

Week 8

Transition Metal Electronic Structure

8.1 Crystal Field Theory and Ligand Field Theory

11/14:

- Questions from the video lectures.
- For descent in symmetry questions, how do we qualitatively describe orbital energy? Do the Mulliken symbols matter?
 - Rationalize based on the molecular orbital parentage.
 - a_1 may or may not be more stable than a_{1g} , for example.
 - We need to visualize the orbitals and see if we're ripping something off from the axial or equatorial positions, for example, because that matters.
 - The labels are just there to make sure that our molecular orbital assignments are correct within the point group. It's mainly a bookkeeping thing. It also helps us create further MOs based on symmetry, though.
 - The Mulliken symbol alone doesn't convey any information about energetics.
- Walsh diagram structure.
 - Curved vs. straight lines is irrelevant; we just need to know whether the energy increases or decreases qualitatively under a procedural change in the molecule.
 - All we need to be able to do is derive the initial and final state and know what maps to what.
- How do you determine straight vs. curved line?
 - No way to do this from first principles; you need a combination of DFT and experiment.
- Comparing orbital energies.
 - If we're deriving from a new set of orbitals (as in the FeBz example), we need SALCs.
 - In an exam setting, though, we won't encounter brand new orbitals.
- We now start in on today's lecture topic.
- **CFT** and **LFT** are frameworks that help us estimate orbital energies.
 - They are quantitative methods for evaluating d -orbital splitting.
- **Crystal field theory:** A method in which one estimates ligands as point charges (which is not great, but ok) and the metal as a cation C. *Also known as CFT.*
 - Established by physicists at the same time as the birth of inorganic chemistry, before chemists came on the scene.

- Ligands are structure-less and orbital-less (think like a physicist; no structure, just charge).
 - Group theory tells us that a d -orbital in an O_h environment has t_{2g} or e_g symmetry regardless of how much structure is present.
 - Even though the transition metal C is positive, the electrons in the d -orbitals are negative, so we get the following visualization.
 - Visualization: In the free ion, you have five degenerate d -orbitals. You then first approximate a negatively charged spherical field around the point charge, which causes the d -orbitals to go-up in energy. Lastly, we split the t_{2g} and e_g set in an O_h field. Essentially, Wuttig redraws Figure VI.11 of Labalme (2022).
- **Crystal field splitting parameter:** The difference in energy between the e_g and t_{2g} set. Denoted by Δ_0 , $10Dq$. Also known as **field strength, field splitting**.
 - $10Dq$ nomenclature will return when we derive Tanabe-Sugano diagrams from scratch.
 - This is a *quantitative* term for the field splitting, different from every other aspect of the course this far (all of which have been qualitative).
 - This is phenomenological parameter that is determined from experiment and cannot be calculated from first principles.
 - The typical experiment is something like UV-Vis spectroscopy.
 - The terms field strength and field splitting are basically synonymous. However, we can also think of “field strength [as] the quantitative number associated with the field splitting.”
 - Main takeaway of CFT: The origin of something called field splitting, and the origin of quantitatively understanding d -orbital energetics.
 - We now move on from CFT to LFT.
 - **Ligand field theory:** The same as CFT; we just put the bonds back in.
 - It's not just point charges any more; there are actual bonds created by overlapping the metal d -orbitals (which we had also previously assumed to exist) with the ligand orbitals (which we are only adding in now).
 - Thus, free d electrons have adjustable interelectronic repulsion parameters. These are important in free ion terms.
 - Field splitting.
 - Origins of ligand field splitting: Molecular symmetry and group theory, as discussed since the beginning of the course.
 - Extent of ligand field splitting: We will talk about this now.
 - Manifestations of ligand field splitting: The colors of TM complexes, reactivity, magnetism, life (bioinorganic chemistry), and many more.
 - There are four main things to consider when determining the magnitude of the field strength.
 1. σ and π effects.
 - We talked a bit about this when we derived the ML_6 MO diagrams from scratch with and without π orbitals.
 - Recall the discussion regarding Figures 6.13, 7.3, and 7.5.
 - We want to compare a π -donor, to a σ -only ligand, to a π -acceptor.
 - In every case, the e_g set will be comprised of $M-L\sigma^*$ orbitals.
 - However, the t_{2g} set will decrease in energy from $M-L\pi^*$ to nb to $M-L\pi$, respectively.
 - It follows that Δ_0 increases as we go from π -donors, to σ -only donors, to π -acceptors.

