Week 6

Enhancing Collision Theory

6.1 Threshold Energy and Line-Of-Centers Model

- 5/2: Picking up from the previous lecture...
 - It follows by plugging in the energy substitutions from last time that

$$u_r f(u_r) du_r = \left(\frac{2}{k_{\rm B}T}\right)^{3/2} \left(\frac{1}{\mu\pi}\right)^{1/2} E_r e^{-E_r/k_{\rm B}T} dE_r$$

- Thus,

$$k = \int_0^\infty du_r f(u_r) k(u_r)$$
$$= \left(\frac{2}{k_B T}\right)^{3/2} \left(\frac{1}{\mu \pi}\right)^{1/2} \int_0^\infty dE_r E_r e^{-E_r/k_B T} \sigma_r(E_r)$$

- Now assume that only those collisions for which the relative kinetic energy exceeds a threshold energy E_0 result in a collision. Thus, define

$$\sigma_r(E_r) = \begin{cases} 0 & E_r < E_0 \\ \pi d_{AB}^2 & E_r \ge E_0 \end{cases}$$

- Consequently,

$$k = \left(\frac{2}{k_{\rm B}T}\right)^{3/2} \left(\frac{1}{\mu\pi}\right)^{1/2} \int_{E_0}^{\infty} dE_r E_r e^{-E_r/k_{\rm B}T} \pi d_{\rm AB}^2$$

$$= \left(\frac{8k_{\rm B}T}{\mu\pi}\right)^{1/2} \pi d_{\rm AB}^2 e^{-E_0/k_{\rm B}T} \left(1 + \frac{E_0}{k_{\rm B}T}\right)$$

$$= \langle u_r \rangle \, \sigma_{\rm AB} e^{-E_0/k_{\rm B}T} \left(1 + \frac{E_0}{k_{\rm B}T}\right)$$

- We can use

$$E_a = k_{\rm B} T^2 \frac{\mathrm{d} \ln k}{\mathrm{d} T}$$

to relate the above to the activation energy.

• Another simplification we've made is that the reaction cross section is not constant, but actually depends on relative speed.

- Accounting for the collision geometry between the two hard spheres gives rise to the line-of-centers model.
- Line-of-centers model: A model for $\sigma_r(E_r)$ in which the cross section depends on the component of the relative kinetic energy that lies along the line that joins the centers of the colliding molecules.

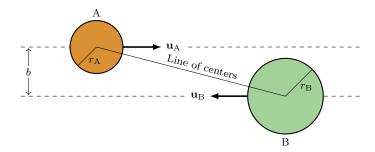


Figure 6.1: Line-of-centers model.

- If we denote the relative kinetic energy along the line of centers by E_{loc} , then we are assuming that a reaction occurs when $E_{loc} > 0$.
- The main thrust of this model is that we are redefining E_r instead of $\sigma_r(E_r)$ overall.
- The line-of-centers model asserts that two molecules will collide only if the **impact parameter** is less than the sum of the radii of the colliding molecules.
 - In particular, we (re)define

$$\sigma_r(E_r) = \begin{cases} 0 & E_r < E_0 \\ \pi d_{AB}^2 \left(1 - \frac{E_0}{E_r}\right) \end{cases}$$

- It follows from math similar to the above that

$$k = \left(\frac{2}{k_{\rm B}T}\right)^{3/2} \left(\frac{1}{\mu\pi}\right)^{1/2} \int_0^{\infty} dE_r E_r e^{-E_r/k_{\rm B}T} \sigma_r(E_r)$$
$$= \left(\frac{8k_{\rm B}T}{\mu\pi}\right)^{1/2} \pi d_{\rm AB}^2 e^{-E_0/k_{\rm B}T}$$
$$= \langle u_r \rangle \sigma_{\rm AB} e^{-E_0/k_{\rm B}T}$$

- Impact parameter: The perpendicular distance between the two dashed lines in Figure 6.1. Denoted by **b**.
- The cross section exhibits a threshold energy.
 - The dependence of the reaction cross section on the relative kinetic energy of the collision is consistent with the line-of-centers model.
- Relating E_0 to the Arrhenius equation parameters.
 - For the activation energy E_a , we have

$$E_a = k_{\rm B} T^2 \frac{\mathrm{d} \ln k}{\mathrm{d}T}$$
$$= k_{\rm B} T^2 \frac{\mathrm{d}}{\mathrm{d}T} \left\{ \ln \left[\left(\frac{8k_{\rm B}T}{\pi \mu} \right)^{1/2} \pi d_{\rm AB}^2 \right] - \frac{E_0}{k_{\rm B}T} \right\}$$

$$= k_{\rm B} T^2 \frac{\mathrm{d}}{\mathrm{d}T} \left\{ \ln T^{1/2} - \frac{E_0}{k_{\rm B}T} + \text{terms not involving } T \right\}$$

$$= E_0 + \frac{1}{2} k_{\rm B} T$$

- Tian wants us to memorize the last line above.
- Considering the line-of-centers collision model and the Arrhenius equation yields

$$A = \langle u_r \rangle \, \sigma_{AB} e^{1/2}$$

• Tian goes through a practice problem.

6.2 Isotropy, Internal Energy, and Center of Mass Assumptions

• Doing away with the assumption that the spheres are isotropic.

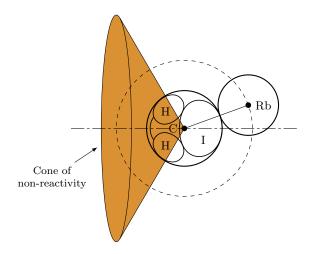


Figure 6.2: Molecules are not isotropic.

