

Modeling and Constraining the Cluster Mass Function to Test Gravity at Large Scales

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Work presented here in collaboration with:

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Constraining alternative theories of gravity

1. Strength from cluster observations: constraining gravity at large scales, either using concrete models or general parameterizations.
2. Worked example: viable $f(R)$ gravity models with the chameleon screening mechanism.
3. For $f(R)$, current prospects for improvement are primarily based on analyses including less massive and low- z objects.

Hu-Sawicki f(R) gravity model (departing from GR)

$$f(R) = -2\Lambda - \frac{f_{R0}}{n} \frac{\bar{R}_0^{n+1}}{R^n}$$

Approximated Hu-Sawicki model in the high curvature regime. In the limit $|f_{R0}| < 10^{-2}$ closely mimics the Λ CDM expansion history

$$\lambda_{C0} \approx 29.9 \sqrt{\frac{|f_{R0}|}{10^{-4}} \frac{n+1}{4-3\Omega_m}} h^{-1} \text{Mpc}$$

Compton wavelength; scales below this present modified gravity until GR is recovered when the fifth force is screened by non-linear effects

$$g(a, k) \equiv -\frac{1}{3} \frac{k^2}{k^2 + m_{f_R}^2 a^2}$$

$$m_{f_R}^{-2} = \lambda_C^2 = 3f_{RR} \quad f_R \equiv \frac{\partial f}{\partial R}$$

Linear growth is different than GR+ Λ CDM

Halo mass function modeling

$$n_{\Delta_v} \equiv \frac{dn}{d \ln M_v} = \frac{\bar{\rho}_m}{M_v} \frac{d \ln \nu}{d \ln M_v} \nu f(\nu)$$

Sheth-Tormen mass function

$$\nu f(\nu) = A \sqrt{\frac{2}{\pi}} a \nu^2 [1 + (a \nu^2)^{-p}] \exp [-a \nu^2 / 2]$$

$$\nu = \delta_c / \sigma(M_v)$$

Peak height

$$\delta_c(\Omega_m, z) = \mathcal{A} \left(1 - \mathcal{B} \log_{10} \left[1 + \frac{\Omega_m^{-1} - 1}{(1+z)^3} \right] \right)$$

Density threshold fitting formula

$$n_{\Delta} = \left(\frac{n_{\Delta}^{f(R)}}{n_{\Delta}^{\text{GR}}} \right) \Big|_{\text{ST}} n_{\Delta}|_{\text{Tinker}}$$

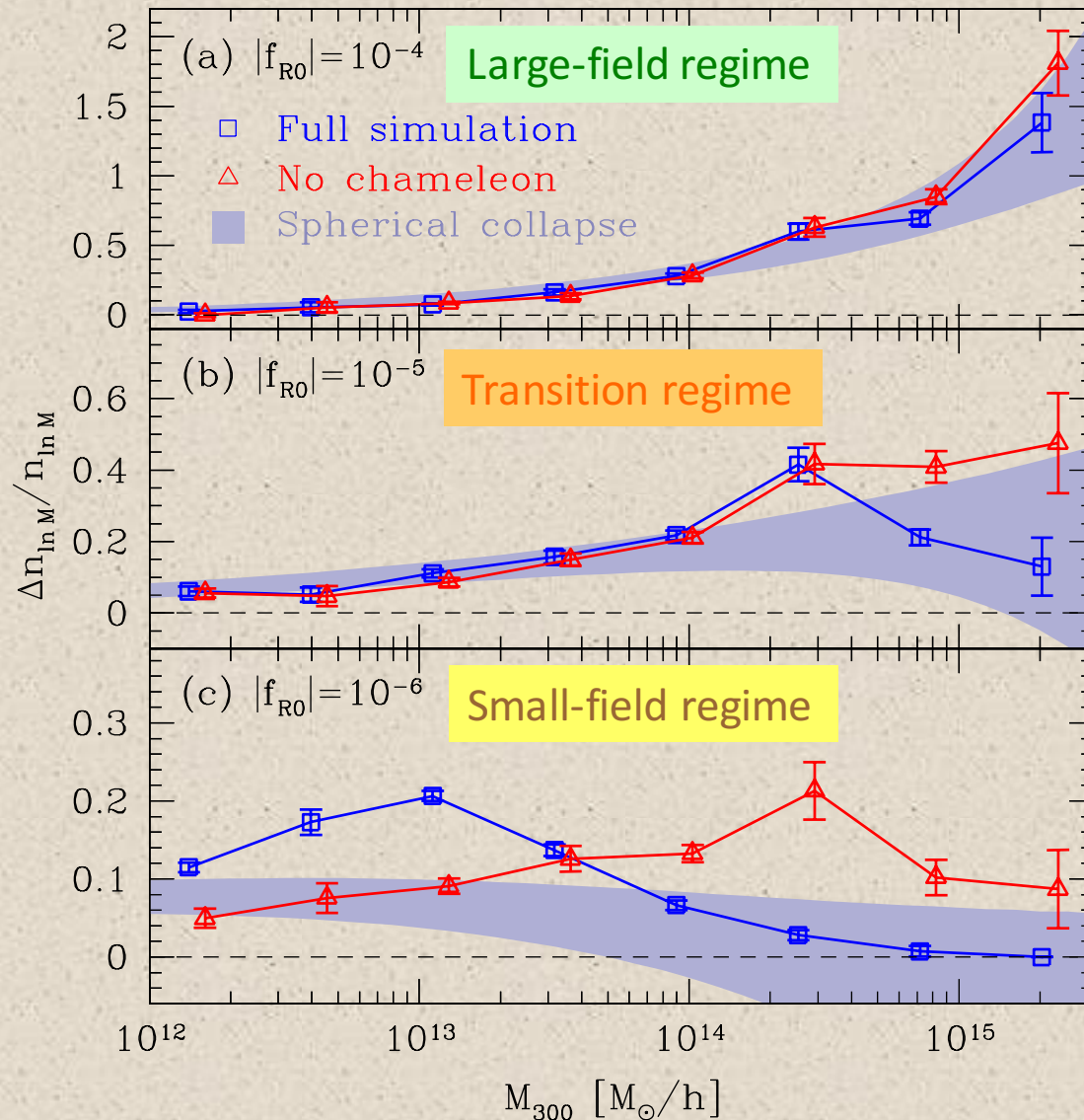
$$\mathcal{A} = 1.6865 \text{ and } \mathcal{B} = 0.0123 \text{ for GR}$$

