

Section 4

We compute the linear representation of $SO(2)$ on \mathcal{T}_2

```
In[1]:= Clear[r];
n = 2;
vars = Table[x[i], {i, 1, n}];
rlist = Flatten[Table[r[i, j, k], {i, 1, n}, {j, 1, n}, {k, 1, n}]];
rename = Table[rlist[[m]] -> Subscript[a, 4 - m], {m, 1, 4}];
f =
  (Sum[r@@Sort[{i, j, k}] x[i] x[j] x[k], {i, 1, n}, {j, 1, n}, {k, 1, n}] /. rename)
Out[6]= a3 x[1]^3 + 3 a2 x[1]^2 x[2] + 3 a1 x[1] x[2]^2 + a0 x[2]^3
```

The last line is the expression for the cubic polynomial associate to a tensor Γ . Note that the coordinates are $x[1]$ and $x[2]$ where the indices are arguments and not subscripts.

```
In[7]:= Clear[sigma];
sigma[theta_] = {{Cos[theta], -Sin[theta]}, {Sin[theta], Cos[theta]}};
MatrixForm[sigma[theta]]
Out[9]//MatrixForm=
  ( Cos[theta]  -Sin[theta] )
  ( Sin[theta]   Cos[theta] )
```

Multiplying by this rotation matrix on the left gives the action of $SO(2)$ on \mathbb{R}^2 where the elements are thought of a column vectors. The action corresponds to rotating 'counter-clockwise' by an angle θ .

The action on Tensors (or equivalently on polynomials) is given by $\sigma \circ f(z) = f(\sigma^{-1} \cdot z)$ for $z \in \mathbb{R}^2$.

```
In[10]:= Substitution = Table[x[i] -> (sigma[-theta].{z[1], z[2]})[[i]], {i, 1, n}]
Out[10]=
  {x[1] -> Cos[theta] z[1] + Sin[theta] z[2], x[2] -> -Sin[theta] z[1] + Cos[theta] z[2]}
```

```
In[11]:= Transformedf = f //. Substitution
Out[11]=
  a0 (-Sin[theta] z[1] + Cos[theta] z[2])^3 +
  3 a1 (-Sin[theta] z[1] + Cos[theta] z[2])^2 (Cos[theta] z[1] + Sin[theta] z[2]) +
  3 a2 (-Sin[theta] z[1] + Cos[theta] z[2]) (Cos[theta] z[1] + Sin[theta] z[2])^2 +
  a3 (Cos[theta] z[1] + Sin[theta] z[2])^3
```

This is a cubic polynomial in $z[1], z[2]$. We can now read off the transformations of the coefficients from $\sigma \circ f(z) = b_3 z[1]^3 + 3 b_2 z[1]^2 z[2] + 3 b_1 z[1] z[2]^2 + b_0 z[2]^3$. We account for the factors of 3 in the coefficients $b[1]$ and $b[2]$ and order the coefficients as a column vector from b_0 to b_3 .

```
In[12]:= newcoeffs = Simplify[DiagonalMatrix[{1, 1/3, 1/3, 1}].
      CoefficientList[Transformedf /. {z[2] → 1}, z[1]]]
```

```
Out[12]= {Cos[θ]3 a0 + Sin[θ] (3 Cos[θ]2 a1 + Sin[θ] (3 Cos[θ] a2 + Sin[θ] a3)),
      1/4 (-4 Cos[θ]2 Sin[θ] a0 + (Cos[θ] + 3 Cos[3 θ]) a1 +
      2 Sin[θ] (a2 + 3 Cos[2 θ] a2 + Sin[2 θ] a3), Cos[θ] Sin[θ]2 a0 +
      1/4 ((Sin[θ] - 3 Sin[3 θ]) a1 + 2 Cos[θ] ((-1 + 3 Cos[2 θ]) a2 + Sin[2 θ] a3),
      -Sin[θ]3 a0 + Cos[θ] (3 Sin[θ]2 a1 + Cos[θ] (-3 Sin[θ] a2 + Cos[θ] a3))}
```

```
In[13]:= Lσ = Grad[newcoeffs, Table[ai-1, {i, 1, 4}]]
```

```
Out[13]= {{Cos[θ]3, 3 Cos[θ]2 Sin[θ], 3 Cos[θ] Sin[θ]2, Sin[θ]3},
      {-Cos[θ]2 Sin[θ], 1/4 (Cos[θ] + 3 Cos[3 θ]), 1/2 (1 + 3 Cos[2 θ]) Sin[θ], 1/2 Sin[θ] Sin[2 θ]},
      {Cos[θ] Sin[θ]2, 1/4 (Sin[θ] - 3 Sin[3 θ]), 1/2 Cos[θ] (-1 + 3 Cos[2 θ]), 1/2 Cos[θ] Sin[2 θ]},
      {-Sin[θ]3, 3 Cos[θ] Sin[θ]2, -3 Cos[θ]2 Sin[θ], Cos[θ]3}}
```

```
In[14]:= MatrixForm[Lσ]
```

```
Out[14]//MatrixForm=
      Cos[θ]3      3 Cos[θ]2 Sin[θ]      3 Cos[θ] Sin[θ]2      Sin[θ]3
      -Cos[θ]2 Sin[θ]  1/4 (Cos[θ] + 3 Cos[3 θ])  1/2 (1 + 3 Cos[2 θ]) Sin[θ]  1/2 Sin[θ] Sin[2 θ]
      Cos[θ] Sin[θ]2  1/4 (Sin[θ] - 3 Sin[3 θ])  1/2 Cos[θ] (-1 + 3 Cos[2 θ])  1/2 Cos[θ] Sin[2 θ]
      -Sin[θ]3      3 Cos[θ] Sin[θ]2      -3 Cos[θ]2 Sin[θ]      Cos[θ]3
```

This is the representation of SO(2) on the space of $2 \times 2 \times 2$ symmetric tensors. This representation is used in Sec. 4 (Pg. 15) of the paper. We can also determine the generator for this action.

Section 2.2

```
In[15]:= L = D[Lσ, θ] /. {θ → 0}
```

```
Out[15]= {{0, 3, 0, 0}, {-1, 0, 2, 0}, {0, -2, 0, 1}, {0, 0, -3, 0}}
```

```
In[16]:= MatrixForm[L]
```

```
Out[16]//MatrixForm=
      0  3  0  0
      -1 0  2  0
      0 -2  0  1
      0  0 -3  0
```

This is the matrix L in Sec. 2.1 (Pg. 11)

```
In[17]:= {vals, vecs} = Eigensystem[L]
Out[17]= {{3 i, -3 i, i, -i}, {{i, -1, -i, 1}, {-i, -1, i, 1}, {-3 i, 1, -i, 3}, {3 i, 1, i, 3}}}
```

```
In[18]:= Δ = DiagonalMatrix[vals]
Out[18]= {{3 i, 0, 0, 0}, {0, -3 i, 0, 0}, {0, 0, i, 0}, {0, 0, 0, -i}}
```

If we treat vecs as a matrix instead of a list of vectors, each eigenvector will be treated as a row. To make them columns, as appropriate for a right eigenvector, we need to take a transpose.

```
In[19]:= L.Transpose[vecs] - Transpose[vecs].Δ
Out[19]= {{0, 0, 0, 0}, {0, 0, 0, 0}, {0, 0, 0, 0}, {0, 0, 0, 0}}
```

Transpose[vecs] gives the right eigenvectors of \mathcal{L} . To get the Left eigenvectors as rows, we need to invert Transpose[vecs]. Below, we include an additional normalization to clear denominators.

```
In[20]:= Lvecs = Sqrt[Det[vecs]] Inverse[Transpose[vecs]]
Out[20]= {{1, -3 i, -3, i}, {-1, -3 i, 3, i}, {-1, i, -1, i}, {1, i, 1, i}}
```

These are the left Eigenvectors of \mathcal{L} . Lets Check

```
In[21]:= Lvecs.L - Δ.Lvecs
Out[21]= {{0, 0, 0, 0}, {0, 0, 0, 0}, {0, 0, 0, 0}, {0, 0, 0, 0}}
```

```
In[22]:= MatrixForm[Lvecs]
```

```
Out[22]//MatrixForm=

$$\begin{pmatrix} 1 & -3 i & -3 & i \\ -1 & -3 i & 3 & i \\ -1 & i & -1 & i \\ 1 & i & 1 & i \end{pmatrix}$$

```

This is the matrix E in Sec. 2.1 of the paper and the corresponding eigenvalues are $\lambda_1=3i$, $\lambda_2=-3i$, $\lambda_3=i$, $\lambda_4=-i$.

For a diagonal tensor $(\beta_1, 0, 0, \beta_2)$ we get $v_i = \beta_1 + i \beta_2 = v_4$. This gives $\text{Exp}[4 i \theta] = v_4(0) / v_1(0) = \frac{(a_0+a_2)+i(a_1+a_3)}{(a_0-3 a_2)+i(-3 a_1+a_3)}$. Subject to the necessary condition $(a_0 + a_2)^2 + (a_1 + a_3)^2 = (a_0 - 3 a_2)^2 + (-3 a_1 + a_3)^2$, we get 4 solutions for θ , say φ , $\varphi+\pi/2$, $\varphi+\pi$ and $\varphi+3\pi/2$.

These solutions define β_1 and β_2 by $\beta_1 + i \beta_2 = \text{Exp}[-i\varphi] ((a_0 + a_2) + i(a_1 + a_3))$. The rotations $\varphi+\pi/2$, $\varphi+\pi$ and $\varphi+3\pi/2$ then correspond, respectively, to the diagonal tensors $(\beta_2, 0, 0, -\beta_1)$, $(-\beta_1, 0, 0, -\beta_2)$ and $(-\beta_2, 0, 0, \beta_1)$ respectively.