- People made this observation by looking at a series of homoleptic O_h complexes, specifically the **spectrochemical series**.
- 2. Oxidation state.
 - As oxidation state increases (the metal ion becomes more positive), ligand field splitting increases, too.
 - This is because orbital overlap increases as ligands are pulled in and hence the amount of ligand field splitting increases, too.
 - Mathematically, increasing the oxidation state increases the overlap integral S .
 - Alternatively, we can think of this as more electropositive transition metals drawing more electrons toward them.
 - Wuttig illustrates this point with example Δ_o values for compounds that are alike in every way except for the oxidation state. For instance,

$$\begin{aligned}\Delta_o(\text{Cr}(\text{H}_2\text{O})_6^{2+}) &= 14\,100 \text{ cm}^{-1} \\ \Delta_o(\text{Cr}(\text{H}_2\text{O})_6^{3+}) &= 17\,000 \text{ cm}^{-1}\end{aligned}$$

- 3. Row of the TM on the periodic table.

- The first row has a lower field strength than the second row, which has a similar field strength to the third row.
- This is because the radial extension of second and third row metals is much higher, leading to better overlap (higher S), and thus higher splitting of bonding/anti-bonding orbitals.
- Example: In the following d^6 complexes, we can observe a drastic increase in Δ_o between the first- and second-row TM and a smaller increase in Δ_o between the second- and third-row TM.

$$\begin{aligned}\Delta_o(\text{Co}(\text{NH}_3)_6^{3+}) &= 22\,870 \text{ cm}^{-1} \\ \Delta_o(\text{Rh}(\text{NH}_3)_6^{3+}) &= 34\,100 \text{ cm}^{-1} \\ \Delta_o(\text{Ir}(\text{NH}_3)_6^{3+}) &= 41\,200 \text{ cm}^{-1}\end{aligned}$$

- 4. Geometry.

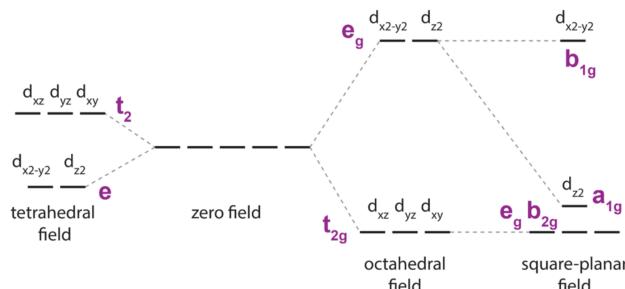
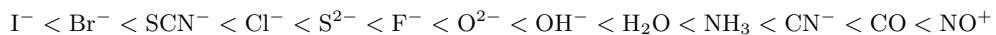


Figure 8.1: Geometry affects field strength.

- The number of ligands around the metal, as well as whether or not certain ligands are on-axis, affects the number of interactions and hence the splitting.
- In a T_d field for example, there are far fewer on-axis repulsive interactions compared to O_h .
- **Spectrochemical series:** A measure of the orbital overlap, electronegativity, and VOLE considerations that determine how big the ligand field splitting is.
 - Some ligands engender higher splitting energies than others, and we can measure this via the electronic transitions of (light absorbed/emitted by) homoleptic complexes.



————— π-donor ————— σ-only ————— π-acceptor —————

Figure 8.2: The spectrochemical series.

- The spectrochemical series was initially just a qualitative observation; it was later explained by MO theory.
- Example: Rationalize the following energy trends

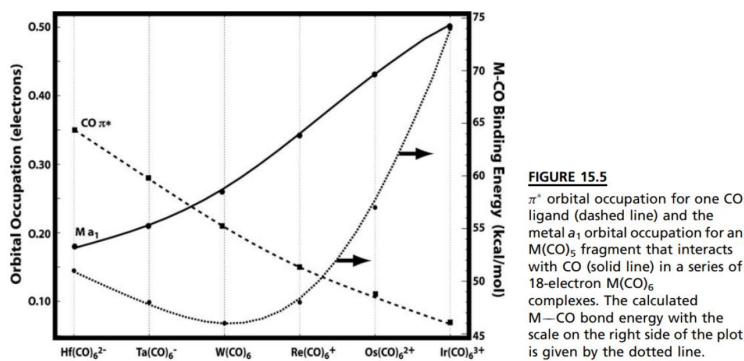


Figure 8.3: Energy trends across a row.