$$\mathcal{A} = 1.7063 \text{ and } \mathcal{B} = 0.0136 \text{ for } f(R)$$

$$n_{\text{ST}}^{f(R)} / n_{\text{ST}}^{\text{GR}}$$

We set this equal to 1 when becomes <1

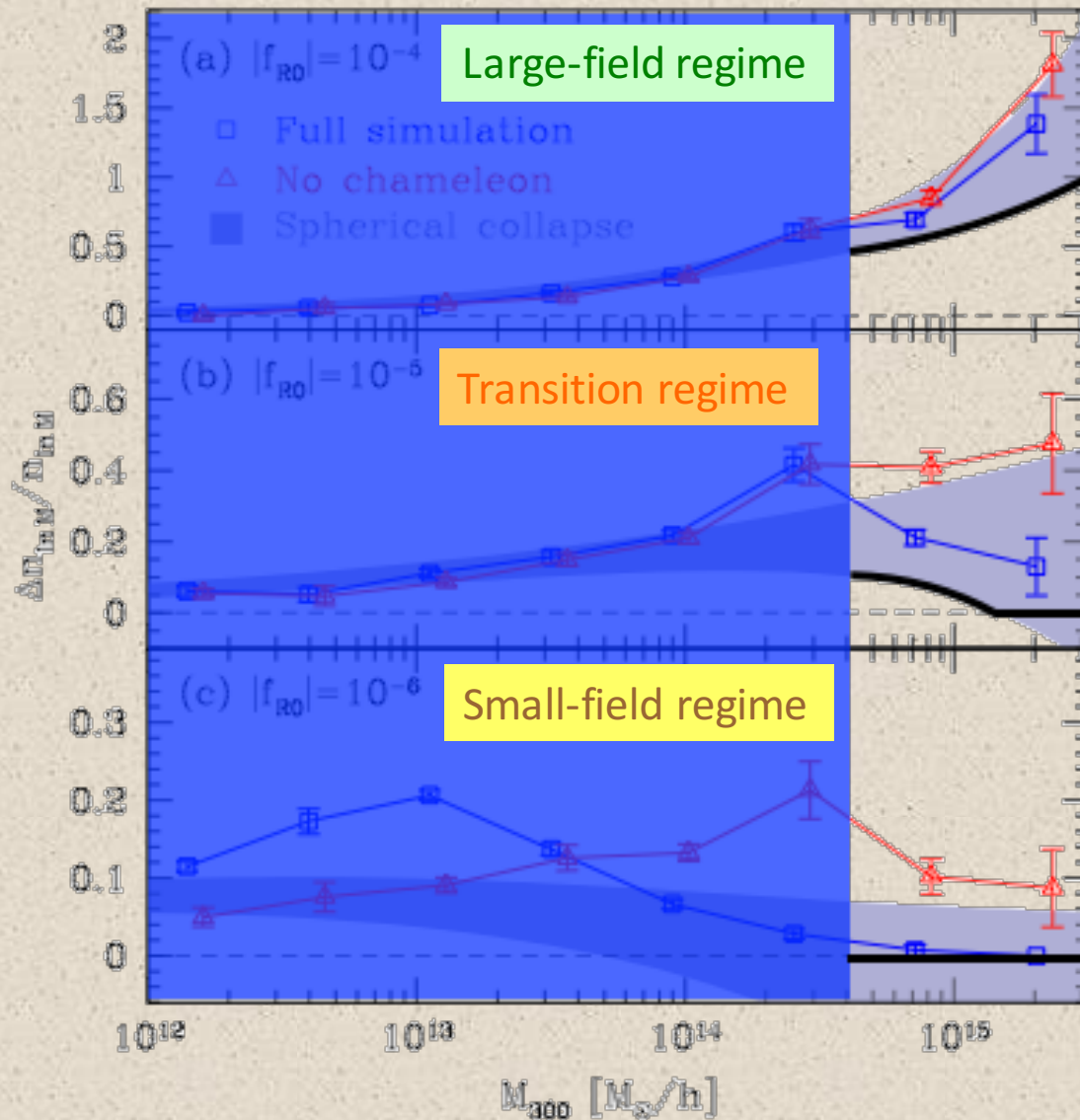
Halo mass function modeling



- Using N-body simulations including the **Chameleon screening mechanism** (Schmidt et al 2009).

- Our **modeling** is based on the **bottom line** of the blue shaded area, which as shown here is **conservative**.