Hilbert series for the SO(2) invariants.

In[23]:= **Simplify**[1 / (2 π) **Integrate**[1 / **Simplify**[**Det**[**IdentityMatrix**[4] - λ **L** σ]], { θ , 0, 2 π }]
 Out[23]=

$$\left\{ \begin{array}{ll} \frac{1+\lambda^4}{(-1+\lambda^2)^3 (1+\lambda^2)} & \text{Abs}[\lambda] > 1 \\ \frac{2+\lambda^{2/3} (-1-\lambda^{2/3}+\lambda^{4/3}) (1+\lambda^{2/3}+\lambda^{4/3}+2\lambda^2)}{3 (-1+\lambda^2)^3 (1+\lambda^2)} & \frac{1}{\text{Abs}[\lambda]^{1/3}} < 1 \text{ if } \text{Abs}[\lambda]^{1/3} \neq 1 \\ -\frac{1+\lambda^4}{(-1+\lambda^2)^3 (1+\lambda^2)} & \text{True} \end{array} \right.$$

We need the result for $\text{Abs}[\lambda] < 1$, so this is the last line in the piecewise defined integral.

In[24]:= **Φ SO₂** = **Simplify** $\left[-\frac{1+\lambda^4}{(-1+\lambda^2)^3 (1+\lambda^2)} \right]$

Out[24]=

$$-\frac{1+\lambda^4}{(-1+\lambda^2)^3 (1+\lambda^2)}$$

Hilbert series for the O(2) invariants.

We now compute the action of O(2) by adding a reflection operator corresponding to $x[1] \rightarrow x[1]$, $x[2] \rightarrow -x[2]$. In terms of the tensor coefficients, this action is given by the matrix

In[25]:= $\mathcal{N} = \{\{\theta, \theta, \theta, 1\}, \{\theta, \theta, 1, \theta\}, \{\theta, 1, \theta, \theta\}, \{1, \theta, \theta, \theta\}\};$

In[26]:= **Simplify**[**Det**[**IdentityMatrix**[4] - λ \mathcal{N} .**L** σ]]
 Out[26]=

$$(-1+\lambda^2)^2$$

This determinant does not depend explicitly on θ , so it is easy to integrate the reciprocal.

In[27]:= **Φ O₂** = **Simplify**[(1 / **Det**[**IdentityMatrix**[4] - λ \mathcal{N} .**L** σ] + **Φ SO₂**) / 2]
 Out[27]=

$$-\frac{1}{(-1+\lambda^2)^3 (1+\lambda^2)}$$

For O(2) covariants corresponding to linear forms, the generating function is 1/2 the generating function for SO(2). Does this make sense?

If we take u and w as the fundamental vector covariants, then other covariants are obtained by taking linear combinations of u and w with coefficients given by O(2) invariants. Consequently, the generating function is

In[28]:= **Simplify**[($\lambda + \lambda^3$) **Φ O₂**]

Out[28]=

$$-\frac{\lambda}{(-1+\lambda^2)^3}$$

Computation of the invariants

In[29]:= **u = Simplify[Grad[Laplacian[f, vars], vars] / 6]**

Out[29]=
 $\{a_1 + a_3, a_0 + a_2\}$

This is the trace vector. We ‘lift’ this vector to form the cubic polynomial f_1

In[30]:= **f₁ = 3 u.vars (vars.vars) / (n + 2)**

Out[30]=

$$\frac{3}{4} ((a_1 + a_3) x[1] + (a_0 + a_2) x[2]) (x[1]^2 + x[2]^2)$$

We can form an additional SO(2) covariant vector by rotating u counterclockwise by $\pi/2$.

In[31]:= **uperp = $\sigma[\pi/2]$.u**

Out[31]=
 $\{-a_0 - a_2, a_1 + a_3\}$

f_1 corresponds to a $2 \times 2 \times 2$ symmetric tensor \mathcal{B}

In[32]:= **$\mathcal{B} = \text{Table}[D[f_1/6, x[i], x[j], x[k]], \{i, 1, n\}, \{j, 1, n\}, \{k, 1, n\}];$**

In[33]:= **Table[MatrixForm[$\mathcal{B}[[i]]$], {i, 1, n}]**

Out[33]=

$$\left\{ \begin{pmatrix} \frac{3}{4} (a_1 + a_3) & \frac{1}{4} (a_0 + a_2) \\ \frac{1}{4} (a_0 + a_2) & \frac{1}{4} (a_1 + a_3) \end{pmatrix}, \begin{pmatrix} \frac{1}{4} (a_0 + a_2) & \frac{1}{4} (a_1 + a_3) \\ \frac{1}{4} (a_1 + a_3) & \frac{3}{4} (a_0 + a_2) \end{pmatrix} \right\}$$

With our normalization, the trace-free part is given by $f_3 = (n+2)f - f_1$

In[34]:= **f₃ = Collect[Expand[(n + 2) (f - f₁)], vars]**

Out[34]=

$$(-3 a_1 + a_3) x[1]^3 + (-3 a_0 + 9 a_2) x[1]^2 x[2] + (9 a_1 - 3 a_3) x[1] x[2]^2 + (a_0 - 3 a_2) x[2]^3$$

f_3 corresponds to a trace-free $2 \times 2 \times 2$ symmetric tensor \mathcal{D} . To eliminate denominators, we multiply by a factor of (n+2), which equals 4 in the case n=2.

In[35]:= **$\mathcal{D} = \text{Table}[\text{Simplify}[D[f_3/6, x[i], x[j], x[k]]], \{i, 1, n\}, \{j, 1, n\}, \{k, 1, n\}];$**

In[36]:= **Table[MatrixForm[$\mathcal{D}[[i]]$], {i, 1, n}]**

Out[36]=

$$\left\{ \begin{pmatrix} -3 a_1 + a_3 & -a_0 + 3 a_2 \\ -a_0 + 3 a_2 & 3 a_1 - a_3 \end{pmatrix}, \begin{pmatrix} -a_0 + 3 a_2 & 3 a_1 - a_3 \\ 3 a_1 - a_3 & a_0 - 3 a_2 \end{pmatrix} \right\}$$

In[37]:= **Dstarsqrd = Simplify[TensorContract[$\mathcal{D} \otimes \mathcal{D}$, {{1, 4}, {2, 5}}]]**

Out[37]=

$$\left\{ \left\{ 2 ((a_0 - 3 a_2)^2 + (-3 a_1 + a_3)^2), 0 \right\}, \left\{ 0, 2 ((a_0 - 3 a_2)^2 + (-3 a_1 + a_3)^2) \right\} \right\}$$

In[38]:= **v = Simplify[TensorContract[($\mathcal{D} \otimes \mathcal{D}$) $\otimes \mathcal{D}$, {{1, 4}, {2, 5}, {3, 7}, {6, 8}}]]**

Out[38]=
 $\{0, 0\}$

```
In[39]:= Simplify[Dstarsqrd.D]
```

```
Out[39]=
```

$$\left\{ \left\{ \left\{ 2(-3a_1 + a_3) \left((a_0 - 3a_2)^2 + (-3a_1 + a_3)^2 \right), 2(-a_0 + 3a_2) \left((a_0 - 3a_2)^2 + (-3a_1 + a_3)^2 \right) \right\}, \right. \right. \\ \left. \left\{ 2(-a_0 + 3a_2) \left((a_0 - 3a_2)^2 + (-3a_1 + a_3)^2 \right), 2(3a_1 - a_3) \left((a_0 - 3a_2)^2 + (-3a_1 + a_3)^2 \right) \right\} \right\}, \\ \left\{ \left\{ 2(-a_0 + 3a_2) \left((a_0 - 3a_2)^2 + (-3a_1 + a_3)^2 \right), 2(3a_1 - a_3) \left((a_0 - 3a_2)^2 + (-3a_1 + a_3)^2 \right) \right\}, \right. \\ \left. \left\{ 2(3a_1 - a_3) \left((a_0 - 3a_2)^2 + (-3a_1 + a_3)^2 \right), 2(a_0 - 3a_2) \left((a_0 - 3a_2)^2 + (-3a_1 + a_3)^2 \right) \right\} \right\} \right\}$$

```
In[40]:= Simplify[Table[Tr[Simplify[Dstarsqrd.D][i]], {i, 1, n}]]
```

```
Out[40]=
```

$$\{0, 0\}$$

```
In[41]:= w = Expand[D.u.u]
```

```
Out[41]=
```

$$\{a_0^2 a_1 - 3a_1^3 + 10a_0 a_1 a_2 + 9a_1 a_2^2 - 3a_0^2 a_3 - 5a_1^2 a_3 + 2a_0 a_2 a_3 + 5a_2^2 a_3 - a_1 a_3^2 + a_3^3, \\ a_0^3 + 5a_0 a_1^2 - a_0^2 a_2 + 9a_1^2 a_2 - 5a_0 a_2^2 - 3a_2^3 + 2a_0 a_1 a_3 + 10a_1 a_2 a_3 - 3a_0 a_3^2 + a_2 a_3^2\}$$

```
In[42]:= TeXForm[w]
```

```
Out[42]//TeXForm=
```

$$\left\{ -3a_1^3 - 5a_3a_1^2 + a_0^2a_1 + 9a_2^2a_1 - a_3^2a_1 + 10a_0a_2a_1 + a_3^3 - 3a_0^2a_3 + 5a_2^2a_3 + 2a_0a_2a_3, a_0^3 - a_2a_0^2 + 5a_1^2a_0 - 5a_2^2a_0 - 3a_2^3 + a_0a_1a_3 + 10a_1a_2a_3 - 3a_0a_3^2 + a_2a_3^2 \right\}$$

```
In[43]:= Simplify[u.w]
```

```
Out[43]=
```

$$a_0^4 - 3a_1^4 - 3a_2^4 - 8a_1^3 a_3 + 24a_1 a_2^2 a_3 + 6a_2^2 a_3^2 + a_3^4 - \\ 8a_0 a_2 (-3a_1^2 + a_2^2 - 3a_1 a_3) + 6a_0^2 (a_1^2 - a_2^2 - a_3^2) + 6a_1^2 (3a_2^2 - a_3^2)$$

```
In[44]:= Simplify[Det[{u, w}]]
```

```
Out[44]=
```

$$4 \left(-3a_0^2 a_1 a_2 + a_0^3 a_3 + a_2 (3a_1^3 + 6a_1^2 a_3 - 2a_2^2 a_3 - 3a_1 (a_2^2 - a_3^2)) + a_0 (2a_1^3 - 6a_1 a_2^2 + 3a_1^2 a_3 - a_3 (3a_2^2 + a_3^2)) \right)$$