- Trends: The M–CO binding energy increases left to right. The orbital occupation of the $\text{CO}(\pi^*)$ decreases left to right. Regardless, the orbital occupation of the $M(a_1)$ (primarily of d_{z^2} parentage) increases left to right.
- The analytical procedure is as follows.
 - Investigate the ligand: CO is a π -acceptor; thus, the d_{z^2} orbital is of $M-\text{L}\sigma^*$ parentage. Recall from the PES of CO that its LUMO (π^* symmetry) lies higher in energy than the d atomic orbitals, which are sandwiched between that and the HOMO (σ symmetry).
 - Investigate the metal center: All molecules are third-row d^6 transition metals.
 - Investigate what changes: The charge (e.g., Ir^{3+} vs. Hf^{2-}) and the electronegativity of the metal center (decreases left to right; does she mean the amount of negative charge?? Seems like she's been using electronegativity and electropositivity potentially wrong a few times??).
 - On the left side, the metal d orbitals are high in energy, so they are close to the high-lying $\text{CO}(\pi^*)$ orbitals in energy, so they donate a lot of electron density to $\text{CO}(\pi^*)$. On the right side, the opposite is true.
 - Additionally, there is a smaller gap between d_{z^2} and $\text{CO}(\sigma)$ on the right side.
- **Electronic spectroscopy:** Electron absorption spectra are broad — occur much more rapidly than molecular vibrations, i.e., we see a snapshot of the molecule in various rotational vibrational states.
- We're done learning about trends and onto quantifying stuff. To understand how to quantify things better, we need to understand the origins of electronic spectra.
- Thus, we'll discuss some basic properties of UV-Vis spectra in relation to the Franck-Condon principle.
- In the first half of the course, we focused on vibration. Now we're looking at electronic excitations and, in particular, electronic excitation wells.

- Franck-Condon Principle.

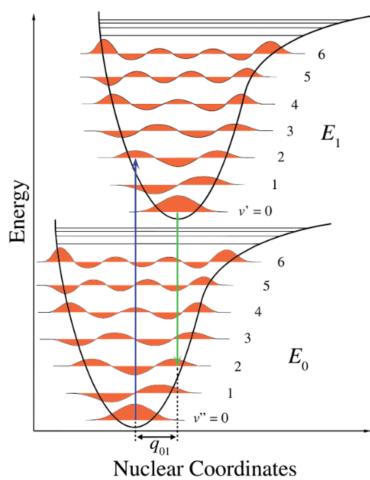


Figure 8.4: Franck-Condon Principle potential wells.

- E_0 is our ground electronic state.
- E_1 is the first electronic excited state.
- All of the little levels in between are the vibrational energy levels (overtones, etc.) that we discussed in the beginning.
- We now assume that an electronic transition is much faster than a vibrational transition, and hence occurs virtually without a change in nuclear coordinates. We just do a vertical jump from one well to the other, without moving nuclear coordinates along a vibration (“electronic transitions are ‘vertical’ with respect to the nuclear coordinate axis”).
- The Franck-Condon principle arises from the Born-Oppenheimer approximation.
- **Born-Oppenheimer approximation:** We can separate the electronic and vibrational wave functions of a given molecule.
- Terms next time.

8.2 Electronic Spectroscopy II

11/16:

- No in-person lecture on Friday; a video will be posted. Watch before Monday’s lecture (last lecture of the class). Wednesday is the review session, Friday is the second exam, no final.
- Last problem set has been posted; due the review session day. Submit virtually (as a PDF).
- Goal: Understand and derive term symbols in the context of electron configuration. Free ions \rightarrow weak field \rightarrow strong field. It’s important to understand where these things come from because that allows us to use the tools in the literature in a more informed fashion.
- Terms allow us to predict the energy of the electronic transition. But where do they come from? How do they vary for ground and excited states?
- **Term:** Energy level of a spectrum.
- Free ion term: d -orbital manifold in the absence of any other geometry. We have 5 d orbitals and, suppose, a d^2 system. We can fill electrons however we want as long as we satisfy the Pauli exclusion principle. But different fillings have different energies. *picture*

