

TESTING GRAVITY 2017

25-28 JANUARY 2017, SFU HARBOUR CENTER, VANCOUVER, BC, CANADA



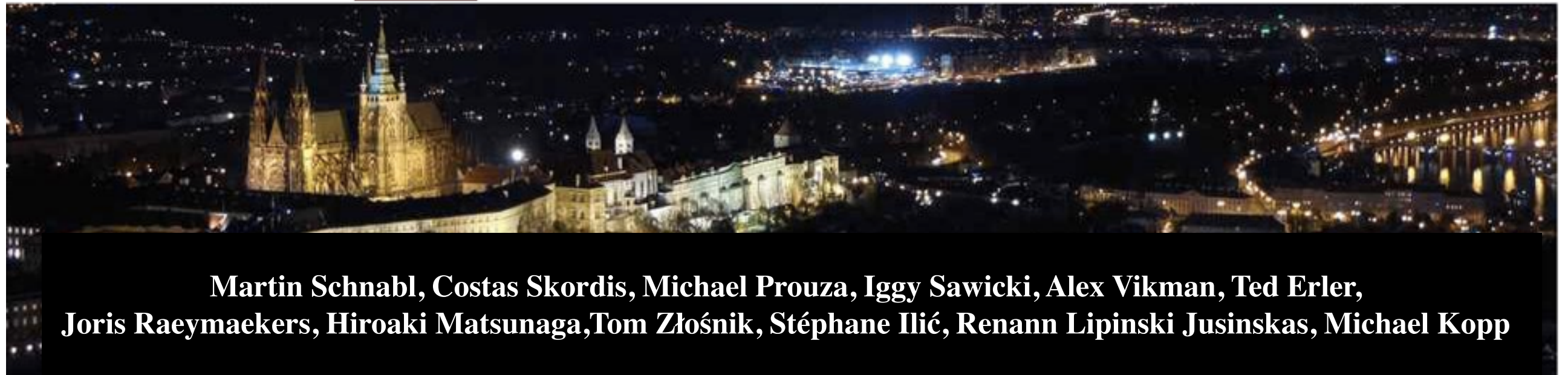
Canonical Exorcism for Cosmological Ghosts

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WE HAVE JOBS!!!

2. 2017-01-24: [String Theory](#) ([Prague, Inst. Phys.](#) - Europe) [Deadline: 2017-02-15]

Postdoc - hep-th

[Detailed record](#)

1 postdoc position

3. 2017-01-23: [Senior Postdoc in Cosmology and Gravitation](#) ([Prague, Inst. Phys.](#) - Europe) [Deadline: 2017-02-15]

Postdoc - astro-ph, gr-qc, hep-th

[Detailed record](#)

1 senior researcher position

4. 2017-01-23: [Cosmology and Gravitation](#) ([Prague, Inst. Phys.](#) - Europe) [Deadline: 2017-02-15]

Postdoc - astro-ph, gr-qc, hep-th

[Detailed record](#)

3 postdoc positions

Special THANKS for Support and Letters of Intent to cooperate with CEICO

Christos Charmousis (LPT Orsay)

Ruth Durrer (Geneva U.)

Andrei Frolov (SFU)

Gregory Gabadadze (NYU)

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Graham Smith (BNL)

Mark Trodden (UPenn)

Tanmay Vachaspati (ASU)