We can now compute SO(2) invariants from u, w and \mathcal{D} .

```
In[45]:= Clear[j, h, l, m]
```

```
Invariants =
```

$$\{j_2 - u.u, h_2 - \text{Tr}[\text{TensorContract}[\mathcal{D} \otimes \mathcal{D}, \{\{1, 4\}, \{2, 5\}\}]], l_4 - w.u, m_4 - \text{Det}[\{u, w\}]\}$$

```
Out[46]=
```

$$\left\{ -(a_0 + a_2)^2 - (a_1 + a_3)^2 + j_2, -(a_0 - 3a_2)^2 - 3(-a_0 + 3a_2)^2 - 3(3a_1 - a_3)^2 - (-3a_1 + a_3)^2 + h_2, \right. \\ \left. - \left((a_0 + a_2) (a_0^3 + 5a_0 a_1^2 - a_0^2 a_2 + 9a_1^2 a_2 - 5a_0 a_2^2 - 3a_2^3 + 2a_0 a_1 a_3 + 10a_1 a_2 a_3 - 3a_0 a_3^2 + a_2 a_3^2) \right) - \right. \\ \left. (a_1 + a_3) (a_0^2 a_1 - 3a_1^3 + 10a_0 a_1 a_2 + 9a_1 a_2^2 - 3a_0^2 a_3 - 5a_1^2 a_3 + 2a_0 a_2 a_3 + 5a_2^2 a_3 - a_1 a_3^2 + a_3^3) + \right. \\ l_4, -8a_0 a_1^3 + 12a_0^2 a_1 a_2 - 12a_1^3 a_2 + 24a_0 a_1 a_2^2 + 12a_1 a_3^2 - 4a_0^3 a_3 - \\ \left. 12a_0 a_1^2 a_3 - 24a_1^2 a_2 a_3 + 12a_0 a_2^2 a_3 + 8a_2^3 a_3 - 12a_1 a_2 a_3^2 + 4a_0 a_3^3 + m_4 \right\}$$

This is the basis for the ideal of polynomials on \mathbb{R}^8 (corresponding to the 4 coefficients a_0, a_1, a_2, a_3 and the 4 invariants j_2, h_2, l_4, m_4

We seek potential relations among the invariants by finding a basis for the ideal generated by the

definitions of the invariants intersected with the polynomials in j_2, h_2, l_4, m_4 that do not depend on a_0, a_1, a_2, a_3 , i.e. we are eliminating the coefficients between the relations defining the invariants.

```
In[47]:= GroebnerBasis[Invariants, {m4, l4, h2, j2},
  {a0, a1, a2, a3}, MonomialOrder -> EliminationOrder]
```

```
Out[47]= {h2 j2^3 - 4 l4^2 - 4 m4^2}
```

We see that there is one identity that allows us to replace m_4^2 by an expression in the other invariants. There are no further relations, so this implies l_4 , h_2 and j_2 are algebraically independent.

Alternate choices for the fundamental invariants are the trace and determinant of Γ^{*2} , which are themselves $O(2)$ invariants, so they can be expressed in terms of the fundamental invariants j_2, h_2 and l_4

```
In[48]:= r = Table[D[f/6, x[i], x[j], x[k]], {i, 1, n}, {j, 1, n}, {k, 1, n}];
Q = Simplify[TensorContract[rTensor, {{1, 4}, {2, 5}}]];
MatrixForm[Q]
```

```
Out[49]//MatrixForm=
( a1^2 + 2 a2^2 + a3^2      a0 a1 + a2 (2 a1 + a3) )
( a0 a1 + a2 (2 a1 + a3)    a0^2 + 2 a1^2 + a2^2 )
```

This is the matrix Γ^{*2} .

```
In[50]:= NewInvariants = {tau - Tr[Q], delta - Det[Q]}
```

```
Out[50]= {tau - a0^2 - 3 a1^2 - 3 a2^2 - a3^2,
  delta - 2 a1^4 + 4 a0 a1^2 a2 - 2 a0^2 a2^2 - a1^2 a2^2 - 2 a2^4 + 2 a0 a1 a2 a3 + 4 a1 a2^2 a3 - a0^2 a3^2 - 2 a1^2 a3^2}
```

```
In[51]:= GroebnerBasis[Join[NewInvariants, Invariants],
  {tau, delta, l4, h2, j2}, {m4, a0, a1, a2, a3}, MonomialOrder -> EliminationOrder]
```

```
Out[51]= {16 tau - h2 - 12 j2, -1024 delta + h2^2 + 8 h2 j2 + 80 j2^2 - 128 l4}
```

This shows that $\text{Tr}[\Gamma^{*2}] = \frac{h_2 + 12 j_2}{16}$, $\text{Det}[\Gamma^{*2}] = \frac{h_2^2 + 8 h_2 j_2 + 80 j_2^2 - 128 l_4}{1024}$

```
In[52]:= DiagSetting = {a3 -> alpha, a2 -> 0, a1 -> 0, a0 -> beta}
```

```
Out[52]= {a3 -> alpha, a2 -> 0, a1 -> 0, a0 -> beta}
```

```
In[53]:= Simplify[Invariants /. DiagSetting]
```

```
Out[53]= {-alpha^2 - beta^2 + j2, -4 (alpha^2 + beta^2) + h2, -alpha^4 + 6 alpha^2 beta^2 - beta^4 + l4, 4 alpha^3 beta - 4 alpha beta^3 + m4}
```

```
In[54]:= GroebnerBasis[{-alpha^2 - beta^2 + j2, -4 (alpha^2 + beta^2) + h2, -alpha^4 + 6 alpha^2 beta^2 - beta^4 + l4, 4 alpha^3 beta - 4 alpha beta^3 + m4},
  {alpha, h2, j2, l4, m4}, {beta}]
```

```
Out[54]= {j2^4 - l4^2 - m4^2, h2 - 4 j2, 8 alpha^4 - 8 alpha^2 j2 + j2^2 - l4}
```

```

In[55]:= u /. DiagSetting
Out[55]=
 $\{\alpha, \beta\}$ 

In[56]:= w /. DiagSetting
Out[56]=
 $\{\alpha^3 - 3 \alpha \beta^2, -3 \alpha^2 \beta + \beta^3\}$ 

In[57]:= Simplify[(Invariants /. DiagSetting) /. {\beta \to \alpha}]
Out[57]=
 $\{-2 \alpha^2 + j_2, -8 \alpha^2 + h_2, 4 \alpha^4 + l_4, m_4\}$ 

In[58]:= GroebnerBasis[{-2 \alpha^2 + j_2, -8 \alpha^2 + h_2, 4 \alpha^4 + l_4, m_4}, {\alpha, h_2, j_2, l_4, m_4}]
Out[58]=
 $\{m_4, j_2^2 + l_4, h_2 - 4 j_2, 2 \alpha^2 - j_2\}$ 

In[59]:= Simplify[(Invariants /. DiagSetting) /. {\beta \to 0}]
Out[59]=
 $\{-\alpha^2 + j_2, -4 \alpha^2 + h_2, -\alpha^4 + l_4, m_4\}$ 

In[60]:= GroebnerBasis[{-\alpha^2 + j_2, -4 \alpha^2 + h_2, -\alpha^4 + l_4, m_4}, {\alpha, h_2, j_2, l_4, m_4}]
Out[60]=
 $\{m_4, j_2^2 - l_4, h_2 - 4 j_2, \alpha^2 - j_2\}$ 

```

Section 7.2

Fully decoupleable $3 \times 3 \times 3$ tensors

```

In[61]:= n = 3; vars = Table[x[i], {i, 1, n}]; f = Sum[\beta_i x[i]^3, {i, 1, n}]
Out[61]=
 $\beta_1 x[1]^3 + \beta_2 x[2]^3 + \beta_3 x[3]^3$ 

```

This is the cubic polynomial corresponding to a fully decoupleable tensor.

```

In[62]:= r = Simplify[Table[D[f, x[i], x[j], x[k]], {k, 1, n}, {j, 1, n}, {i, 1, n}] / 6];
u = Simplify[Table[D[Laplacian[f, vars], x[i]], {i, 1, n}] / 6];
f3 = (n + 2) f - 3 (u.vars) (vars.vars)
Out[64]=
 $-3 (\beta_1 x[1] + \beta_2 x[2] + \beta_3 x[3]) (x[1]^2 + x[2]^2 + x[3]^2) + 5 (\beta_1 x[1]^3 + \beta_2 x[2]^3 + \beta_3 x[3]^3)$ 

