

# LSS Probes of Cosmic Acceleration

Matteo Fasiello

Stanford University



MF, Z. Vlah

Phys.Rev. D94 (2016) no.6, 063516

MF, Z. Vlah

ArXiv:1611.00542

MF, Z. Vlah &..

+ In progress

January 28th, 2017, “Testing Gravity 2017”, Vancouver

# Beyond $\Lambda$ CDM

A dynamical mechanism driving acceleration usually comes with additional degree(s) of freedom

## Extra Scalar



$$\mathcal{L} = -\frac{1}{2} Z^{\mu\nu}(\phi, \partial\phi, \partial^2\phi, \dots) \partial_\mu\phi\partial_\nu\phi - V(\phi) + g(\phi)T$$

Screening where GR extremely well-tested, e.g. Solar system

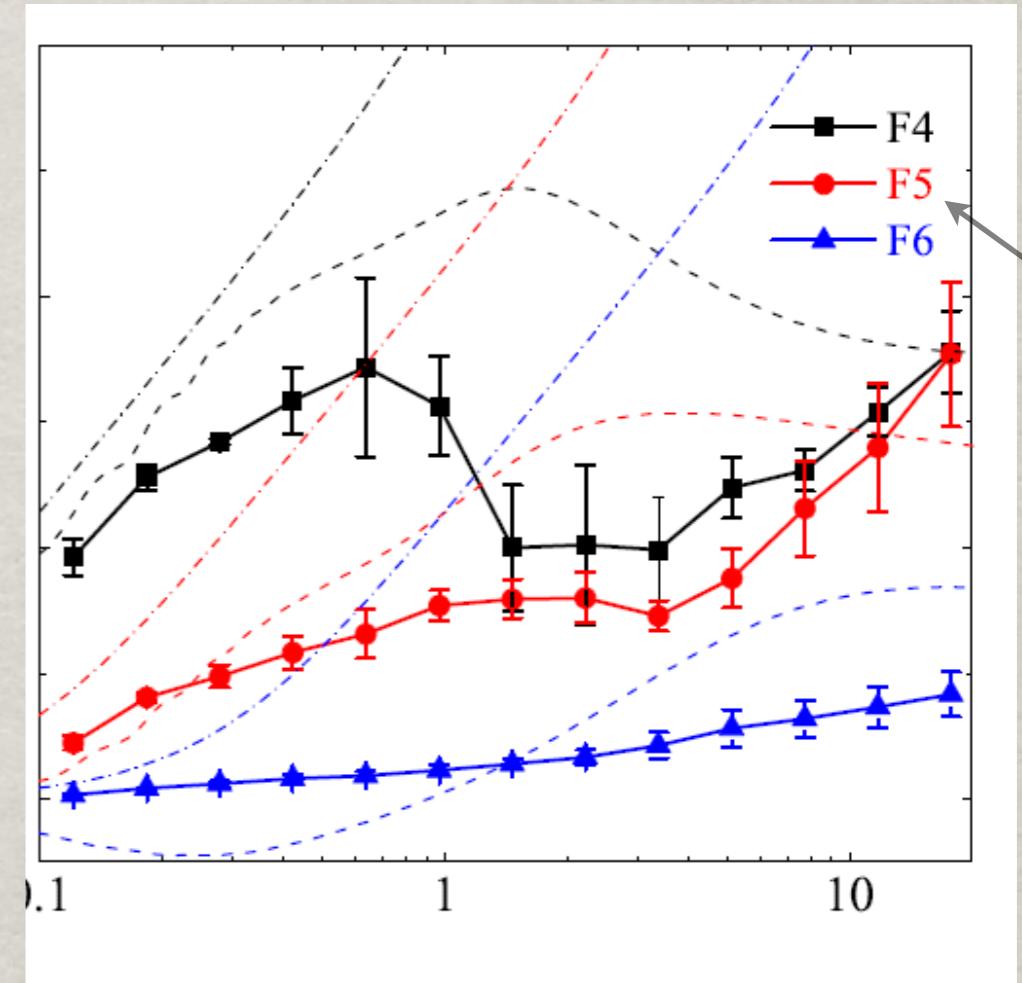
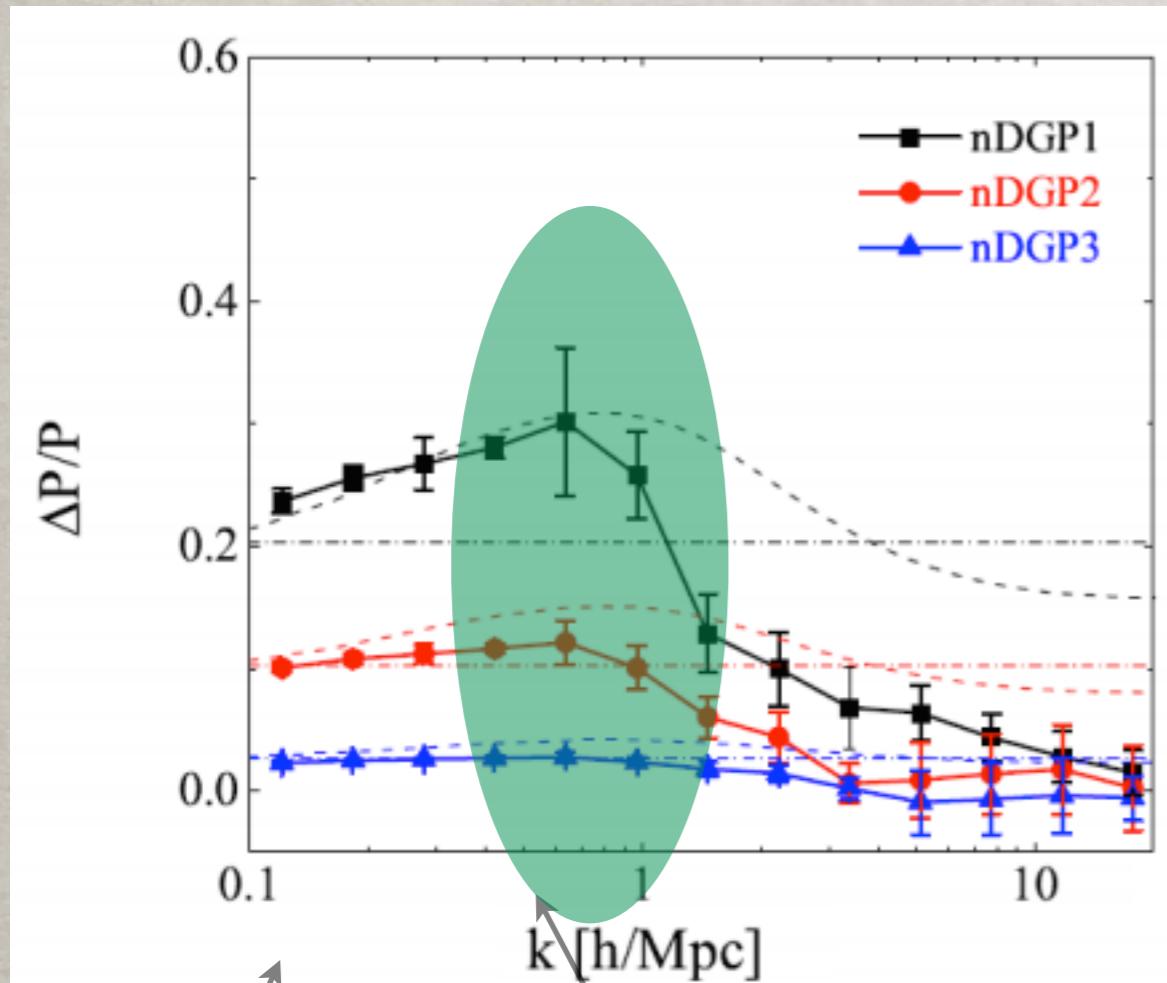
$$V(r) \sim -\frac{g^2(\phi)}{Z(\phi)} \frac{e^{-\frac{m(\phi)}{\sqrt{Z(\phi)}}r}}{4\pi r} \mathcal{M}$$

