Monenton Equation:

$$Ro\left[\frac{\partial \hat{u}}{\partial \hat{t}} + \hat{u} - \nabla \hat{u}\right] + \hat{f} \times \hat{u} = -\hat{\nabla} \hat{\chi}$$

$$\hat{u} = \hat{u}_0 + Ro \hat{u}_1 + Ro \hat{u}_2 + \cdots$$

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Sub in:

$$R_{0} \left[\begin{array}{cccc} \frac{\partial \hat{u}_{0}}{\partial \hat{t}} & + R_{0} & \frac{\partial \hat{u}_{1}}{\partial \hat{t}} & + \dots & + & \left(\hat{u}_{0} + R_{0} \hat{u}_{1}^{2}, + \dots \right) \cdot \tilde{Y} \left(\hat{u}_{0} + R_{0} \hat{u}_{1}^{2} + \dots \right) \\ + & \left(1 + R_{0} & \hat{\beta} \hat{y} \right) \hat{z}_{2}^{2} \times \left(\hat{u}_{0} + R_{0} \hat{u}_{1}^{2} + \dots \right) \\ = & - & \tilde{Y} \left(\hat{\eta}_{0} + R_{0} \hat{\eta}_{1}^{2}, + \dots \right) \end{array}$$

* Equate all the terms that depend on Ro (ie. no dependence on Ro)

$$\frac{1}{2} \times \frac{1}{4} = -\frac{1}{2} \text{M}_0$$
 (terms at $O(R_0^\circ)$).

Geostrophic balance emerges at o (Ro°).

* Equate tems that depend on Ro :

$$\frac{\partial \hat{\Omega}_0}{\partial t} + \hat{\Omega}_0 \cdot \nabla \hat{\Omega}_0 + \frac{1}{2} \times \hat{\Omega}_1 + \hat{\beta} \hat{\gamma} \hat{z} \times \hat{\Omega}_0 = -\hat{\nabla} \hat{\gamma}_1$$

$$+ ems at o(Ro^{\frac{1}{2}}).$$

Eliminate of, by taking cut of this equation.

Name use of:

This gives:

$$\frac{1}{2} \times \left(\frac{\sqrt{3}}{\sqrt{3}} \times \sqrt{3} \right) = \left(\frac{1}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) - \left(\frac{\sqrt{3}}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) \\
- \frac{1}{2} \times \left(\frac{\sqrt{3}}{\sqrt{3}} \times \sqrt{3} \right) = \left(\frac{1}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) + \frac{1}{2} \cdot \left(\frac{\sqrt{3}}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) \\
- \frac{1}{2} \times \left(\frac{\sqrt{3}}{2} \times \sqrt{3} \right) = \left(\frac{\sqrt{3}}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) + \frac{1}{2} \cdot \left(\frac{\sqrt{3}}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) \\
- \frac{1}{2} \times \left(\frac{\sqrt{3}}{2} \times \sqrt{3} \right) = \left(\frac{\sqrt{3}}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) + \frac{1}{2} \cdot \left(\frac{\sqrt{3}}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) \\
- \frac{1}{2} \times \left(\frac{\sqrt{3}}{2} \times \sqrt{3} \right) = \left(\frac{\sqrt{3}}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) + \frac{1}{2} \cdot \left(\frac{\sqrt{3}}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) \\
- \frac{1}{2} \times \left(\frac{\sqrt{3}}{2} \times \sqrt{3} \right) = \left(\frac{\sqrt{3}}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) + \frac{1}{2} \cdot \left(\frac{\sqrt{3}}{2} \cdot \frac{\sqrt{3}}{\sqrt{3}} \right) \\
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- \frac{1}{2} \times \left(\frac{\sqrt{3}}{2} \times \sqrt{3} \right) = \frac{1}{2} \times \left(\frac{\sqrt{3}}{2} \times \sqrt{3} \right) + \frac{1}{2} \cdot \left(\frac{\sqrt{3}}{2} \times \sqrt{3} \right) + \frac{1}{2} \cdot \left(\frac{\sqrt{3}}{2} \times \sqrt{3} \right) \\
- \frac{1}{2} \times \left(\frac{\sqrt{3}}{2} \times \sqrt{3} \right) + \frac{1}{2} \times \left(\frac{\sqrt{3}}{2}$$