```

This is the trace-free part

```

In[65]:= d = Simplify[Table[D[f3, x[i], x[j], x[k]], {k, 1, n}, {j, 1, n}, {i, 1, n}] / 6];
v = Simplify[TensorContract[(d \otimes d) \otimes d, {{1, 4}, {2, 5}, {3, 7}, {6, 8}}]];
w = Simplify[d.u.u];

```

We have now computed the vectors u, v and w and the trace-free tensor \mathcal{D} , which are the ingredients needed to compute the Integrty basis given by Olive and Auffray.


```

In[68]:= Clear[H, J, K, L, M, Q];
Q = Simplify[TensorContract[D⊗D, {{1, 4}, {2, 5}}]];
γuu = w;
Coeffs = Table[βi, {i, 1, n}];
Trivialize = Table[Coeffs[[j]] → 0, {j, 1, Length[Coeffs]}];
IntegrityPolys =
  {H[2] - Simplify[Tr[Q]], H[4] - Simplify[Tr[Q.Q]], J[2] - Simplify[u.u],
    L[4] - Simplify[γuu.u], H[6] - Simplify[v.v], H[10] - Simplify[D.v.v.v],
    J[4] - Simplify[u.Q.u], K[4] - Simplify[Tr[Q.(D.u)]],
    J[6] - Simplify[(u.Q).γuu], K[6] - Simplify[v.w], L[6] - Simplify[(u.Q).v],
    M[6] - Simplify[γuu.γuu], H[8] - Simplify[(u.Q).(Q.v)]};
OAINvariants = IntegrityPolys /. Trivialize;

```

The Ideal IntegrityPolys is generated by the polynomials defining the Integrity basis elements in terms of the coefficients of a fully decoupled tensor, with the labels of the Integrity invariants as the slack variables.

```

In[74]:= IntegrityPolys

```

```

Out[74]=

```

$$\begin{aligned}
&\{H[2] - 10 (\beta_1^2 + \beta_2^2 + \beta_3^2), H[4] - 44 \beta_1^4 - 44 \beta_2^4 - 58 \beta_2^2 \beta_3^2 - 44 \beta_3^4 - 58 \beta_1^2 (\beta_2^2 + \beta_3^2), \\
&J[2] - \beta_1^2 - \beta_2^2 - \beta_3^2, L[4] - 2 (\beta_1^4 + \beta_2^4 - 3 \beta_2^2 \beta_3^2 + \beta_3^4 - 3 \beta_1^2 (\beta_2^2 + \beta_3^2)), \\
&H[6] - 4 (\beta_3^2 (\beta_1^2 + \beta_2^2 - 4 \beta_3^2)^2 + \beta_2^2 (\beta_1^2 - 4 \beta_2^2 + \beta_3^2)^2 + \beta_1^2 (-4 \beta_1^2 + \beta_2^2 + \beta_3^2)^2), \\
&H[10] - 8 (128 \beta_1^{10} + 128 \beta_2^{10} - 60 \beta_2^8 \beta_3^2 - 95 \beta_2^6 \beta_3^4 - 95 \beta_2^4 \beta_3^6 - \\
&\quad 60 \beta_2^2 \beta_3^8 + 128 \beta_3^{10} - 60 \beta_1^8 (\beta_2^2 + \beta_3^2) + \beta_1^6 (-95 \beta_2^4 + 60 \beta_2^2 \beta_3^2 - 95 \beta_3^4) + \\
&\quad \beta_1^4 (-95 \beta_2^6 + 90 \beta_2^4 \beta_3^2 + 90 \beta_2^2 \beta_3^4 - 95 \beta_3^6) - 30 \beta_1^2 (2 \beta_2^8 - 2 \beta_2^6 \beta_3^2 - 3 \beta_2^4 \beta_3^4 - 2 \beta_2^2 \beta_3^6 + 2 \beta_3^8)), \\
&J[4] - 2 (3 \beta_1^4 + 3 \beta_2^4 + \beta_2^2 \beta_3^2 + 3 \beta_3^4 + \beta_1^2 (\beta_2^2 + \beta_3^2)), K[4] - 8 \beta_1^4 - 8 \beta_2^4 + 4 \beta_2^2 \beta_3^2 - 8 \beta_3^4 + 4 \beta_1^2 (\beta_2^2 + \beta_3^2), \\
&J[6] - 12 \beta_1^6 - 12 \beta_2^6 + 19 \beta_2^4 \beta_3^2 + 19 \beta_2^2 \beta_3^4 - 12 \beta_3^6 + 19 \beta_1^4 (\beta_2^2 + \beta_3^2) + \beta_1^2 (19 \beta_2^4 + 18 \beta_2^2 \beta_3^2 + 19 \beta_3^4), \\
&K[6] - 2 (8 \beta_1^6 + 8 \beta_2^6 - 11 \beta_2^4 \beta_3^2 - 11 \beta_2^2 \beta_3^4 + 8 \beta_3^6 - 11 \beta_1^4 (\beta_2^2 + \beta_3^2) + \beta_1^2 (-11 \beta_2^4 + 18 \beta_2^2 \beta_3^2 - 11 \beta_3^4)), \\
&L[6] - 6 (8 \beta_1^6 + 8 \beta_2^6 - \beta_2^4 \beta_3^2 - \beta_2^2 \beta_3^4 + 8 \beta_3^6 - \beta_1^4 (\beta_2^2 + \beta_3^2) - \beta_1^2 (\beta_2^2 + \beta_3^2)^2), \\
&M[6] - \beta_3^2 (3 \beta_1^2 + 3 \beta_2^2 - 2 \beta_3^2)^2 - \beta_2^2 (3 \beta_1^2 - 2 \beta_2^2 + 3 \beta_3^2)^2 - (2 \beta_1^3 - 3 \beta_1 (\beta_2^2 + \beta_3^2))^2, \\
&H[8] - 4 (72 \beta_1^8 + 18 \beta_1^6 (\beta_2^2 + \beta_3^2) - 11 \beta_1^4 (3 \beta_2^4 + \beta_2^2 \beta_3^2 + 3 \beta_3^4) + \\
&\quad \beta_1^2 (18 \beta_2^6 - 11 \beta_2^4 \beta_3^2 - 11 \beta_2^2 \beta_3^4 + 18 \beta_3^6) + 3 (24 \beta_2^8 + 6 \beta_2^6 \beta_3^2 - 11 \beta_2^4 \beta_3^4 + 6 \beta_2^2 \beta_3^6 + 24 \beta_3^8)) \}
\end{aligned}$$

Characteristic Polynomial coefficients

```

In[75]:= rstarsqrd = Simplify[TensorContract[r⊗r, {{1, 4}, {2, 5}}]];

```

```

In[76]:= Clear[q, ξ];

```

```

ξ = Rest[Reverse[CoefficientList[Simplify[Det[λ IdentityMatrix[3] + rstarsqrd]], λ]], λ]]

```

```

Out[76]=

```

$$\{\beta_1^2 + \beta_2^2 + \beta_3^2, \beta_1^2 \beta_2^2 + \beta_1^2 \beta_3^2 + \beta_2^2 \beta_3^2, \beta_1^2 \beta_2^2 \beta_3^2\}$$

As expected, these are the elementary symmetric polynomials of the quantities β_i^2 .

```
In[77]:= DiagInvars = Table[qi -  $\xi$ [[i]], {i, 1, 3}]
```

```
Out[77]= {q1 -  $\beta_1^2$  -  $\beta_2^2$  -  $\beta_3^2$ , q2 -  $\beta_1^2 \beta_2^2$  -  $\beta_1^2 \beta_3^2$  -  $\beta_2^2 \beta_3^2$ , q3 -  $\beta_1^2 \beta_2^2 \beta_3^2$ }
```

This is the basis of invariants for the group $G_R = S_3 \times (\mathbb{Z}_2)^3$. Since all the Olive and Auffray invariants, when restricted to fully decoupled tensors, are also G_R invariants, we can express them in terms of the quantities q_i

```
In[78]:= Table[GroebnerBasis[Join[{IntegrityPolys[[i]], DiagInvars},
Join[{OAIInvariants[[i]], {q1, q2, q3}}, { $\beta_1$ ,  $\beta_2$ ,  $\beta_3$ },
MonomialOrder  $\rightarrow$  EliminationOrder] [[1]], {i, 1, Length[OAIInvariants]}]
```

```
Out[78]= {H[2] - 10 q1, -H[4] + 44 q12 - 30 q2, J[2] - q1,
-L[4] + 2 q12 - 10 q2, -H[6] + 64 q13 - 220 q1 q2 + 300 q3,
-H[10] + 1024 q15 - 5600 q13 q2 + 5800 q1 q22 + 7600 q12 q3 - 7000 q2 q3,
-J[4] + 6 q12 - 10 q2, -K[4] + 8 q12 - 20 q2, -J[6] + 12 q13 - 55 q1 q2 + 75 q3,
-K[6] + 16 q13 - 70 q1 q2 + 150 q3, -L[6] + 48 q13 - 150 q1 q2 + 150 q3,
-M[6] + 4 q13 - 15 q1 q2 + 75 q3, -H[8] + 288 q14 - 1080 q12 q2 + 300 q22 + 1300 q1 q3}
```

This is the ideal corresponding to the relations in Eq. (7.1).

