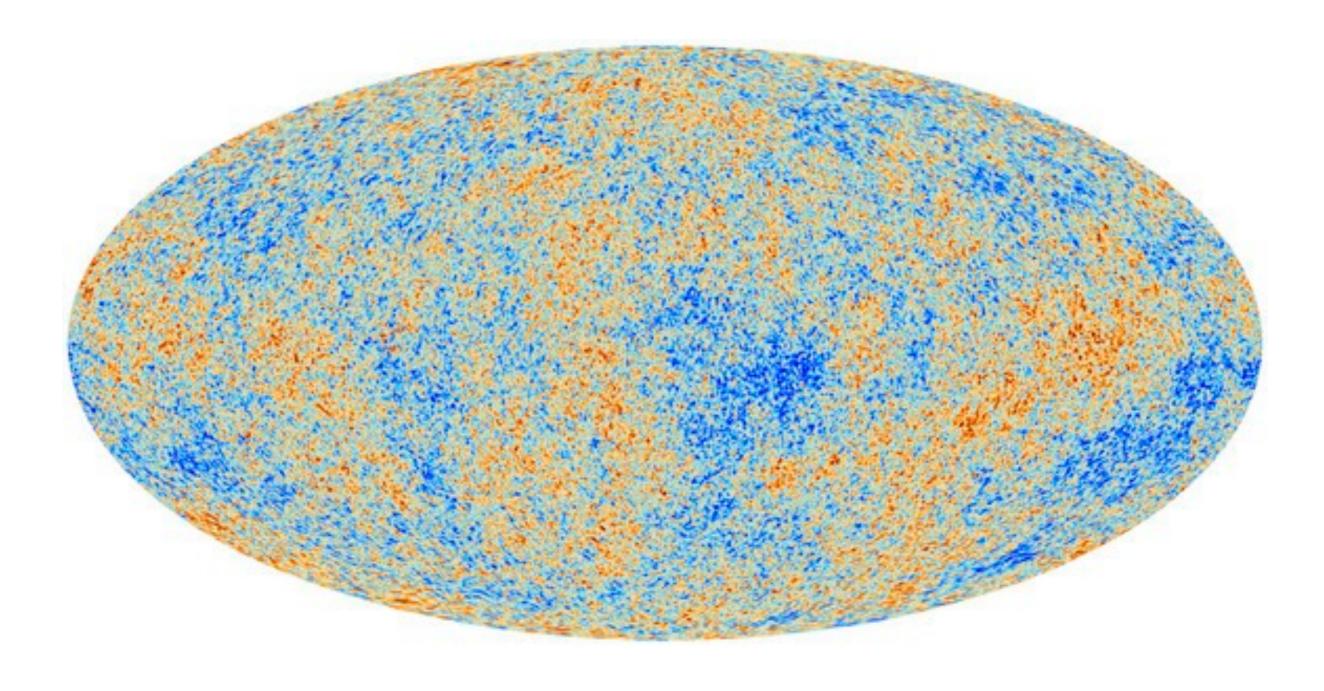
Gravity from Cosmology

Pedro G. Ferreira
University of Oxford

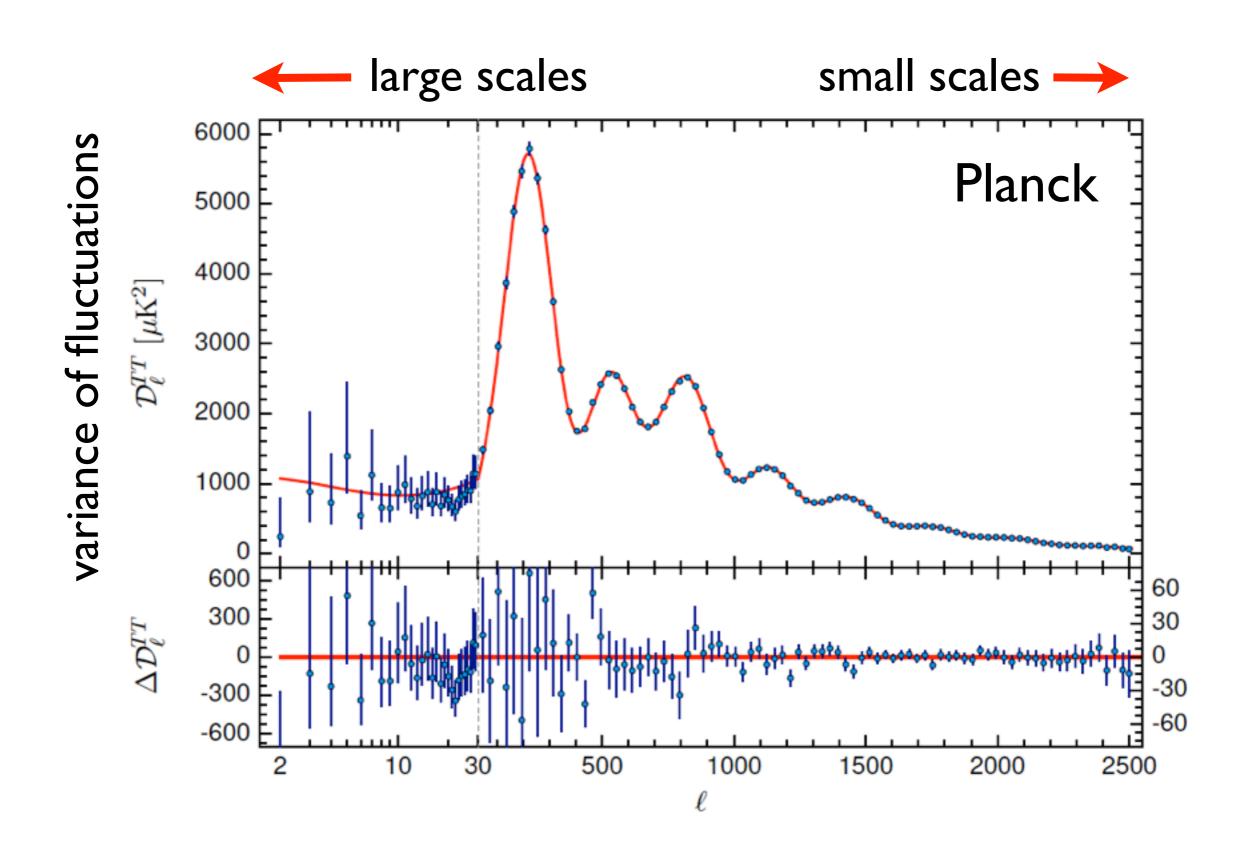
The state of General Relativity in 1957

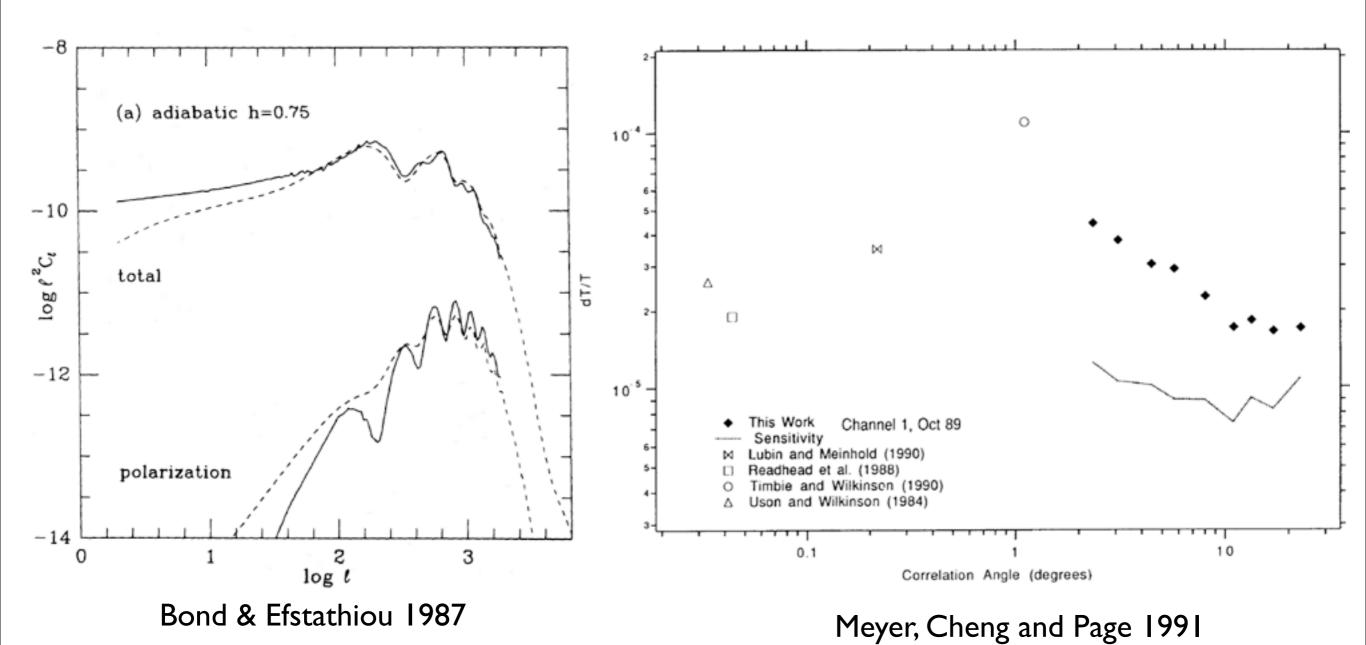
"There exists... one serious difficulty, and that is the lack of experiments. Furthermore, we are not going to get any experiments, so we have to take the viewpoint of how to deal with the problems where no experiments are available. ... the best viewpoint is to pretend that there are experiments and calculate. In this field we are not pushed by experiments but pulled by imagination."

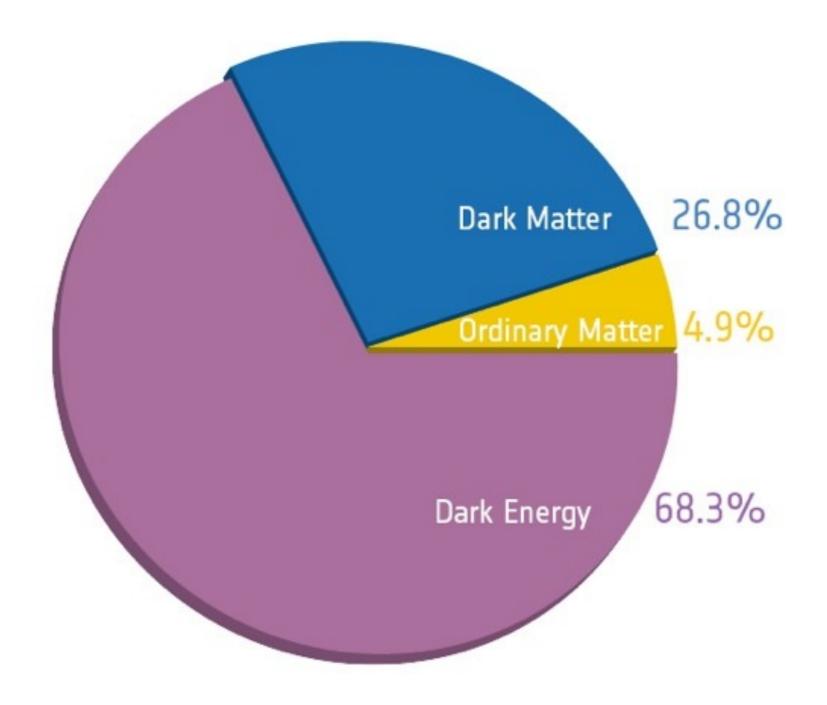
R. Feynman, Chapel Hill workshop on gravitation (1957)



Planck 2015

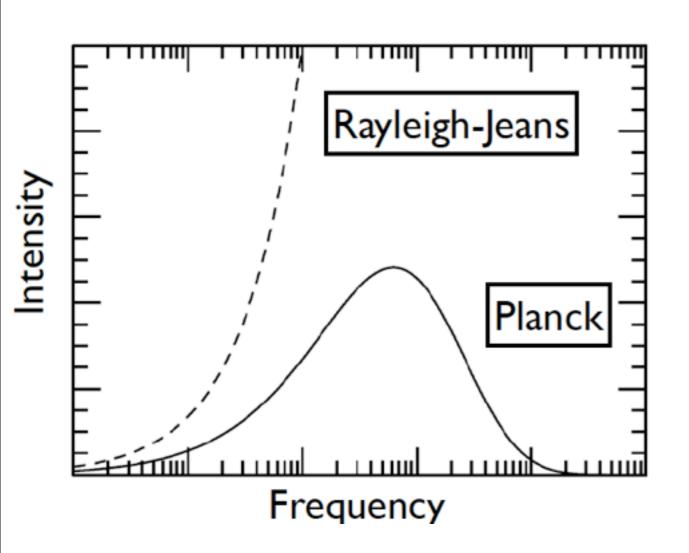






