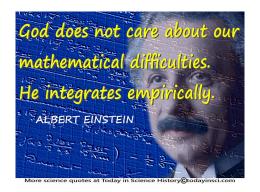
#### MA 207 - Differential Equations-II

#### Amiya Kumar Pani

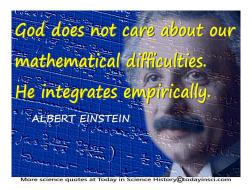
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November 13, 2020

#### Start with ......



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#### Outline of the lecture

- Vibrating string-wave equation
- Method of separation of variables
- D'Alembert's Method
- Method of Fourier transforms

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BC:  $u(0, t) = 0, \ u(l, t) = 0$ 

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- $c^2 > 0 = T/\rho$  is a constant dependent on material property,
- T: tension and  $\rho$ : density of the string.
- f(x) and g(x) are initial displacement and velocity respectively.

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$$u(x,t) = X(x)T(t)$$

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(Note that the first expression is a function of x alone, the second expression is a function of t alone, since x and t are independent variables, hence both sides must be equal to a constant.)

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That is, X(x) is the solution of the eigenvalue problem defined by:

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The cases  $\lambda = 0$  and  $\lambda < 0$  give only trivial solutions to the EVP.

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Hence the eigenvalues are  $\lambda = (\frac{n\pi}{I})^2$ 

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Hence the eigenvalues are  $\lambda = (\frac{n\pi}{I})^2$  and eigenfunctions are

$$\lambda = (\frac{n\pi}{I})^2$$

$$X_n(x) = B_n \sin(\frac{n\pi x}{I}).$$

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 $u(x,0) = f(x), \ u_t(x,0) = g(x)$ 

 $\implies T_n(t) = C_n \cos(\frac{n\pi ct}{I}) + D_n \sin(\frac{n\pi ct}{I})$ ,  $C_n$  and  $D_n$  being arbitrary constants.

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$$u(x,0) = f(x), \ u_t(x,0) = g(x)$$

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$$u_n(x,t) = X_n(x)T_n(t)$$

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$$u_n(x,t) = X_n(x)T_n(t)$$

$$= \left[a_n \cos(\frac{n\pi ct}{l}) + b_n \sin(\frac{n\pi ct}{l})\right] \sin(\frac{n\pi x}{l})$$

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Under suitable assumptions on the initial data (which will be specified later),

$$u(x,t) = \sum_{n=1}^{\infty} u_n(x,t)$$
 also satisfies the wave equation.

$$u(x,0) = f(x) \Longrightarrow \sum_{n=1}^{\infty} a_n \sin(\frac{n\pi x}{l}) = f(x)$$

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$$a_n = \frac{2}{I} \int_0^I f(x) \sin(\frac{n\pi x}{I}) dx;$$

$$u(x,0) = f(x) \Longrightarrow \sum_{n=1}^{\infty} a_n \sin(\frac{n\pi x}{l}) = f(x)$$

$$u_t(x,0) = g(x) \Longrightarrow \sum_{n=1}^{\infty} \left(\frac{n\pi c}{l}\right) b_n \sin(\frac{n\pi x}{l}) = g(x)$$

$$a_n = \frac{2}{l} \int_0^l f(x) \sin(\frac{n\pi x}{l}) dx;$$

$$b_n = \frac{2}{l} \times \frac{l}{n\pi c} \int_0^l g(x) \sin(\frac{n\pi x}{l}) dx$$

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The formal solution of the wave equation (1D- space) with Dirichlet BC & IC is given by

$$u(x,t) = \sum_{n=1}^{\infty} \left[ a_n \cos(\frac{n\pi ct}{l}) + b_n \sin(\frac{n\pi ct}{l}) \right] \sin(\frac{n\pi x}{l}), \text{ with the}$$

Fourier coefficients  $a_n$  and  $b_n$  defined as above.

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② 
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$$g(x)$$
,  $g'(x)$  are continuous in  $[0, I]$ ,

$$g(0) = g(1) = 0$$

then, the formal solution is in fact a solution.