- Consider the reaction

$$Rb(g) + CH_3I(g) \Longrightarrow RbI(g) + CH_3(g)$$

- The rubidium atom must collide with the iodomethane in the vicinity of the iodine atom for a reaction to occur.
- Indeed, many molecules have a **cone of non-reactivity**.
- Additionally, the internal energy of the reactants can affect the cross section of a reaction.
 - Consider the reaction

$$H_2^+(g) + He(g) \Longrightarrow HeH^+(g) + H(g)$$

- As the reactant molecule ${\rm H_2}^+$ passes through different vibrational states, its reaction cross section changes.
- We only need to understand that other types of energy can have an effect qualitatively; we do not need to work with the shape of the curves quantitatively.
- A reactive collision can be described in a center-of-mass coordinate system.
 - Consider the collision and subsequent scattering process for the bimolecular reaction

$$A(g) + B(g) \Longrightarrow C(g) + D(g)$$

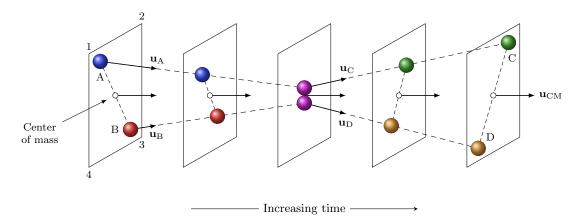


Figure 6.3: Center-of-mass coordinate system.

- Before the collision, A and B are traveling with velocities \mathbf{u}_{A} and \mathbf{u}_{B} , respectively.
- The collision generates molecules C and D, which then move away from each other with velocities \mathbf{u}_{C} and \mathbf{u}_{D} , respectively.
- $-\mathbf{R}$, the location of the center of mass, is given by

$$\mathbf{R} = \frac{m_{\mathrm{A}}\mathbf{r}_{\mathrm{A}} + m_{\mathrm{B}}\mathbf{r}_{\mathrm{B}}}{M} \qquad M = m_{\mathrm{A}} + m_{\mathrm{B}}$$

– The velocity $\mathbf{u}_{\rm cm}$ of the center of mass is the time derivative of the position vector. Therefore, it is given by

$$\mathbf{u}_{\rm cm} = \frac{m_{\rm A}\mathbf{u}_{\rm A} + m_{\rm B}\mathbf{u}_{\rm B}}{M}$$

- We assume that this is an elastic collision (thus, energy is conserved).
- The total kinetic energy is given by

$$\mathrm{KE}_{\mathrm{react}} = \frac{1}{2} m_{\mathrm{A}} u_{\mathrm{A}}^2 + \frac{1}{2} m_{\mathrm{B}} u_{\mathrm{B}}^2$$

- Combining the fact that the relative speed of the two molecules is given by $\mathbf{u}_{r} = \mathbf{u}_{A} - \mathbf{u}_{B}$ with the definition of \mathbf{u}_{cm} yields

$$\mathbf{u}_{\mathrm{A}} = \mathbf{u}_{\mathrm{cm}} + \frac{m_{\mathrm{B}}}{M} \mathbf{u}_{\mathrm{r}}$$
 $\mathbf{u}_{\mathrm{B}} = \mathbf{u}_{\mathrm{cm}} - \frac{m_{\mathrm{A}}}{M} \mathbf{u}_{\mathrm{r}}$

- Note that the change in plus to minus sign between the two above forms hails from our definition of relative speed as A minus B and not the other way around (as we could also very well define it). In other words, it's just a convention thing, and all that matters is that we're consistent.
- It follows that

$$\begin{aligned} \mathrm{KE}_{\mathrm{react}} &= \frac{m_{\mathrm{A}}}{2} \left(\mathbf{u}_{\mathrm{cm}} + \frac{m_{\mathrm{B}}}{M} \mathbf{u}_{\mathrm{r}} \right)^2 + \frac{m_{\mathrm{B}}}{2} \left(\mathbf{u}_{\mathrm{cm}} - \frac{m_{\mathrm{A}}}{M} \mathbf{u}_{\mathrm{r}} \right)^2 \\ &= \frac{1}{2} M u_{\mathrm{cm}}^2 + \frac{1}{2} \mu u_{\mathrm{r}}^2 \end{aligned}$$

- Thus, the kinetic energy is composed of two contributions: one due to the motion of the center of mass, and one due to the relative motion of the two colliding molecules.
- We can do a similar analysis for the products to determine that

$$KE_{prod} = \frac{1}{2}Mu_{cm}^2 + \frac{1}{2}\mu'u_r'^2$$

- Note that momentum is conserved, i.e.,

$$m_{\rm A}\mathbf{u}_{\rm A} + m_{\rm B}\mathbf{u}_{\rm B} = m_{\rm C}\mathbf{u}_{\rm C} + m_{\rm D}\mathbf{u}_{\rm D}$$

- This implies that $u_{\rm cm}$ does not change from reactants to products.
- The energy associated with the motion of the center of mass is therefore constant, and we will ignore its constant contribution to the total kinetic energy.

$$E_{\text{react,int}} + \frac{1}{2}\mu u_{\text{r}}^2 = E_{\text{prod,int}} + \frac{1}{2}\mu' u_{\text{r}}^{\prime 2}$$

- $E_{\text{react,int}}$ and $E_{\text{prod,int}}$ are the total internal energies of the reactants and products, respectively.
- This internal energy takes into account all the degrees of freedom other than translation.