Big puzzles

Small inconsistencies



- Lamb shift
- Wu parity violation experiment
- Fitch-Cronin CP violation
- Precession of perihelion of Mercury

The Theory

Einstein Gravity: a field theory view

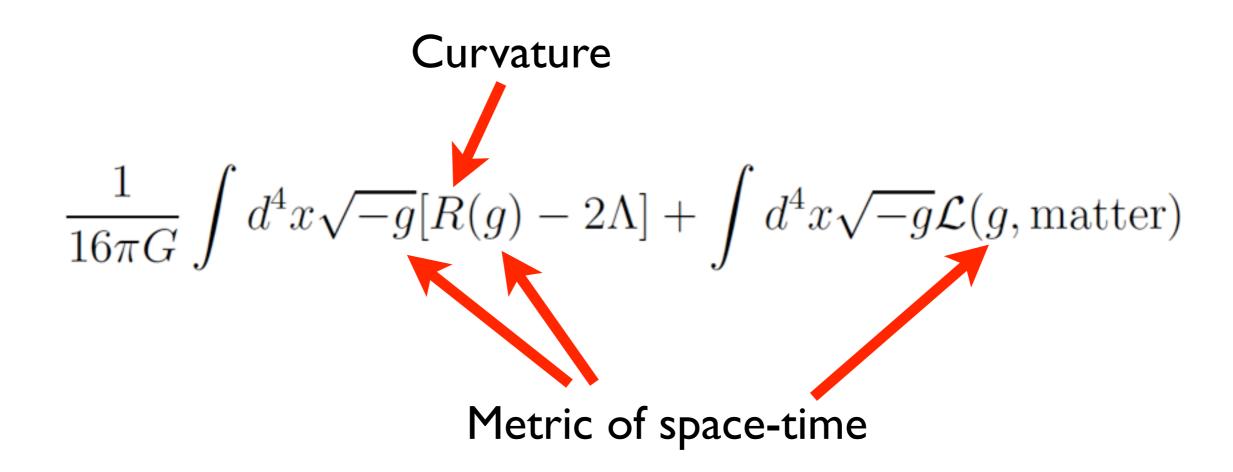
Feynman/Weinberg Theorem

$$g_{\alpha\beta}=\eta_{\alpha\beta}+h_{\alpha\beta}$$
 Spin-2 field
$$\Box h_{\mu\nu}=16\pi G(T_{\mu\nu}^M+T_{\mu\nu}^G)$$

Feynman (1963) Weinberg (1965) Deser (1970)

Unique non-linear completion is GR...

Einstein Gravity: a relativists view



Lovelock's theorem (1971): "The only second-order, local gravitational field equations derivable from an action containing solely the 4D metric tensor (plus related tensors) are the Einstein field equations with a cosmological constant."