Halo mass function modeling

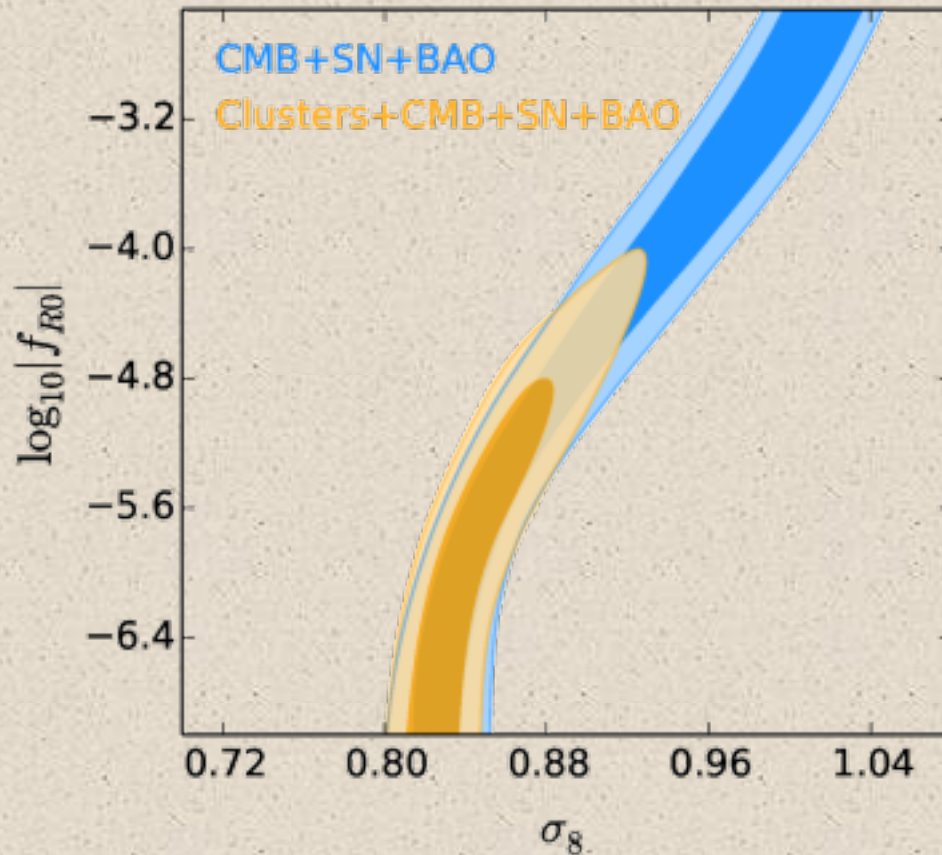


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HS $n=1$ $f(R)$ growth + flat Λ CDM

Cataneo et al 15 (PRD 2015, 92, 044009)



$$|f_{R0}| < 1.62 \times 10^{-5} (All)$$

Clusters: XLF: BCS+REFLEX+MACS ($z < 0.5$) 224 survey with 94 X-ray follow-up (Mantz et al 2015) + 50 weak lensing (Weighing the Giants; von der Linden et al 2014) + cluster f_{gas} (Mantz et al 2014) CMB (Planck collaboration 2014; SPT, Story et al 2013; ACT, Das et al 2014) + SNIa (Union 2.1, Suzuki et al 2012) + BAO (6dF, Beutler et al 2011; SDSS, Padmanabhan et al 2012, Anderson et al 2014; WiggleZ, Blake et al 2011)

