Week 8

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8.1 Midterm 2 Review

11/14: • Still 3 problems total and 5 points each.

- The problems will be calculations based on the basic concepts.
- Figure out the stable and unstable subspaces of some finite systems.
- Figure out whether or not a system is stable.
- Prove whether or not a function is planar linear
- Starting with the classification of planar linear autonomous systems.
 - We have y' = Ay where A is a 2×2 real matrix.
 - As a result of the realness, the eigenvalues behave regularly, i.e., there are only finitely many types
 of eigenvalues. These are...
 - 1. Real, nonzero, same sign. Depending on the sign, we'll either have a source or a sink. The orbits will be a distorted graph of a power function. If asked to investigate the phase portrait, then we need to figure out the stable and unstable subspaces and clearly indicate a basis. If asked to draw, we need to clearly indicate which subspaces are stable and unstable. We also need to clearly indicate the direction of the phase lines. First case: Everything is stable; second case: Everything is unstable. We draw the eigenspaces as well with arrows on the "axes." Figure 5.3a-5.3b.
 - 2. Real, different sign. One stable and one unstable subspace. We need to clearly indicate how the axes are tilted. Figure 5.3c.
 - 3. A is similar to the Jordan block with zero eigenvalues and 1 in the upper right hand corner.

 Then

$$A \sim \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} e^{tA} = \begin{pmatrix} 1 & t \\ 0 & 1 \end{pmatrix}$$

- 4. Purely imaginary eigenvalues. These must appear in a conjugate pair. The phase diagram will be concentric ellipses, and we essentially have the harmonic oscillator equation. If we have to sketch, we must show how the ellipses are tilted.
- 5. Complex eigenvalues $\sigma \pm i\beta$. Either we have a spiral source or a spiral sink. It's meaningless to indicate how the spiral tilts here, so don't bother trying. Determining whether they spin clockwise or counterclockwise. If

$$\begin{pmatrix} x \\ y \end{pmatrix}' = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

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then our fundamental solution is

$$\begin{pmatrix} x \\ y \end{pmatrix}' = \begin{pmatrix} \cos t & \sin t \\ -\sin t & \cos t \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

and we rotate counterclockwise. Since $A^2 = -\mu^2 I_2$, $e^{tA} = I_2 \cos \mu t + \mu^2 I_2 \sin t$. Negative reverses everything. Harmonic oscillator goes counterclockwise.

- There is an online website that gives us phase portraits for an equation. We can use this to help develop intuition.
- If you have a set of eigenvectors, how do you know how to tilt it?
 - Shao goes over examples of eigenvalues and eigenvectors.
- This is not something you need to memorize, but something you need to be able to recover.
- This is not a course for math majors; thus, there will not be proofs concerning the contraction mapping principle. We will not be asked to show existence, uniqueness, continuous difference, or differentiability with respect to parameters.
- We do need to know Grönwall's inequality, however.
- Grönwall's inequality: If $\phi:[p,T]\to\mathbb{R}$ and

$$\phi(t) \le b + a \int_0^t \phi(\tau) d\tau$$

then

$$\phi(t) \le b e^{at}$$

- Usually stated in the integral form, and we usually only need a special case.
- We may need to prove this; the proof mimics the derivation of the Duhamel formula.
- $-a,b \in \mathbb{R}$.
- We need to memorize the proof.
- We also need to be able to recognize when we can and should use it. Let $\phi(t) = \Phi'(t) \le b + a\Phi(t)$, $\Phi(0) = 0$. Then $\phi(t) \le b + a \int_0^t \varphi(\tau) d\tau$.
- Use it when we want to bound a function that satisfies either an integral or a differential quantity.
- This is the only proof in the theory of ODE systems we need to memorize.
- We need to master the methods to compute perturbation series.
 - Suppose our IVP depends on a parameter μ differentiably.

$$\frac{\mathrm{d}y}{\mathrm{d}t} = f(t; y(t; \mu); \mu), y(t_0) = x(\mu), \mu \approx 0$$

- If the parameter is close to zer, then you should be able to compute the μ -derivative with respect to the parameter.
- By Taylor expanding with respect to the parameter, you should be able to recover solutions that are close to the actual.

$$y(t; \mu) = y_0(t) + y_1(t)\mu + y_2(t)\mu^2 + O(\mu^3)$$

- We are typically satisfied with approximations to the second order.
- We expand our ODE into a Taylor series of μ . The differentiability with respect to parameters theorem (see Lecture 6.2 or Theorem 2.11 in Teschl (2012)) tells us that this is legitimate.

$$\frac{\mathrm{d}}{\mathrm{d}t}(y_0(t)) = f(t; y_0(t); 0), y_0(t_0) = x(0)$$

$$\frac{\mathrm{d}}{\mathrm{d}t}(y_1(t)) = \frac{\partial f}{\partial z}y_1(t) + \frac{\partial f}{\partial u}, y_1(t) = \frac{\partial x}{\partial u}$$

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 Just know the basic Taylor expansions (trig ones and exponential functions; usually we'll stick to polynomials, though).

