# Exercise Sheet 2

## Exercise 1. Comparing clocks

In mathematical terms, explain why the clock of a freely falling observer runs faster than the clock of any other observer. How does this relate to the twin paradox?

#### **Exercise 2. Newtonian coordinate systems**

Let M be a Minkowski spacetime. Recall that a *Lorentz* or *inertial* coordinate system is a (local) coordinate system  $\xi = (t, x, y, z) \colon M \to \mathbb{R}^4_1$  such that  $\xi$  is an isometry.

Given a freely falling observer o with normalized proper time  $\tau$ , a Newtonian coordinate system is an inertial coordinate system  $\xi = (t, x, y, z)$  such that  $\xi(o(\tau)) = (\tau, 0, 0, 0)$ .

- (1) Given a freely falling observer o, how unique is a Newtonian coordinate system?
- (2) What about existence?
- (3) Does it matter that the observer is freely falling? Does it matter that *M* is a Minkowski spacetime and not any Lorentzian manifold?

#### **Exercise 3. Failure of simultaneity**

Let o be a freely falling observer in a Minkowski spacetime M. Two events p and q are considered *simultaneous* from the point of view of o if they have the same t-coordinate in any/some Newtonian coordinate system relative to o.

- (1) Show that two events p and q are simultaneous from the point of view of o if and only if  $\vec{pq}$  is orthogonal to o (i.e. p and q are in the same restspace relative to o).
- (2) Show that the notion of simultaneity is not independent of the observer. In fact, give a necessary and sufficient condition for two observers to agree on simultaneity.

## **Exercise 4. Length contraction**

Consider an unaccelerated train travelling at a fraction v of the speed of light relative to some observer o. If d is the length of the train from the point of view of a sitting passenger, what is the length of the spaceship from the point of view of o? Answer the questions below.

- (1) Let  $o_1$  and  $o_2$  be two freely falling particles in a Minkowski spacetime M, representing the endpoints of the train. Show that  $o_1$  and  $o_2$  are parallel in M if and only if they are parallel in the restspace of some/any freely falling observer.
- (2) From now on we assume that  $o_1$  and  $o_2$  are parallel. Assume that o sits on the train. Show that the distance L between  $o_1$  and  $o_2$  in the restspace of o is constant and equal to their distance in M.
- (3) Now let o be any freely falling observer. Assume that o,  $o_1$ , and  $o_2$  are coplanar in M. What does this assumption mean physically? From now on we thus assume that dim M = 2.
- (4) Let  $\theta$  denote the hyperbolic angle between o and  $o_1$ . Draw o,  $o_1$ , and  $o_2$  in an inertial coordinate system relative to o. Show that the distance  $L_o$  between  $o_1$  and  $o_2$  is constant, and by doing hyperbolic trigonometry in the right triangle, verifies  $L = L_o \cosh \theta$ . Conclude that  $L_o = L\sqrt{1-v^2}$  where v is the relative speed of the train with respect to o.

## Exercise 5. Velocity-addition formula for collinear motions

(1) (a) For  $\theta \in \mathbb{R}$ , consider the hyperbolic rotation matrix

$$R_{\theta} = \begin{pmatrix} \cosh \theta & \sinh \theta & 0 & 0 \\ \sinh \theta & \cosh \theta & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}.$$

Show that  $R_{\theta}$  is an isometry of Minkowski space  $\mathbb{R}^4_1$ , in fact show that it is a Lorentz boost: cf Ex. 2 in Exercise Sheet #1. What is  $R_{\theta_2} \circ R_{\theta_1}$ ?

