{z} \oplus \underline{a} \cdot \underline{z}} \right] |z_1 \dots z_n b_1 \dots b_n\rangle \\ &= \sum_{\substack{z_1, \dots, z_n : \\ a_1 z_1 \oplus \dots \oplus a_n z_n = 0}} |z_1 \dots z_n b_1 \dots b_n\rangle \end{aligned}$$

## Exploring the Versatility of the The Hadamard Gate

**Problem :** Prepared State :  $x_1 \oplus a_1 = y_1, \dots, x_n \oplus a_n = y_n$

$$|\phi\rangle = \frac{1}{\sqrt{2}}(|x_1 \dots x_n \ b_1 \dots b_n\rangle + |y_1 \dots y_n \ b_1 \dots b_n\rangle)$$

$x_1, \dots, x_n, a_1, \dots, a_n$  unknown. Find  $a_1, \dots, a_n$ .

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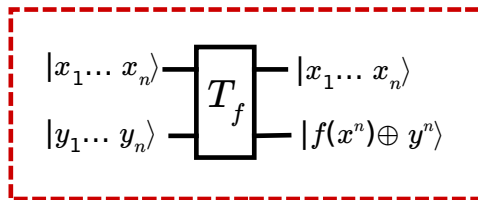
$$\begin{aligned} |\phi\rangle &\mapsto \sum_{z_1, \dots, z_n \in \{0,1\}^n} \left[ (-1)^{\underline{x} \cdot \underline{z}} + (-1)^{\underline{x} \cdot \underline{z} \oplus \underline{a} \cdot \underline{z}} \right] |z_1 \dots z_n \ b_1 \dots b_n\rangle \\ &= \sum_{\substack{z_1, \dots, z_n : \\ a_1 z_1 \oplus \dots \oplus a_n z_n = 0}} |z_1 \dots z_n \ b_1 \dots b_n\rangle \end{aligned}$$

For any  $(a_1, \dots, a_n)$  there are  $2^{n-1}$  terms in above sum.

Step 2: Apply Measurement :  $\{|0\dots 0\rangle \langle 0\dots 0| \otimes I, \dots, |1\dots 1\rangle \langle 1\dots 1| \otimes I\}$ .

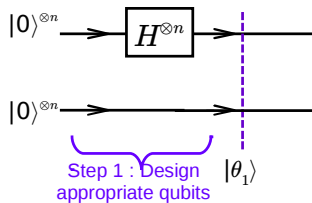
Outcome provides one choice of  $z_1, z_2, \dots, z_n$  for which  $a_1 z_1 \oplus \dots \oplus a_n z_n = 0$ .

## Quantum Oracle for our period finding function



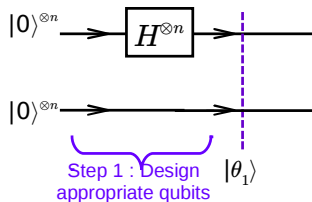
$$|x_1 \dots x_n \ y_1 \dots y_n\rangle \xrightarrow{T_f} |x_1 \dots x_n \ f(x^n) \oplus (y_1 \dots y_n)\rangle$$

## Simon's circuit for period finding



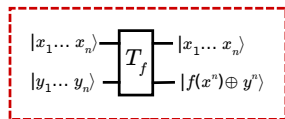
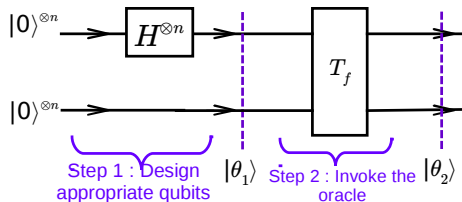
$$|\theta_1\rangle = H^{\otimes(n)} |0\rangle^{\otimes n} \otimes I |0\rangle^{\otimes n}$$

## Simon's circuit for period finding



$$\begin{aligned} |\theta_1\rangle &= H^{\otimes(n)} |0\rangle^{\otimes n} \otimes I |0\rangle^{\otimes n} \quad \text{Ignoring } \frac{1}{\sqrt{2}} \text{ factors} \\ &= \sum_{x^n \in \{0,1\}^n} |x_1 \cdots x_n \ 0 \cdots 0\rangle \end{aligned}$$