See also Hojman, Kuchar & Teitelboim (1976)

Jordan-Brans-Dicke Theory

One free parameter

$$S = \int d^4x \sqrt{-g} \left[\phi R - \frac{\omega_{\rm BD}}{\phi} \nabla_{\mu} \phi \nabla^{\mu} \phi + V + \mathcal{L}_M[g_{\mu\nu}, \phi] \right\}$$

Cassini (Bertotti et al 2003) $\omega_{
m BD} > 40,000$

Extra degrees of freedom

metric
$$\longrightarrow$$
 add ϕ , A^{μ} , $f_{\alpha\beta}$ etc.

4D — e.g. in 5 dimensions:

$$g_{AB} = \left(\begin{array}{cc} g_{\alpha\beta} & A_{\alpha} \\ A_{\beta} & \phi \end{array}\right)$$

2nd order
$$\longrightarrow$$
 e.g. if $\int d^4x \sqrt{-g} f(R)$ define $\phi = \frac{df}{dR}$.

Local
$$\longrightarrow$$
 e.g. $\phi = \frac{R}{\Box}$.

All transform differently under diffeomorphisms

Background $ds^2 = g^{(0)}_{\mu\nu} dx^{\mu} dx^{\nu} = -dt^2 + a^2(t)(d\vec{x})^2$

$$G_{\alpha\beta} = 8\pi G T_{\alpha\beta} + U_{\alpha\beta}$$

where
$$U_{\alpha\beta}(a,\dot{a},\rho_M,P_M,\phi,\cdots)$$

$$P_X = \underline{w}\rho_X$$

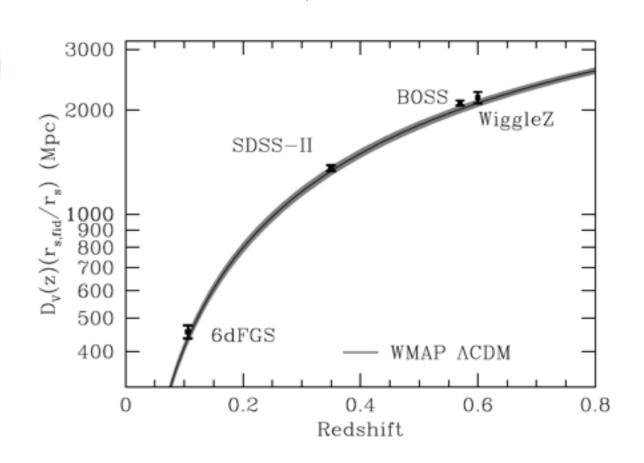
Homogeneity and isotropy

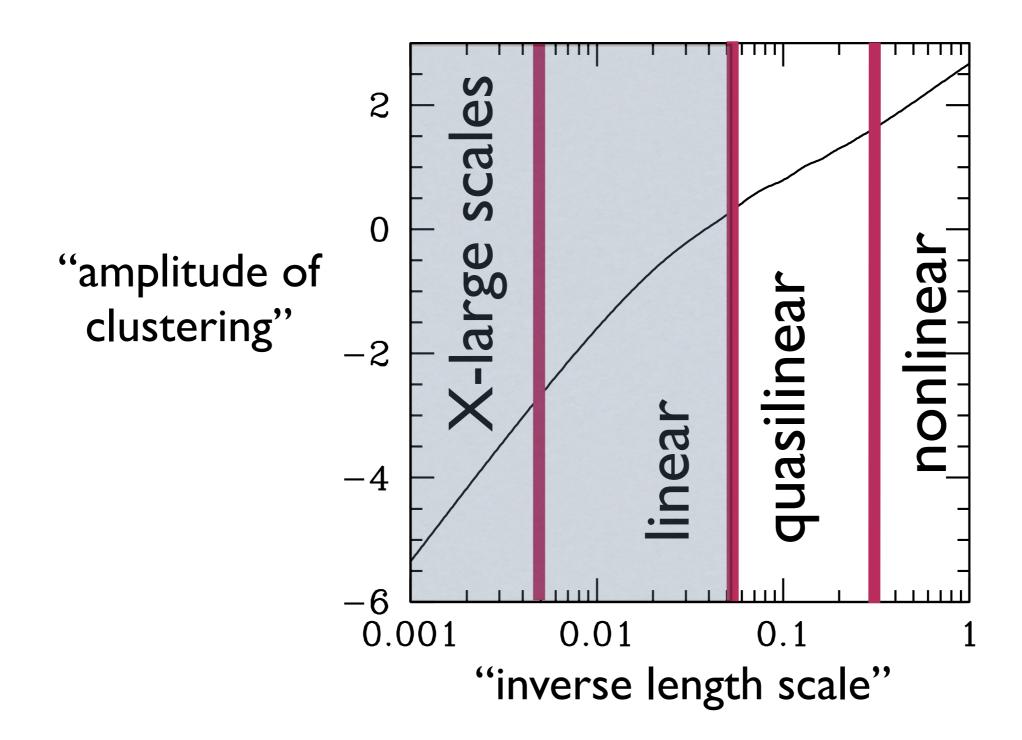
$$U_{\alpha\beta} = 8\pi G \begin{pmatrix} \rho_X & 0\\ 0 & a^2 P_X \delta_{ij} \end{pmatrix}$$

Bianchi identities

$$\nabla^{\alpha}[8\pi G T_{\alpha\beta} + U_{\alpha\beta}] = 0$$

BOSS, Anderson et al 2013.





$$N(k) \propto k^3$$

More statistical power -

Linear Perturbations

$$g_{\alpha\beta} = g_{\alpha\beta}^{(0)} + h_{\alpha\beta}$$

$$\rho_M = \bar{\rho}_M (1 + \delta_M)$$

$$\phi = \phi_0 + \delta\phi$$

Construct most general quadratic action which has:

- upto 2nd order in time derivatives
- $h_{\alpha\beta} \to h_{\alpha\beta} + \nabla_{\alpha}\xi_{\beta} + \nabla_{\beta}\xi_{\alpha}$ where $x^{\alpha} \to x^{\alpha} + \xi^{\alpha}$
- inherits symmetries of the background