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- g(x), g'(x) are continuous in [0, I],
- g(0) = g(1) = 0

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#### Exercise:

1. Is the solution unique?

If the function f(x) & g(x) satisfy the properties:

- f(x), f'(x), f''(x) are continuous in [0, I],
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- 2. Tut. Sheet 5: 4, 5, 6, 10 (ii), (iii)

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#### Exercise:

- 1. Is the solution unique?
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Consider  $u_{tt} = 4u_{xx}, \ 0 < x < 30, \ t > 0$  with u(0,t) = 0 = u(30,t) and initial position  $u(x,0) = f(x) := \begin{cases} x/10, & 0 \le x \le 10, \\ (30-x)/20, & 10 < x \le 30. \end{cases}$  with initial velocity  $u_t(x,0) = g(x) := 0$ .

Consider  $u_{tt} = 4u_{xx}$ , 0 < x < 30, t > 0 with u(0,t) = 0 = u(30,t) and initial position  $u(x,0) = f(x) := \begin{cases} x/10, & 0 \le x \le 10, \\ (30-x)/20, & 10 < x \le 30. \end{cases}$  with initial velocity  $u_t(x,0) = g(x) := 0$ . The solution u(x,t) is now given with  $b_n = 0$ , c = 2 and l = 30 by

$$u(x,t) = \sum_{n=1}^{\infty} a_n \cos(\frac{2n\pi t}{30}) \sin(\frac{n\pi x}{30}),$$

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$$u(x,t) = \sum_{n=1}^{\infty} a_n \cos(\frac{2n\pi t}{30}) \sin(\frac{n\pi x}{30}),$$

with the Fourier coefficients  $a_n$  defined as

$$a_n = \frac{2}{30} \int_0^{10} \frac{x}{10} \sin(\frac{n\pi x}{30}) dx + \frac{2}{30} \int_{10}^{30} \frac{30 - x}{20} \sin(\frac{n\pi x}{30}) dx$$

Consider  $u_{tt} = 4u_{xx}$ , 0 < x < 30, t > 0 with u(0, t) = 0 = u(30, t) and initial position  $\begin{cases} x/10 & 0 < x < 0 \end{cases}$ 

$$u(x,0) = f(x) := \begin{cases} x/10, & 0 \le x \le 10, \\ (30-x)/20, & 10 < x \le 30. \end{cases}$$

with initial velocity  $u_t(x,0)=g(x):=0$ . The solution u(x,t) is now given with  $b_n=0$ , c=2 and l=30 by

$$u(x,t) = \sum_{n=1}^{\infty} a_n \cos(\frac{2n\pi t}{30}) \sin(\frac{n\pi x}{30}),$$

with the Fourier coefficients  $a_n$  defined as

$$a_n = \frac{2}{30} \int_0^{10} \frac{x}{10} \sin(\frac{n\pi x}{30}) dx + \frac{2}{30} \int_{10}^{30} \frac{30 - x}{20} \sin(\frac{n\pi x}{30}) dx$$
$$= \frac{9}{n^2 \pi^2} \sin(\frac{n\pi}{3}), \quad n = 1, 2, \dots$$



# Jean-Baptiste le Rond d'Alembert (16 November 1717 - 29 October 1783)

French mathematician, mechanician, physicist, philosopher, and music theorist.

Until 1759 he was, together with Denis Diderot, a co-editor of the Encyclopédie.

D'Alembert's formula for obtaining solutions to the wave equation is named after him. The wave equation is sometimes referred to as d'Alembert's equation,

The fundamental theorem of algebra is known as the d'Alembert/Gauss theorem, as an error in d'Alembert's proof was caught by Gauss. He also created his ratio test, a test to see if a series converges.

The D'Alembert operator, which first arose in D'Alembert's analysis of vibrating strings, plays an important role in modern theoretical physics.



Consider the wave equation  $u_{tt} - c^2 u_{xx} = 0$ .

$$u_{tt}-c^2u_{xx}=0.$$

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$$\left(\omega^{\pm}=\pm c,\; \frac{dx}{dt}=\pm c,\Longrightarrow x-ct=c_1,\; x+ct=c_2\right)$$

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Note that the transformations  $\xi \equiv x - ct = c_1$ ,  $\eta \equiv x + ct = c_2$  converts the equation into canonical form.