Symmetron

Vainshtein

Chameleon

# N-body



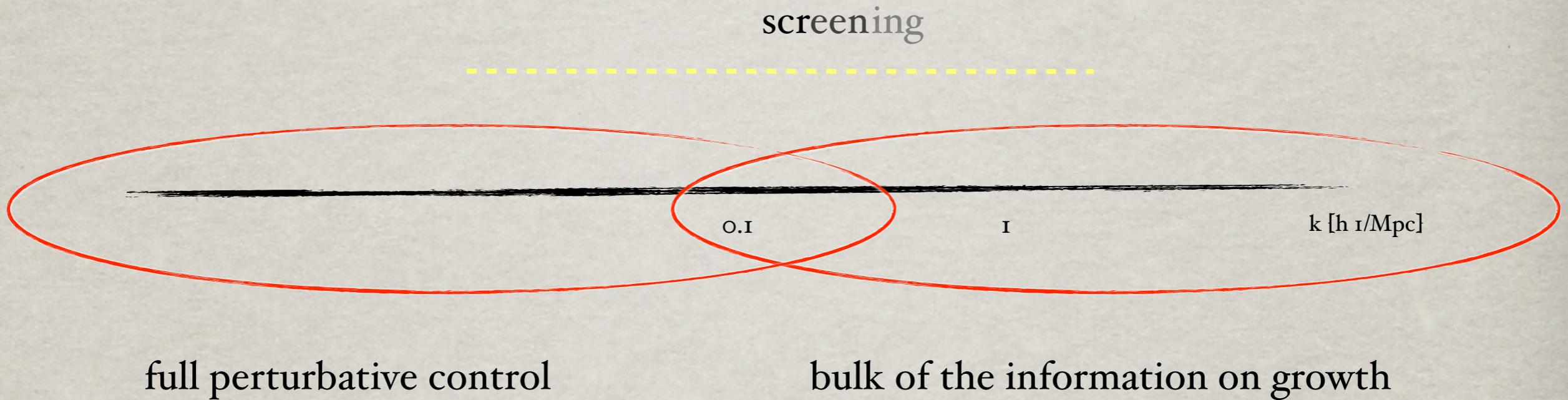
Falck, Koyama, Zhao (2015)