$$[F_{NN} \quad \nabla \times (\underline{a} \times \underline{b}) = (\underline{v} \cdot \underline{\nabla}) \cdot \underline{a} + \underline{a} (\underline{\nabla} \cdot \underline{b}) - \underline{b} (\underline{\nabla} \cdot \underline{a}) - (\underline{a} \cdot \underline{\nabla}) \cdot \underline{b}]$$

$$\mathring{\nabla} \times (\mathring{\beta} \mathring{y} \stackrel{?}{\underline{a}} \times \mathring{u}_{0}) = (\mathring{u}_{0} \cdot \mathring{\underline{\nabla}}) \mathring{\beta} \mathring{y} \stackrel{?}{\underline{a}} + \mathring{\beta} \mathring{y} \stackrel{?}{\underline{a}} (\underline{\nabla} \cdot \mathring{u}_{0})$$

$$-\dot{u_0}\left(\dot{\nabla}\cdot\dot{\beta}\dot{y}\,\hat{z}\right)-\dot{\beta}\dot{\gamma}\,\hat{z}\cdot\dot{\nabla}\dot{u_0}$$

$$=\left(\dot{u_0}\cdot\dot{\nabla}\right)\left(\dot{\beta}\dot{\gamma}\,\hat{z}\right)$$

$$\hat{\nabla} \times (-\hat{\nabla}\hat{\gamma}_{i}) = 0$$

Finally find:

$$\frac{\partial}{\partial t} \hat{S}_{0} \hat{2} - \hat{\nabla} \times (\hat{\Omega}_{0} \times \hat{S}_{0} \hat{2}) + \hat{2} (\hat{\nabla}_{0} \cdot \hat{\Omega}_{0}) + (\hat{\Omega}_{0} \cdot \hat{\Omega}_{0}) (\hat{\beta} \hat{Q} \hat{2}) = 0$$

$$= \hat{\Sigma}_{0} \hat{S}_{0} \hat{S}_{0$$

which finally gives

$$\frac{\partial}{\partial t} \hat{y}_{0} \hat{z} + \hat{u}_{0} \cdot \hat{\nabla} \left(\hat{y}_{0} + \hat{y}_{0} \hat{y}_{0} \right) \hat{z} = -\hat{z} \left(\hat{\nabla} \cdot \hat{u}_{0} \right)$$

Take the 2-component of this:

$$\frac{\partial \hat{S}_{0}}{\partial t} + \hat{\Omega}_{0} \cdot \hat{\nabla} (\hat{S}_{0} + \hat{P}\hat{S}_{0}) = -\hat{\nabla} \cdot \hat{\Omega}_{0} \cdot \hat{\Omega}_{0}$$

Vorkcity equation at o(Ro)

Continuity equation

$$Ro \cdot F \cdot \left[\frac{\partial \hat{\gamma}}{\partial t} + \hat{u} \cdot \hat{\nabla} \hat{\gamma} \right] + \left[1 + Ro F \hat{\gamma} \right] \hat{\nabla} \cdot \hat{u} = 0$$

$$o(1).$$

Sulo this in:

incompressibility. at o(Ro). (already satisfied by geospophic

$$O(P_0'): F\left[\frac{\partial \hat{\gamma}_0}{\partial \hat{\tau}} + \hat{\gamma}_0 \cdot \hat{\nabla} \hat{\gamma}_0\right] + \hat{\nabla} \cdot \hat{\gamma}_0 + F \hat{\gamma}_0 \hat{\nabla} \cdot \hat{\gamma}_0 = 0$$

divergence of in,

Use continuity equation (to o(Ro1) to eliminate the $\vec{\nabla} \cdot \vec{u}$, in the vorticity equation (to o(Ro1):

$$\frac{\partial \hat{S}_{0}}{\partial \hat{t}} + \hat{u}_{0} \cdot \hat{\nabla} \left(\hat{S}_{0} + \hat{\beta} \hat{y} \right) = - \hat{\nabla} \cdot \hat{u}_{0}'$$