Lagos et al 2016

Linear Perturbations

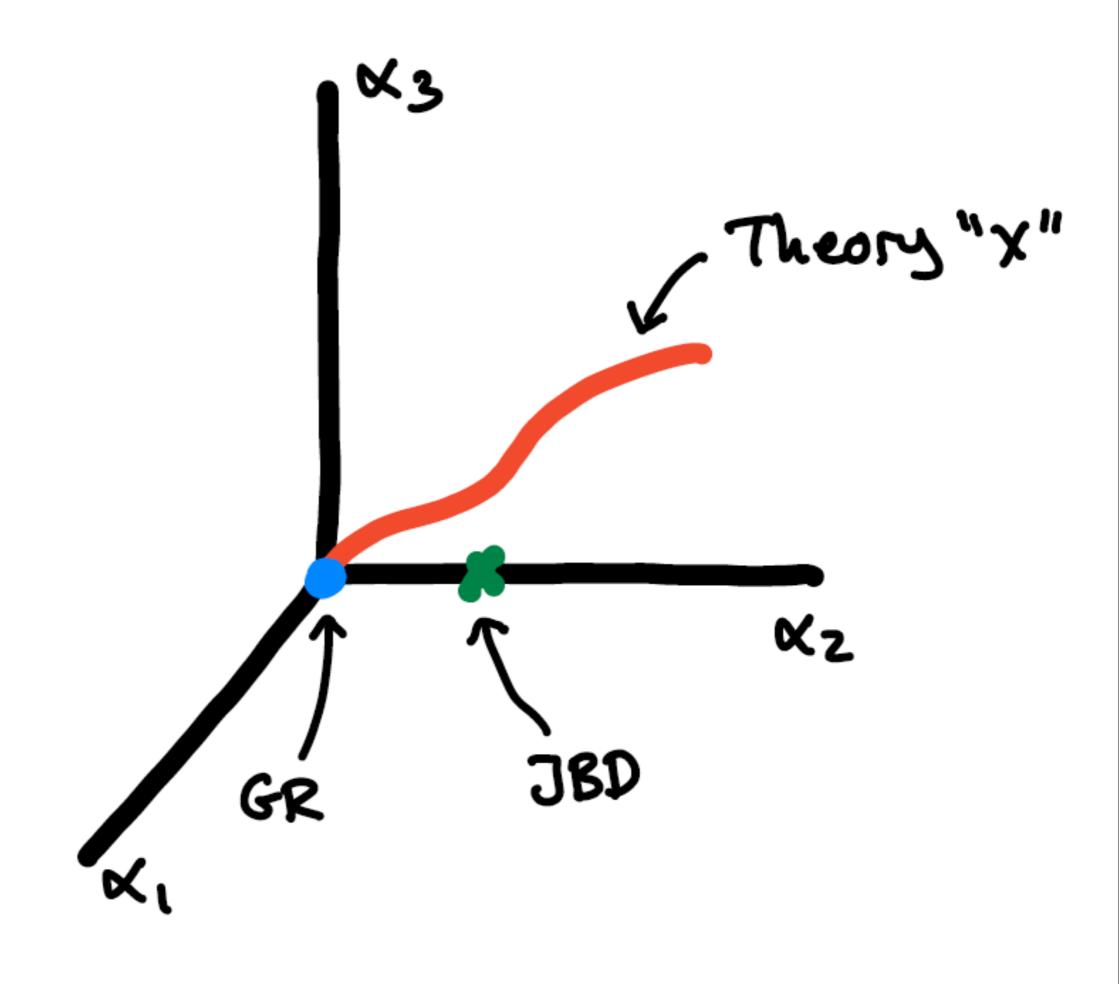
$$S = \int d^4x \sqrt{-g^{(0)}} \left[\alpha_1 \nabla^{\mu} h^{\alpha\beta} \nabla_{\mu} h_{\alpha\beta} + \alpha_2 \nabla^{\mu} h^{\alpha\beta} \nabla_{\alpha} h_{\mu\beta} \right]$$
$$+ \alpha_3 \nabla^{\mu} h \nabla^{\alpha} h_{\mu\alpha} + \alpha_4 \nabla^{\mu} h \nabla_{\mu} h + \alpha_5 \nabla^{\mu} h_{\mu\alpha} \nabla^{\alpha} \delta \phi + \cdots \right]$$

Properties:

- $(\alpha_1, \alpha_2, \cdots)$ are functions of t
- $\alpha_X(t)$ depend on transf. props of extra fields
- clear mapping theory \longleftrightarrow $\alpha_X(t)$
- clear physical interpretation of each $\alpha_X(t)$

Examples:

- Scalar-tensor (Horndeski): five $\alpha_X(t)$
- Vector-tensor (Einstein-Aether, Proca): nine $\alpha_X(t)$
- Tensor-tensor (Bigravity, massive gravity): three $\alpha_X(t)$



The Data

A preferred length scale- the horizon

$$\mathcal{H}^{-1} \equiv \left(\frac{\dot{a}}{a}\right)^{-1} \propto \tau \simeq 3000h^{-1}\mathrm{Mpc}$$

Most surveys $\ll \tau$ so that $k\tau \gg 1$

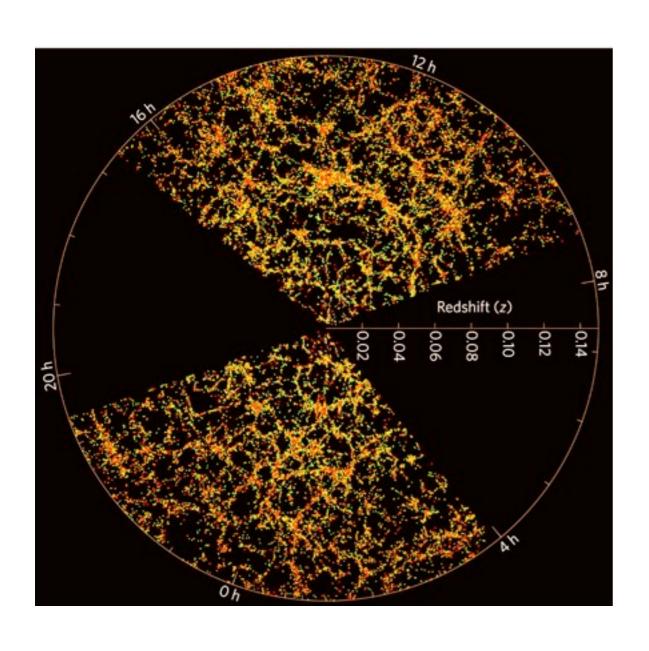
Newtonian potentials:
$$h_{\alpha\beta}=2\left(\begin{array}{cc}\Phi&0\\0&a^2\Psi\delta_{ij}\end{array}\right)$$

Einstein equations:
$$-k^2\Phi = 4\pi G\mu a^2\rho\Delta$$

$$\gamma\Psi = \Phi$$

 (μ, γ) are rational functions of $\alpha_X(t)$ and k^2

We measure matter and light.