Linear

Vainshtein Onset

$f(R)$

# Perturbation Theory



A lot going on to conquer the quasi-linear scales

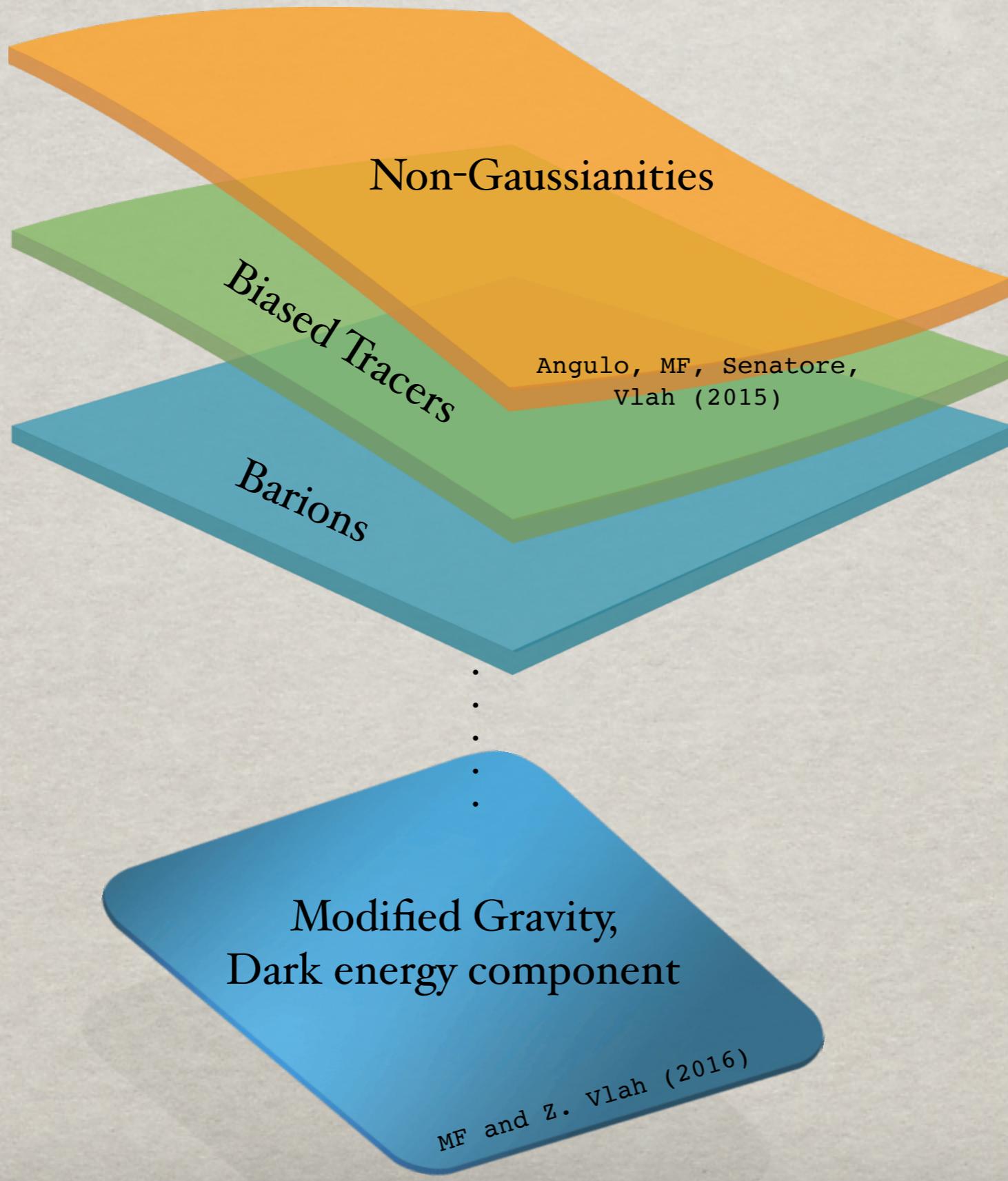
RPT (Crocce, Scoccimarro)  
TRG (Matarrese, Pietroni; Pietroni)  
TSPT(Blas et al)

....  
Lagrangian approach  
(Matsubara; Porto et al; Vlah et al)

...

EFT of LSS

# Layers of physics



# Why Perturbation Theory?

Underlying physical principles

More clear cut map with DE/MG model space

Symmetries (e.g. their role in consistency conditions)

Easy to include additional layers (e.g. non-Gaussianity, baryons..)