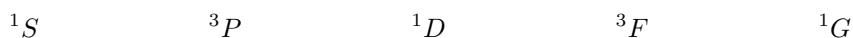
- Multi-electron transition metal ions have multiple energy levels even when *d*-orbital degeneracy is not lifted by a ligand field.
- **Configuration:** The assignment of a given set of electrons to a certain set of orbitals.
 - Applies to *d* orbital setups that are *not* free ion, i.e., degeneracy has been lifted.
 - Each configuration gives rise to a number of energy levels and the number of energy levels have to have an even number of terms.
- Example:
 - 2 electrons in *d* orbitals. (t_{2g}^2) , for instance, implies that we have 2 electrons in 2 t_{2g} orbitals.
 - 4 electrons in *d* orbitals. t_{2g}^4 , for instance, implies that all electrons are in the t_{2g} set. $(t_{2g}^2)(t_{2g}^2)$, for instance, implies that electrons are in different orbitals; this describes an excited state.
- Energy ordering of the terms.
 - To construct suitable energy level diagrams for various states in a transition metal ion, we consider the last three terms (IER – interelectronic repulsion, SOC – spin-orbit coupling, and LFI – ligand field interaction).
 - The interaction between these three terms determines whether we're in a weak or strong field.
 - Case 1: IER > LFI > SOC gives rise to the weak field approximation (ligand field induces only a small perturbation of the energy levels).
 - Case 2: LFI > IER > SOC gives rise to the strong field approximation (implies more low spin configurations).
- We've now started understanding the underlying principles that change the energy ordering.
- Consider only weak field approximation. We first approximate the weakest field possible: Where there is no field. World's most boring electronic structure, but helps us understand how to build correlation diagrams. Let's determine the number of microstates or the ways we can put a single electron into *d* orbitals.
 - d^1 case: No effects from other electrons, yet there are still numerous microstates.
 - A *d* orbital can be described by $\ell = 2$ (the angular momentum quantum number), $m_\ell \in \{-2, -1, 0, 1, 2\}$ (the magnetic quantum number; takes integer values that span $-\ell$ to ℓ), and $m_s = \pm 1/2$ (the spin quantum number).
 - Thus, there are 5 degenerate *d* orbitals (one for each m_ℓ ; these labels can be thought of as alternatives to $z^2, x^2 - y^2, \dots$).
 - The one electron can occupy any *d* orbital and take spin $\pm 1/2$ for 10 total degenerate sets of microstates.
 - We can designate the collection of microstates with a term symbol.
- **Term symbol:** An abbreviation describing a selection of microstates.
- How do I know that I need to get 10 microstates?
 - The number ℓ^q of microstates for a given configuration, where ℓ is the angular momentum quantum number and q is the number of electrons, can be calculated for a number of equivalent electrons. The formula is
$$n = \frac{(4\ell + 2)!}{q!(4\ell + 2 - q)!}^{[1]}$$

¹Wuttig puts $4\ell - 2$ in the denominator.

- If we do this for $\ell = 2$, $q = 1$, we get

$$n = \frac{10!}{9!} = 10$$

- How do I know what the term symbol is?
 - Term symbols have the form $^{2S+1}L_J$.
 - $M_L = \sum m_\ell$ for a given microstate.
 - $M_S = \sum m_s$ for a given microstate.
 - S , the spin multiplicity, is the total spin angular momentum (values $0, 1/2, 1, 3/2, 2, 5/2$) or the maximum M_S value.
 - L , the total orbital angular momentum, is the maximum M_L value.
 - J , the L-S/Russell-Saunders/spin-orbit coupling, is only relevant for multielectron systems where the presence of an electron in the same shell diminishes the coupling of another. J is the total angular momentum. $J = L + S, L + S - 1, L + S - 2, \dots, L - S - 1, L - S$ ^[2].
- Example: Term symbol for d^1 .
 - We must have $S = 1/2$ (the smallest possible value of S , and the only one achievable by one electron) and $L = 2$ (degenerate set of orbitals that span -2 to 2). Thus, 2D .
 - No spin-orbit coupling to worry about because there's only one electron in the system (coupling inherently requires two or more electrons).
- Multielectron atom procedure.
 1. Determine the total number of microstates (m_ℓ, m_s values).
 2. Determine the possible values of $M_L = \sum m_\ell$ and $M_S = \sum m_s$.
 3. Construct a Pauli-allowed chart of microstates.
 4. Group microstates into terms.
 5. Assign ground state term symbol.
- Let's tackle d^2 .
 - By the formula, there are 45 microstates.
 - Goes through the microstate table and rectangles (we won't be asked to derive this, though).
 - The maximum value of L is 4. This means the two electrons are in the same orbital, so we have $S = 0$. Thus, we get 1G for our term symbol. We now slash off all microstates for which $L = 4$ and $S = 0$, i.e., microstates for which $M_L = 0, \pm 1, \pm 2, \pm 3, \pm 4$ and $M_S = 0$.
 - Next up is $L = 3$ for maximum possible L . Maximum possible S in this case is 1. Thus, we slash off all states for which $M_L = 0, \pm 1, \pm 2, \pm 3$ and $M_S = 0, \pm 1$ except the ones we've already counted^[3]. Now that we have $L = 3$, the electrons can be in different orbitals, so $S = 1$.
 - My big question: What do the microstates that we group together into one term symbol have in common??
- Takeaway: We have multiple term symbols/energy levels for free ion states.
- For d^2 , we show in handout that we have 5 terms:



- But how do we order the energy? Use Hund's rules:

²This is probably wrong, too; should iterate down to $|L - S|$.