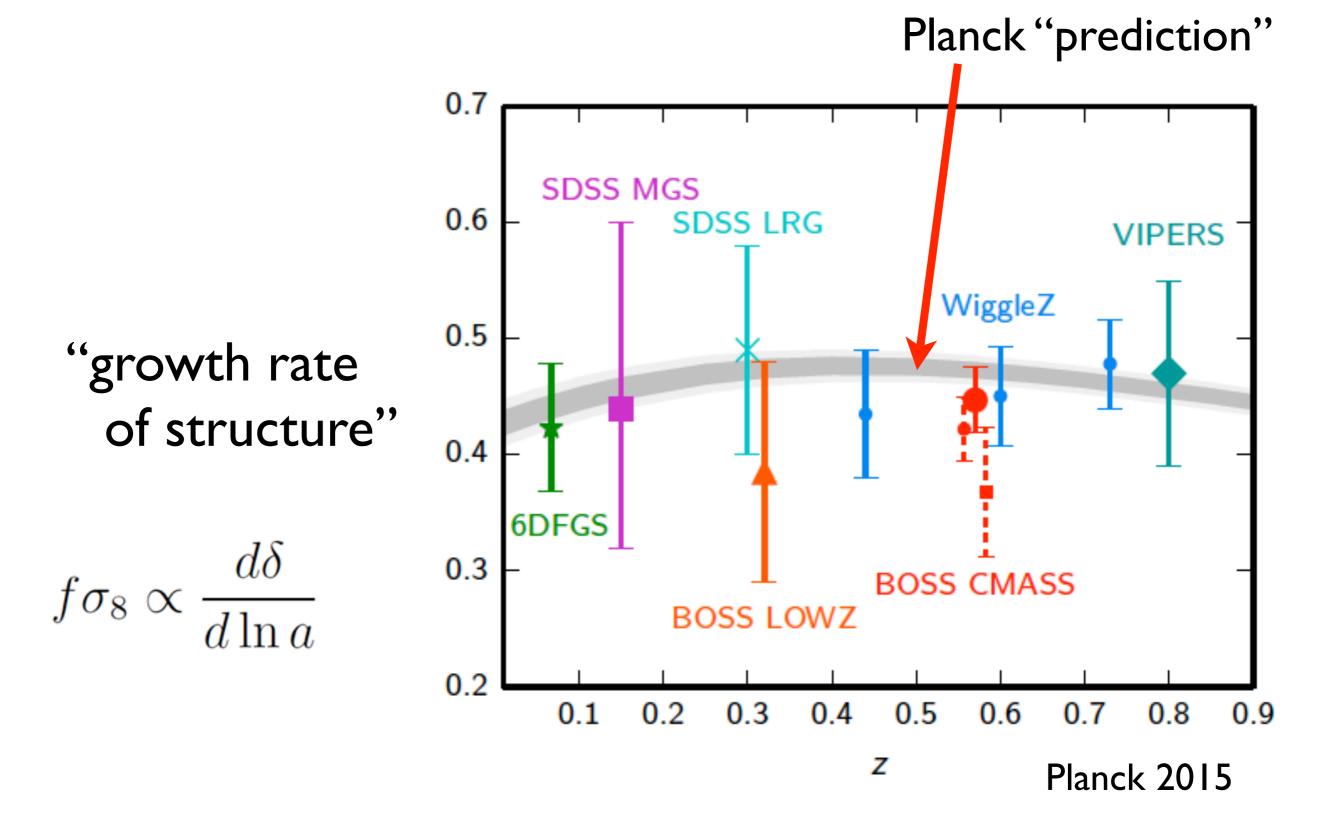
Growth rate

$$f(k,a) = \frac{d \ln \delta_M(k,a)}{d \ln a} \qquad f$$

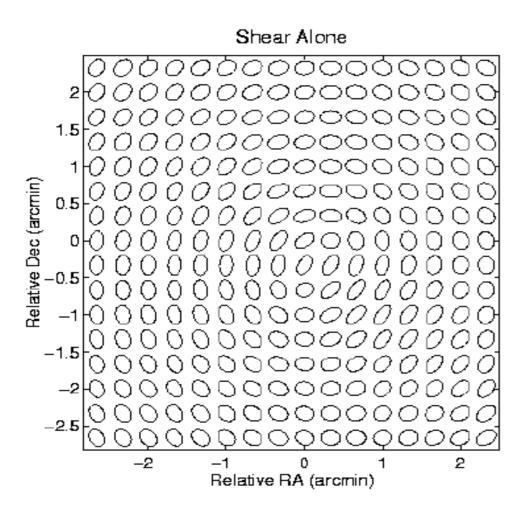
f satisfies a simple ODE

$$\frac{df}{d\ln a} + qf + f^2 = \frac{3}{2}\Omega_M \xi$$

with
$$q=rac{1}{2}[1-3w(1-\Omega_M)]$$
 and $\xi=rac{\mu}{\gamma}$



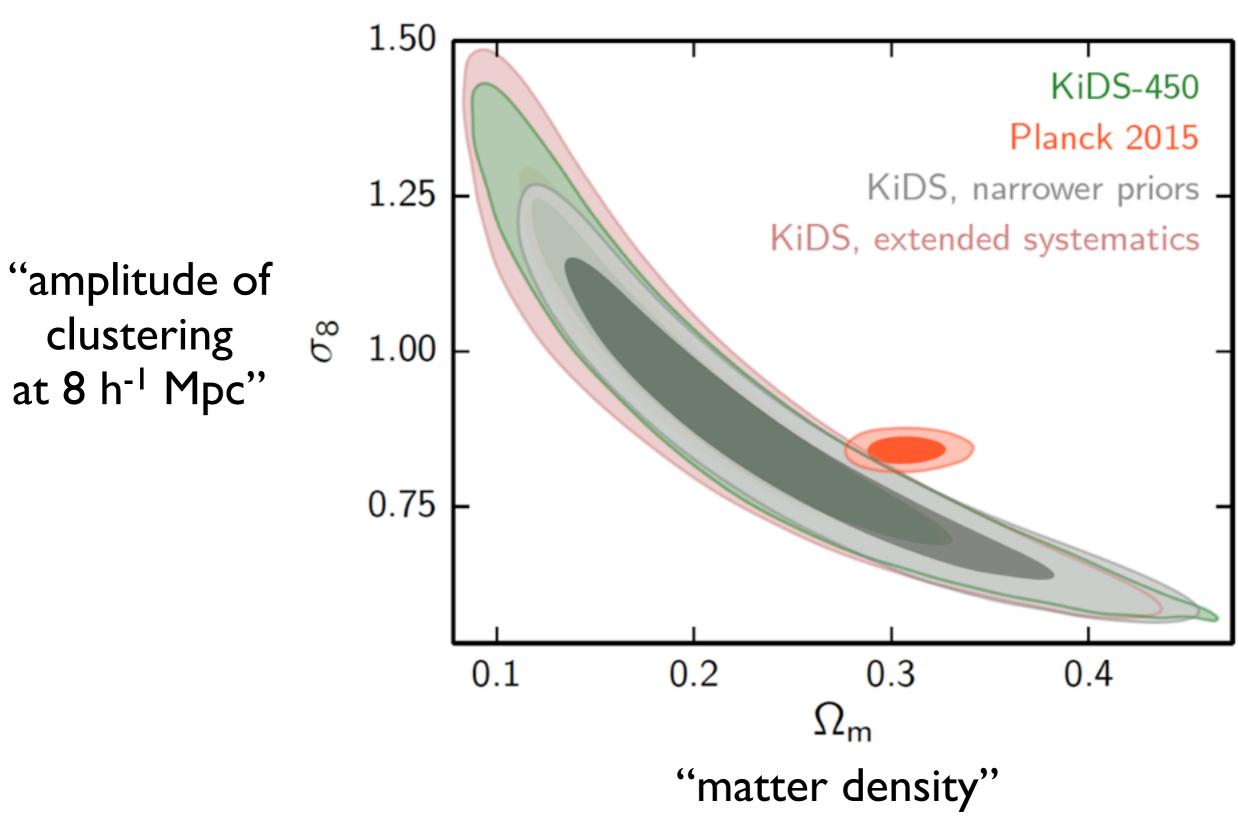
Weak Lensing

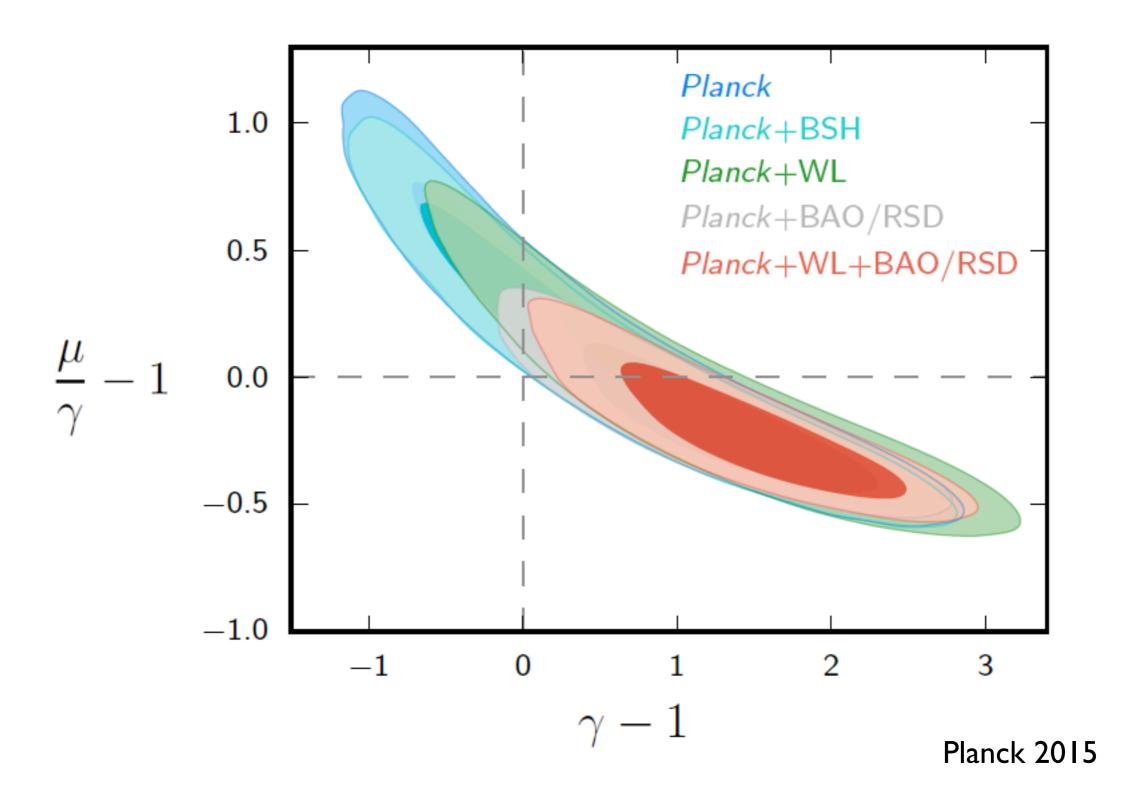


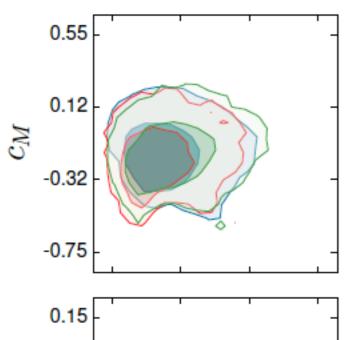
$$\begin{split} & \operatorname{shear} \simeq \int_0^\chi \nabla_\perp^2 [\Phi + \Psi](\chi') \left[\chi' \left(1 - \frac{\chi'}{\chi} \right) \right] d\chi' \\ & \operatorname{shear} \, \sim \, \Sigma \equiv \mu \left(1 + \frac{1}{\gamma} \right) \end{split}$$