Section 7.3

Partially decoupleable $3 \times 3 \times 3$ tensors

```
In[79]:= n = 3;
vars = Table[x[[i]], {i, 1, n}];
f = 3  $\alpha$  x[[1]] (x[[1]]2 + x[[2]]2) +
 $\gamma_1$  (3 x[[2]]2  $\times$  x[[1]] - x[[1]]3) +  $\gamma_2$  (3 x[[1]]2  $\times$  x[[2]] - x[[2]]3) +  $\beta_3$  x[[3]]3
```

```
Out[79]= 3  $\alpha$  x[[1]] (x[[1]]2 + x[[2]]2) +  $\gamma_1$  (-x[[1]]3 + 3 x[[1]] x[[2]]2) +  $\gamma_2$  (3 x[[1]]2 x[[2]] - x[[2]]3) +  $\beta_3$  x[[3]]3
```

This is the canonical form corresponding to a partially decoupleable tensor.

```
In[80]:= Coeffs = { $\alpha$ ,  $\gamma_1$ ,  $\gamma_2$ ,  $\beta_3$ };
```

This is the list of tensor coefficients in the canonical form.

```
In[81]:=  $\Gamma$  = Simplify[Table[D[f, x[[i]], x[[j]], x[[k]], {k, 1, n}, {j, 1, n}, {i, 1, n}] / 6]
```

```
Out[81]= {{ {3  $\alpha$  -  $\gamma_1$ ,  $\gamma_2$ , 0}, { $\gamma_2$ ,  $\alpha$  +  $\gamma_1$ , 0}, {0, 0, 0}},
{{  $\gamma_2$ ,  $\alpha$  +  $\gamma_1$ , 0}, { $\alpha$  +  $\gamma_1$ , - $\gamma_2$ , 0}, {0, 0, 0}}, {{ {0, 0, 0}, {0, 0, 0}, {0, 0,  $\beta_3$ }}}
```

```
In[82]:= Table[MatrixForm[ $\Gamma$ [[i]], {i, 1, n}]
```

```
Out[82]= { {  $\begin{pmatrix} 3\alpha - \gamma_1 & \gamma_2 & 0 \\ \gamma_2 & \alpha + \gamma_1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$ ,  $\begin{pmatrix} \gamma_2 & \alpha + \gamma_1 & 0 \\ \alpha + \gamma_1 & -\gamma_2 & 0 \\ 0 & 0 & 0 \end{pmatrix}$ ,  $\begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \beta_3 \end{pmatrix}$  }
```

```
In[83]:= u = Simplify[Table[D[Laplacian[f, vars], x[i]], {i, 1, n}]/6];
f3 = (n + 2) f - 3 (u.vars) (vars.vars)
```

```
Out[84]= -3 (4 α x[1] + β3 x[3]) (x[1]^2 + x[2]^2 + x[3]^2) +
5 (3 α x[1] (x[1]^2 + x[2]^2) + γ1 (-x[1]^3 + 3 x[1] x[2]^2) + γ2 (3 x[1]^2 x[2] - x[2]^3) + β3 x[3]^3)
```

This is the trace-free part

```
In[85]:= D = Simplify[Table[D[f3, x[i], x[j], x[k]], {k, 1, n}, {j, 1, n}, {i, 1, n}]/6];
v = Simplify[TensorContract[(D⊗D)⊗D, {{1, 4}, {2, 5}, {3, 7}, {6, 8}}]];
w = Simplify[D.u.u];
```

The calculations follow the same steps as in the fully decoupled case.

```
In[88]:= Clear[H, J, K, L, M, Q];
v = Simplify[TensorContract[(D⊗D)⊗D, {{1, 4}, {2, 5}, {3, 7}, {6, 8}}]];
w = Simplify[D.u.u];
Q = Simplify[TensorContract[D⊗D, {{1, 4}, {2, 5}}]];
γuu = w;
Trivialize = Table[Coeffs[j] → 0, {j, 1, Length[Coeffs]}];
```

In what follows, we use the decomposition $\mathcal{I} = \mathcal{I}^+ \cup \mathcal{I}'$, corresponding to the ‘Fundamental’ and ‘Secondary’ invariants. This is not standard terminology!!

```
In[94]:= FundamentalRelations = {H[2] - Simplify[Tr[Q]],
H[4] - Simplify[Tr[Q.Q]], J[2] - Simplify[u.u], L[4] - Simplify[γuu.u]};
FundamentalInvariants = FundamentalRelations /. Trivialize;
SecondaryRelations = {H[6] - Simplify[v.v], H[10] - Simplify[D.v.v.v],
J[4] - Simplify[u.Q.u], K[4] - Simplify[Tr[Q.(D.u)]],
J[6] - Simplify[(u.Q).γuu], K[6] - Simplify[v.w], L[6] - Simplify[(u.Q).v],
M[6] - Simplify[γuu.γuu], H[8] - Simplify[(u.Q).(Q.v)]};
SecondaryInvariants = SecondaryRelations /. Trivialize;
PartialDecoupleRelations = Join[FundamentalRelations, SecondaryRelations];
```

```
In[99]:= G = Table[Γ[i, j, k], {i, 1, 2}, {j, 1, 2}, {k, 1, 2}]
```

```
Out[99]= {{ {3 α - γ1, γ2}, {γ2, α + γ1}}, {{γ2, α + γ1}, {α + γ1, -γ2}}}
```

This is the $2 \times 2 \times 2$ block of Γ

```
In[100]:= GStarSqrD = Expand[TensorContract[G⊗G, {{1, 4}, {2, 5}}]]
```

```
Out[100]= {{ {10 α^2 - 4 α γ1 + 2 γ1^2 + 2 γ2^2, 4 α γ2}, {4 α γ2, 2 α^2 + 4 α γ1 + 2 γ1^2 + 2 γ2^2}}}
```

```
In[101]:= MatrixForm[GStarSqrD]
```

```
Out[101]//MatrixForm=
( 10 α^2 - 4 α γ1 + 2 γ1^2 + 2 γ2^2      4 α γ2
   4 α γ2      2 α^2 + 4 α γ1 + 2 γ1^2 + 2 γ2^2 )
```

```
In[102]:=
ured = TensorContract[G, {1, 2}]
```

```
Out[102]:=
{4 α, 0}
```

This is the trace of the $2 \times 2 \times 2$ block

```
In[103]:=
Clear[q];
defns = {q1 - β3^2, q2 - ured.ured, q4 - Tr[GStarSqr], q3 - Det[GStarSqr]}
```

```
Out[104]:=
{q1 - β3^2, -16 α^2 + q2, -12 α^2 + q4 - 4 γ1^2 - 4 γ2^2,
-20 α^4 + q3 - 32 α^3 γ1 - 8 α^2 γ1^2 - 4 γ1^4 - 8 α^2 γ2^2 - 8 γ1^2 γ2^2 - 4 γ2^4}
```

These are the expressions of the $O(2) \times \mathbb{Z}_2$ invariants (the quantities q_i) in terms of the parameters defining the Canonical form.

```
In[105]:=
EliminateCoeffs = Table[
  GroebnerBasis[Join[{FundamentalRelations[[i]]}, defns], Join[FundamentalInvariants,
    {q3, q4, q2, q1}], Coeffs, MonomialOrder → EliminationOrder][[1]], {i, 1, 4}]
```

```
Out[105]:=
{H[2] - 10 q1 + 15 q2 - 25 q4,
-H[4] + 44 q1^2 - 42 q1 q2 + 144 q2^2 - 30 q3 + 100 q1 q4 - 420 q2 q4 + 320 q4^2,
J[2] - q1 - q2, -2 L[4] + 4 q1^2 - 12 q1 q2 + 4 q2^2 - 20 q3 - 5 q2 q4 + 5 q4^2}
```

These are the relations right after (7.2) expressing the elements in \mathcal{I}^+ in terms of the $O(2) \times \mathbb{Z}_2$ invariants q_i . We can invert these relations and solve for the quantities q_i .

```
In[106]:=
TriangularSystem =
  GroebnerBasis[EliminateCoeffs, {q4, q3, q2, q1, H[2], H[4], J[2], L[4]}]
```

```
Out[106]:=
{-H[2]^2 + 2 H[4] + 3 H[2] × J[2] - 6 J[2]^2 - 6 L[4] + 9 H[2] q1 - 90 J[2] q1,
-J[2] + q1 + q2, 8 H[2]^2 - 25 H[4] - 60 H[2] × J[2] - 1500 J[2]^2 + 1200 L[4] +
11 250 J[2] q1 - 11 250 q1^2 + 11 250 q3, -H[2] - 15 J[2] + 25 q1 + 25 q4}
```

```
In[107]:=
Substitutions = Table[Solve[TriangularSystem[[i]] == 0, qi][[1, 1]], {i, 1, 4}]
```

```
Out[107]:=
{q1 →  $\frac{H[2]^2 - 2 H[4] - 3 H[2] \times J[2] + 6 J[2]^2 + 6 L[4]}{9 (H[2] - 10 J[2])}$ , q2 → J[2] - q1,
q3 →  $\frac{-8 H[2]^2 + 25 H[4] + 60 H[2] \times J[2] + 1500 J[2]^2 - 1200 L[4] - 11 250 J[2] q1 + 11 250 q1^2}{11 250}$ ,
q4 →  $\frac{1}{25} (H[2] + 15 J[2] - 25 q1)$ }
```

These are the substitutions implied by Eq. (7.3).