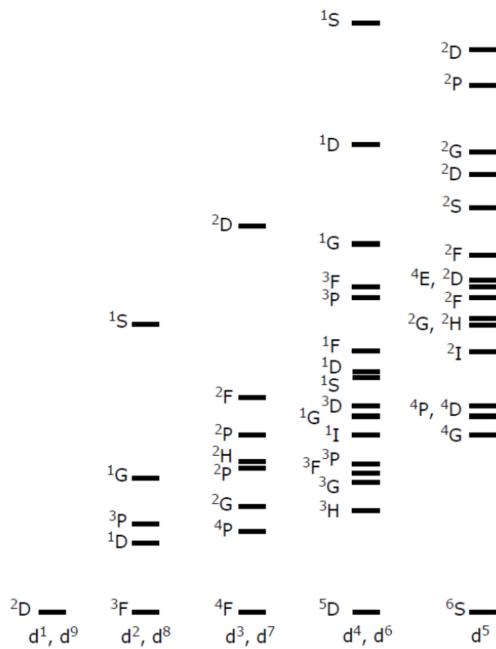
³Note that since there are $7 = 2L + 1$ possible M_L values and $3 = 2S + 1$ possible M_S values, we will have $7 \cdot 3 = 21$ microstates of this type. This is where Talapin's $(2L + 1)(2S + 1)$ formula comes from!

1. Maximize the spin multiplicity (i.e., terms with higher multiplicity will have lower energy).
 2. Maximize the total angular momentum (i.e., terms with higher L will have lower energy).
- Thus, the energy ordering is $^3F, ^3P, ^1G, ^1D, ^1S$.
 - Electron hole formalism: d^n and d^{10-n} have equivalent weak field energy orderings.
 - Will be asked about in our homework.

8.3 Electronic Spectroscopy III and TS Diagrams

11/18:

- Goals:
 - Building Tanabe-Sugano diagrams from scratch.
 - More about term symbols (deriving and ordering them; ground states v. excited states).
- Last time: We derived the term symbols and ordering for d^1, d^9 and d^2, d^8 . These term symbols have different energies; these must be spectroscopically determined.

Figure 8.5: d -orbital weak field splittings.

- The weak field energy orderings are shown without spin-orbit coupling (the J term).
- But recall that S-O coupling can exist.
- Example: 3F S-O coupling for d^2 .
 - 3F means $L = 3$ and $S = 1$.
 - We can split it into $(L + S) - (L - S) + 1 = 2S + 1 = 3$ terms^[4], namely those with $L + S = 4$, $L + S - 1 = 3$, and $L - S = 2$. Thus, we get 3F_4 as the highest energy of these terms, 3F_3 in the middle, and 3F_2 at the bottom.

⁴This is the spin multiplicity, as expected.

- Additional energy ordering rule: If the free ion term is less than half-filled, the lowest J has the lowest energy; if the free ion term is more than half-filled, the highest J has the lowest energy!
- We now start to include a ligand field.
- Splitting of energy levels can be described by the character table.
- Example: Consider the 1D term in the free ion for d^2 in an O_h field.
 - In a similar fashion to the d orbitals in an O_h field, the D states arising from groups of electrons have the same 5-fold degeneracy, so splitting will be the same.
 - Thus, we split 1D into 1E_g and $^1T_{2g}$.
 - Notice that the spin multiplicity is conserved under splitting. This is because the chemical environment does not directly interact with spin, so all the states into which a particular term is split have the same spin multiplicity as the parent term.
- Example: Consider the 3F term in the free ion for d^2 in an O_h field (without S-O coupling).
 - f orbitals transform as the cubic functions, i.e., as a_{2u}, t_{1u}, t_{2u} (see Table 6.8).
 - However, since this is a d^n configuration and all d orbitals have a center of symmetry, we change the u to a g .
 - Thus, the splitting is $^3A_{2g}$, $^3T_{1g}$, and $^3T_{2g}$.
- We can now undertake the full construction of energy level diagrams.
- Continue with the d^2 example.
 - At this point, we know how to get the weak field energy ordering and split each free ion term into its relative weak field splitting. This latter splitting gives the **weak free ion**.
 - Thus, we can now derive the left (weak field) side of the correlation diagram for d^2 in an O_h environment.
- **Weak free ion:** The interaction with the O_h chemical environment by assuming that the electron in the d -orbital manifold is repulsed by the negative charge around it.
 - This is essentially CFT.
- Next task: Strong field terms.
 - Given a large splitting of d^2 in O_h , the following configurations are possible: $t_{2g}^2, t_{2g}e_g, e_g^2$.
 - The latter two are excited states.
 - These are analogous to the initial term symbols, in that these states only exist in an infinitely strong ligand field.
 - As the ligand field weakens...
 1. The electrons will feel each other's presence and start to couple.
 2. The symmetry properties of these states can be determined by taking the direct products of the representation of a single electron.
 - Example:

$$\begin{aligned}t_{2g}^2 &= t_{2g} \times t_{2g} = A_{1g} + E_g + T_{1g} + T_{2g} \\t_{2g} \times e_g &= T_{1g} + T_{2g} \\e_g \times e_g &= A_{1g} + A_{2g} + E_g\end{aligned}$$

- Remember that both explicit computation and a direct product table can lead to the right answer in computations such as these.