Sarah Bridle lectures (2003)



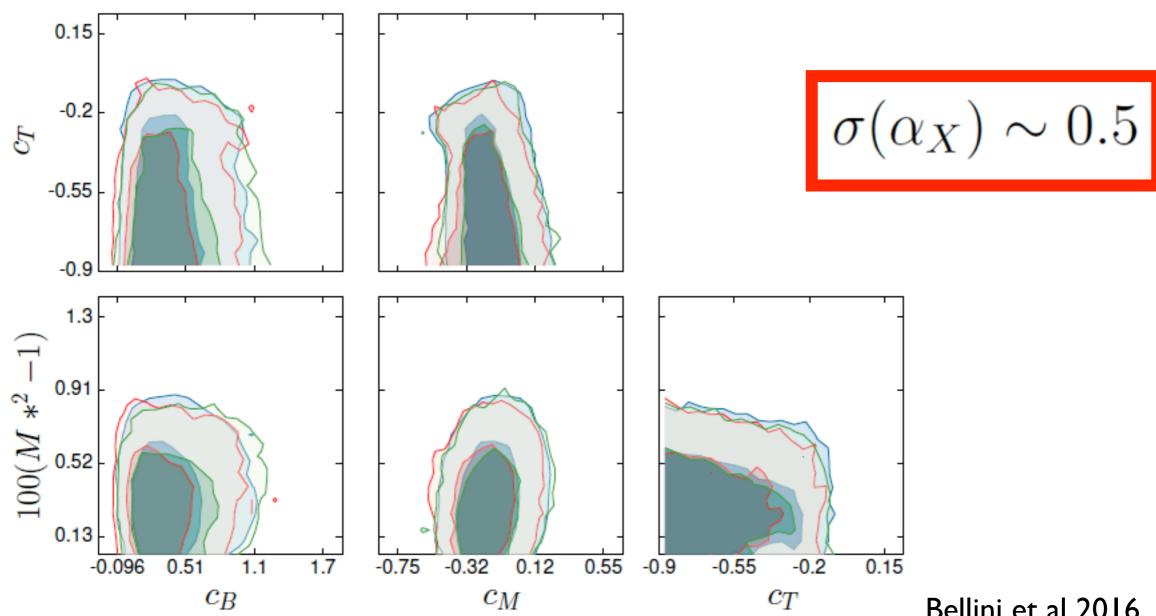






Constrain $\alpha_X(t)$ in scalar-tensor theories

Parametrize:
$$\alpha_X = b_X + c_X \frac{\Omega_{\rm DE}(z)}{\Omega_{\rm DE}(z=0)}$$



Jordan-Brans-Dicke Theory

One free parameter

$$S = \int d^4x \sqrt{-g} \left[\phi R - \frac{\omega_{\rm BD}}{\phi} \nabla_{\mu} \phi \nabla^{\mu} \phi + V + \mathcal{L}_M[g_{\mu\nu}, \phi] \right\}$$

Cassini (Bertotti et al 2003)

$$\omega_{\rm BD} > 40,000$$

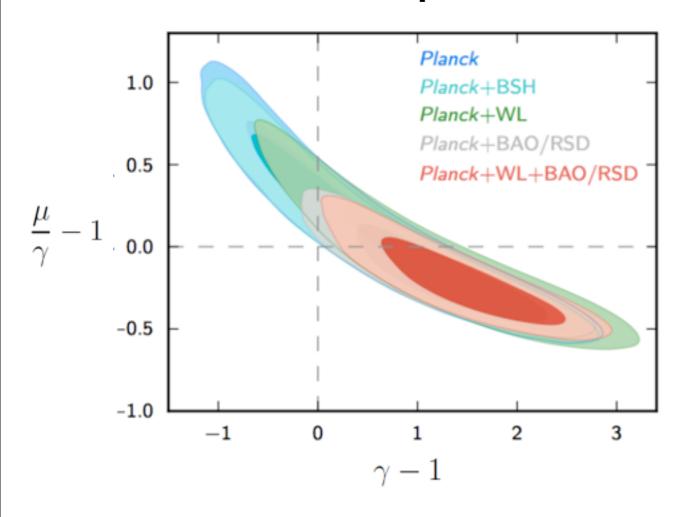
Planck (Avilez & Skordis 2015)

