

Math 228B Lec 7

plan:

- Boundary conditions in homework 1
- $\|B^n\|_{L_h^2} = \|f^n\|_{L^2(-\pi, \pi)}$
- 2-norm of a multiplication operator ($\|G^n\|_{L^2} = \|G\|_\infty^n$)
- Z-transform of a sampled function
(Fourier transform, Poisson summation formula)

Boundary conditions



For Dirichlet conditions

$(L=1 \text{ in homework})$

$$u(0, t) = 0, \quad u(L, t) = 0$$

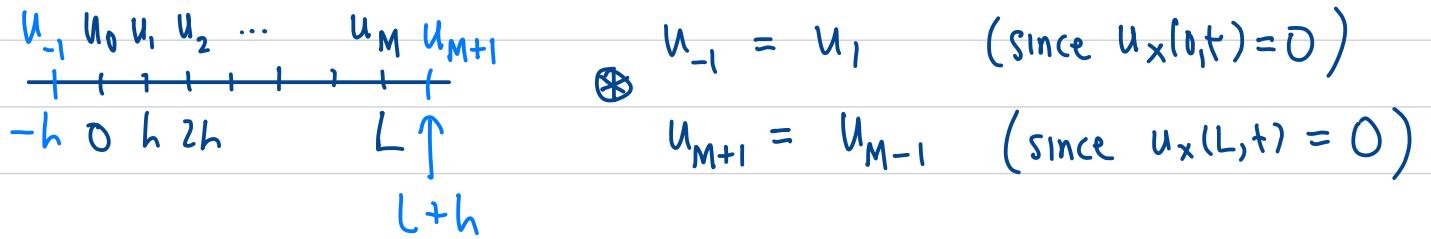
$$D_t^+ u = \alpha D_x^+ D_x^- u$$

update:
$$\begin{cases} u_j^{n+1} = \nu u_{j-1}^n + (1-2\nu) u_j^n + \nu u_{j+1}^n & 1 \leq j \leq M-1 \\ u_0^{n+1} = 0, \quad u_M^{n+1} = 0 & \end{cases}$$

(or don't store u_0^n, u_M^n and use a modified formula
for $j=1, j=M-1$)

Neumann b.c.'s $u_x(0, t) = 0, \quad u_x(L, t) = 0$

- include u_0, u_M among the unknowns
- use "ghost nodes" at $x_{-1} = -h, x_{M+1} = L+h$



Often the exact solution satisfies these reflective symmetries

$$u(-x, t) = u(x, t), \quad u(L+x, t) = u(L-x, t)$$

- eliminate u_{-1}, u_{M+1} from the update using \otimes

$$\text{e.g. } u_0^{n+1} = v u_{-1}^n + (1-2v) u_0^n + v u_1^n$$

$$= (1-2v) u_0^n + 2v u_1^n$$

↑
final update for u_j^{n+1} with $j=0$

(similar formula for $j=M$)

$$(\text{next: } \|B^n\|_{L_h^2} = \|f^n\|_{L^2(-\pi, \pi)} = \|G\|_\infty^n)$$

Observation: The norm of a unitary operator $Q: X \rightarrow Y$ is 1.

X, Y Hilbert spaces, $Q^{-1} = Q^*$ $\Rightarrow Q^* Q = I$ so

$$\|Qx\| = \sqrt{\langle Qx, Qx \rangle} = \sqrt{\langle Q^* Qx, x \rangle} = \sqrt{\langle x, x \rangle} = \|x\| \quad \forall x \in X.$$

$$\|Q\| = \sup_{x \neq 0} \frac{\|Qx\|}{\|x\|} = \sup_{x \neq 0} \frac{\|x\|}{\|x\|} = \sup_{x \neq 0} 1 = 1$$

We saw last time that the \mathcal{Z} transform is unitary

and diagonalizes any finite difference operator

up to a constant factor

$$\begin{array}{ccc} \mathbb{L}_h & \xrightarrow{\mathcal{B}} & \mathbb{L}_h \\ \mathcal{Z} \downarrow & \xrightarrow{G} & \downarrow \mathcal{Z} \\ \mathbb{L} & \xrightarrow{\Lambda} & \mathbb{L} \end{array}$$
$$\mathcal{B} u_j = \sum_j c_m u_{j+m}$$
$$G(\xi) = \sum_m c_m e^{im\xi}$$
$$G f(\xi) = G(\xi) f(\xi)$$