- Full correlation diagram.

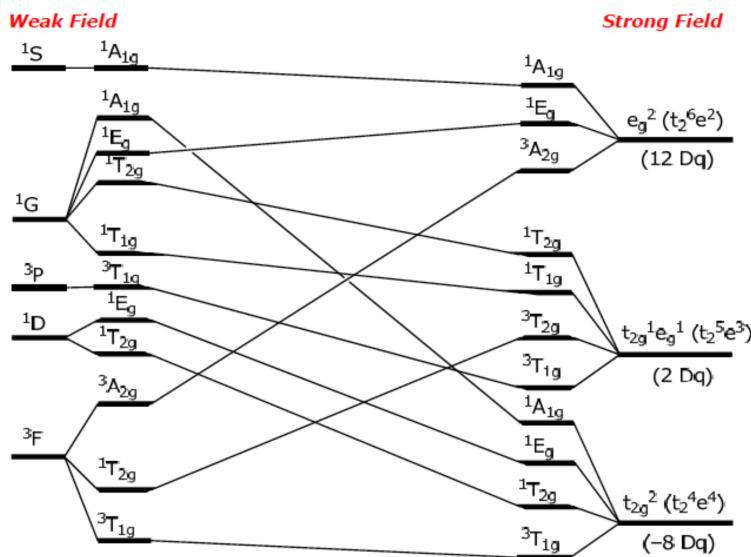


Figure 8.6: Correlation diagram for a d^2 configuration.

- Axioms:
 1. There must exist a 1-to-1 correspondence between the states at the 2 extremes of the abscissa.
 2. As the strength of the interaction changes, states of the same spin degeneracy do not cross.
- The multiplicities of the strong field splitting terms are difficult to determine using the principles of *this* course. Determined by descent in symmetry to reveal population and split degenerate states. Beyond the scope of this course. The multiplicities will always be given to us.
- This concludes the construction of correlation diagrams.
- **Tanabe-Sugano diagrams:** Correlation diagrams that are useful in interpretation of electronic spectra of coordination compounds.
 - Answer the question of the precise scaling of the energies of these states with the ligand field strength.
 - Dreamt up by two physicists in 1954.
- Example: d^1 case (simplest; e.g., Ti^{3+} in an O_h field).

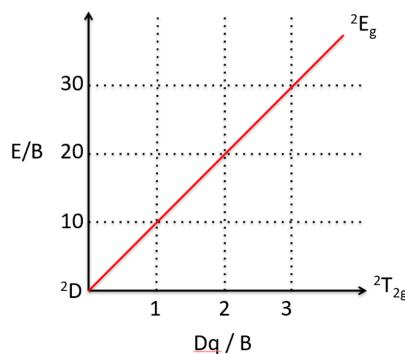


Figure 8.7: Tanabe-Sugano diagram for a d^1 configuration.

- We have one free ion term 2D .
- Our strong field terms are $^2T_{2g}$ and 2E_g (possible one electron configurations are t_{2g}^{-1}, e_g^{-1}).
 - 2E_g is the excited state in the octahedron.
 - $^2T_{2g}$ is the lowest energy on the horizontal, so the vertical distance between the x -axis and the sloping line is a measure of the energy of the excited state above the ground state (and hence is directly proportional to Δ_o).
- E describes the energy of the optical transition.
- B is the **Racah parameter**.
- Interpretation, a d^1 TM complex (such as a Ti^{3+} complex) should have one absorption.
- **Racah parameter:** A measure of the repulsion between terms of the same multiplicity. *Denoted by B . Given by*
$$B = \frac{e^2}{r_{ij}}$$

 - Directly correlated to the degree of electron repulsion (see its definition).
 - Originally defined as a purely free ion parameter, but does change upon coordination (**nephelauxetic effect**).
 - Important trends:
 - Complexes have smaller B than free ions.
 - More covalent M–L bonds implies lower B .

- Example (continued): Does the spectrum for $Ti(OH_2)_6^{3+}$ reflect our intuition that there should only be one absorption?

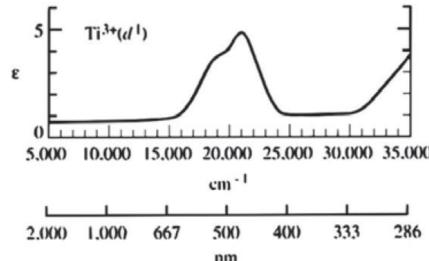


Figure 8.8: Spectrum for $Ti(OH_2)_6^{3+}$.