$$\omega_{\rm BD} > 1,000$$

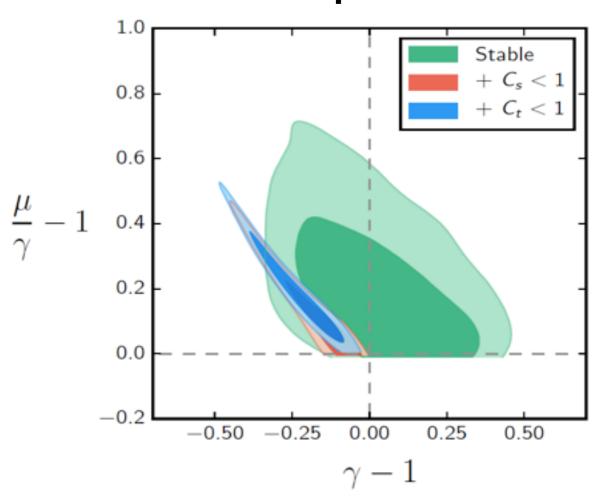
The Challenge

Priors

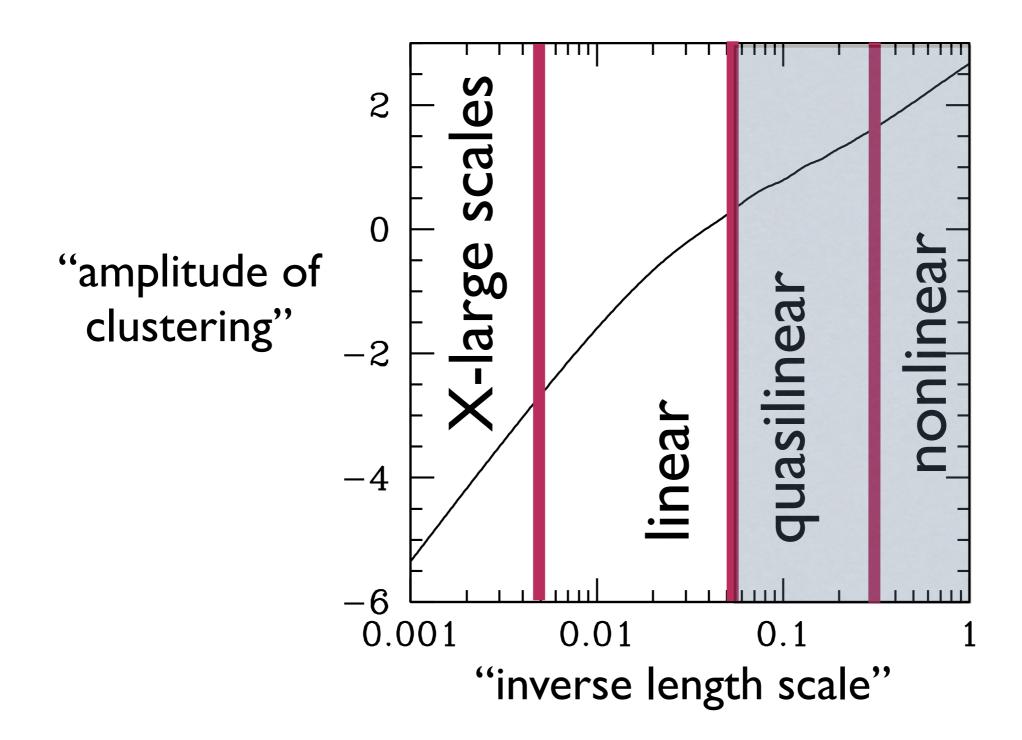
without priors



with priors



Salvatelli et al 2016



$$N(k) \propto k^3$$

More statistical power —

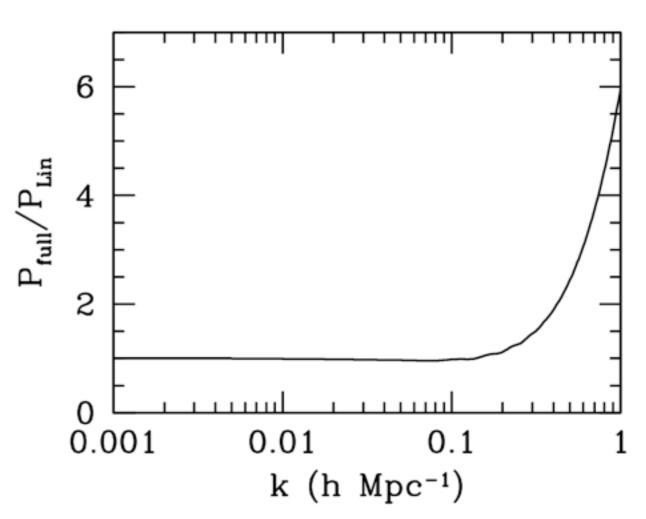
Non-linear physics

baryonic feedback

λ (Mpc/h) 1.0 10.0 0.1 REF/DMONLY 1.2 P(k)/P(k)_{DARK} 0.8 σ_8 =0.84, f_{gas} =0.175 --- σ₈=0.74, f_{oas}=0.120 0.6 0.1 1.0 10.0 100.0 k (h/Mpc)

Sembloni et al 2012

non-linear growth



Screening

Newtonian potential

$$\Phi = -\frac{GM}{r}$$

$$\ddot{\vec{r}} = -\nabla[\Phi + \phi]$$

Fifth force

$$\phi = -\frac{GM}{r}e^{-mr}$$

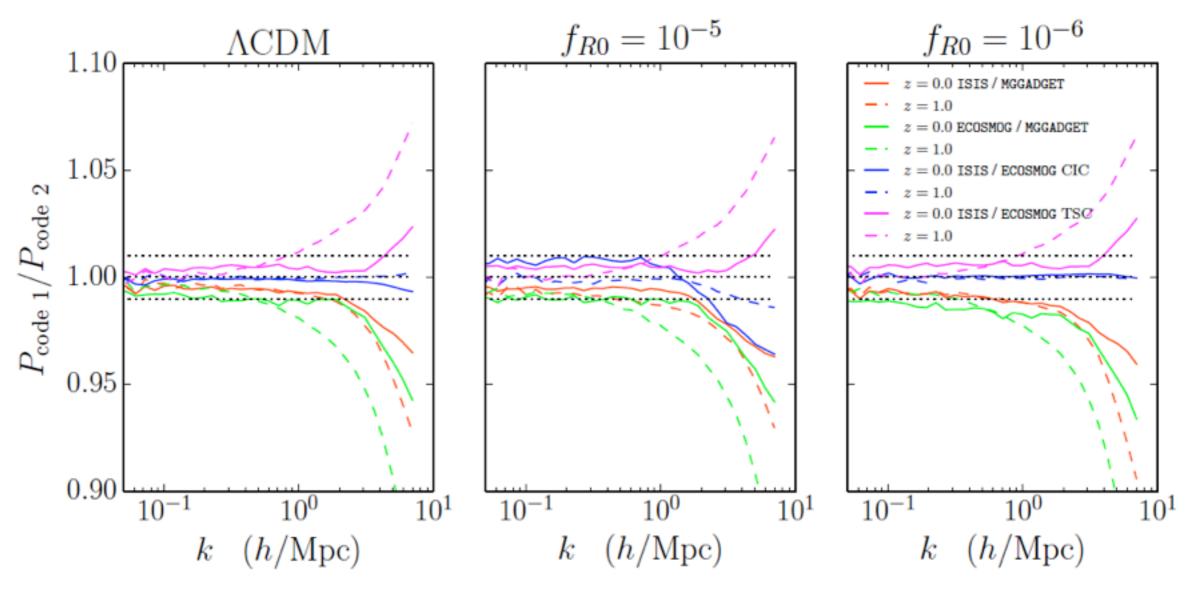
Chameleon: $m=m(\rho)$ $m\to\infty$ when $\rho\gg\rho_S$

Vainshtein: $\tilde{G}
ightarrow 0$ when $r \ll r_V$

Non-linear effect...

Non-linear physics

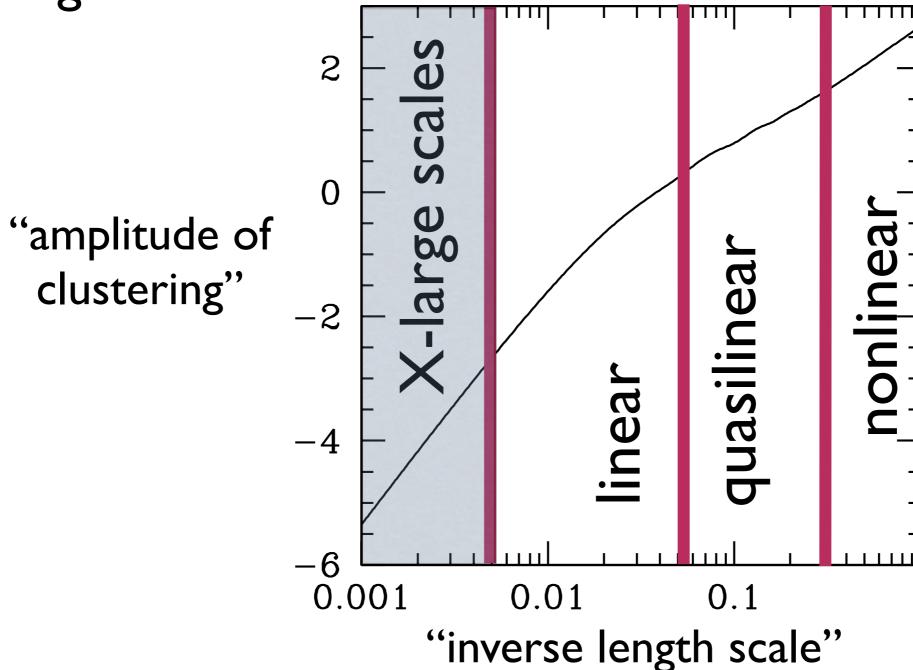
Comparison of N-body codes: 1% out to k~1



Winther et al 2015

Model specific ———— Model independent?