Gold: all data sets combined

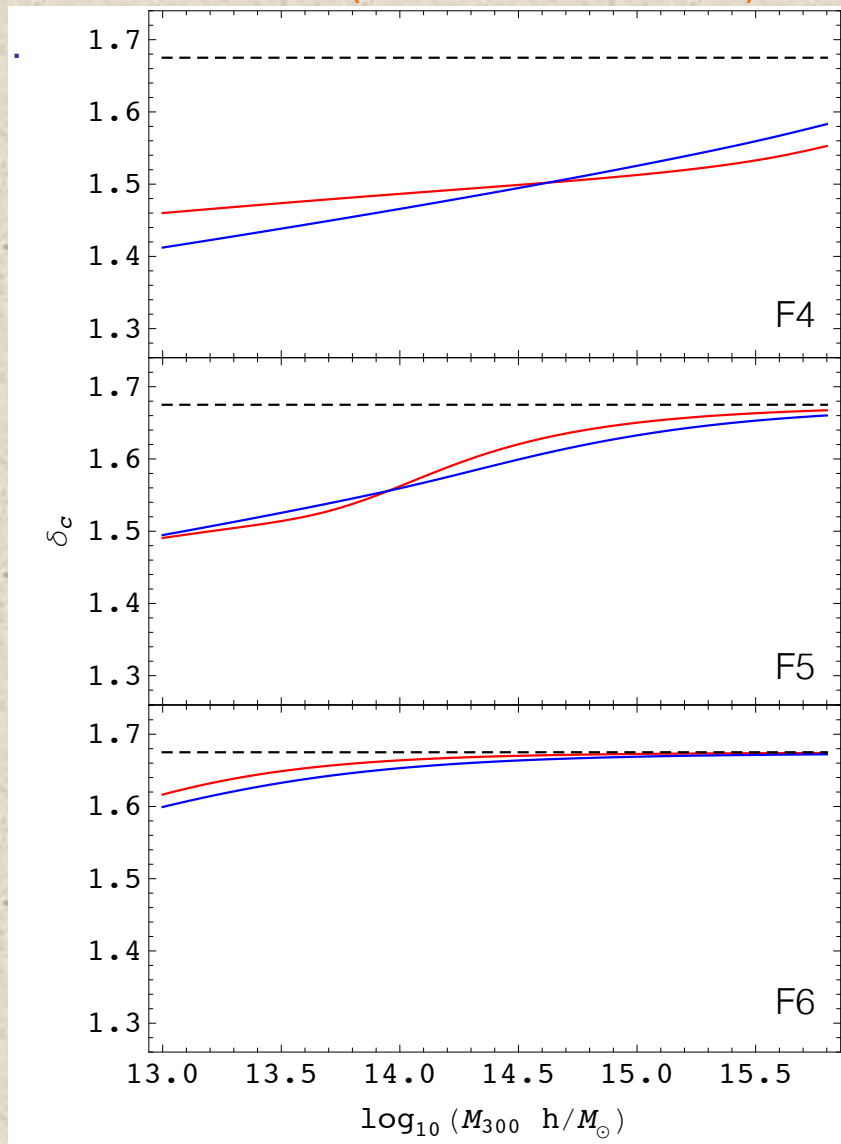
For General Relativity $|f_{R0}|=0$

Correlation between σ_8 and $|f_{R0}|$

$$\rho = 0.73 (All) \quad \rho = 0.90 (CMB)$$

Chameleon screening refinement

Cataneo et al 16, (JCAP 2016, 12, 024)

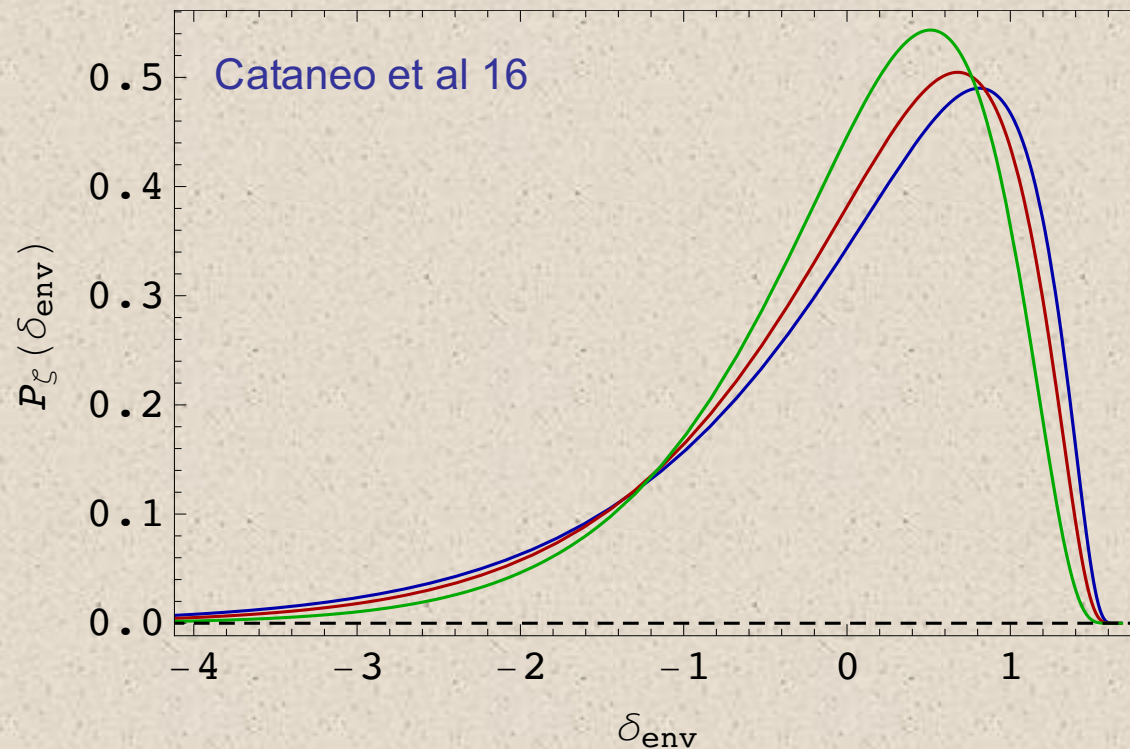


- Modeling the effect of the Chameleon screening mechanism on the mass function using high-resolution N-body simulations to obtain a more accurate mass function (Cataneo, Rapetti, Lombriser, Li, 2016)
- Spherical collapse thresholds at $z=0$ ($\Omega_m=0.281$). Lombriser et al 2014 calculations at the peak of the environmental density distribution in blue and our corrected/calibrated δ_c in red to account for self-screening and environmental screening mechanisms.

Chameleon screening refinement

$$\delta_c^{\text{eff}} = \epsilon(M \mid M_{\text{th}}^{(1)}, \mu, M_{\text{th}}^{(2)}, \nu, \alpha) \times \delta_c(\delta_{\text{env}}^{\text{peak}})$$

