

FDP 2022 Classical and quantum mechanics - Finals

24 April, 2021

The duration of the test is **3 hours**, and it will account for 40% of your final grade. There are 2 exercises and 3 problems in total, with uniform weightage.

Exercises

1) Lagrangian and Newtonian formalisms

Write down the Lagrangians and recover Newton's equations (without solving them) for the following situations:

(a) 1D simple harmonic oscillator: $F = -kx$

(b) Charged particle in electromagnetic field: $\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$, where:

$$\vec{E} = -\vec{\nabla}\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \quad (1)$$

$$\vec{B} = \vec{\nabla} \times \vec{A}. \quad (2)$$

2) Properties of Gaussian wavepackets

Consider a 1D Gaussian wavepacket centered at the origin, with momentum p_0 :

$$\psi(x) = e^{ip_0x/\hbar} \frac{e^{(-x'^2/(2\Delta^2))}}{(\pi\Delta^2)^{1/4}}. \quad (3)$$

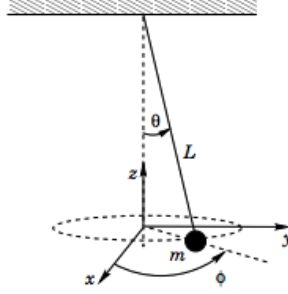
Using expressions for Gaussian integrals,

- (a) Find the probability distribution $P_\psi(x)$ of the wavepacket, and show that it is normalized to unity.
- (b) Show that the packet has a mean position $\langle X \rangle_\psi = 0$, mean momentum $\langle P \rangle_\psi = p_0$, position uncertainty $(\Delta X)_\psi = \Delta/2^{1/2}$, and momentum uncertainty $(\Delta P)_\psi = \hbar/(2^{1/2}\Delta)$.
- (c) Show that this packet saturates the lower bound given by Heisenberg's uncertainty relation.

Problems

1) Spherical pendulum

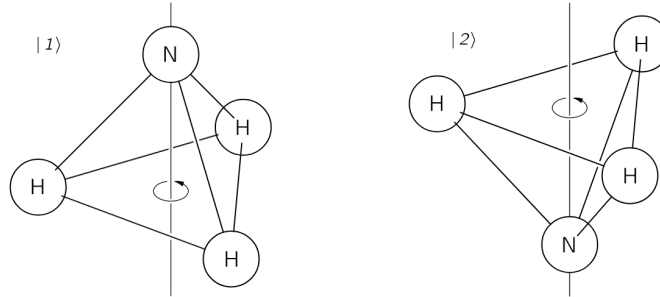
A pendulum with mass m and a massless stick of fixed length L moves in three dimensions under the influence of a constant gravitational field, with gravitational acceleration g :



- Working with spherical coordinates (r, θ, ϕ) as labelled in the diagram, write down the Lagrangian of the problem. How many degrees of freedom does the problem have?
- Derive the pendulum's equations of motion.
- Show that the magnitude of the angular momentum about the z axis, L_z , is equal to the generalized momentum, p_ϕ . Why is the generalized momentum conserved in time?
- Solve the equations of motion in the special case where $\dot{\phi} = 0$, and θ is small. What situation does this correspond to?

2) Ammonia molecule

Consider a free ammonia molecule at rest, which can be described by a vector in a two-dimensional Hilbert space. The orthonormal basis $\{|1\rangle, |2\rangle\}$ formed by basis vectors $|1\rangle$ and $|2\rangle$ correspond to the following geometrical configurations of the molecule:



The Hamiltonian of the molecule in this basis is represented as:

$$\hat{H} = \begin{pmatrix} E_0 & -A \\ -A & E_0 \end{pmatrix}, \quad (4)$$

where E_0 and A are positive real numbers. In this problem, we study the stability of these geometric configurations of ammonia.

- Show that the eigenvalues of \hat{H} are $E_0 + A$ and $E_0 - A$, with corresponding eigenvectors $|E_1\rangle = \frac{1}{\sqrt{2}}(|1\rangle - |2\rangle)$ and $|E_2\rangle = \frac{1}{\sqrt{2}}(|1\rangle + |2\rangle)$.
- Construct the propagator, $\hat{U}(t)$.
- Suppose the initial state $|\psi(0)\rangle$ of the molecule is $|1\rangle$. Determine its time evolution, i.e. find $|\psi(t)\rangle$.
- Determine the probability of finding the state to be $|1\rangle$ at time t .
- Sketch the probability as a function of time. What can you say about the stability of the initial state $|1\rangle$? Give an example of a state that is stable.

3) Canonical brackets

- (a) Show the *canonical commutation relation* between the position and momentum operators, which is:

$$[\hat{x}, \hat{p}] = i\hbar \hat{1}, \quad (5)$$

where $[\hat{A}, \hat{B}] \equiv \hat{A}\hat{B} - \hat{B}\hat{A}$ is the commutator between two arbitrary operators \hat{A} and \hat{B} . (Hint: The commutator is an operator. Work in the position basis and apply this operator to an arbitrary state $|\psi\rangle$.)

- (b) The Heisenberg-Robertson inequality for arbitrary operators \hat{A} and \hat{B} is given by:

$$(\Delta A)_\rho^2 (\Delta B)_\rho^2 \geq \frac{1}{4} \langle i[\hat{A}, \hat{B}] \rangle_\rho^2. \quad (6)$$

Recover Heisenberg's uncertainty relation using this inequality.

- (c) The quantum mechanical canonical commutation relation has an analogue in the Hamiltonian formalism of classical mechanics. Instead of the commutator $[\hat{A}, \hat{B}]$ between two operators, the *Poisson brackets* between two classical dynamical variables $\omega(x, p)$ and $\lambda(x, p)$ are defined as:

$$\{\omega, \lambda\} \equiv \frac{\partial \omega}{\partial x} \frac{\partial \lambda}{\partial p} - \frac{\partial \omega}{\partial p} \frac{\partial \lambda}{\partial x}. \quad (7)$$

Show that:

$$\{x, p\} = 1. \quad (8)$$

- (d) Show that Hamilton's canonical equations of motion can be written in terms of the Poisson brackets:

$$\dot{x} = \{x, H\} \quad (9)$$

$$\dot{p} = \{p, H\}. \quad (10)$$