This talk is based on

 arXiv: **1702.XXXXX**

PREPARED FOR SUBMISSION TO JCAP

YITP-

Canonical Exorcism for Cosmological Ghosts

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Cosmological Ghosts

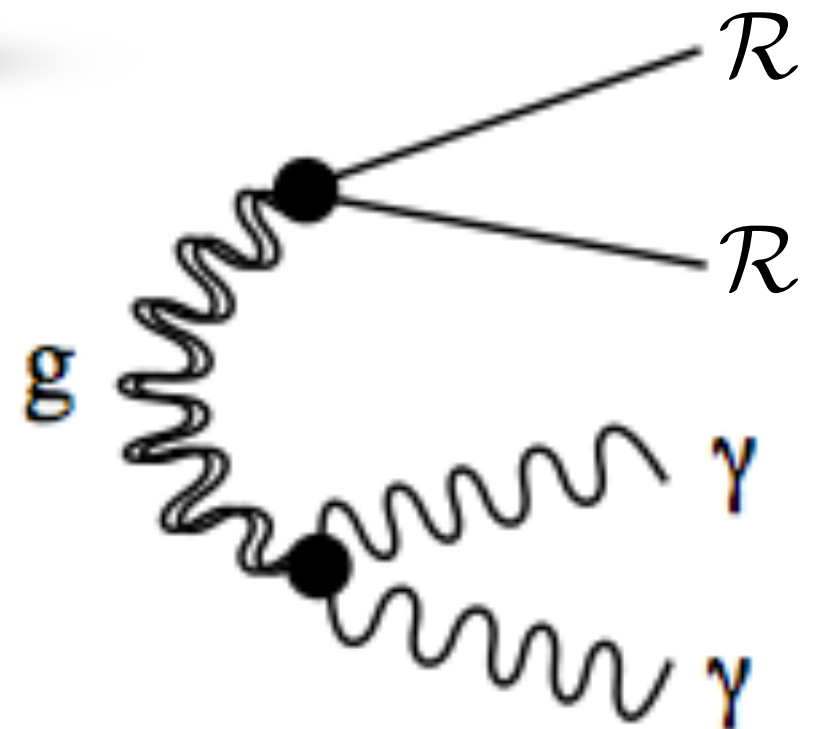


$$S = \frac{1}{2} \int d\eta d^3x Z \left((\mathcal{R}')^2 - c_S^2 (\partial_i \mathcal{R})^2 \right), \quad \mathcal{R} = \Phi + H \frac{\delta\varphi}{\dot{\varphi}},$$

Ghost $Z(t) < 0$

$$H_{\mathbf{k}} = \frac{|P_{\mathbf{k}}|^2}{2Z} + \frac{Z c_S^2 k^2 |\mathcal{R}_{\mathbf{k}}|^2}{2} < 0$$

*ghosts - modes (oscillators)
with the **negative mass***



$$\Gamma_{0 \rightarrow 2\gamma 2\phi} \sim \frac{\Lambda^8}{M_{\text{Pl}}^4}$$

Cline, Jeon, Moore, (2003)

*Can one change
the sign of the Hamiltonian
by a canonical transformation?*

Canonical Transformations

$$(q, p, H) \rightarrow (\theta, \pi, \mathcal{H})$$

preserve Poincaré-Cartan integral invariant:

$$I = \oint p dq - H dt = \oint \pi d\theta - \mathcal{H} dt$$



generating function: $p dq - H dt - (\pi d\theta - \mathcal{H} dt) = dF$



$$p = \frac{\partial F}{\partial q}, \quad \pi = -\frac{\partial F}{\partial \theta}, \quad \mathcal{H} = H + \frac{\partial F}{\partial t}$$

preserve Poisson brackets: $\{q, p\} = \{\theta, \pi\} = 1$

Motion -canonical transformation

$$(q, p, H) \rightarrow (Q_0, P_0, 0)$$

with the generating function which is the on-shell action- i.e. Hamilton principal function $S(q, Q_0)$

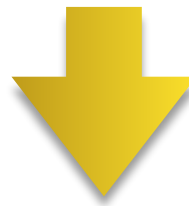
$$\frac{\partial S}{\partial t} + H \left(\frac{\partial S}{\partial q}, q, t \right) = 0$$

Hamilton-Jacobi equation

Transformation in the Heisenberg picture

$$\hat{\mathcal{O}}' = \hat{U} \hat{\mathcal{O}} \hat{U}^\dagger$$

*unitary time-dependent transformation of the **canonical variables***



$$\hat{H}'(\hat{\mathcal{O}}', t) = \hat{U}^\dagger \hat{H}(\hat{\mathcal{O}}', t) \hat{U} - i\hat{U}^\dagger \left(\frac{\partial \hat{U}}{\partial t} \right)_{\mathcal{O}'}$$

*the Hamiltonian transforms as a
connection in the non-Abelian field theories!*

Time-Dependent Linear Bogolyubov Transformations

$$\begin{aligned}\hat{q} &= \alpha \hat{\theta} + \beta \hat{\pi}, \\ \hat{p} &= \gamma \hat{\theta} + \delta \hat{\pi},\end{aligned}$$

$$\alpha(t), \beta(t), \gamma(t), \delta(t)$$

canonical: $\alpha\delta - \beta\gamma = 1$

$$F(q, \theta, t) = -\frac{1}{\beta} \left(\theta q - \frac{\delta}{2} q^2 - \frac{\alpha}{2} \theta^2 \right)$$

Change in Hamiltonian for an Oscillator

$$H_- = -\frac{1}{2} (p^2 + \omega^2 q^2) \Rightarrow H_n = \frac{1}{2} \pi^2 A + \pi \theta C + \frac{1}{2} \theta^2 B$$

$$A = \beta^2 \left(\frac{d}{dt} \left(\frac{\delta}{\beta} \right) - \left(\frac{\delta}{\beta} \right)^2 - \omega^2 \right)$$