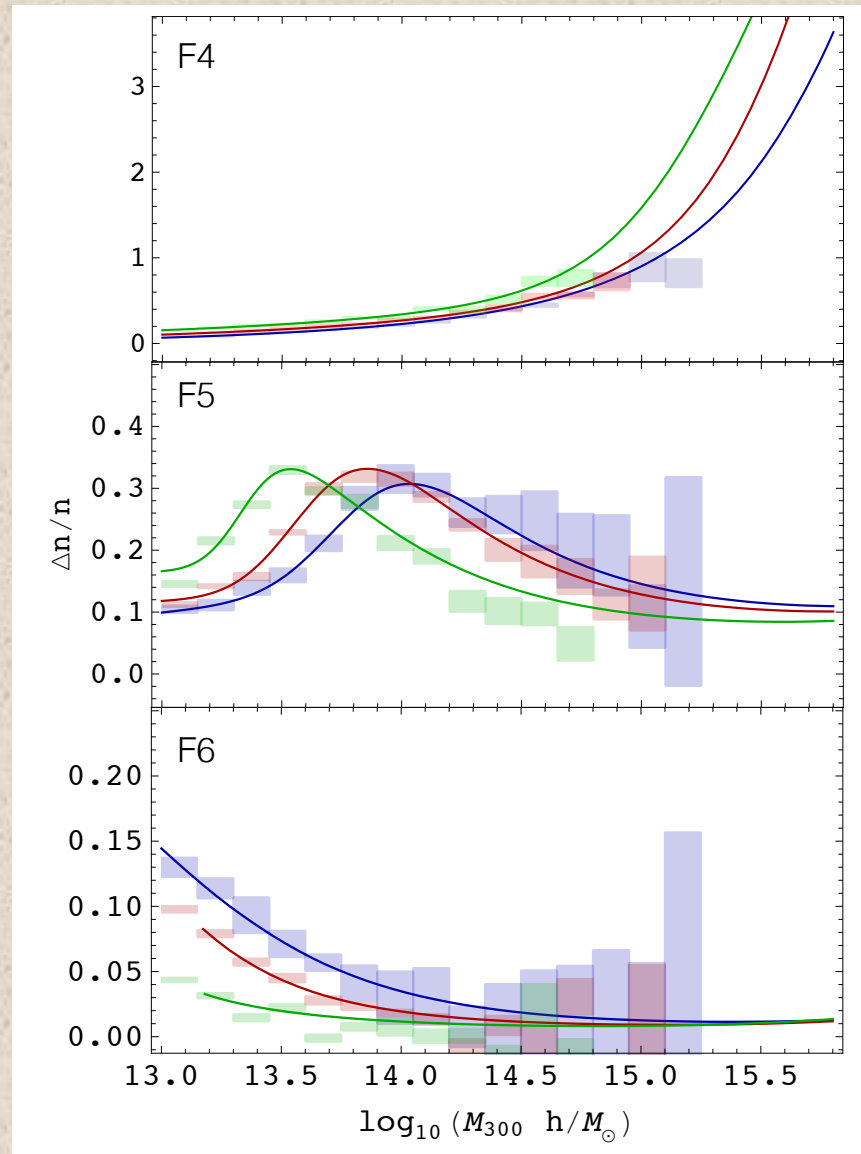
$$\epsilon = \frac{1 + (M/M_{\text{th}}^{(1)})^\eta (\delta_c^\Lambda / \delta_c^{f(R)})^\chi + (M/M_{\text{th}}^{(2)})^\vartheta (\delta_c^{f(R)} / \delta_c^\Lambda)}{1 + (M/M_{\text{th}}^{(1)})^\eta + (M/M_{\text{th}}^{(2)})^\vartheta}$$



GR/LCDM environmental density probability distribution (Lam & Sheth 08); $z=0$, 0.2 , 0.5 ; $\Omega_m=0.281$