# Simon's circuit for period finding



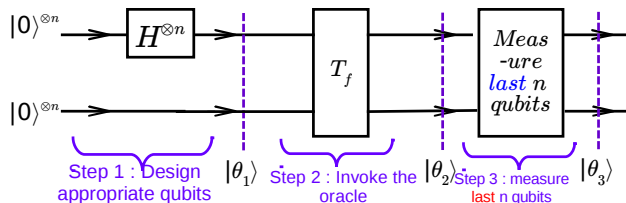
$$|\theta_1\rangle = H^{\otimes(n)} |0\rangle^{\otimes n} \otimes I |0\rangle^{\otimes n} \quad \text{Ignoring } \frac{1}{\sqrt{2}} \text{ factors}$$

$$= \sum_{x^n \in \{0,1\}^n} |x_1 \dots x_n \ 0 \dots 0\rangle$$

$$|\theta_2\rangle = \sum_{x^n \in \{0,1\}^n} |x_1 \dots x_n \ f(x_1 \dots x_n)\rangle$$



## Simon's circuit for period finding



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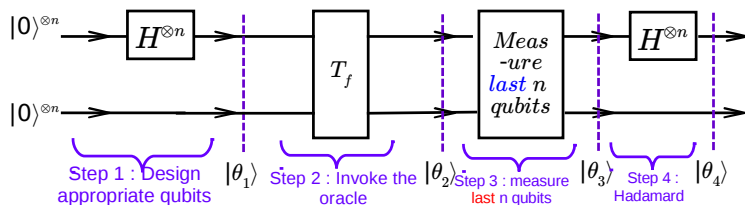
$$= \sum_{x^n \in \{0,1\}^n} |x_1 \cdots x_n \ 0 \cdots 0\rangle$$

$$|\theta_2\rangle = \sum_{x^n \in \{0,1\}^n} |x_1 \cdots x_n \ f(x_1 \cdots x_n)\rangle$$

Suppose outcome of the measurement were  $b_1, \dots, b_n$

$$|\theta_3\rangle = \frac{1}{\sqrt{2}} (|x_1 \cdots x_n \ b_1 \cdots b_n\rangle + |(x_1 \oplus a_1) \cdots (x_n \oplus a_n) \ b_1 \cdots b_n\rangle)$$

# Simon's circuit for period finding



$$|\theta_1\rangle = H^{\otimes(n)} |0\rangle^{\otimes n} \otimes I |0\rangle^{\otimes n} \quad \text{Ignoring } \frac{1}{\sqrt{2}} \text{ factors}$$

$$= \sum_{x^n \in \{0,1\}^n} |x_1 \cdots x_n \ 0 \cdots 0\rangle$$

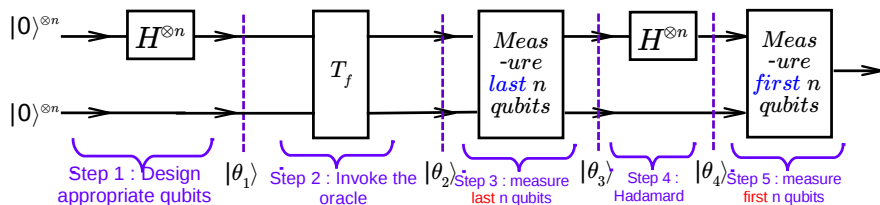
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$$|\theta_4\rangle = \sum_{\substack{z_1, \dots, z_n: \\ a_1 z_1 \oplus \cdots \oplus a_n z_n = 0}} |z_1 \cdots z_n \ b_1 \cdots b_n\rangle$$

# Simon's circuit for period finding



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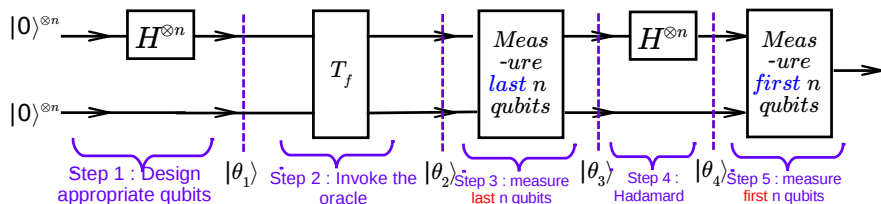
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## Simon's circuit for period finding