# Adding a MG or DDE component

$$\left\{ \begin{array}{l} \frac{\partial \delta_m}{\partial \tau} + \vec{\nabla} \cdot [(1 + \delta_m) \vec{v}] = 0 \\ \frac{\partial \delta_Q}{\partial \tau} - 3\omega \mathcal{H} \delta_Q + \vec{\nabla} \cdot [(1 + \omega + \delta_Q) \vec{v}] = 0 \\ \frac{\partial \vec{v}}{\partial \tau} + \mathcal{H} \vec{v} + (\vec{v} \cdot \vec{\nabla}) \vec{v} = -\nabla \Phi \end{array} \right.$$

Sefusatti, Vernizzi (2011);  
Anselmi et al (2011);  
D'Amico, Sefusatti (2011);

clustering quintessence,  $c_s=0$

$$\nabla^2 \Phi = \frac{3}{2} \mathcal{H}^2 \Omega_m \left( \delta_m + \delta_Q \frac{\Omega_Q}{\Omega_m} \right)$$

$\delta_T$

known exactly only up to quadratic order

# All orders, integral & differential solutions

MF, Vlah (2016)

$$\delta_{\mathbf{k}}(\eta) = \sum_{n=1}^{\infty} F_n^s(\mathbf{q}_1.. \mathbf{q}_n, \eta) D_+^n(\eta) \delta_{\mathbf{q}_1}^{\text{in}} .. \delta_{\mathbf{q}_n}^{\text{in}}$$

$$\Theta_{\mathbf{k}}(\eta) = \sum_{n=1}^{\infty} G_n^s(\mathbf{q}_1.. \mathbf{q}_n, \eta) D_+^n(\eta) \delta_{\mathbf{q}_1}^{\text{in}} .. \delta_{\mathbf{q}_n}^{\text{in}}$$

$$F_n(\eta) = \int_{-\infty}^{\eta} \frac{d\tilde{\eta}}{C(\tilde{\eta})} \left\{ e^{(n-1)(\tilde{\eta}-\eta)} \frac{\tilde{f}_+}{\tilde{f}_+ - \tilde{f}_-} \left[ \left( \tilde{h}_{\beta}^{(n)} - \frac{\tilde{f}_-}{\tilde{f}_+} \tilde{h}_{\alpha}^{(n)} \right) + e^{\tilde{\eta}-\eta} \frac{D_-(\eta)}{\tilde{D}_-(\eta)} \left( \tilde{h}_{\alpha}^{(n)} - \tilde{h}_{\beta}^{(n)} \right) \right] \right\}$$

$$G_n(\eta) = \int_{-\infty}^{\eta} \frac{d\tilde{\eta}}{C(\tilde{\eta})} \left\{ e^{(n-1)(\tilde{\eta}-\eta)} \frac{\tilde{f}_+}{\tilde{f}_+ - \tilde{f}_-} \left[ \left( \tilde{h}_{\beta}^{(n)} - \frac{\tilde{f}_-}{\tilde{f}_+} \tilde{h}_{\alpha}^{(n)} \right) + e^{\tilde{\eta}-\eta} \frac{f_-}{f_+} \frac{D_-(\eta)}{\tilde{D}_-(\eta)} \left( \tilde{h}_{\alpha}^{(n)} - \tilde{h}_{\beta}^{(n)} \right) \right] \right\}$$