New halo mass function

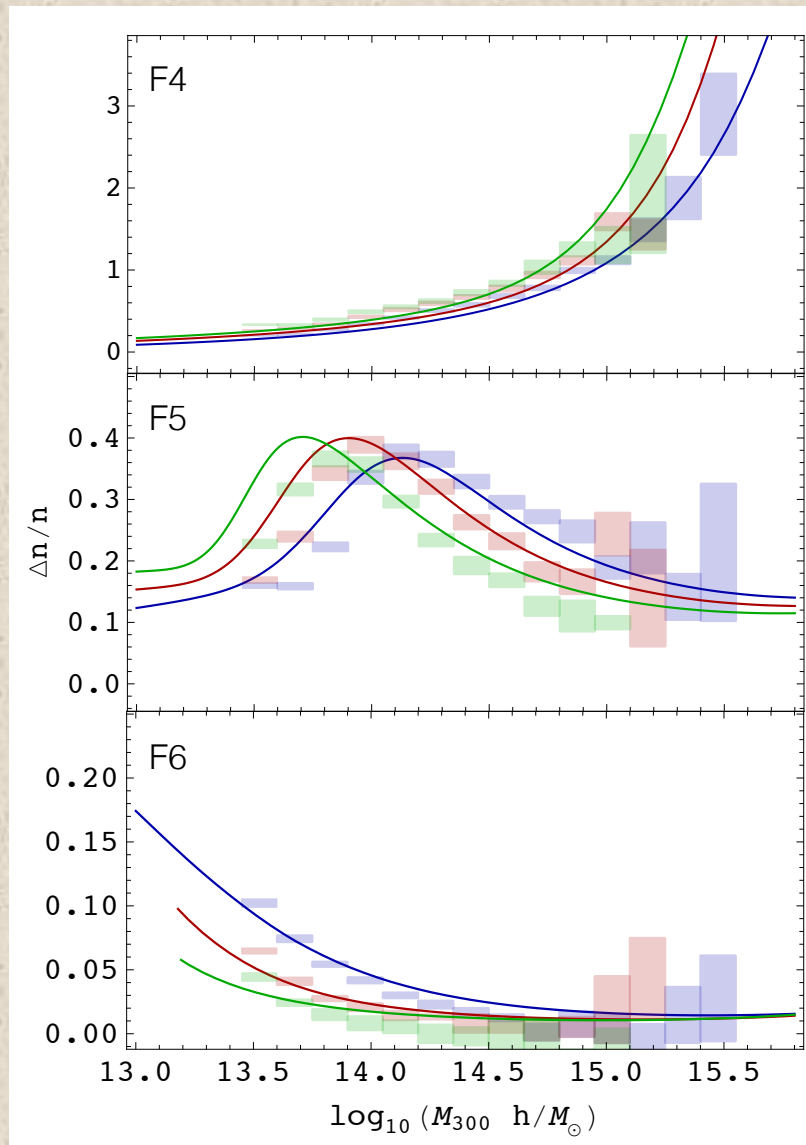
Cataneo et al 16



- High-resolution: $L_{\text{box}}=1024 \text{ Mpc}/h$; $N_{\text{particles}}=1024$; force resolution = $15.3 \text{ kpc}/h$; $N_{\text{realizations}} = 1$; $z = 0$ (blue), 0.1, 0.2 (red), 0.3, 0.4, 0.5 (green); $n_s=0.971$, $\Omega_m=0.281$, $H_0=69.7$, $T_{\text{CMB}}=2.7255$, $Y_{\text{He}}=0.24$, $N_{\nu}=3$, $\sigma_8=0.82$ (flat Λ CDM background). Current fits:
5% for $10^{13.5} \leq M_{300m}(M_{\odot}/h)^{-1} \leq 10^{15.5}$
 $10^{-6} \leq |f_{R0}| \leq 10^{-4}$ and $0 \leq z \leq 0.5$
- Rockstar halo finder with spherical overdensity masses with average density equals $300\rho_m$; simulations divided in octants; uncertainties on the HMF $f(R)/\text{GR}$ ratios propagated with the jackknife method with halos in mass bins of $\Delta \log_{10}(M)=0.15$. (We only keep mass bins with $N_{\text{halos}} > 20$.)

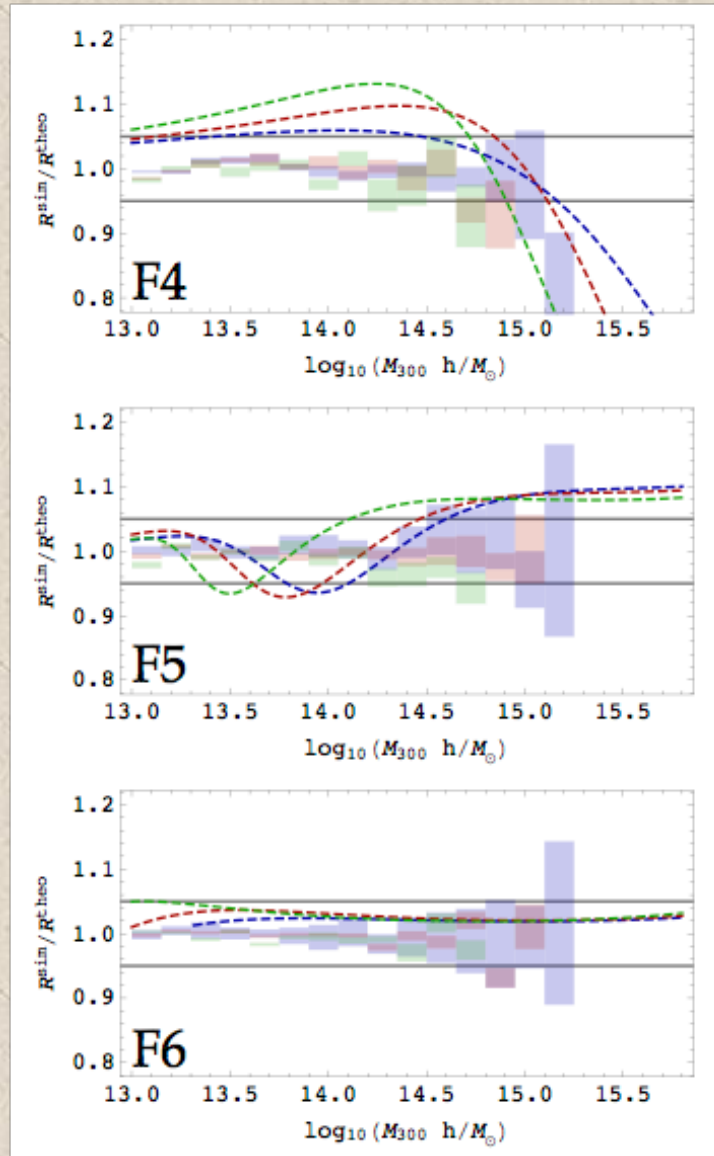
New halo mass function

Cataneo et al 16



- Low-resolution: $L_{\text{box}}=1.5 \text{ Gpc}/h$; $N_{\text{particles}}=1024$; force resolution = $22.9 \text{ kpc}/h$; $N_{\text{realizations}} = 6$; $z = 0$ (blue), 0.11, 0.25 (red), 0.43 (green); $n_s=0.958$, $\Omega_m=0.24$, $H_0=73$, $T_{\text{CMB}}=2.7255$, $Y_{\text{He}}=0.24$, $N_{\nu}=3$, $\sigma_8=0.796$ (flat Λ CDM background).
- Rockstar halo finder with spherical overdensity masses with average density equals $300\rho_m$; uncertainties on the HMF $f(R)/\text{GR}$ ratios propagated with the jackknife method using the six realizations with halos in mass bins of $\Delta \log_{10}(M)=0.15$. (We only keep mass bins with $N_{\text{halos}} > 20$.)