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Let  $v(\xi, \eta) = u(x, t) (= u(\xi(x, t), \eta(x, t))).$ 

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 $U_X$ 

Consider the wave equation  $u_{tt} - c^2 u_{xx} = 0$ .

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$$u_x = v_\xi \cdot \xi_x + v_\eta \cdot \eta_x$$

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$$= \frac{\partial}{\partial \xi} [v_{\xi} + v_{\eta}] \frac{\partial \xi}{\partial x} + \frac{\partial}{\partial \eta} [v_{\xi} + v_{\eta}] \frac{\partial \eta}{\partial x}$$

Consider the wave equation  $u_{tt} - c^2 u_{xx} = 0$ .

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Let  $v(\xi, \eta) = u(x, t) (= u(\xi(x, t), \eta(x, t))).$ 

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$$= \frac{\partial}{\partial \xi} [v_{\xi} + v_{\eta}] \frac{\partial \xi}{\partial x} + \frac{\partial}{\partial \eta} [v_{\xi} + v_{\eta}] \frac{\partial \eta}{\partial x}$$

That is,  $u_{xx} = v_{\xi\xi} + 2v_{\xi\eta} + v_{\eta\eta}$ .

$$u_{tt} = (\underbrace{c(-v_{\xi}+v_{\eta})}_{u_t})_t$$

$$u_{tt} = (\underbrace{c(-v_{\xi} + v_{\eta})}_{u_t})_t = \frac{\partial}{\partial t}[c(-v_{\xi} + v_{\eta})]$$

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That is,  $u_{tt} = c^2(v_{\xi\xi} - 2v_{\xi\eta} + v_{\eta\eta})$ .

$$u_{tt} = (\underbrace{c(-v_{\xi} + v_{\eta})}_{u_{t}})_{t} = \frac{\partial}{\partial t}[c(-v_{\xi} + v_{\eta})]$$
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That is,  $u_{tt} = c^2(v_{\xi\xi} - 2v_{\xi\eta} + v_{\eta\eta}).$ 

Substituting the expressions for  $u_{xx}$  and  $u_{yy}$  in the wave equation, we obtain

$$c^{2}(v_{\xi\xi}-2v_{\xi\eta}+v_{\eta\eta})=c^{2}(v_{\xi\xi}+2v_{\xi\eta}+v_{\eta\eta})$$

$$u_{tt} = (\underbrace{c(-v_{\xi} + v_{\eta})}_{u_{t}})_{t} = \frac{\partial}{\partial t}[c(-v_{\xi} + v_{\eta})]$$
$$= \frac{\partial}{\partial \xi}[c(-v_{\xi} + v_{\eta})]\frac{\partial \xi}{\partial t} + \frac{\partial}{\partial \eta}[c(-v_{\xi} + v_{\eta})]\frac{\partial \eta}{\partial t}$$

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Substituting the expressions for  $u_{xx}$  and  $u_{yy}$  in the wave equation, we obtain

$$c^{2}(v_{\xi\xi} - 2v_{\xi\eta} + v_{\eta\eta}) = c^{2}(v_{\xi\xi} + 2v_{\xi\eta} + v_{\eta\eta})$$

$$\implies \boxed{v_{\xi\eta} = 0}$$

$$v_{\xi\eta}=0\Longrightarrowrac{\partial^2 v}{\partial \xi\partial\eta}=0\Longrightarrowrac{\partial}{\partial \xi}igg(rac{\partial v}{\partial\eta}igg)=0\Longrightarrowrac{\partial v(\xi,\eta)}{\partial\eta}=G(\eta)\,.$$

$$v_{\xi\eta} = 0 \Longrightarrow rac{\partial^2 v}{\partial \xi \partial \eta} = 0 \Longrightarrow rac{\partial}{\partial \xi} \left(rac{\partial v}{\partial \eta}\right) = 0 \Longrightarrow rac{\partial v(\xi,\eta)}{\partial \eta} = G(\eta)$$