Lemma 7.3. Finding real solutions for the parameters in terms of the q tildes and

the associated inequalities that are needed for solvability

```
In[108]:= GroebnerBasis[defs, Join[{β3}, Table[qi, {i, 1, 4}]], {α, γ1, γ2}]
```

```
Out[108]= {-q1 + β32}
```

We need $q_1 \geq 0$ to solve for a real β_3 .

```
In[109]:= GroebnerBasis[defs, Join[{γ2}, Table[qi, {i, 1, 4}]], {α, γ1, β3}]
```

```
Out[109]= {q24 - 4 q22 q3 + 16 q32 - 2 q23 q4 + 8 q2 q3 q4 + 2 q22 q42 - 8 q3 q42 - 2 q2 q43 + q44 + 4 q23 γ22}
```

```
In[110]:= Collect[%[[1]], γ2]
```

```
Out[110]= q24 - 4 q22 q3 + 16 q32 - 2 q23 q4 + 8 q2 q3 q4 + 2 q22 q42 - 8 q3 q42 - 2 q2 q43 + q44 + 4 q23 γ22
```

We get a linear equation for γ_2^2 .

```
In[111]:= Eqnγ2 = (% /. {γ22 → γ2sqrd}) == 0
```

```
Out[111]= 4 γ2sqrd q23 + q24 - 4 q22 q3 + 16 q32 - 2 q23 q4 + 8 q2 q3 q4 + 2 q22 q42 - 8 q3 q42 - 2 q2 q43 + q44 == 0
```

```
In[112]:= Solve[Eqnγ2, γ2sqrd][[1]]
```

```
Out[112]= {γ2sqrd →  $\frac{-q_2^4 + 4 q_2^2 q_3 - 16 q_3^2 + 2 q_2^3 q_4 - 8 q_2 q_3 q_4 - 2 q_2^2 q_4^2 + 8 q_3 q_4^2 + 2 q_2 q_4^3 - q_4^4}{4 q_2^3}$ }
```

To get a real solution, we therefore need $q_2 \neq 0$ and the fraction (or equivalently the product of the numerator and the denominator in the above expression) is non-negative. We see below that $q_2 > 0$ is necessary, and along with this condition, we will need that the numerator be greater than or equal to zero.

```
In[113]:= Eqnsααγ = GroebnerBasis[defs, Join[{α, γ1}, Table[qi, {i, 1, 4}]], {γ2, β3}]
```

```
Out[113]= {-q24 + 16 q22 q3 - 64 q32 + 4 q23 q4 - 32 q2 q3 q4 - 8 q22 q42 + 32 q3 q42 + 8 q2 q43 - 4 q44 + 16 q23 γ12,  
α q22 - 8 α q3 - 2 α q2 q4 + 2 α q42 + q22 γ1},  
q23 - 8 q2 q3 - 4 q22 q4 + 16 q3 q4 + 6 q2 q42 - 4 q43 + 128 α q3 γ1 - 32 α q42 γ1 - 16 q22 γ12,  
q22 - 8 q3 - 2 q2 q4 + 2 q42 + 16 α q2 γ1}, 16 α2 - q2}}
```

```
In[114]:= Eqnsααγ[[5]] == 0
```

```
Out[114]= 16 α2 - q2 == 0
```

To find a real solution, we need $q_2 \geq 0$. This, along with the earlier requirement $q_2 \neq 0$ implies that $q_2 > 0$.

In[115]:=

Eqnsaα[[2]] == 0

Out[115]=

$$\alpha q_2^2 - 8 \alpha q_3 - 2 \alpha q_2 q_4 + 2 \alpha q_4^2 + q_2^2 \gamma_1 == 0$$

In[116]:=

Simplify[Solve[Eqnsaα[[2]] == 0, γ₁][[1]]]

Out[116]=

$$\left\{ \gamma_1 \rightarrow -\frac{\alpha (q_2^2 - 8 q_3 - 2 q_2 q_4 + 2 q_4^2)}{q_2^2} \right\}$$

We get no further conditions from the solvability for γ_1

Example 7.4

In[117]:=

```
n = 3;
vars = Table[x[i], {i, 1, n}];
f = Sum[2 i x[i]^3, {i, 1, n}] + (3 x[1]^2 x[2] - x[2]^3) - 12 x[1] x[2] x[3];
```

Out[117]=

$$2 x[1]^3 + 3 x[1]^2 x[2] + 3 x[2]^3 - 12 x[1] x[2] x[3] + 6 x[3]^3$$

This is an explicit numerical example.

In[118]:=

```
r = Simplify[Table[D[f, x[i], x[j], x[k]], {k, 1, n}, {j, 1, n}, {i, 1, n}] / 6];
u = Simplify[Table[D[Laplacian[f, vars], x[i]], {i, 1, n}] / 6];
f3 = (n + 2) f - 3 (u.vars) (vars.vars);
d = Simplify[Table[D[f3, x[i], x[j], x[k]], {k, 1, n}, {j, 1, n}, {i, 1, n}] / 6];
Clear[H, J, K, L, M, Q];
v = Simplify[TensorContract[(d⊗d)⊗d, {{1, 4}, {2, 5}, {3, 7}, {6, 8}}]]];
w = Simplify[d.u.u];
Q = Simplify[TensorContract[d⊗d, {{1, 4}, {2, 5}}]]];
γuu = w;
FundamentalValues = {H[2] → Simplify[Tr[Q]],
                    H[4] → Simplify[Tr[Q.Q]], J[2] → Simplify[u.u], L[4] → Simplify[γuu.u]};
SecondaryValues =
{ H[6] → Simplify[v.v], H[10] → Simplify[d.v.v.v], J[4] → Simplify[u.Q.u],
  K[4] → Simplify[Tr[Q.(d.u)]], J[6] → Simplify[(u.Q).γuu], K[6] → Simplify[v.w],
  L[6] → Simplify[(u.Q).v], M[6] → Simplify[γuu.γuu], H[8] → Simplify[(u.Q).(Q.v)]};
```

In[129]:=

FundamentalValues

Out[129]=

$$\{H[2] \rightarrow 1060, H[4] \rightarrow 518384, J[2] \rightarrow 56, L[4] \rightarrow -4528\}$$

In[130]:=

Specialization = FundamentalRelations /. FundamentalValues

Out[130]=

$$\left\{ 1060 - 10 \left(6 \alpha^2 + \beta_3^2 + 10 \gamma_1^2 + 10 \gamma_2^2 \right), 518384 - 32 \alpha^2 \beta_3^2 - \left(32 \alpha^2 + 6 \beta_3^2 \right)^2 - \right. \\ \left. 800 \alpha^2 \gamma_2^2 - 4 \left(13 \alpha^2 + \beta_3^2 - 10 \alpha \gamma_1 + 25 \gamma_1^2 + 25 \gamma_2^2 \right)^2 - 4 \left(\beta_3^2 + (\alpha + 5 \gamma_1)^2 + 25 \gamma_2^2 \right)^2, \right. \\ \left. 56 - 16 \alpha^2 - \beta_3^2, -4528 - 2 \left(-48 \alpha^2 \beta_3^2 + \beta_3^4 + 32 \alpha^3 (3 \alpha - 5 \gamma_1) \right) \right\}$$

In[131]:=

GroebnerBasis[Specialization, Coeffs]

Out[131]=

$$\left\{ 332 + 15 \beta_3^2, 3173103609 + 125768785 \gamma_2^2, \right. \\ \left. -52993421209 + 1509225420 \gamma_1^2, 230203 \alpha - 85849 \gamma_1 \right\}$$

Section 7.5

Expressing the invariants in \mathcal{I}' in terms of \mathcal{I}^+

In[132]:=

```
n = 3; Clear[x];
vars = Table[x[i], {i, 1, n}];
f = 3 α x[1] (x[1]^2 + x[2]^2) +
γ1 (3 x[2]^2 x[1] - x[1]^3) + γ2 (3 x[1]^2 x[2] - x[2]^3) + β3 x[3]^3
```

Out[132]=

$$3 \alpha x[1] \left(x[1]^2 + x[2]^2 \right) + \gamma_1 \left(-x[1]^3 + 3 x[1] x[2]^2 \right) + \gamma_2 \left(3 x[1]^2 x[2] - x[2]^3 \right) + \beta_3 x[3]^3$$

In[133]:=

r = Simplify[Table[D[f, x[i], x[j], x[k]], {k, 1, n}, {j, 1, n}, {i, 1, n}] / 6]

Out[133]=

$$\left\{ \left\{ 3 \alpha - \gamma_1, \gamma_2, 0 \right\}, \left\{ \gamma_2, \alpha + \gamma_1, 0 \right\}, \left\{ 0, 0, 0 \right\} \right\}, \\ \left\{ \left\{ \gamma_2, \alpha + \gamma_1, 0 \right\}, \left\{ \alpha + \gamma_1, -\gamma_2, 0 \right\}, \left\{ 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0 \right\}, \left\{ 0, 0, 0 \right\}, \left\{ 0, 0, \beta_3 \right\} \right\} \right\}$$