$$C = 1 + (1 + \omega) \frac{\Omega_Q(\eta)}{\Omega_m(\eta)}$$

iteratively derived, first recursion are usual

$$\alpha(\mathbf{q}_1, \mathbf{q}_2), \beta(\mathbf{q}_1, \mathbf{q}_2)$$

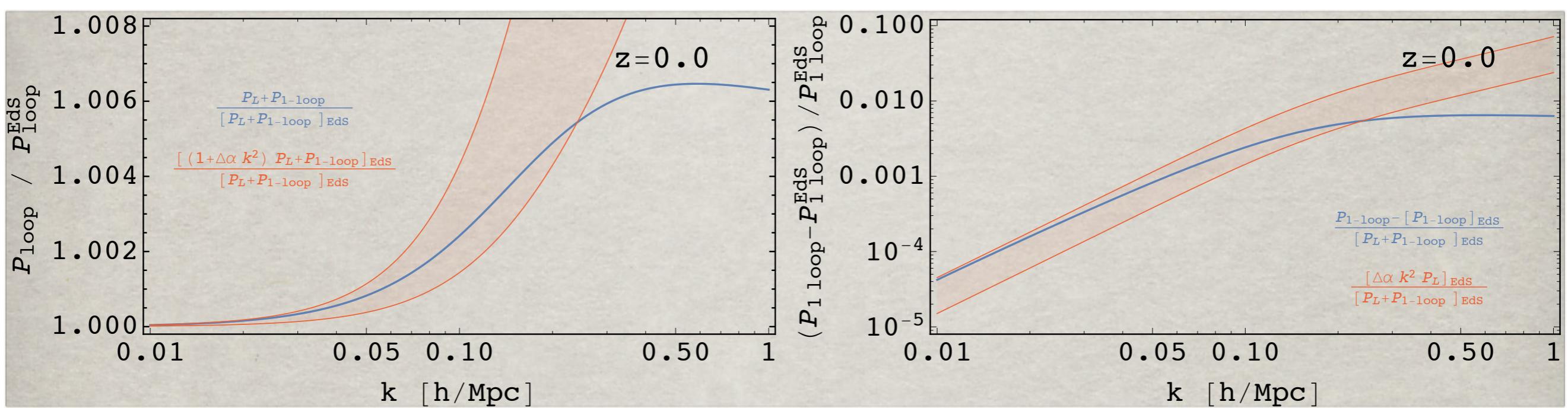
related to

$\propto$  linear growth rate

# Observables

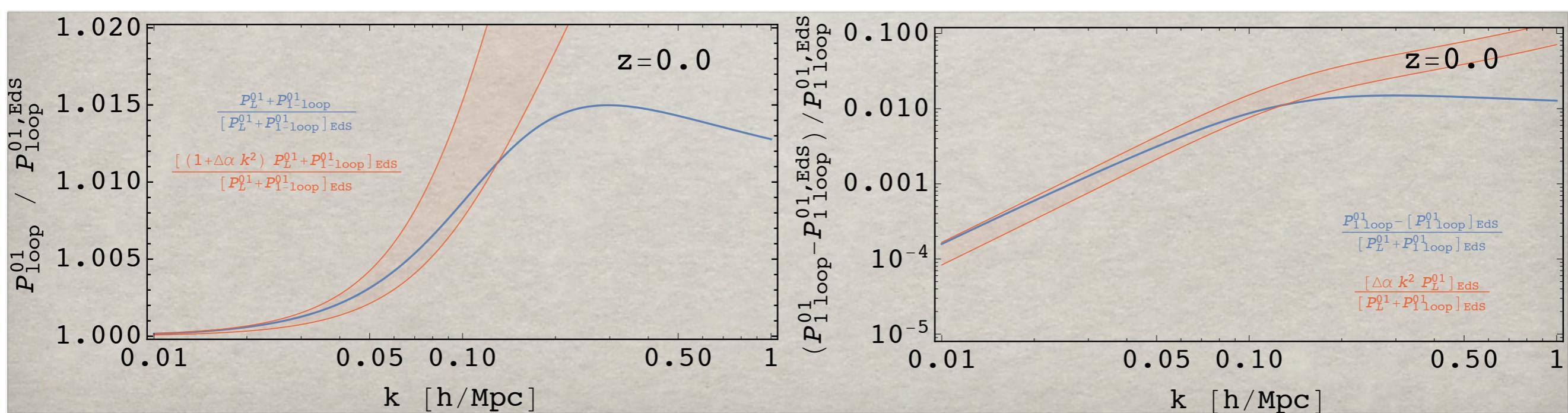
MF, Vlah (2016)