$\mathcal{B} = \mathcal{Z}^{-1} G \mathcal{Z}$ ← just like the eigen-decomposition
 $A = Q \Lambda Q^{-1}$ of a normal matrix

G plays the role of the diagonal matrix Λ

$$\begin{aligned} \therefore \| \mathcal{B}^n \| &\leq \| \mathcal{Z}^{-1} \| \cdot \| G \| \cdot \| \mathcal{Z} \| \\ &= \sqrt{\frac{h}{2\pi}} \| G \| \cdot \sqrt{\frac{2\pi}{h}} = \| G \| \end{aligned}$$

$\| \mathcal{Z} \| = 1$ unitary

using $G = \mathcal{Z} \mathcal{B}^n \mathcal{Z}^{-1}$ we also have

$$\| G \| \leq \underbrace{\| \mathcal{Z} \|}_{\sqrt{\frac{2\pi}{h}}} \cdot \| \mathcal{B}^n \| \cdot \underbrace{\| \mathcal{Z}^{-1} \|}_{\sqrt{\frac{h}{2\pi}}} = \| \mathcal{B}^n \|$$

$$\text{so } \| \mathcal{B}^n \|_{\mathbb{L}_h} = \| G \|_{\mathbb{L}(-\pi, \pi)}$$

Note: $G(\xi) = \sum_m c_m e^{im\xi}$ is continuous (it's a finite sum)

$$\underline{\text{Claim}}: \|Gf\|_2 = \|G\|_\infty = \max_{-\pi \leq \xi \leq \pi} |G(\xi)|$$

proof: let C be the RHS: C

step 1: show $\|Gf\|_2 \leq C \|f\|_2$ for all $f \in L^2(-\pi, \pi)$

$$\begin{aligned}\|Gf\|^2 &= \int_{-\pi}^{\pi} |Gf(\xi)|^2 d\xi = \int_{-\pi}^{\pi} |G(\xi)f(\xi)|^2 d\xi \\ &\leq C^2 \int_{-\pi}^{\pi} |f(\xi)|^2 d\xi = C^2 \|f\|^2\end{aligned}$$

↑ key step: $|G(\xi)|^2 \leq C^2$ for every $\xi \in [-\pi, \pi]$,

step 2: show that if $K < C$ then $\exists f$ s.t. $\|Gf\| > K \|f\|$ (i.e. no smaller constant than C will work)

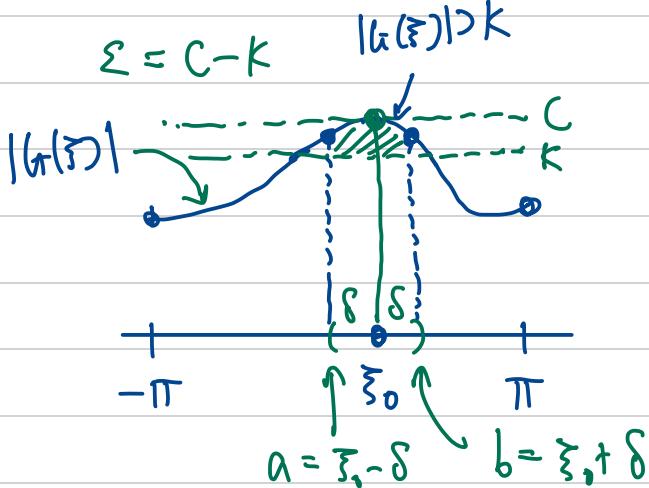
Since $|G(\xi)|$ is a continuous function, it achieves its

maximum at some point

$\xi_0 \in [-\pi, \pi]$, and there is
a neighborhood (a, b)

containing ξ_0 so that

$|G(\xi)| > K$ for $a \leq \xi \leq b$.