$$F\left[\frac{\partial \hat{y_0}}{\partial \hat{t}} + \hat{y_0} \cdot \hat{\nabla} \hat{y_0}\right] = - \hat{\nabla} \cdot \hat{y_0}$$

$$=) \frac{\partial \vec{y}_{0}}{\partial \vec{t}} + \hat{u}_{0} \cdot \hat{\nabla} (\hat{y}_{0} + \hat{\beta}\hat{y}) - F \frac{\partial \vec{y}_{0}}{\partial \hat{t}} - F \hat{u}_{0}^{*} \cdot \hat{\nabla} \hat{y}^{*} = 0$$

But py has no t-dependence, so can include it

$$\frac{\partial f}{\partial t}$$
 $\left(\frac{\partial}{\partial s} + \hat{\beta}\hat{y} - F\hat{y}_{0}\right) = 0$ consenation of spotential vorticity advection by

advection by geostophic flow is.

Looks like potential valicity (ignaring to, H)

Can express this in terms of geostrophic streamfunction $\hat{\varphi}_0$ = $\nabla^2 \hat{\psi}_0$, $\hat{\psi}_0 = \left(-\partial \hat{\psi}_0 - \partial \hat{\psi}_0\right)$, $\hat{\chi}_0 = \hat{\psi}_0$

Put back dinensional scales:

$$\frac{D_o}{Dt} = \frac{\partial}{\partial t} + \underline{u}_o \cdot \underline{\nabla} = \frac{1}{T} \frac{D_o}{D\hat{t}} = \frac{1}{T} \frac{D_o}{D\hat{t}} = \frac{1}{T} \frac{D_o}{D\hat{t}}$$

Sub into non-din PU equation:

$$\frac{D_o}{D_f^2} \left(\hat{S}_o + \hat{\beta} \hat{\hat{y}} - F \hat{q}_o \right) = 0$$

$$\frac{\partial}{\partial t} = 0 \qquad Q = \frac{1}{2} + \frac{1}{2} + \frac{1}{2}$$

Divensional agpl conservation.

Q: quasigeostrophic potential vothicity

Can also expres geostophie streamfunction in tens of y: $\frac{2}{2} \times \frac{1}{10} = -\frac{7}{2} y_0$

$$\Rightarrow \hat{\mathbf{u}}_{0} = -\hat{\mathbf{z}} \times (\hat{\mathbf{z}} \times \hat{\mathbf{u}}_{0}) = \hat{\mathbf{z}} \times \hat{\mathbf{v}}_{0} = \hat{\mathbf{v}}_{0}$$

Sunnay: derivation et QG system.

* non-dineusionalized the momentum e conservation equations -> two non-dimensional numbers.

* In the limit of Rock I but Fro(1) did an as yn ptoke expansion of is, if, if

* Collected terms in each equation at o(ho):

- =) geostophic belance =) geostophic flow: 40, 40, 40

* Collected tems in each equation at o(Ro'):

- =) correction to geostrophic balance
- =) "ageostrophic flow": û,, ñ,

* Took the curl of the O (Ro') movember equation and combined it with O (Ro') continuity equation

Grandineusianal QG potential vorticity equation.

* Re-dinensionalize the QQ potential nothicity equation

A non-négarous derivation: (c.f. Assignment 3)

SW potential nothicity:

$$Q_{SW} = \frac{f+5}{4+4}$$

$$= \frac{f+5}{4+5}$$

This is the same as elq potential varticity except for constants like fo, H.

Rossby waves in the QG system

(Include effects of stratification though lottex-stretching tem)

* Single-larger QG shallow water model: PU conservation

votex stretching

Introduce le = L' = deformation wavenunlost.

* Base state: uniform zonal flow

=> sloping upper surface: $4 = \frac{9\overline{y}}{f_0} \Rightarrow \overline{y}(y) = -\frac{fv}{g}.y.$

* PV of the base stable is

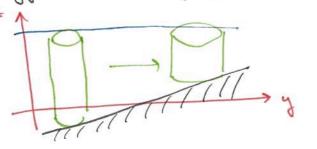
gradient in PV in y-direction.