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$$v_{\eta} = G(\eta) \implies v_{\eta}(\xi, \eta) = Q'(\eta)$$

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$$egin{aligned} v_{\eta} &= G(\eta) &\implies v_{\eta}(\xi,\eta) = Q'(\eta) \ &\implies rac{\partial}{\partial \eta} igg[ v(\xi,\eta) - Q(\eta) igg] = 0 \end{aligned}$$

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$$v_{\eta} = G(\eta) \implies v_{\eta}(\xi, \eta) = Q'(\eta)$$

$$\implies \frac{\partial}{\partial \eta} \left[ v(\xi, \eta) - Q(\eta) \right] = 0$$

$$\implies v(\xi, \eta) - Q(\eta) = P(\xi) \implies v(\xi, \eta) = P(\xi) + Q(\eta)$$

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Hence, u(x,t) = P(x-ct) + Q(x+ct).

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Hence, u(x,t) = P(x-ct) + Q(x+ct).

Now we need to determine P and Q using the initial conditions.

$$v_{\xi\eta} = 0 \Longrightarrow \frac{\partial^2 v}{\partial \xi \partial \eta} = 0 \Longrightarrow \frac{\partial}{\partial \xi} \left( \frac{\partial v}{\partial \eta} \right) = 0 \Longrightarrow \frac{\partial v(\xi, \eta)}{\partial \eta} = G(\eta)$$

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Now, 
$$u(x,0) = f(x) \Longrightarrow f(x) = P(x) + Q(x)$$
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$$v_{\xi\eta} = 0 \Longrightarrow \frac{\partial^2 v}{\partial \xi \partial \eta} = 0 \Longrightarrow \frac{\partial}{\partial \xi} \left( \frac{\partial v}{\partial \eta} \right) = 0 \Longrightarrow \frac{\partial v(\xi, \eta)}{\partial \eta} = G(\eta)$$

$$v_{\eta} = G(\eta) \implies v_{\eta}(\xi, \eta) = Q'(\eta)$$

$$\implies \frac{\partial}{\partial \eta} \left[ v(\xi, \eta) - Q(\eta) \right] = 0$$

$$\implies v(\xi, \eta) - Q(\eta) = P(\xi) \implies v(\xi, \eta) = P(\xi) + Q(\eta)$$

Hence, u(x,t) = P(x-ct) + Q(x+ct).

Now we need to determine P and Q using the initial conditions.

Now, 
$$u(x,0) = f(x) \Longrightarrow f(x) = P(x) + Q(x)$$
.  
Also,  $u_t(x,0) = g(x)$ .

$$v_{\xi\eta} = 0 \Longrightarrow \frac{\partial^2 v}{\partial \xi \partial \eta} = 0 \Longrightarrow \frac{\partial}{\partial \xi} \left( \frac{\partial v}{\partial \eta} \right) = 0 \Longrightarrow \frac{\partial v(\xi, \eta)}{\partial \eta} = G(\eta)$$

$$v_{\eta} = G(\eta) \implies v_{\eta}(\xi, \eta) = Q'(\eta)$$

$$\implies \frac{\partial}{\partial \eta} \left[ v(\xi, \eta) - Q(\eta) \right] = 0$$

$$\implies v(\xi, \eta) - Q(\eta) = P(\xi) \implies v(\xi, \eta) = P(\xi) + Q(\eta)$$

Hence, u(x,t) = P(x-ct) + Q(x+ct).

Now we need to determine P and Q using the initial conditions.

Now, 
$$u(x,0) = f(x) \Longrightarrow f(x) = P(x) + Q(x)$$
.

$$\left. \frac{\partial u}{\partial t} \right|_{t=0} = \left. \frac{\partial}{\partial t} \left[ P(x-ct) + Q(x+ct) \right] \right|_{t=0}$$

$$v_{\xi\eta} = 0 \Longrightarrow rac{\partial^2 v}{\partial \xi \partial \eta} = 0 \Longrightarrow rac{\partial}{\partial \xi} \left(rac{\partial v}{\partial \eta}\right) = 0 \Longrightarrow rac{\partial v(\xi, \eta)}{\partial \eta} = G(\eta)$$

$$v_{\eta} = G(\eta) \implies v_{\eta}(\xi, \eta) = Q'(\eta)$$

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Hence, u(x,t) = P(x-ct) + Q(x+ct).

