
Consistent Energy Preserving Neural Networks

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Abstract

Unlike vanilla Artificial Neural Networks, Neural Networks embedded with physically-informed priors achieve remarkable results in accurately learning and predicting non-linear dynamical systems. Despite this success, their generalization performance is often limited to short-range trajectories that run for less than the final training time. The performance degradation is due to state and energy drift induced by errors accumulating over time and it poses a critical challenge to learning highly non-linear and chaotic long-range dynamics precisely when data for only a limited number of short-range trajectories are available. One proposed solution to this challenge has been to incorporate symplectic integrators that aim to preserve the symplectic flow. Here, we present an alternative, an L1 penalty of the Hamiltonian derivative with respect to time: $d\mathcal{H}(p, q, t)/dt$. The inclusion of this term, a straightforward addition to the network, enforces the Hamiltonian to remain constant in time for a given trajectory. We empirically illustrate the stabilising effect this penalty has on our network. In addition, we use it to solve problems in non-linear dynamics such as heinon-heiles, 3-body problem and double pendulum.

1 Introduction

- Physicists have long been interested in learning the governing equations of complex dynamical systems.
- Neural networks, as universal function approximators, have shown resounding success across a host of

domains. - Naturally, physicists decided to use neural networks to learn dynamical systems. - However, their performance in learning physical systems has often been limited. - Physicists, as a consequence, have not been encouraged by the initial excitement of machine learning. - However, new research aimed at *scientific machine learning* - a branch that tackles science problems with domain-specific ML, is paving a way to address numerous challenges. - One crucial method has been to incorporate prior theoretical information into the network, such as hamiltonian mechanics. - This excitement has spurred others to work with lagrangians, ODEs and even graphs in order to tackle dynamical systems. - Despite their widespread adoption, Hamiltonian Neural Networks still struggle to predict long-range trajectories precisely. This is in part due to the energy drift that ensues as we accumulate errors over time. - A major topic of discussion to address this issue has been to adopt symplectic integrators that preserve the symplectic flow. - However, we identify an additional component, often neglected when stating Hamilton's equations, that significantly stabilises learning. - We show that the addition of the dh/dt penalty induces the network to conserve energy across time for a given trajectory, a facet not yet invoked by neural networks. - We extensively benchmark this addition across multiple datasets and consistently find the inclusion to be of benefit. - Furthermore, we emphasise that the constraint is an easy plug-and-play addition to existing networks and illustrate how existing networks such as HNN, Symp ODE and Hnets benefit from its inclusion. - Most importantly, we take the HNN model used for Heinon-Heiles and include a symplectic integrator with dhd .

2 Background

2.1 Hamiltonian Neural Networks

Recently, [?] demonstrated that dynamic predictions through time can be improved using Hamiltonian Neural Networks (HNNs) which endow models with a Hamiltonian constraint. The Hamiltonian is an im-

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portant representation of a dynamical system because it is one of two approaches that generalizes classical mechanics. The Hamiltonian \mathcal{H} is a scalar function of position $\mathbf{q} = (q_1, q_2, \dots, q_M)$ and momentum $\mathbf{p} = (p_1, p_2, \dots, p_M)$. In representing physical systems with a Hamiltonian, one can simply extract the time derivatives of the inputs by differentiating the Hamiltonian with respect to its inputs (see Eqn. 1.)

$$\frac{d\mathbf{q}}{dt} = \frac{\partial \mathcal{H}}{\partial \mathbf{p}}, \quad \frac{d\mathbf{p}}{dt} = -\frac{\partial \mathcal{H}}{\partial \mathbf{q}} \quad (1)$$

As a consequence, it is noted in [?] that by accurately learning a Hamiltonian, the system’s dynamics can be naturally extracted through backpropagation. This information allows us to build two 1st-order differential equations which can be used to update the state space, (\mathbf{q}, \mathbf{p}) . Equation 2 shows this integral, in which we define the symplectic gradient $\mathbf{S} = \left[\frac{\partial \mathcal{H}}{\partial \mathbf{p}}, -\frac{\partial \mathcal{H}}{\partial \mathbf{q}} \right]$:

$$(\mathbf{q}, \mathbf{p})_{t+1} = (\mathbf{q}, \mathbf{p})_t + \int_t^{t+1} \mathbf{S}(\mathbf{q}, \mathbf{p}) dt \quad (2)$$

It can be shown that the Hamiltonian in many systems also represents the total energy of the system. Therefore, the Hamiltonian is a powerful inductive bias that can be utilised to evolve a physical state while maintaining energy conservation.

3 Method

The time-derivative of a Hamiltonian $\mathcal{H}(\mathbf{q}, \mathbf{p}, t)$ can be obtained using the chain rule:

$$\frac{d\mathcal{H}}{dt} = \frac{\partial \mathcal{H}}{\partial \mathbf{q}} \frac{\partial \mathbf{q}}{\partial t} + \frac{\partial \mathcal{H}}{\partial \mathbf{p}} \frac{\partial \mathbf{p}}{\partial t} + \frac{\partial \mathcal{H}}{\partial t} \quad (3)$$

The underlying fact in energy preserving systems is that $H = E$ is a constant. As such, the time derivative should be set to zero. By learning the Hamiltonian and differentiating it with respect to \mathbf{q}, \mathbf{p} we obtain the time derivatives of the state vectors which when replaced in the above equation yield:

$$\frac{d\mathcal{H}}{dt} = -\frac{\partial \mathbf{p}}{\partial t} \frac{\partial \mathbf{q}}{\partial t} + \frac{\partial \mathbf{q}}{\partial t} \frac{\partial \mathbf{p}}{\partial t} + \frac{\partial \mathcal{H}}{\partial t} = 0 \quad (4)$$

which leaves us with:

$$\frac{d\mathcal{H}}{dt} = \frac{\partial \mathcal{H}}{\partial t} = 0 \quad (5)$$

To enforce this additional constraint in our networks, all we need to do is simply provide the time to the network, which extends the input dimension by 1, and allow the time component to be differentiable. We

then compute the gradient and use an L1 penalty such that the final loss is:

$$\mathcal{L}_{EPNN} = \left\| \frac{\partial \mathcal{H}_\theta}{\partial \mathbf{q}} + \frac{\partial \mathbf{p}}{\partial t} \right\| + \left\| \frac{\partial \mathcal{H}}{\partial \mathbf{p}} - \frac{\partial \mathbf{q}}{\partial t} \right\| + \left| \frac{\partial \mathcal{H}}{\partial t} \right| \quad (6)$$

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Figure 1: Sample Figure Caption

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Use one line space before the table title, one line space after the table title, and one line space after the table. The table title must be initial caps and each table numbered consecutively.

Table 1: Sample Table Title

PART	DESCRIPTION
Dendrite	Input terminal
Axon	Output terminal
Soma	Cell body (contains cell nucleus)

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If you need to include additional appendices during submission, you can include them in the supplementary material file. You can submit a single file of additional supplementary material which may be either

²Sample of the second footnote.

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```

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of the authors of your paper, all
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Acknowledgements

All acknowledgments go at the end of the paper, including thanks to reviewers who gave useful comments, to colleagues who contributed to the ideas, and to funding agencies and corporate sponsors that provided financial support. To preserve the anonymity, please include acknowledgments *only* in the camera-ready papers.

References

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