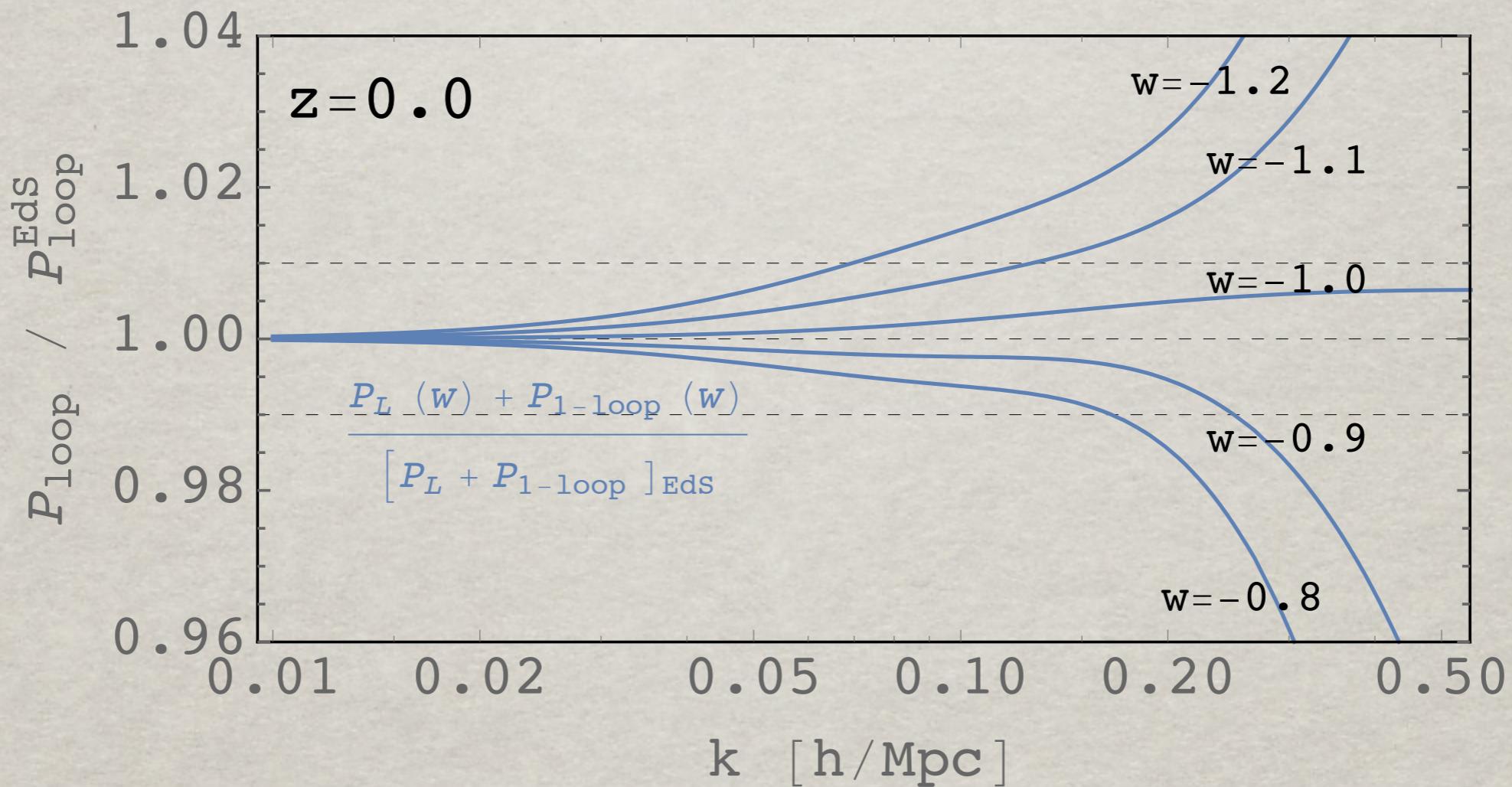
$$P_{1\text{-loop}}(k, a) = P_L(k, a) + P_{22}(k, a) + 2P_{13}(k, a) + P_{\text{c.t.}}(k, a)$$

