

Chapter 18

Thermal Properties of Matter

18.1 Macroscopic Derivation of the Ideal Gas Law

- 8/19:
- Atoms have a nucleus (composed of protons and neutrons) orbited by electrons.
 - **Atomic number:** The number of protons. *Denoted by Z .*
 - **Atomic mass:** Essentially the number of protons plus the number of neutrons. *Denoted by A , μ .*
Units amu.
 - $1 \text{ amu} = \frac{1}{12} m(^{12}_6\text{C})$.
 - **Avogadro's number:** The number of molecules per mole of a substance, i.e., 6.02×10^{23} . *Denoted by N_A .*
 - N_A carbon-12 atoms weighs 12 g.
 - We define $1 \text{ mol} = N_A$ of something.
 - **Boyle's Law:** The product of the pressure and volume of a gas is a constant (that depends on the gas at hand).
 - **Ideal gas law:** The relation
$$pV = nRT$$
relating the pressure, volume, number of moles, and temperature of a dilute gas to a constant (that is not specific to any particular gas).
 - **Universal gas constant:** The constant $8.314 \frac{\text{J}}{\text{mol K}}$. *Denoted by R .*
 - Thus, we can think of temperature as being a reflection of a few macroscopic properties of gasses (e.g., pressure, volume, and number of moles).

18.2 Microscopic Derivation of the Ideal Gas Law

- Covers the derivation of the ideal gas law from KMT, as described Chapter 5 of Labalme (2020).
- Important addition:

$$p = \frac{1}{3} \rho \bar{v}^2$$

where p is pressure, ρ is density, and \bar{v} is the average velocity of the molecules.

- This relates a macroscopic and a microscopic quantity.

- Thus, we can relate the average speed of the molecules to the measurable pressure via

$$v_{\text{rms}} = \sqrt{\frac{3p}{\rho}}$$

- We can calculate from the above equation that the root mean square velocity of hydrogen gas at STP is about 1 800 m/s. For nitrogen gas, it's about 450 m/s.
- **Boltzmann constant:** The quotient of the universal gas constant and Avogadro's constant, having value 1.38×10^{-23} J/K. *Denoted by k .*
- It follows that

$$\overline{KE}_{\text{molecule}} = \frac{3}{2}kT$$

- Additionally, we have that

$$v_{\text{rms}} = \sqrt{\frac{3kT}{\mu}}$$

- This property can be taken advantage of for **diffusion separation of isotopes**.
- How to separate ^{238}U from ^{235}U :
 - Create UF_6 , a gas.
 - The lighter molecules will effuse slightly faster out of a box with a hole.
 - If you apply the cycle over and over again, you will enrich it a little bit each time.
 - Eventually, you will have a large proportion of $^{235}\text{UF}_6$, from which the ^{235}U can be extracted.

18.3 Thermodynamic Work

8/20:

- Final on Wednesday.
 - Posted at 10 AM CT, due on Canvas at 12:20 PM CT.
 - 2 hours.
 - Quantitative questions on HW 1-5 material, qualitative questions on Monday/Tuesday lecture material.
 - All chapters.
 - Will weight our grades with more/less emphasis on the final and take the higher of the two.
- Assuming that gas molecules have no internal structure with which to store energy (i.e., they can only store kinetic energy of motion), we have that

$$\begin{aligned} E_{\text{int}} &= nN_A \overline{KE}_{\text{molecule}} \\ &= \frac{3}{2}n(N_A k)T \\ &= \frac{3}{2}nRT \end{aligned}$$

- This implies that temperature alone determines the internal energy of a gas.
- This is a good approximation for monoatomic gasses, but we may need other formulas for gasses with more complex molecular structures.
- **Heat bath:** A bath that maintains a constant temperature.
 - Lake Michigan is a good example: Adding or removing heat from it will not significantly affect its temperature.

- Consider a container in a heat bath of temperature T .
 - The container is made of conducting walls and filled with gas of pressure p . It is also capped by a piston with cross-sectional area A , pushing down on the gas with force F .
 - If the piston moves up by a tiny distance dx , then the work dW exerted by the gas on the piston is given by

$$\begin{aligned} dW &= F \cdot dx \\ &= pA \, dx \\ &= p \, dV \end{aligned}$$

- Similarly, the work exerted by the piston on the gas is given by $dW = -p \, dV$.
- Thus, pushing down on the piston raises the internal energy of the gas. But since the container is in a heat bath, temperature stays the same, i.e., internal energy stays the same. Consequently, work being done on the gas by the piston must cause heat to flow out of the gas.
- If the piston moves from a to b , then

$$\begin{aligned} W &= \int_a^b -p \, dV \\ &= -nRT \int_a^b \frac{dV}{V} \\ &= -nRT \ln \left(\frac{V_b}{V_a} \right) \end{aligned}$$

- Do remember that this equation only holds in isothermal conditions.
 - If we have an expansion, then $V_b > V_a$, so $W < 0$.
 - Similarly, if we have a compression, then $V_b < V_a$, so $W > 0$.
- Convention: We will typically talk about work *done on the gas* and heat *added to the gas*.
- **Isothermal process:** A process during which temperature is held constant.
- **Isotherm:** A line of constant temperature in a pressure vs. volume graph.
- **Isochoric process:** A process during which volume is held constant.
- **Isobaric process:** A process during which pressure is held constant.
- Changes in pressure and volume are pathway independent.
- Consider an insulated container.
 - As before, though, the container is filled with gas of pressure p . It is also capped by a piston with cross-sectional area A , pushing down on the gas with force F .
 - If you compress the gas by pushing the piston down, you're doing positive work on the gas and E_{int} increases.
 - Temperature increases so much, actually, that you can ignite a gasoline-soaked cotton ball in a fire piston.
 - In the piston's frame, a gas molecule colliding with a moving piston leaves with the same velocity it came in with.
 - In the lab's frame, a colliding gas molecule leaves the collision with extra velocity (think of a collision between a very light and a very heavy object).
- **Adiabatic process:** A process with no transfer of heat.