Suppose outcome of the measurement were  $b_1, \dots, b_n$

$$|\theta_4\rangle = \sum_{\substack{z_1, \dots, z_n: \\ a_1 z_1 \oplus \dots \oplus a_n z_n = 0}} |z_1 \dots z_n b_1 \dots b_n\rangle$$

Measure the first  $n$  registers with measurement operators

$$\{|0\dots 00\rangle\langle 0\dots 00| \otimes I_2^{\otimes n}, |0\dots 01\rangle\langle 0\dots 01| \otimes I_2^{\otimes n}, |1\dots 11\rangle\langle 1\dots 11| \otimes I_2^{\otimes n}\}$$

Every outcome gives you one linear equation  $o_1 a_1 \oplus \dots \oplus o_n a_n = 0$  where  $(o_1, \dots, o_n)$  is your outcome.

Need  $n$  linear independent eqns to solve for  $a_1, \dots, a_n$ . Repeat whole apparatus  $k$  times.

# Analysis of Simon's period finding algorithm

# Factoring a composite integer

Every composite integer is a product of **powers** of **primes**.

## Example

$$66 = 2 \cdot 3 \cdot 11$$

## Example

$$275 = 5 \cdot 5 \cdot 11$$

## Example

$$277 = ??$$

Given  $n$ -bit integer  $N$ , find **primes**  $p_1, \dots, p_m$  and integers  $q_1, \dots, q_m$  s.t

$$N = p_1^{q_1} \cdots p_m^{q_m}.$$

Given  $n$ -bit integer  $N$ , need a quantum algorithm that identifies prime factors in run-time  $n^k$  for some  $k$ .

# Towards Shor's Algorithm for Prime Factorization

Goal : Design a **polynomial-time** quantum algorithm that can identify the prime factors of a  $n$ -bit composite number  $N$ .

Break the task down.

Efficiently identify **non-trivial** factor of  $N$ .

Find  $\alpha$  such that  $\alpha|N$  and  $\alpha \neq 1$  and  $\alpha \neq N$ .

$N$  =

$\alpha_1 \cdot \alpha_2$

No. Factors    No. Computations

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$$\begin{aligned} N &= \alpha_1 \cdot \alpha_2 \\ &= \alpha_{11} \cdot \alpha_{12} \cdot \alpha_{21} \cdot \alpha_{22} \\ &= \alpha_{111} \alpha_{112} \alpha_{121} \alpha_{122} \cdots \alpha_{211} \alpha_{212} \alpha_{221} \alpha_{222} \end{aligned}$$

No. Factors    No. Computations



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No. Factors    No. Computations

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		No. Factors	No. Computations
$N$	$= \alpha_1 \cdot \alpha_2$	2	$n^k$
	$= \alpha_{11} \cdot \alpha_{12} \cdot \alpha_{21} \cdot \alpha_{22}$	4	$2(n-1)^k$
	$= \alpha_{111} \alpha_{112} \alpha_{121} \alpha_{122} \cdots \alpha_{211} \alpha_{212} \alpha_{221} \alpha_{222}$	8	$4(n-2)^k$
	$= \vdots$	$\dots$	$\vdots$
	$= p_1^{q_1} \cdot p_2^{q_2} \cdot \cdots p_{m-1}^{q_{m-1}} \cdot p_m^{q_m}$	$2^l$	$2^{l-1}(n-l)^k$

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	$= \vdots$	$\cdots$	$\vdots$
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$$l \text{ steps} \Rightarrow \text{No. Computations} \leq n^k + 2n^k + 4n^k + \cdots 2^{l-1}n^k \leq 2^l n^k$$

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$$l \text{ steps} \Rightarrow \text{No. Computations} \leq n^k + 2n^k + 4n^k + \cdots 2^{l-1}n^k \leq 2^l n^k$$

$$\text{Since } p_i \geq 2, \quad \text{No. of factors } 2^l \leq \log_2 N \Rightarrow \text{No. Computations} \leq n^k \log_2 N \leq n^{k+1}.$$

## The factorization Problem

Suffices to efficiently identify **non-trivial** factor of  $n$ -bit integer  $N$ .