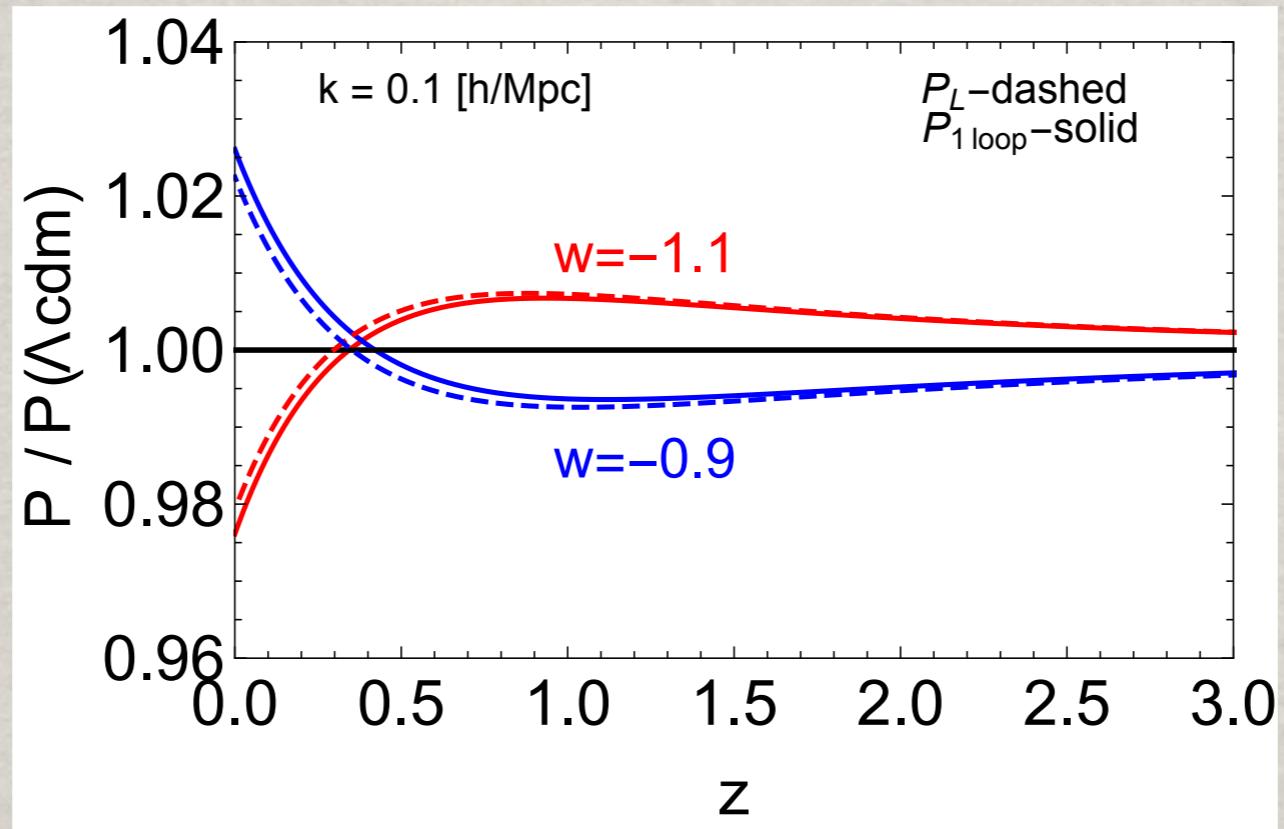


$$C(\eta) = 1$$

test with  $\Lambda$ CDM







All the way to Biased tracers

$$\delta_h(\vec{x}, t) \simeq \int^t H(t') \left[ c_{\delta_T}(t') \frac{\delta_T(\vec{x}_{\text{fl}}, t')}{H(t')^2} + c_{\delta_{\text{d.e.}}}(t') \delta_{\text{d.e.}}(\vec{x}_{\text{fl}}) + c_{\partial v_c}(t') \frac{\partial_i v_c^i(\vec{x}_{\text{fl}}, t')}{H(t')} + c_{\partial v_{\text{d.e.}}}(t') \frac{\partial_i v_{\text{d.e.}}^i(\vec{x}_{\text{fl}}, t')}{H(t')} + c_{\epsilon_c}(t') \epsilon_c(\vec{x}_{\text{fl}}, t') + c_{\epsilon_{\text{d.e.}}}(t') \epsilon_{\text{d.e.}}(\vec{x}_{\text{fl}}, t') + c_{\partial^2 \delta_T}(t') \frac{\partial_{x_{\text{fl}}}^2 \delta_T(\vec{x}_{\text{fl}}, t')}{k_M^2 H(t')^2} + \dots \right].$$

# What's next

MF, z. vlah &..



small deformations  
within CQ,  $c_s \neq 0$

Richer dynamics

Vainshtein-screened  
theories in perturbative setup

new scale  $k_V$  w.r.t.  $k_{NL}$

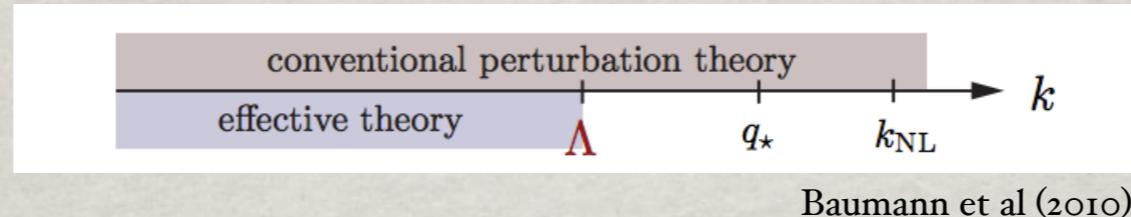
Consistency conditions breaking and squeezed signal enhancement  
typical cases + intriguing exceptions (Galileons)

Thank you!