- Use the ansatz $y(t; \mu) = y_0(t) + y_1(t)\mu + y_2(t)\mu^2 + O(\mu^3)$.
- Substitute $y(t; \mu)$ into $f(t, y(t; \mu); \mu)$. Expand $f(t; y(t; \mu); \mu)$ into a Taylor series of μ . Balance the coefficients of $\mu^0, \mu^1, \mu^2, \dots$
- Then you will get a series of equations that is theoretically solvable. Then a sequence of ODEs for $y_0(t), y_1(t), y_2(t), \dots$
- Your ODEs for $y_1, y_2, ...$ should not involve μ (because they are coefficients in the Taylor expansion with respect to μ . Coefficients of a Taylor series shouldn't involve the argument); if it does, there is something going wrong.
- As for the initial value, $y_0(t_0) + y_1(t_0)\mu + y_2(t_0)\mu^2 + \cdots$. This implies that something equals $x(\mu)$. The Taylor coefficients of $x(\mu)$ at $\mu = 0$.
- These are the general steps you use to find the perturbative series expansion.
- The computations on the exam will not be too heavy.
- If you're still unclear on the calculation, look through the HW answer keys.
- Conclusion: The Grönwall's inequality is something we need to remember from the theory; the perturbative procedure is something we need to be able to do.
- Why do we expand with respect to μ ?
 - We do it with respect to μ because our function is a function of μ . Differentiability and smallness imply we can use the Taylor series.
- Shao reiterates: Definitely read through the key to HW5!!! All the steps you will need to do are done completely and in detail.
- There will be things that are in HW6 (the one due Friday) that will appear on the exam because we have discussed these things in lecture.
- The definitions of Lyapunov stability and asymptotic stability. These will appear in the exam. We need to *clearly* remember the definitions.
- Consider y' = f(y), $f(x_0) = 0$ (an autonomous system with a fixed point; we can transform our system via $(y x_0)' = f(x_0 + (y x_0))$ to translate our fixed point to zero; implies $y' = f(x_0 + y)$, y = 0 is a fixed point). We should be able to determine the asymptotic stability near x_0 by computing the linearization (i.e., the Jacobian $f'(x_0)$) at the fixed point.
 - Regarding determining stability near x_0 , remember the following theorem.
 - Theorem: If all eigenvalues of $f'(x_0)$ have negative real parts, then x_0 is asymptotically stable. If at least one eigenvalue has real part greater than zero, then x_0 is not Lyapunov stable.
 - We should be able to apply the above criterion in practice.
 - We should also be able to reproduce the proof of the first part of Lyapunov's theorem (related to a question in HW6).
 - Lyapunov functions: $f(x_0) = 0$. Definition:
 - 1. L(x) is C^1 near x_0 , $L(x_0) = 0$. L(x) > 0 for x near x_0 .
 - 2. $\nabla L(x) \cdot f(x) \leq 0$ for x near x_0 iff $L(\phi_t(x)) \leq L(x)$, $t \geq 0$. If $L(\phi_t(x))$ is always strictly decreasing, then it is a strict Lyapunov function.
 - Theorem (Lyapunov's theorem): Usually, we can explicitly determine a Lyapunov function:
 - 1. If there is a Lyapunov function near the fixed point, then it is Lyapunov stable. For trajectories starting at nearby points, the trajectory can never excape nearby points.
 - 2. If there is a strict Lyapunov function, then it is asymptotically stable.

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- We need to be able to apply this theorem in practice; we don't need to know the proof.
- Examples of Lyapunov functions: Newton's second law.
 - Suppose you have a particle moving within a potential field with potential function U, i.e.,

$$mx'' = -U'(x)$$

- Then by a standard process, you can convert it to a planar linear system by introducing the variable v (the velocity), i.e.,

$$\begin{pmatrix} x \\ v \end{pmatrix}' = \begin{pmatrix} v \\ -U(x)'/m \end{pmatrix}$$

- Then $E(x,v) = \frac{m}{2}v^2 + U(x)$ is constant along the orbits, that is,

$$\nabla E(x, v) \cdot \begin{pmatrix} v \\ -U'(x)/m \end{pmatrix} = 0$$

- The gradient of the energy function is orthogonal to the vector field.
- -E(x,v) is a Lyapunov function (global). This happens and induces a fixed point exactly where the velocity is zero and the function takes on a critical value.
- Linearization at the fixed point $(x_0, 0)$ is

$$\begin{pmatrix} 0 & 1 \\ -\frac{U''(x_0)}{m} & 0 \end{pmatrix}$$

So $E(x_0, v) > E(x_0, 0)$ for $x \sim x_0$, $v \sim 0$ iff U takes a minimum at x_0 . The energy function cannot always stay larger than the energy at the fixed point. Satisfies second Lyapunov condition, but not the first.

One question: Classification of planar linear autonomous systems, one on Grönwall, one on qualitative asymptotic analysis using Lyapunov. Three questions total. There will also be some questions (parts of questions, I guess) on perturbative series.