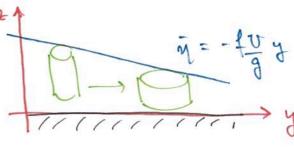
planetary

vortex stretching/ squeezing due to sloping upper surface

→ can get a PV gradient even when \$=0!

* Analogy with topographic B-effect:





Add small petubations:

$$4 = -Uy + \tilde{4}$$
 $\Rightarrow u = U - \frac{3\tilde{4}}{3x}$

$$V = \frac{3\tilde{4}}{3x}$$

Then

$$\frac{\partial Q}{\partial t} = \frac{\partial Q}{\partial t} + \underline{u} \cdot \underline{\nabla} Q$$

$$= \frac{\partial}{\partial t} \left(\nabla^2 \widetilde{4} - k \partial^2 \widetilde{4} \right) + \left(\upsilon - \partial \widetilde{4} \right) \frac{\partial}{\partial x} \left(\overline{Q}(y) + \nabla^2 \widetilde{4} - k \partial^2 \widetilde{4} \right)$$

$$+ \frac{\partial \widetilde{4}}{\partial x} \frac{\partial}{\partial y} \left(\overline{Q}(y) + \nabla^2 \widetilde{4} - k \partial^2 \widetilde{4} \right) = 0$$

Neglect continear tems:

$$\left(\frac{\partial}{\partial t} + U \frac{\partial}{\partial x}\right) \left(\nabla^2 \ddot{4} - k_0^2 \ddot{4}\right) + \left(\beta + k_0^2 U\right) \frac{\partial \ddot{4}}{\partial x} = 0$$

linearited QG PV equation.

Cs PDE with constant coefficients.

Plane wave solutions:
$$\tilde{Y} = \tilde{Y} = i(kx + ly - wt)$$

 $(-i\omega + ikU)(-(k^2+l^2) - ko^2)\tilde{Y}$
 $+ik(p + ko^2U)\tilde{Y} = 0$

non-thinial solvs I to

$$= 1 \quad \omega = k U - le \quad \beta + le_0^2 U$$
 Dispersion rela.
$$le^2 + l^2 + le_0^2$$

Two ney differences from SW:

* kg = Lo' \Rightarrow ettech of stratification.

* U \Rightarrow effect of mean flow

Case 1:
$$k_0 = 0$$
 ($L_0 \rightarrow \infty$)
$$\omega = kU - \beta k$$

$$R^2 + l^2 \qquad \text{where } k_0 = 0$$
Rossley wave

(as before, but with U)

Effect of U:

- changes the frequency by let of Dopples - changed the Cph by J. Shift.

Case 2: V = 0

$$\omega = -\frac{\beta k}{k^2 + l^2 + lkp^2}$$
new tem.

les \$0: "barodinic Rossby wave"

Rewrite as

Treat a as a constant (find contour of constant w)

Completing the square (check!):

$$\left(k + \frac{\beta}{2\omega}\right)^2 + l^2 = \left(\frac{\beta}{2\omega}\right)^2 - ko^2$$

$$\left(\frac{k_D}{k_D} + \frac{\beta}{2\omega k_D}\right)^2 + \left(\frac{k_D}{k_D}\right)^2 = \left(\frac{\beta}{2\omega k_D}\right)^2 - 1.$$

=) Circle centred at
$$(k_0, l_0) = (-\frac{\beta}{2\omega k_0}, 0)$$

with radius
$$\sqrt{\left(\frac{\beta}{2\omega k_B}\right)^2 - 1}$$

This gives a condition on w:

$$\left(\frac{\beta}{2\omega k_0}\right)^2 - 1$$
 >0

=)
$$\omega < \frac{\beta}{2k_D} = \omega_{max} \leftarrow maximum$$

frequency

Then
$$\left(\frac{R}{R_D} + \frac{\omega_{max}}{\omega}\right)^2 + \left(\frac{l}{R_D}\right)^2 = \left(\frac{\omega_{max}}{\omega}\right)^2 - 1$$
.

