

Liouville's Theorem
and the application to
propagation of Cosmic Rays

Stefan Hackstein

stefan.hackstein@hs.uni-hamburg.de

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Liouville's Theorem



derived 1838 by Joseph Liouville (1809 - 1882)

Theorem

The phase-space distribution function $\rho(q_i, p_i; t)$ is constant along trajectory of a conservative system

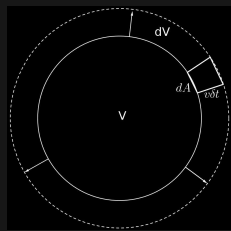
$$\frac{d\rho}{dt} = \frac{\partial \rho}{\partial t} + \sum_i \left(\frac{\partial \rho}{\partial q_i} \dot{q}_i + \frac{\partial \rho}{\partial p_i} \dot{p}_i \right) = 0$$

i. e. the density of system points is constant in the vicinity of a given system point traveling through phase-space.

Liouville's Theorem

Theorem

The Volume enclosed by a closed surface in phase-space is constant as surface moves through phase-space



Proof for Hamiltonian systems:

$$dV = \int_S dA \hat{n} \vec{v} \delta t \Rightarrow \frac{dV}{dt} = \int_S dA \hat{n} \vec{v} = \int_V \nabla \cdot \vec{v} dV$$

$$\vec{v} = (\dot{q}, \dot{p}) \Rightarrow \nabla \cdot \vec{v} = \partial_q \dot{q} + \partial_p \dot{p} = \partial_q (\partial_p H) + \partial_p (-\partial_q H) = 0 = \frac{dV}{dt}$$

Examples

Harmonic Oscillator

Pendulum

Liouville's Theorem

Theorem

Orbits in phase-space are unique.

i. e. trajectories in phase-space cannot cross each other.

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application to Cosmic Rays?

Liouville's Theorem for Cosmic Rays

from Parker 1967: Dynamics of the Magnetosphere:

Consider a static magnetic field B . Density of particles at point s on a field line of B with pitch angle θ at time t : $\Psi(s, \theta, t)$.

"Liouville said that [...] the particle density in phase-space is preserved along the particle trajectory. That is to say the total derivative of the density function Ψ is zero in moving with the particle along the line of force."

$$0 = \frac{d\Psi}{dt} = \frac{\partial\Psi}{\partial t} + \frac{\partial\Psi}{\partial s} \frac{ds}{dt} + \frac{\partial\Psi}{\partial\theta} \frac{d\theta}{dt}$$

with velocity along B -line: $\frac{ds}{dt} = v \sin \theta$ and invariance of magnetic moment $\frac{mv_{\perp}^2}{2B} = \text{const.}$ this differential equation can be written as

$$0 = \frac{\partial\Psi}{\partial t} + v \cos \theta \frac{\partial\Psi}{\partial s} + \frac{v \sin \theta}{2B} \frac{dB}{ds} \frac{\partial\Psi}{\partial\theta}$$

Liouville's Theorem for Cosmic Rays

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assuming small variation in time, $\frac{\partial \Psi}{\partial t} \simeq 0$, the general solution has the form $\Psi(s, \theta) = f\left(\frac{\sin^2 \theta}{B}\right)$ with arbitrary function f . More conveniently one can write

$$\Psi = A \left(\frac{\sin^2 \theta}{B} \right)^\alpha$$

Liouville's Theorem for Cosmic Rays

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Assume: particle density Ψ is isotropic at any point along the trajectory. This demands $\alpha = 0$ and therefore $\Psi = A = \text{const}$ along the trajectory. Also Ψ is no longer a function of B .

Theorem

In presence of a static or slowly changing magnetic field, if the particle density is isotropic at one point in space it is isotropic everywhere.

This implies that **magnetic fields cannot change the flux of cosmic rays measured at earth**. Especially, no anisotropy can be introduced purely by deflection in B

q.e.d

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