Now define $f(\xi) = \begin{cases} 0 & \xi < a \\ 1 & a \leq \xi \leq b \\ 0 & \xi > b \end{cases}$

$$\text{Then } \|yf\|_2^2 = \int_{-\pi}^{\pi} |G(\xi)f(\xi)|^2 d\xi = \int_a^b |G(\xi)|^2 d\xi$$

$$\geq \int_a^b K^2 d\xi = K^2(b-a)$$

$$\text{and } \|f\|^2 = \int_{-\pi}^{\pi} |f(\xi)|^2 d\xi = \int_a^b 1 d\xi = b-a$$

$$\text{So } \|yf\| \geq K \|f\|$$

Conclusion: The 2-norm of a finite difference operator is the maximum absolute value of the amplification factor $G(\xi)$.

$$\text{Also, } y^n f(\xi) = G(\xi)^n f(\xi) \quad \text{and} \quad \|G^n\|_\infty = \|G\|_\infty^n$$

$$\text{so } \|B^n\|_{L_h^2} = \|y^n\|_2 = \|G\|_\infty^n$$

von Neumann analysis also handles implicit methods since

$$B^{-1} = Z^{-1} G^{-1} Z, \quad G^{-1} f(\xi) = \frac{1}{G(\xi)} f(\xi)$$

$$\text{and } \|B^{-1}\|_{L_h^2} = \|G^{-1}\|_2 = \|\frac{1}{G}\|_\infty = \max_{-\pi \leq \xi \leq \pi} \left| \frac{1}{G(\xi)} \right|$$

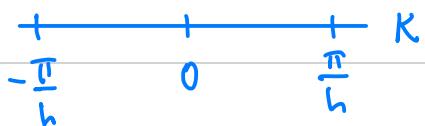
So far we are using the λ^2 version of the Z-transform, which is the easiest way to compute norms.

$$\hat{u}(\xi) = \sum_j u_j e^{-ij\xi} \quad \leftarrow \begin{matrix} \text{no change with} \\ \text{different grid spacings} \end{matrix}$$

To interpret this physically, it is useful to introduce a variant that is tailored to the grid:

def: $\tilde{u}(k) = h \sum_j u_j e^{-ijk}$

$k = \text{wave number}$
 $\uparrow (\text{units: } \frac{1}{\text{length}})$
 kappa



$$\xi = hk = \text{dimensionless wave number}$$

$$\psi_j(k) = e^{-ijk} \quad \text{basis functions}, \quad \tilde{u}(k) = \sum_j h u_j \psi_j(k)$$

orthogonality: $\frac{h}{2\pi} \int_{-\pi/h}^{\pi/h} \overline{\psi_j(k)} \psi_l(k) dk = \delta_{jl} = \begin{cases} 1 & j=l \\ 0 & j \neq l \end{cases}$

$$\sum_j |hu_j|^2 = \frac{h}{2\pi} \int_{-\pi/h}^{\pi/h} |\tilde{u}(k)|^2 dk \quad \leftarrow \begin{matrix} \text{Parseval's identity} \\ (\text{just the Pythagorean}) \end{matrix}$$

$$\therefore h \sum_j |u_j|^2 = \frac{1}{2\pi} \int_{-\pi/h}^{\pi/h} |\tilde{u}(k)|^2 dk$$

- Observations:
- $\tilde{U}(K) = h \hat{U}(hK)$
 - $\hat{U}(\xi)$ and $\tilde{U}(K)$ are periodic with periods 2π and $\frac{2\pi}{h}$, respectively

Suppose we sample a continuous function $U(x)$, $x \in \mathbb{R}$:

$$u_j = U(jh), \quad -\infty < j < \infty$$

Physical Z-transform of the sampled sequence u :

$$\tilde{U}(K) = h \sum_j u_j e^{-ijhK} = h \sum_j U(jh) e^{-i(jh)K} = h \sum_j V(jh)$$

$$V(x) = U(x) e^{-ikx} \quad \left(K \text{ fixed, different } K \text{'s give different } V \text{'s} \right)$$

Fourier transform of U, V : $\hat{U}(K) = \int_{-\infty}^{\infty} U(x) e^{-ikx} dx$

$$\hat{V}(r) = \int_{-\infty}^{\infty} V(x) e^{-irx} dx = \int_{-\infty}^{\infty} U(x) e^{-i(K+r)x} dx = \hat{U}(K+r)$$

Poisson summation formula

$$h \sum_{j=-\infty}^{\infty} V(jh) = \sum_{m=-\infty}^{\infty} \hat{V}\left(-\frac{2\pi}{h} m\right)$$

also true without the minus sign

$$\therefore \tilde{U}(K) = \sum_m \hat{U}\left(K - \frac{2\pi}{h} m\right) \quad \leftarrow \begin{array}{l} \text{significance will be} \\ \text{explained next time...} \end{array}$$