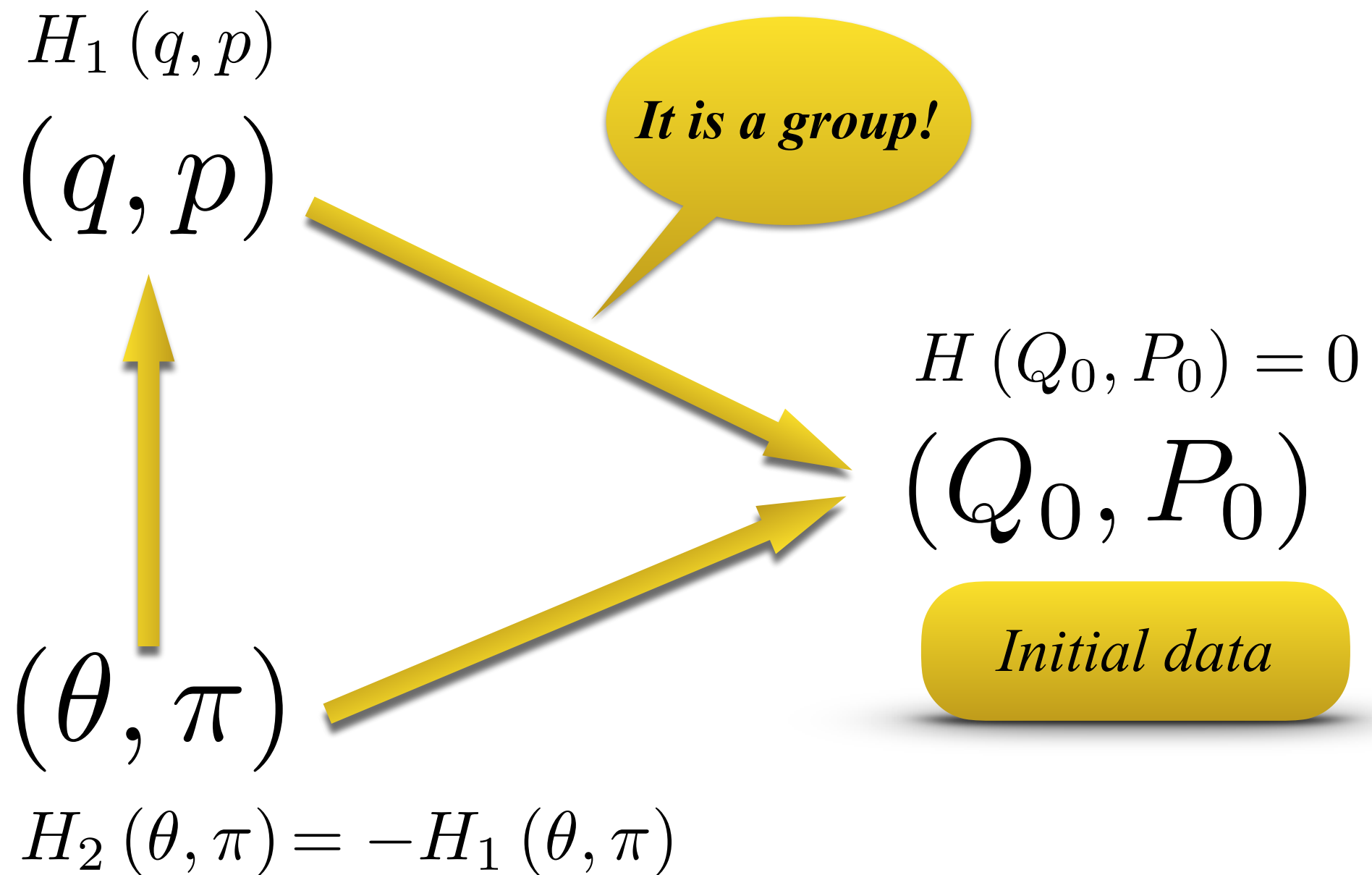
$$C = \frac{1}{\beta} \left[\frac{d\beta}{dt} + \delta + \alpha A \right]$$

$$B = \alpha^2 \left[\frac{d}{dt} \left(\frac{\gamma}{\alpha} \right) - \left(\frac{\gamma}{\alpha} \right)^2 - \omega^2 \right],$$

***Strong time
dependence***

*Can one always solve these
nonlinear equations and
change the sign of energy?*

Canonical map through zero Hamiltonian



Harmonic Oscillator Example

$$H_- = -\frac{1}{2} (p^2 + \omega^2 q^2) \quad \Rightarrow \quad H_+ = \frac{1}{2} (\pi^2 + \omega^2 \theta^2)$$

$$q(t) = q_0 \cos \omega t - \frac{p_0}{\omega} \sin \omega t$$

$$p(t) = q_0 \omega \sin \omega t + p_0 \cos \omega t,$$

$$p = -\dot{q}$$

$$\theta(t) = q_0 \cos \omega t + \frac{p_0}{\omega} \sin \omega t$$

$$\pi(t) = -q_0 \omega \sin \omega t + p_0 \cos \omega t,$$

$$\pi = \dot{\theta}$$

$$q = \theta \cos 2\omega t - \frac{\pi}{\omega} \sin 2\omega t,$$

$$p = \theta \omega \sin 2\omega t + \pi \cos 2\omega t.$$

$$\{q, p\} = \{\theta, \pi\} = 1$$

Conclusions

- There are many *canonical* variables for cosmological perturbations
- Time-dependent linear canonical, unitary (at least mode by mode) transformations always allow to find such *canonical* variables which are *not ghosty*

Thanks a lot for attention!