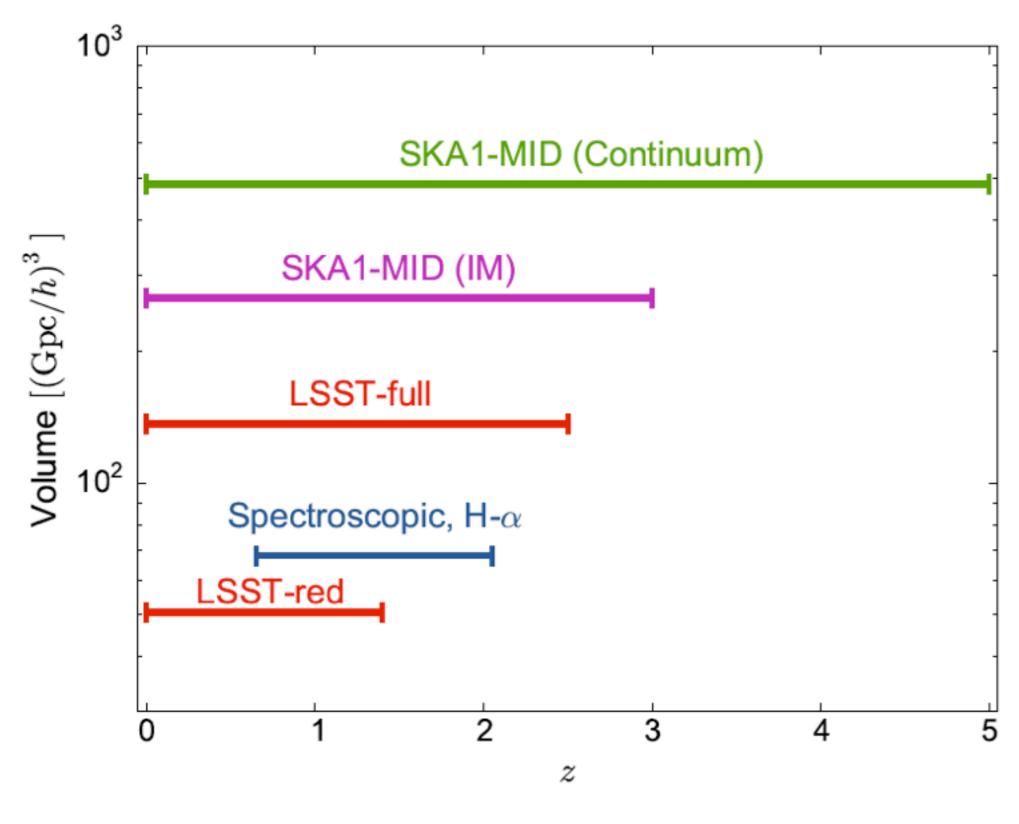
Ultra-large scales



$$N(k) \propto k^3$$

More statistical power ——

Ultra-large scales

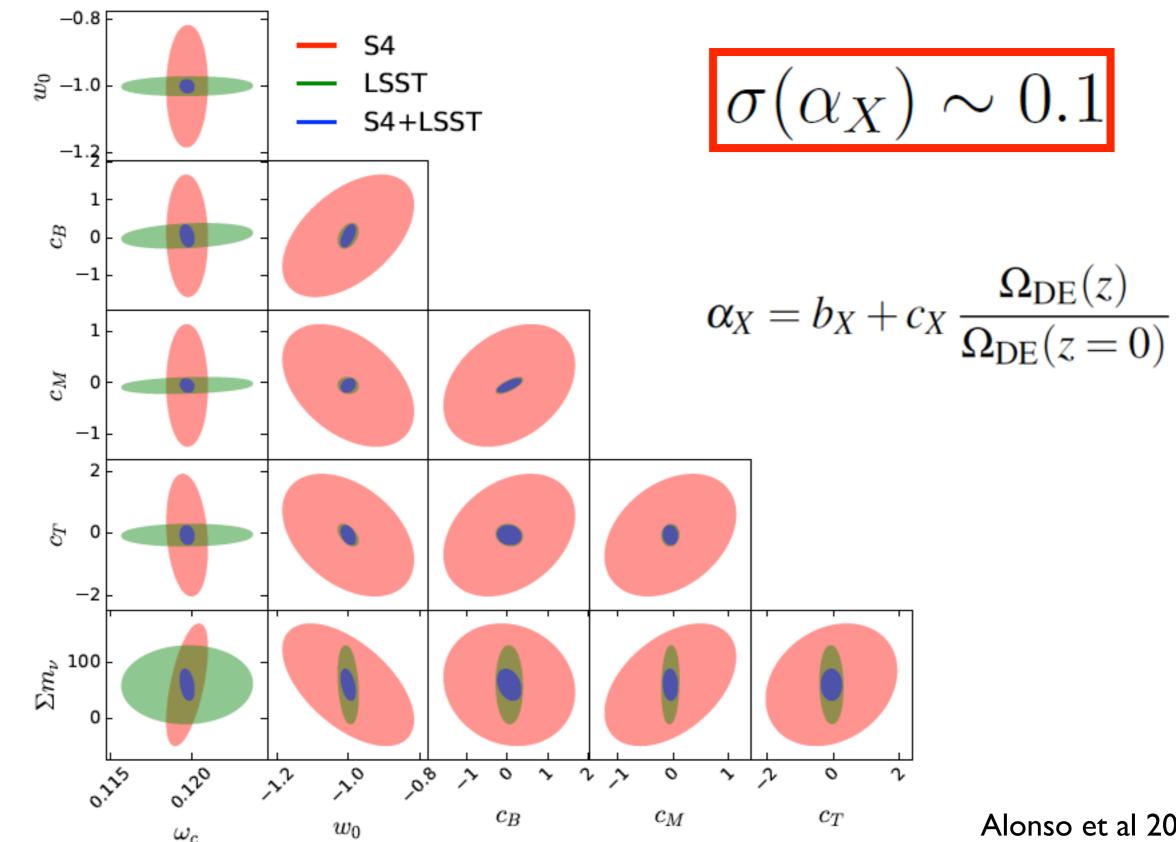


Plot by D. Alonso

The Future

The Future is now

Data Type	Now	Soon	Future
Photo-z:LSS (weak lensing)	DES, RCS, KIDS	HSC	LSST, Euclid, SKA, WFIRST
Spectro-z (BAO, RSD,)	BOSS	DESI,PFS,HETDEX, Weave	Euclid, SKA
SN Ia	HST, Pan-STARRS, SCP, SDSS, SNLS	DES, J-PAS	JWST,LSST
CMB/ISW	WMAP, Planck	AdvACT	Simons Array, Stage IV, LiteBird
sub-mm, small scale lensing, SZ	ACT, SPT,Planck, ACTPol,SPTPol,	PolarBear,Spider, Vista	CCAT, SKA
X-Ray clusters	ROSAT, XMM, Chandra	XMM, XCS, eRosita	
HI Tomography	GBT	Meerkat, Baobab, Chime, Kat 7	SKA



Alonso et al 2016

Jordan-Brans-Dicke Theory

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Cassini (Bertotti et al 2003)

$$\omega_{\rm BD} > 40,000$$

Planck (Avilez & Skordis 2015)

$$\omega_{\rm BD} > 1,000$$

$$\omega_{\rm BD} > 20,000$$

"An important contribution of the general theory of relativity to cosmology has been to keep out theologians by a straightforward application of tensor analysis."

E. Schucking

D. Alonso, L. Amendola, M. Amin, T. Baker, R. Bean, E. Bellini, C. Blake, P. Brax, P. Bull, C. Burrage, S. Daniels, A. Davies, P. G. Ferreira, J. Gleyzes, G. Gubitosi, W. Hu, L. Hui, C, Heymans, S. Joudaki, B. Jain, J. Khoury, K. Koyama, M. Kunz, M. Lagos, D. Langlois, D. Leonard, E. Linder, L. Lombrisier, D. Mota, A. Narimani, J. Noller, J. Peacock, W. Percival, F. Piazza, L. Pogosian, D. Sapone, D. Scott, I. Sawicki, A. Silvestri, F. Simpson, A. Taylor, F. Vernizzi, H. Winther, J. Zuntz, ...