Goal : Given  $N$ , find  $1 < x < N$  s.t,  $\text{GCD}(x, N) > 1$ .

# Some Number Theoretic Preliminaries

Goal : Given  $N$ , find  $1 < x < N$  s.t,  $\text{GCD}(x, N) > 1$ .

$\text{co-pr}(N) = \{a : 1 < a < N - 1, \text{ s.t } \text{GCD}(a, N) = 1, \text{ i.e., } a, N \text{ are co-prime}\}$

1.  $\text{co-pr}(N)$  is a finite group under  $\text{mod } N$  multiplication.

Need  $b$  s.t :  $ab = 1 \text{ mod } N$ . As you sweep  $b$ ,  $ab$ 's are distinct.

2. Being a finite group, each element of  $\text{co-pr}(N)$  has a **finite order**.

$\text{ord}(a) = \min\{k : a^k = 1 \text{ mod } N\} = \text{period of the fn. } f_{a,N}(k) = a^k \text{ mod } N$ .

3. Suppose  $r = \text{ord}(a)$  for  $a \in \{1, \dots, N - 1\}$ . Then

$$a^r = \theta N + 1 \Rightarrow N \mid (a^r - 1) \text{ and } N \nmid (a^{\frac{r}{2}} - 1)$$

**Case 1:**  $r$  is even.

$$N \mid (a^r - 1) = (a^{\frac{r}{2}} - 1)(a^{\frac{r}{2}} + 1)$$

If  $N \nmid (a^{\frac{r}{2}} - 1)$ , then we are done. Indeed,  $a^{\frac{r}{2}} - 1$  and  $a^{\frac{r}{2}} + 1$  have non-trivial common factors with  $N$ , i.e.,  $\text{GCD}(a^{\frac{r}{2}} - 1, N) > 1$  and  $\text{GCD}(a^{\frac{r}{2}} + 1, N) > 1$ .

## Chances of this happening are HIGH

### Theorem

Suppose  $N = p_1^{q_1} \cdots p_m^{q_m}$  is the prime factorization. Let  $X \in \text{co-pr}(N)$  be chosen uniformly at random, Let  $R = \text{ord}(X)$ . Then

$$P(R \text{ is even and } N \nmid (X^{\frac{R}{2}} - 1)) \geq 1 - \frac{1}{2^m}.$$

Suppose we can efficiently compute

$$\text{ord}(a) = \min\{k : a^k = 1 \pmod N\} = \text{period of the fn. } f_{a,N}(k) = a^k \pmod N.$$

Pick  $X_1, \dots, X_l$  uniformly at random, compute  $R_1 = \text{ord}(X_1), \dots, R_l = \text{ord}(X_l)$  and obtain a non-trivial factor of  $N$  with high probability.



# Algorithm for Identifying Prime Factors

Efficiently identify **non-trivial** factor of  $n$ -bit integer  $N$ .

## Algorithm

**Inputs:** Composite  $n$ -bit number  $N$

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Step 5 : If  $r$  is even and  $N \nmid (x^{\frac{r}{2}} + 1)$ , then compute  $\text{GCD}(x^{\frac{r}{2}} + 1, N)$ ,  $\text{GCD}(x^{\frac{r}{2}} - 1, N)$ . **Return if either is non-trivial factor**. If none is non-trivial factor **return FAILURE**

## Period Finding is $\mathbb{Z}_{2^n}$ is fundamental to Factorization

Simon's algorithm utilized the **Hadamard** transform to provide us period in  $(\mathbb{Z}_2)^n$ .

Suppose  $f : \{0, 1, \dots, 2^n - 1\} \rightarrow \{0, 1, \dots, 2^m - 1\}$  is a periodic function in  $\mathbb{Z}_{2^n}$ , i.e.,

$$f(x) = f(x + r) \text{ for some } 0 < r < 2^n - 1 \text{ and } \forall x \text{ valid.}$$

$\exists$  efficient algo. to compute  $r$  with high prob.

$\Rightarrow$

$\exists$  efficient algo. to FACTOR composite integer  $N$  with high. prob.

**Quantum Fourier Transform** in place of **Hadamard transform** yield period in  $\mathbb{Z}_{2^n}$ .