- The spectrum for the titanium complex does (mostly) reflect this intuition.
- There is a slight shoulder (pH dependent) due to deprotonation of an aqua group resulting in some $Ti(OH_2)_5(OH)^{2+}$ in solution.
- The **molar extinction coefficient** ϵ is very low ($\epsilon \approx 4 M^{-1} cm^{-1}$).
 - The transition is weak because it is Laporte forbidden ($g \rightarrow g$).
 - Selection rules will be covered in greater detail on Monday.
- **Electron hole formalism:** A d^{10-n} configuration will behave in the same way, at all points along the abscissa of the energy level diagram, as the corresponding d^n configuration except that all energies of the interaction with the environment will have the opposite sign.
 - n holes in a d shell can be treated like n positrons. n positrons interact with each other the same way as electrons with each other, but with the environment, they will behave as attractive forces because these positive entities are now interacting with a negative ligand field.

- An application of the electron hole formalism.
 - The TS diagram in Figure 8.7 also treats the d^9 case!
- Example: Does the spectrum for $\text{Cu}(\text{OH}_2)_6^{2+}$ reflect our intuition that there should only be one absorption?

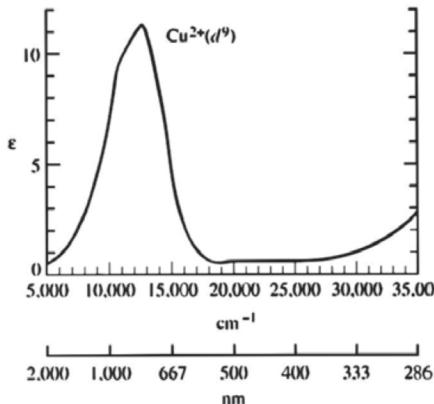


Figure 8.9: Spectrum for $\text{Cu}(\text{OH}_2)_6^{2+}$.

- Mostly yes, and we have the same shoulder again.
- This shoulder can also come from a Jahn-Teller distortion, yielding transitions from the the bottom two d orbitals to the top one. (Why isn't there a transition from the third from the bottom to the top??)
- Note that J-T is not seen in d^1 because the t_{2g} set is primarily nonbonding for the σ -only aquo ligands.
 - Review J-T distortion!!
 - Similarly small ε . Owing to Laporte forbiddenness again.
 - Note that since we're in the hole formalism, the TS diagram labels need to be switched since the hole is moving up.
 - Overall same TS diagram as Figure 8.7, but the ${}^2T_{2g}$ and 2E_g labels are flipped.
- Example: d^2 case (increasingly complex).

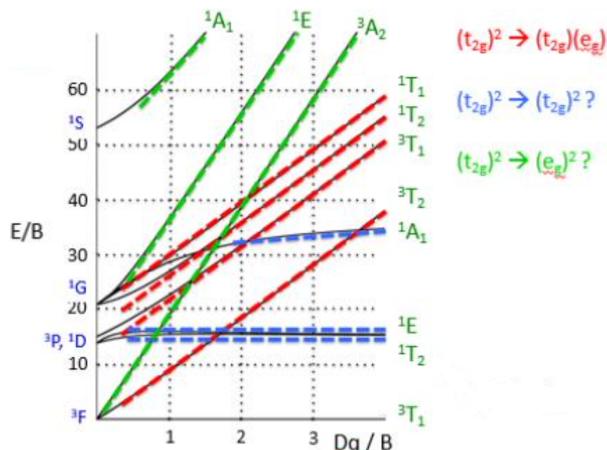


Figure 8.10: Tanabe-Sugano diagram for a d^2 configuration.

- Once again, free ion terms on the left and strong field terms on the right.
 - Note that sometimes, g labels are left out (see the strong field terms).
- Red lines: $t_{2g}^2 \rightarrow t_{2g}e_g$ transitions.
- Blue lines: $t_{2g}^2 \rightarrow t_{2g}^2$ transitions (e.g., spin flip).
- Green lines: $t_{2g}^2 \rightarrow e_g^2$ transitions (two photon transitions have a low probability).
- Curvature: Caused by mixing of states of the same symmetry.
- Spin-allowed transitions include ${}^3T_1 \rightarrow {}^3T_2$ and ${}^3T_1 \rightarrow {}^3T_1$ (both red-line transitions). These should be particularly noticeable.
- Example (continued): Assigning transitions to peaks in a V^{3+} (d^2) complex's spectrum.