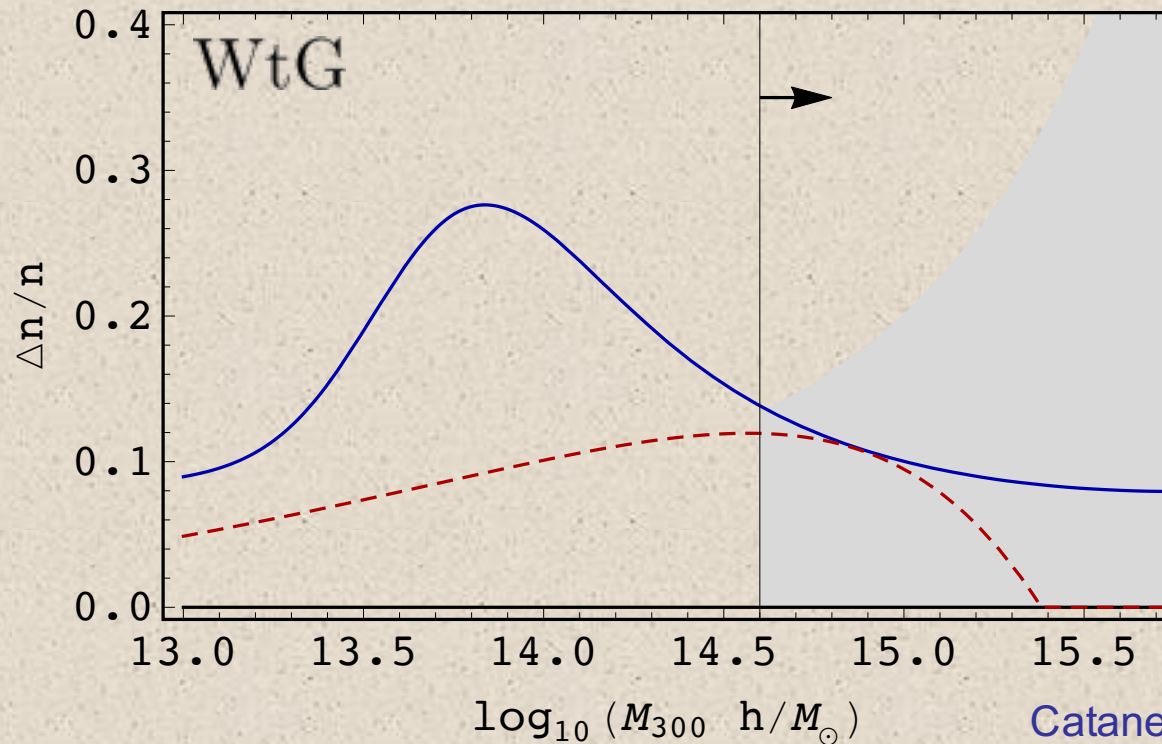
New halo mass function



$$\mathcal{R}(M, z, \bar{f}_{R0}, \boldsymbol{\theta}) = \frac{n_{\ln M}^{f(R)}}{n_{\ln M}^{\text{GR}}}$$

- Dashed lines show the effect of neglecting the correction,
- which is only weakly dependent on standard cosmological parameters
- and potentially insensitive to baryonic physics and massive neutrinos

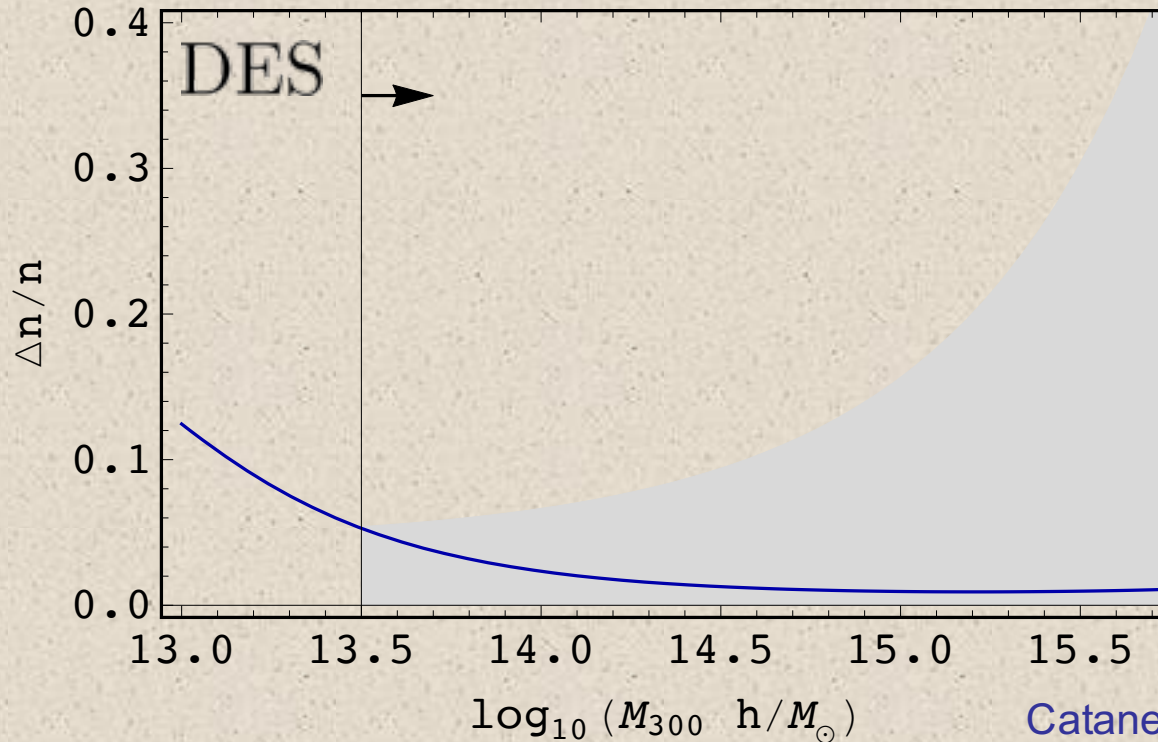
Comparison of current and prospective constraints



$$|\bar{f}_{R0}| \lesssim 5 \times 10^{-6}$$

Gray band, 7% mass calibration from WtG with the corresponding lower mass limit; red line, current 2-sigma level from the WtG constraints using the previous mass function; blue line, projected constraints for the next WtG constraints with the new $f(R)$ HMF (\sim a factor of 2 improvement).