In[134]:=

```

u = Simplify[Table[D[Laplacian[f, vars], x[i]], {i, 1, n}]/6];
f3 = f - 3 (u.vars) (vars.vars) / 5;
D = Simplify[5 Table[D[f3, x[i], x[j], x[k]], {k, 1, n}, {j, 1, n}, {i, 1, n}]/6];
Clear[H, J, K, L, M, Q];
v = Simplify[TensorContract[(D⊗D)⊗D, {{1, 4}, {2, 5}, {3, 7}, {6, 8}}]];
w = Simplify[D.u.u];
Q = Simplify[TensorContract[D⊗D, {{1, 4}, {2, 5}}]];
γuu = w;
Coeffs = {α, γ1, γ2, β3};
Eliminated = {α, γ1, γ2};
Trivialize = Table[Coeffs[[j]] → 0, {j, 1, Length[Coeffs]}];
FundamentalRelations =
  {H[2] - Simplify[Tr[Q]], J[2] - Simplify[u.u], L[4] - Simplify[γuu.u]};

FundamentalRelations = {H[2] - Simplify[Tr[Q]],
  J[2] - Simplify[u.u], H[4] - Simplify[Tr[Q.Q]], L[4] - Simplify[γuu.u]};
FundamentalInvariants = Join[{β3}, (FundamentalRelations /. Trivialize)];
SecondaryRelations = {J[4] - Simplify[u.Q.u],
  K[4] - Simplify[Tr[Q.(D.u)]], H[6] - Simplify[v.v], J[6] - Simplify[(u.Q).γuu],
  K[6] - Simplify[v.w], L[6] - Simplify[(u.Q).v], M[6] - Simplify[γuu.γuu],
  H[8] - Simplify[(u.Q).(Q.v)], H[10] - Simplify[D.v.v.v]};
SecondaryInvariants = SecondaryRelations /. Trivialize;
PartialDecoupleRelations = Join[FundamentalRelations, SecondaryRelations];
OAINvariantsList = Join[FundamentalInvariants, SecondaryInvariants];

```

In[152]:=

```
GroebnerBasis[FundamentalRelations, FundamentalInvariants, Eliminated]
```

Out[152]=

$$\{-H[2]^2 + 2 H[4] + 3 H[2] \times J[2] - 6 J[2]^2 - 6 L[4] + 9 H[2] \beta_3^2 - 90 J[2] \beta_3^2\}$$

In[153]:=

```

NeccSuffRelations =
  Join[GroebnerBasis[FundamentalRelations, FundamentalInvariants, Eliminated],
    Table[GroebnerBasis[Join[{SecondaryRelations[[i]]}, FundamentalRelations], Join[
      {SecondaryInvariants[[i]]}, FundamentalInvariants], Eliminated][[2]], {i, 1, 9}]]

```

Out[153]=

$$\begin{aligned}
 & \{-H[2]^2 + 2H[4] + 3H[2] \times J[2] - 6J[2]^2 - 6L[4] + 9H[2] \beta_3^2 - 90J[2] \beta_3^2, \\
 & H[2]^2 - 2H[4] - 12H[2] \times J[2] + 24J[2]^2 + 18J[4] - 12L[4], \\
 & 2H[2]^2 - 4H[4] - 15H[2] \times J[2] + 66J[2]^2 + 9K[4] - 6L[4], 13H[2]^3 - 26H[2] \times H[4] + \\
 & 27H[6] + 7H[2]^2 J[2] - 146H[4] \times J[2] + 156H[2] J[2]^2 - 372J[2]^3 - 30H[2] \times L[4] + \\
 & 924J[2] \times L[4] - 234H[4] \beta_3^2 + 7272J[2]^2 \beta_3^2 + 1512L[4] \beta_3^2 + 8100J[2] \beta_3^4 - 8100\beta_3^6, \\
 & H[2]^3 - 2H[2] \times H[4] - 6H[4] \times J[2] - 12H[2] J[2]^2 + 72J[2]^3 + 36J[6] - \\
 & 12H[2] \times L[4] - 18H[4] \beta_3^2 + 144J[2]^2 \beta_3^2 + 324L[4] \beta_3^2 + 2700J[2] \beta_3^4 - 2700\beta_3^6, \\
 & 19H[2]^3 - 38H[2] \times H[4] - 14H[2]^2 J[2] - 86H[4] \times J[2] - 96H[2] J[2]^2 + \\
 & 744J[2]^3 + 162K[6] - 48H[2] \times L[4] + 744J[2] \times L[4] - 342H[4] \beta_3^2 + \\
 & 8136J[2]^2 \beta_3^2 + 3456L[4] \beta_3^2 + 24300J[2] \beta_3^4 - 24300\beta_3^6, \\
 & 19H[2]^3 - 38H[2] \times H[4] - 23H[2]^2 J[2] - 149H[4] \times J[2] + 174H[2] J[2]^2 + \\
 & 204J[2]^3 - 48H[2] \times L[4] + 852J[2] \times L[4] + 81L[6] - 261H[4] \beta_3^2 + 7488J[2]^2 \beta_3^2 + \\
 & 1998L[4] \beta_3^2 + 12150J[2] \beta_3^4 - 12150\beta_3^6, -H[2]^3 + 2H[2] \times H[4] + 14H[2]^2 J[2] - \\
 & 22H[4] \times J[2] - 120H[2] J[2]^2 + 552J[2]^3 - 6H[2] \times L[4] - 420J[2] \times L[4] + \\
 & 324M[6] + 18H[4] \beta_3^2 - 5544J[2]^2 \beta_3^2 + 2376L[4] \beta_3^2 + 24300J[2] \beta_3^4 - 24300\beta_3^6, \\
 & -70H[2]^4 + 361H[2]^2 H[4] - 442H[4]^2 + 1458H[8] - 149H[2]^3 J[2] - \\
 & 674H[2] \times H[4] \times J[2] - 830H[2]^2 J[2]^2 + 7216H[4] J[2]^2 + 8868H[2] J[2]^3 - \\
 & 30792J[2]^4 - 30H[2]^2 L[4] - 3828H[4] \times L[4] - 408H[2] \times J[2] \times L[4] + 3624J[2]^2 L[4] + \\
 & 11088L[4]^2 - 5796H[4] \times J[2] \beta_3^2 - 158832J[2]^3 \beta_3^2 + 206928J[2] \times L[4] \beta_3^2 + \\
 & 40500H[4] \beta_3^4 + 356400J[2]^2 \beta_3^4 - 121500L[4] \beta_3^4 - 1895400J[2] \beta_3^6, \\
 & -70H[2]^5 + 879H[2]^3 H[4] - 1478H[2] H[4]^2 + 2187H[10] + 1161H[2]^4 J[2] - \\
 & 2866H[2]^2 H[4] \times J[2] - 2506H[4]^2 J[2] - 110H[2]^3 J[2]^2 - 12212H[2] \times H[4] J[2]^2 + \\
 & 159160H[2]^2 J[2]^3 - 230888H[4] J[2]^3 - 607632H[2] J[2]^4 + 1226976J[2]^5 + \\
 & 630H[2]^3 L[4] - 2040H[2] \times H[4] \times L[4] - 204H[2]^2 J[2] \times L[4] + \\
 & 26160H[4] \times J[2] \times L[4] + 125412H[2] J[2]^2 L[4] + 423096J[2]^3 L[4] - 2448H[2] L[4]^2 + \\
 & 35928J[2] L[4]^2 - 10782H[4]^2 \beta_3^2 + 95256H[4] J[2]^2 \beta_3^2 + 9237600J[2]^4 \beta_3^2 + \\
 & 71496H[4] \times L[4] \beta_3^2 + 505440J[2]^2 L[4] \beta_3^2 + 35640L[4]^2 \beta_3^2 + 251100H[4] \times J[2] \beta_3^4 + \\
 & 4017600J[2]^3 \beta_3^4 + 777600J[2] \times L[4] \beta_3^4 - 510300H[4] \beta_3^6 + 5832000J[2]^2 \beta_3^6\}
 \end{aligned}$$

In[154]:=

```
Map[Length, NeccSuffRelations, {1}]
```

Out[154]=

```
{7, 6, 6, 14, 12, 14, 14, 14, 22, 32}
```

In[155]:=

```
SecondaryInvariants = Join[{H[4]}, SecondaryInvariants]
```

Out[155]=

```
{H[4], J[4], K[4], H[6], J[6], K[6], L[6], M[6], H[8], H[10]}
```

Redefine Secondary invariants to include H_4

In[156]:=

```
Normalizations = Table[D[NeccSuffRelations[[i]], SecondaryInvariants[[i]], {i, 1, 10}]
```

Out[156]=

```
{2, 18, 9, 27, 36, 162, 81, 324, 1458, 2187}
```

In[157]:=

```
SolveNeccSuff =
```

```
Table[Solve[NeccSuffRelations[[i]] == 0, SecondaryInvariants[[i]][[1, 1]], {i, 1, 10}] /.  
  { $\beta_3^{k_-} \rightarrow q_1^{k/2}$ }
```