- (b) Let V be a normalized timelike vector in Minkowski space  $\mathbb{R}^4_1$ . Assume V is contained in the tx-plane. Call  $\theta$  the hyperbolic angle  $\theta = \angle(e_0, V)$ . Show that  $R_\theta$  is the only orthochronous isometry  $\mathbb{R}^4_1$  fixing the yz-plane such that  $V = R_\theta(e_0)$ .
- (2) Let M be a Minkowski spacetime. Consider three observers o,  $o_1$ , and  $o_2$ , all freely falling, through some event  $p \in M$ .
  - (a) From now on we assume that o,  $o_1$ , and  $o_2$  are coplanar. Show that this amounts to saying that seen from o, the Newtonian motions of  $o_1$  and  $o_2$  are collinear.
  - (b) Explain why, by choosing an appropriate inertial coordinate system, one can assume that  $M = \mathbb{R}^4_1$ , p = 0,  $o(\tau) = (\tau, 0, 0, 0)$ , and  $o_1$  and  $o_2$  are contained in the tx-plane.
  - (c) Call V,  $V_1$ , and  $V_2$  the initial tangent vectors to o,  $o_1$ , and  $o_2$  respectively. Denote the following oriented hyperbolic angles:  $\theta_1 = \angle(V, V_1)$ ,  $\theta_2 = \angle(V_1, V_2)$ , and  $\theta = \angle(V, V_2)$ . Show that  $R_{\theta_1}(o) = o_1$ ,  $R_{\theta_2}(o_1) = o_2$ , and  $R_{\theta}(o) = o_2$ . Show that  $\theta = \theta_1 + \theta_2$ .
  - (d) Call  $v_1$  (resp.  $v_2$ ) the signed relative speed of  $o_1$  (resp.  $o_2$ ) from the point of view of o (resp.  $o_1$ ). Show that the signed relative speed of  $o_2$  from the point of view of o is:

$$v = \frac{v_1 + v_2}{1 + v_1 v_2} =: v_1 \oplus v_2 .$$

(e) Is  $\oplus$  a commutative operation on  $\mathbb{R}$ ? Is it associative? Is it  $\mathbb{R}$ -linear?

#### Exercise 6. An accelerated observer

Let  $M = \mathbb{R}^2_1$  and let  $o(\tau) = (\tau, 0)$  be a freely falling observer. Consider the curve  $\alpha(\tau) = (\sinh \tau, \cosh \tau)$ .

- (1) Show that  $\alpha$  is a material particle. Is it freely falling? Show that the norm of its acceleration is constant.
- (2) Show that a light beam emitted from  $\alpha$  towards o reaches o in finite proper time, but never reaches back  $\alpha$  after it is reflected. Hence o never appears on  $\alpha$ 's radar.
- (3) Which events can be picked up on  $\alpha$ 's radar?
- (4) Consider a photon being emitted by  $\alpha$  at  $\tau = 0$ , moving in the same direction as  $\alpha$ . Show that the speed of the photo relative to  $\alpha$  is  $e^{\tau}$ . Does that contradict the speed of light being constant?

#### **Exercise 7. Energy-momentum**

Let  $\alpha$  be a material particle of mass m>0 in Minkowski spacetime. We recall that the energy-momentum of  $\alpha$  is the vector  $P=m\frac{\mathrm{d}\alpha}{\mathrm{d}\tau}$ .

Let o be a freely falling observer and let  $\xi = (t, x, y, z)$  be an inertial coordinate system relative to o. One calls *energy of*  $\alpha$  *relative to* o  $E_o$  the time component of P in the coordinate system  $\xi$  and *momentum of*  $\alpha$  *relative to* o the space component  $\vec{P}_o$  of P in the coordinate system  $\xi$ .

- (1) Find the expressions of  $E_o$  and  $\vec{P}_o$  in terms of the relative speed v. Show that for small values of v,  $E_o$  and  $\vec{P}_o$  almost coincide with the Newtonian kinetic energy and momentum.
- (2) The *scalar momentum* of  $\alpha$  relative to o is  $p_o := \|\vec{P}_o\|^2$ . Show that E, m, and p are related by  $E^2 = m^2 + p^2$  and E = pv.
- (3) Derive from the previous question that if a lightlike particle has a well-defined energy-momentum, then it should have zero mass. By considering a lightlike particle with energy *E* as the limit of a sequence of freely falling particles with constant energy *E* and increasing speeds, propose a definition for the energy-momentum of a lightlike particle.