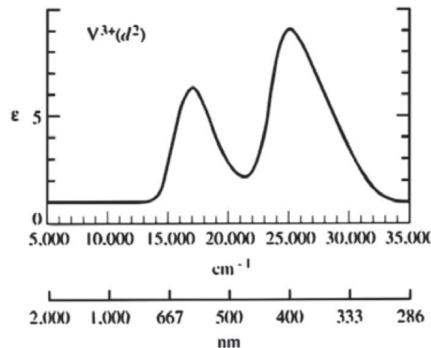


Figure 8.11: Spectrum for a V^{3+} (d^2) complex.

- We assign the lower energy transition (${}^3T_1 \rightarrow {}^3T_2$) to the peak at a lower wavenumber, and the higher energy transition (${}^3T_1 \rightarrow {}^3T_1$) to the peak at the higher wavenumber.
- Example (continued): Estimating B and Dq .

- Take the ratio of the two transitions' wavenumbers in Figure 8.11.

$$\frac{25\,600\,\text{cm}^{-1}}{17\,200\,\text{cm}^{-1}} = 1.49$$

- Looking at the d^2 TS diagram (Figure 8.10), we can see that if we draw a vertical line at $Dq/B = 2.8$, then $E/B = 25.9$ for the lower transition and $E/B = 38.7$ for the higher transition. These energy values also have a ratio of about 1.49^[5].
- Use the equations $E/B_1 = 25.9 = 17200/B_1$ and $E/B_2 = 38.7 = 25600/B_2$ to solve for two values of B that we can average to get a decent estimate for B . Indeed,

$$B_1 = \frac{17\,200\,\text{cm}^{-1}}{25.9} = 664\,\text{cm}^{-1} \quad B_1 = \frac{25\,600\,\text{cm}^{-1}}{38.7} = 661\,\text{cm}^{-1}$$

so

$$B = \frac{664\,\text{cm}^{-1} + 661\,\text{cm}^{-1}}{2} = 663\,\text{cm}^{-1}$$

- We can take this one step further now and estimate the crystal field splitting parameter Dq :

$$\begin{aligned} \frac{Dq}{B} &= 2.8 \\ Dq &= 2.8 \cdot 663\,\text{cm}^{-1} \\ &= 1860\,\text{cm}^{-1} \end{aligned}$$

⁵We can see that moving to the left would increase the ratio, diverging to $+\infty$, and moving to the right would decrease the ratio, converging to 1. Thus, 1.49 is the sweet spot. Note that given linear equations for these two essentially parallel lines, we could also explicitly calculate the sweet spot.

- Some d -electron counts have high- and low-spin cases.

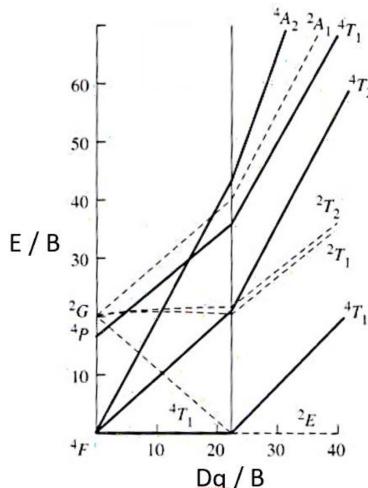


Figure 8.12: Tanabe-Sugano diagram for a d^7 configuration.

- As the field gets stronger, we will eventually have a sharp transition from high spin to low spin.
 - Essentially two different TS diagrams on the two sides of the divide.
- The HS-LS transition is directly proportional to $10Dq = e^2/r_{ij}$.
- In the HS case, we have a negatively sloped line representing a transition from the e_g set to the t_{2g} set that technically forms an excited state since the complex is high spin.

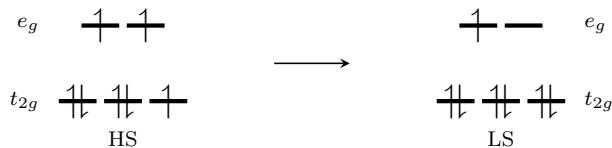


Figure 8.13: Explaining the negatively sloped transition.

- In particular, the transition requires energy because it involves spin pairing, which is costly in a weak field.
- All of the lines on the TS diagram informs us of possible transitions. But are they probable?
 - We've started to make some loose conclusions about what is allowed: Transitions between terms of the same spin multiplicity are spin-allowed, two-photon transitions have low probability, $g \rightarrow g$ transitions have low probability.
 - We will discuss these selection rules next time!
 - The probability of a transition is correlated to its intensity and directly proportional to oscillator strength.
 - We will discuss some applications of high-probability transitions, e.g., photovoltaic cells.
- Conclude the class on Wednesday with a review of bonding and electronic spectroscopy.