Out[157]=

$$\begin{aligned}
& \left\{ H[4] \rightarrow \frac{1}{2} \left(H[2]^2 - 3 H[2] \times J[2] + 6 J[2]^2 + 6 L[4] - 9 H[2] q_1 + 90 J[2] q_1 \right), \right. \\
& J[4] \rightarrow \frac{1}{18} \left(-H[2]^2 + 2 H[4] + 12 H[2] \times J[2] - 24 J[2]^2 + 12 L[4] \right), \\
& K[4] \rightarrow \frac{1}{9} \left(-2 H[2]^2 + 4 H[4] + 15 H[2] \times J[2] - 66 J[2]^2 + 6 L[4] \right), \\
& H[6] \rightarrow \frac{1}{27} \left(-13 H[2]^3 + 26 H[2] \times H[4] - 7 H[2]^2 J[2] + 146 H[4] \times J[2] - \right. \\
& \quad 156 H[2] J[2]^2 + 372 J[2]^3 + 30 H[2] \times L[4] - 924 J[2] \times L[4] + \\
& \quad 234 H[4] q_1 - 7272 J[2]^2 q_1 - 1512 L[4] q_1 - 8100 J[2] q_1^2 + 8100 q_1^3 \Big), \\
& J[6] \rightarrow \frac{1}{36} \left(-H[2]^3 + 2 H[2] \times H[4] + 6 H[4] \times J[2] + 12 H[2] J[2]^2 - 72 J[2]^3 + \right. \\
& \quad 12 H[2] \times L[4] + 18 H[4] q_1 - 144 J[2]^2 q_1 - 324 L[4] q_1 - 2700 J[2] q_1^2 + 2700 q_1^3 \Big), \\
& K[6] \rightarrow \frac{1}{162} \left(-19 H[2]^3 + 38 H[2] \times H[4] + 14 H[2]^2 J[2] + 86 H[4] \times J[2] + \right. \\
& \quad 96 H[2] J[2]^2 - 744 J[2]^3 + 48 H[2] \times L[4] - 744 J[2] \times L[4] + \\
& \quad 342 H[4] q_1 - 8136 J[2]^2 q_1 - 3456 L[4] q_1 - 24300 J[2] q_1^2 + 24300 q_1^3 \Big), \\
& L[6] \rightarrow \frac{1}{81} \left(-19 H[2]^3 + 38 H[2] \times H[4] + 23 H[2]^2 J[2] + 149 H[4] \times J[2] - \right. \\
& \quad 174 H[2] J[2]^2 - 204 J[2]^3 + 48 H[2] \times L[4] - 852 J[2] \times L[4] + \\
& \quad 261 H[4] q_1 - 7488 J[2]^2 q_1 - 1998 L[4] q_1 - 12150 J[2] q_1^2 + 12150 q_1^3 \Big), \\
& M[6] \rightarrow \frac{1}{324} \left(H[2]^3 - 2 H[2] \times H[4] - 14 H[2]^2 J[2] + 22 H[4] \times J[2] + \right. \\
& \quad 120 H[2] J[2]^2 - 552 J[2]^3 + 6 H[2] \times L[4] + 420 J[2] \times L[4] - \\
& \quad 18 H[4] q_1 + 5544 J[2]^2 q_1 - 2376 L[4] q_1 - 24300 J[2] q_1^2 + 24300 q_1^3 \Big), \\
& H[8] \rightarrow \frac{1}{1458} \left(70 H[2]^4 - 361 H[2]^2 H[4] + 442 H[4]^2 + 149 H[2]^3 J[2] + 674 H[2] \times H[4] \times J[2] + \right. \\
& \quad 830 H[2]^2 J[2]^2 - 7216 H[4] J[2]^2 - 8868 H[2] J[2]^3 + 30792 J[2]^4 + \\
& \quad 30 H[2]^2 L[4] + 3828 H[4] \times L[4] + 408 H[2] \times J[2] \times L[4] - 3624 J[2]^2 L[4] - \\
& \quad 11088 L[4]^2 + 5796 H[4] \times J[2] q_1 + 158832 J[2]^3 q_1 - 206928 J[2] \times L[4] q_1 - \\
& \quad 40500 H[4] q_1^2 - 356400 J[2]^2 q_1^2 + 121500 L[4] q_1^2 + 1895400 J[2] q_1^3 \Big), \\
& H[10] \rightarrow \frac{1}{2187} \left(70 H[2]^5 - 879 H[2]^3 H[4] + 1478 H[2] H[4]^2 - 1161 H[2]^4 J[2] + \right. \\
& \quad 2866 H[2]^2 H[4] \times J[2] + 2506 H[4]^2 J[2] + 110 H[2]^3 J[2]^2 + 12212 H[2] \times H[4] J[2]^2 - \\
& \quad 159160 H[2]^2 J[2]^3 + 230888 H[4] J[2]^3 + 607632 H[2] J[2]^4 - 1226976 J[2]^5 - \\
& \quad 630 H[2]^3 L[4] + 2040 H[2] \times H[4] \times L[4] + 204 H[2]^2 J[2] \times L[4] - \\
& \quad 26160 H[4] \times J[2] \times L[4] - 125412 H[2] J[2]^2 L[4] - 423096 J[2]^3 L[4] + 2448 H[2] L[4]^2 - \\
& \quad 35928 J[2] L[4]^2 + 10782 H[4]^2 q_1 - 95256 H[4] J[2]^2 q_1 - 9237600 J[2]^4 q_1 - \\
& \quad 71496 H[4] \times L[4] q_1 - 505440 J[2]^2 L[4] q_1 - 35640 L[4]^2 q_1 - 251100 H[4] \times J[2] q_1^2 - \\
& \quad \left. 4017600 J[2]^3 q_1^2 - 777600 J[2] \times L[4] q_1^2 + 510300 H[4] q_1^3 - 5832000 J[2]^2 q_1^3 \right\}
\end{aligned}$$

In[158]:=

```
FormatAsEquations = Table[Normalizations[[i]] × SecondaryInvariants[[i]] → Collect[
  (Normalizations[[i]] × SecondaryInvariants[[i]] /. SolveNeccSuff), q1], {i, 1, 10}]
```

Out[158]=

```
{2 H[4] → H[2]2 - 3 H[2] × J[2] + 6 J[2]2 + 6 L[4] + (-9 H[2] + 90 J[2]) q1,
 18 J[4] → -H[2]2 + 2 H[4] + 12 H[2] × J[2] - 24 J[2]2 + 12 L[4],
 9 K[4] → -2 H[2]2 + 4 H[4] + 15 H[2] × J[2] - 66 J[2]2 + 6 L[4],
 27 H[6] → -13 H[2]3 + 26 H[2] × H[4] - 7 H[2]2 J[2] + 146 H[4] × J[2] -
 156 H[2] J[2]2 + 372 J[2]3 + 30 H[2] × L[4] - 924 J[2] × L[4] +
 (234 H[4] - 7272 J[2]2 - 1512 L[4]) q1 - 8100 J[2] q12 + 8100 q13,
 36 J[6] → -H[2]3 + 2 H[2] × H[4] + 6 H[4] × J[2] + 12 H[2] J[2]2 - 72 J[2]3 +
 12 H[2] × L[4] + (18 H[4] - 144 J[2]2 - 324 L[4]) q1 - 2700 J[2] q12 + 2700 q13,
 162 K[6] → -19 H[2]3 + 38 H[2] × H[4] + 14 H[2]2 J[2] + 86 H[4] × J[2] +
 96 H[2] J[2]2 - 744 J[2]3 + 48 H[2] × L[4] - 744 J[2] × L[4] +
 (342 H[4] - 8136 J[2]2 - 3456 L[4]) q1 - 24 300 J[2] q12 + 24 300 q13,
 81 L[6] → -19 H[2]3 + 38 H[2] × H[4] + 23 H[2]2 J[2] + 149 H[4] × J[2] -
 174 H[2] J[2]2 - 204 J[2]3 + 48 H[2] × L[4] - 852 J[2] × L[4] +
 (261 H[4] - 7488 J[2]2 - 1998 L[4]) q1 - 12 150 J[2] q12 + 12 150 q13,
 324 M[6] → H[2]3 - 2 H[2] × H[4] - 14 H[2]2 J[2] + 22 H[4] × J[2] +
 120 H[2] J[2]2 - 552 J[2]3 + 6 H[2] × L[4] + 420 J[2] × L[4] +
 (-18 H[4] + 5544 J[2]2 - 2376 L[4]) q1 - 24 300 J[2] q12 + 24 300 q13,
 1458 H[8] → 70 H[2]4 - 361 H[2]2 H[4] + 442 H[4]2 + 149 H[2]3 J[2] + 674 H[2] × H[4] × J[2] +
 830 H[2]2 J[2]2 - 7216 H[4] J[2]2 - 8868 H[2] J[2]3 + 30 792 J[2]4 +
 30 H[2]2 L[4] + 3828 H[4] × L[4] + 408 H[2] × J[2] × L[4] - 3624 J[2]2 L[4] -
 11 088 L[4]2 + (5796 H[4] × J[2] + 158 832 J[2]3 - 206 928 J[2] × L[4]) q1 +
 (-40 500 H[4] - 356 400 J[2]2 + 121 500 L[4]) q12 + 1 895 400 J[2] q13,
 2187 H[10] → 70 H[2]5 - 879 H[2]3 H[4] + 1478 H[2] H[4]2 - 1161 H[2]4 J[2] +
 2866 H[2]2 H[4] × J[2] + 2506 H[4]2 J[2] + 110 H[2]3 J[2]2 + 12 212 H[2] × H[4] J[2]2 -
 159 160 H[2]2 J[2]3 + 230 888 H[4] J[2]3 + 607 632 H[2] J[2]4 - 1 226 976 J[2]5 -
 630 H[2]3 L[4] + 2040 H[2] × H[4] × L[4] + 204 H[2]2 J[2] × L[4] - 26 160 H[4] × J[2] × L[4] -
 125 412 H[2] J[2]2 L[4] - 423 096 J[2]3 L[4] + 2448 H[2] L[4]2 - 35 928 J[2] L[4]2 +
 (10 782 H[4]2 - 95 256 H[4] J[2]2 - 9 237 600 J[2]4 - 71 496 H[4] × L[4] - 505 440 J[2]2 L[4] -
 35 640 L[4]2) q1 + (-251 100 H[4] × J[2] - 4 017 600 J[2]3 - 777 600 J[2] × L[4])
 q12 + (510 300 H[4] - 5 832 000 J[2]2) q13}
```