Week 3

Kinetic Theory of Gases

3.1 Maxwell-Boltzmann Distribution

1/24: • Applying the molecular partition function to the heat capacity of a water molecule.

- A water molecule has three vibrational modes, which we will denote by ν_1, ν_2, ν_3 (corresponding to symmetric stretch, antisymmetric stretch, and bend).
- Main takeaway: Heat capacity can change with temperature.
- After a while (at several thousand kelvin), it will level off (see Figure 18.7).
- Considers CO₂'s vibrational modes, too.
 - The infrared absorption of the bending mode is what's associated with the Greenhouse Effect.
 - The symmetric stretch is IR inactive due to its lack of change of dipole moment.
 - Raman active: Change in the polarizability of the molecule.
- The Maxwell-Boltzmann distribution.
 - Maxwell derived it long before Boltzmann, but Boltzmann's thermodynamic derivation is much easier.
 - We know from the boltzmann factor that $p(E) \propto e^{-E/k_BT}$.
 - Thus, to get the probability p(v) of some speed v, we should have $p(v) \propto e^{-mv^2/2k_BT}$ times a constant giving the number of molecules of each speed? This yields

$$p(v) = A4\pi v^2 e^{-mv^2/2k_BT}$$

where A is a normalization constant.

- The Maxwell-Boltzmann distribution is such that

$$\begin{split} 1 &= \int_0^\infty p(v) \, \mathrm{d}v \\ &= A \int_0^\infty 4\pi v^2 \mathrm{e}^{-mv^2/2k_B T} \, \mathrm{d}v \\ &= A \int_0^\infty 4\pi \left(\frac{2k_B T}{m}\right)^{3/2} u^2 \mathrm{e}^{-u^2} \, \mathrm{d}u \\ &= A 4\pi \left(\frac{2k_B T}{m}\right)^{3/2} \int_0^\infty u^2 \mathrm{e}^{-u^2} \, \mathrm{d}u \\ &= A 4\pi \left(\frac{2k_B T}{m}\right)^{3/2} \frac{\sqrt{\pi}}{4} \end{split}$$

$$A = \left(\frac{m}{2\pi k_B T}\right)^{3/2}$$

- Therefore,

$$p(v) = 4\pi \left(\frac{m}{2\pi k_B T}\right)^{3/2} v^2 e^{-mv^2/2k_B T}$$

- Any distribution that doesn't look like this isn't in thermal equilibrium.
- A system with all particles having v = 0 is at thermal equilibrium with T = 0 K.
- A system with all particles having constant velocity in the same direction is at thermal equilibrium with $T = 0 \,\mathrm{K}$.
 - Think relativity; if you're moving with them, it looks like they're not moving and thus this case is the same as the last one because you're movement doesn't affect the thermodynamics of that system.
- A system with all particles having constant velocity in different directions is not at thermal equilibrium since it does not fit the bell curve but is rather a spike.

3.2 The First Law of Thermodynamics

1/26: • See Labalme (2021) for background on/content of today's lecture.

- Joule best quantified how we think about work and energy.
- System: Part of the world being investigated. It can contain energy, a number of particles, etc.
- The Newtonian way to change the energy into a system is to do work (mechanical, electrical, etc.) on the system. In chemistry, δw is positive if work is done on the system.
 - We have that

$$\delta w = -P \, \mathrm{d}V$$

- A system at thermal equilibrium has a given temperature characteristic of the system. Some property of the system indicates how hot or cold it is (e.g., volume of mercury, etc.)
- Measure heat transfer using a calorimeter and a thermomenter.
 - Convention: Heat put into a system δq is positive and this heat is transferred if the temperature is lower than another system or surroundings.
 - The heat capacity of a system times the change in temperature is equal to the heat put into the system. It is always positive as heat put into the system raises the temperature.
- Example molar heat capacities.
 - For water vapor at low pressure and 20 °C, $\overline{C}_V = 3R = 25 \,\mathrm{J}\,\mathrm{mol}^{-1}\,\mathrm{K}^{-1}$.
 - For liquid water, it's higher (hydrogen bonding).
 - For ice, it's lower.
- First law of thermodynamics: The internal energy of a system changes with heat put into the system and work done on the system.

$$dU = \delta q + \delta w$$

- Note that in engineering, $dU = \delta q \delta w$.
- State variable: A property that describes the system.

- For example, a system of gas molecules has a state defined by the state variables T, P, V, and n.
- State function: A property that depends only upon the state of the system.
 - For example, some equations of state for an ideal gas are PV = nRT or PV = 2U/3.
 - The internal energy is a state function.
 - Heat and work are not state functions because they do not depend uniquely on the values at equilibrium.
 - They also depend on the way you do something.
- Reversible process: A process that can be represented as a path along state variables, e.g., a line on a PV diagram. This implies that it is also a path where all state variables are known, and is therefore a path where the system is always in quasi-equilibrium.
 - Isothermal, isochoric, isobaric, and adiabatic changes are reversible.
 - All of these processes are analyzed exactly as in Labalme (2021).
- Irreversible process: A process that cannot be drawn on a PV diagram.
- Experiment to measure γ (the ratio of specific heats):
 - 1. Let sit at P_0T_0 .
 - 2. Pump in a little gas (add Δn) and let sit, measure $P_0 + \Delta P_1, T_0$.
 - 3. Open the valve to air quickly to P_0 . Adiabatic expansion (cools down).
 - 4. Let sit to measure the new pressure $P + \Delta P_2$ when T is back at T_0 .
 - 5. γ is determined from $\Delta P_1, \Delta P_2$ (this will be a homework problem).
 - In the second step, we add some molecules into the container. We can show that $\Delta P_1/P_0 = \Delta N_1/n_0$.
 - In the third step, we let out the air, and we can show that $\Delta n_2/n_0 = \gamma \Delta P_1/P_0$.
 - In step 4, we have in the container $(n_0 + \Delta n_1 \Delta n_2)RT_0 = (P_0 + \Delta P_2)V_0$.
 - This implies that $\Delta P_2/\Delta P_1 = 1 1/\gamma$.