Prospective constraints for a lower mass/z survey

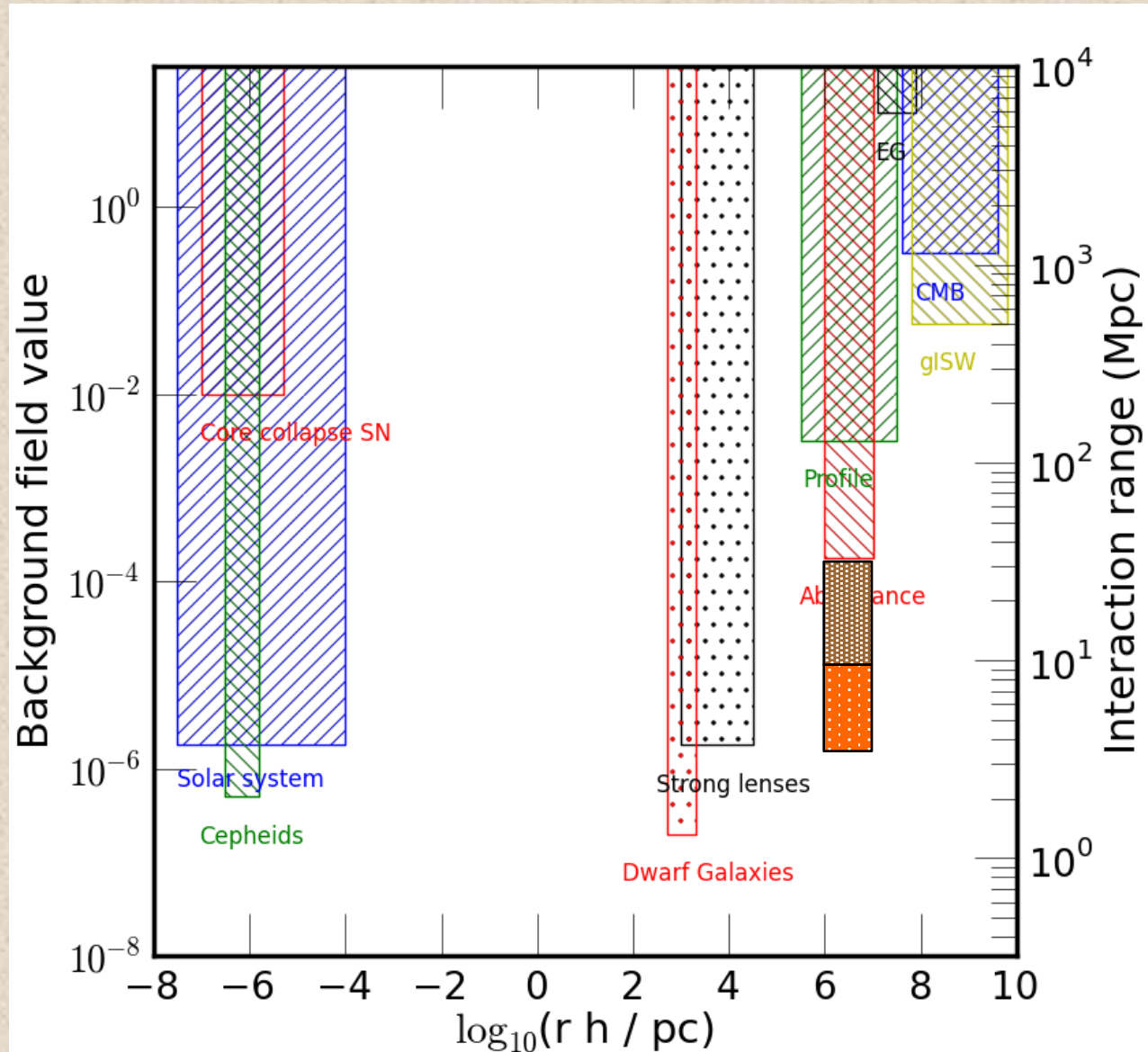


Cataneo et al 16

$$|\bar{f}_{R0}| \lesssim 10^{-6}$$

Gray band, 5% mass calibration from DES with the corresponding lower mass limit; blue line, projected constraints for DES constraints with the new $f(R)$ HMF (~an order of magnitude improvement).

Comparison of current and prospective constraints



Adapted from Joyce,
Jain, Khoury & Trodden
2014 (review)

Cataneo et al 15
 $\lambda_c < \sim 10 \text{ Mpc}$ $|f_{R0}| < \sim 10^{-5}$

Cataneo et al 17 (in pr.),
a factor of ~ 2 better

DES forecast
 $\lambda_c < \sim 3 \text{ Mpc}$ $|f_{R0}| < \sim 10^{-6}$
(Cataneo et al 2016)