# Perturbative approaches to LSS

## conquering quasi-linear scales

### EFT of LSS

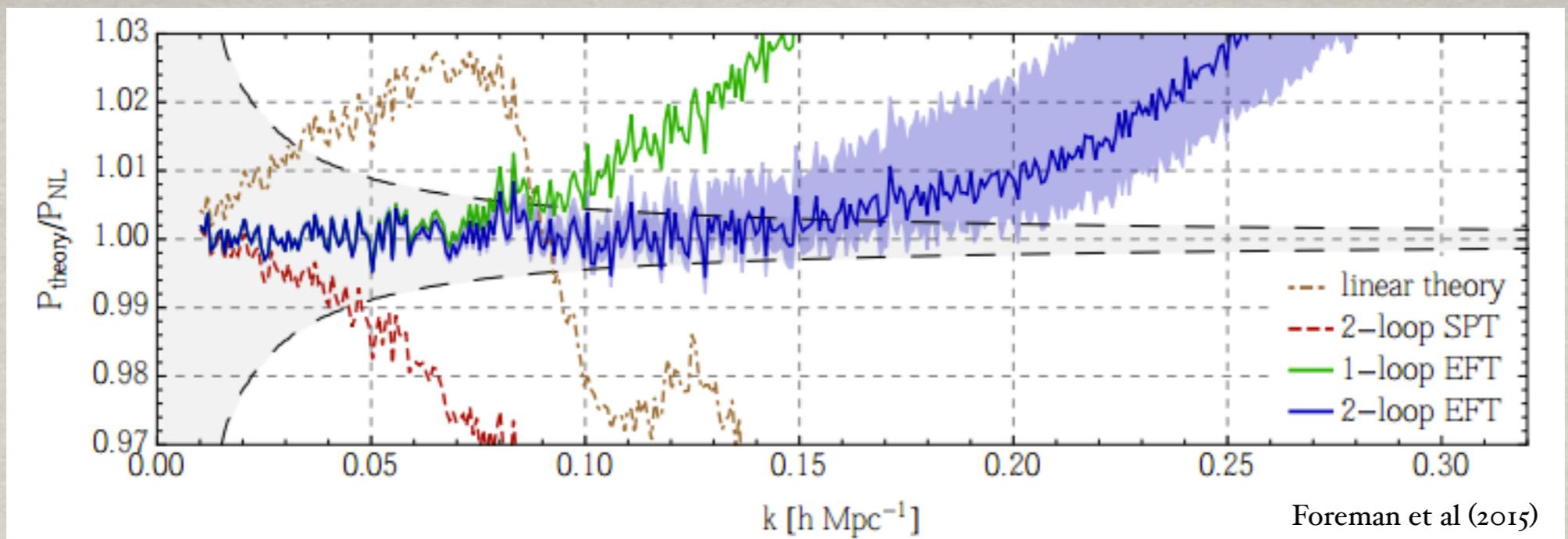


For dark matter (or more), use fluid description

$$\dot{\theta}_\ell + \mathcal{H}\theta_\ell + \frac{3}{2}\Omega_m \mathcal{H}^2 \delta_\ell = -\frac{1}{\rho_\ell} \nabla_i \nabla_j \langle \tau_{ij} \rangle$$

the “EFTness” of the approach is in the fact one describes long-wavelength dynamic informed by a few UV inputs

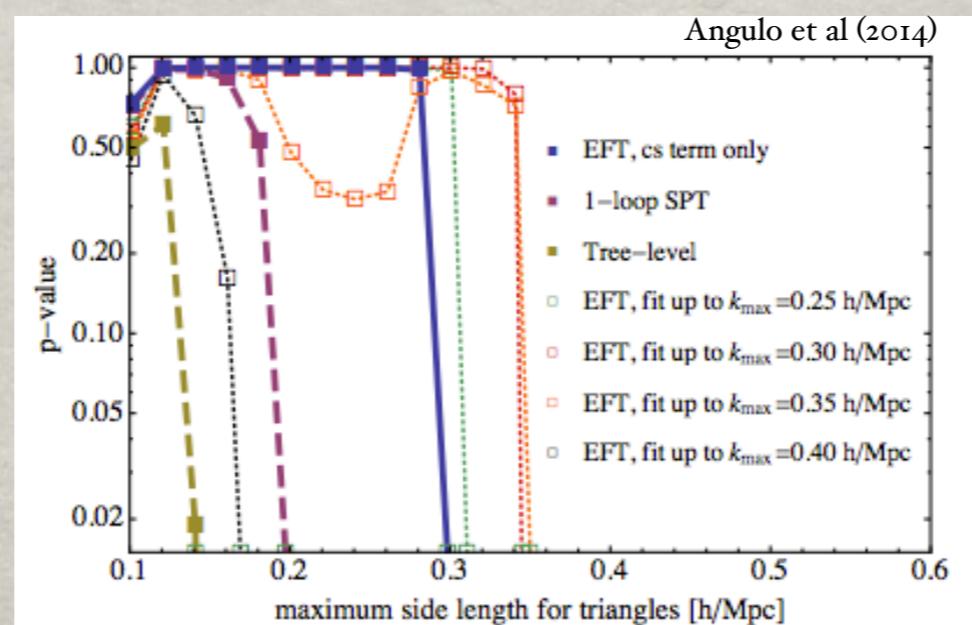
$$\begin{aligned} \langle \tau_{ij} \rangle &= \rho \left[ c_1 \left( \frac{\langle v_s^2 \rangle \partial^2}{\mathcal{H}^2} \right)_{ij} + c_2 \left( \frac{\langle v_s^2 \rangle \partial^2}{\mathcal{H}^2} \right)_{ij}^2 + \dots \right] \phi_\ell + \\ &+ \rho \left[ \left( d_1^{(n)} \left( \frac{\partial^2}{\Lambda^2} \right) + d_2^{(n)} \left( \frac{\partial^2}{\Lambda^2} \right)^2 + \dots \right) \{ v_\ell^2, \delta_\ell \phi_\ell, \dots \} \right]_{ij} \end{aligned}$$



$$P_{\text{EFT-1-loop}}(k, z) = D_1(z)^2 P_{11}(k) + D_1(z)^4 P_{1-\text{loop}}^{\text{usual}}(k) - 2(2\pi)c_{s(1)}^2(z)D_1(z)^2 \frac{k^2}{k_{\text{NL}}^2} P_{11}(k)$$

fit it

Once “Calibrate”  $c_{s(1)}$  on the power spectrum



Move on to other observables