Now we need to determine P and Q using the initial conditions.

Now, 
$$u(x,0) = f(x) \Longrightarrow f(x) = P(x) + Q(x)$$
.

$$\frac{\partial u}{\partial t}\Big|_{t=0} = \frac{\partial}{\partial t} \left[ P(x-ct) + Q(x+ct) \right] \Big|_{t=0}$$
$$= \left[ P'(x-ct)(-c) + Q'(x+ct)(c) \right] \Big|_{t=0}$$

$$v_{\xi\eta} = 0 \Longrightarrow rac{\partial^2 v}{\partial \xi \partial \eta} = 0 \Longrightarrow rac{\partial}{\partial \xi} \left(rac{\partial v}{\partial \eta}\right) = 0 \Longrightarrow rac{\partial v(\xi, \eta)}{\partial \eta} = G(\eta)$$

$$v_{\eta} = G(\eta) \implies v_{\eta}(\xi, \eta) = Q'(\eta)$$

$$\implies \frac{\partial}{\partial \eta} \left[ v(\xi, \eta) - Q(\eta) \right] = 0$$

$$\implies v(\xi, \eta) - Q(\eta) = P(\xi) \implies v(\xi, \eta) = P(\xi) + Q(\eta)$$

Hence, u(x,t) = P(x-ct) + Q(x+ct).

Now we need to determine P and Q using the initial conditions.

Now, 
$$u(x,0) = f(x) \Longrightarrow f(x) = P(x) + Q(x)$$
.

$$\frac{\partial u}{\partial t}\Big|_{t=0} = \frac{\partial}{\partial t} \left[ P(x-ct) + Q(x+ct) \right] \Big|_{t=0}$$

$$= \left[ P'(x-ct)(-c) + Q'(x+ct)(c) \right] \Big|_{t=0} = -cP'(x) + cQ'(x)$$

$$v_{\xi\eta} = 0 \Longrightarrow rac{\partial^2 v}{\partial \xi \partial \eta} = 0 \Longrightarrow rac{\partial}{\partial \xi} \left(rac{\partial v}{\partial \eta}\right) = 0 \Longrightarrow rac{\partial v(\xi, \eta)}{\partial \eta} = G(\eta)$$

$$v_{\eta} = G(\eta) \implies v_{\eta}(\xi, \eta) = Q'(\eta)$$

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Hence, u(x,t) = P(x-ct) + Q(x+ct).

Now we need to determine P and Q using the initial conditions.

Now, 
$$u(x,0) = f(x) \Longrightarrow f(x) = P(x) + Q(x)$$
.

$$\frac{\partial u}{\partial t}\Big|_{t=0} = \frac{\partial}{\partial t} \left[ P(x-ct) + Q(x+ct) \right] \Big|_{t=0}$$

$$= \left[ P'(x-ct)(-c) + Q'(x+ct)(c) \right] \Big|_{t=0} = -cP'(x) + cQ'(x)$$

Hence, 
$$g(x) = c[-P'(x) + Q'(x)]$$
.  
Integrate from 0 to  $x$  with respect to  $x$ , to obtain

$$\int_0^x g(s) \ ds = c \int_0^x -P'(x) \ dx + c \int_0^x Q'(x) \ dx$$

$$\int_{0}^{x} g(s) ds = c \int_{0}^{x} -P'(x) dx + c \int_{0}^{x} Q'(x) dx$$
$$\int_{0}^{x} g(s) ds = -c [P(x)]_{0}^{x} + c [Q(x)]_{0}^{x}$$

$$\int_{0}^{x} g(s) ds = c \int_{0}^{x} -P'(x) dx + c \int_{0}^{x} Q'(x) dx$$

$$\int_{0}^{x} g(s) ds = -c [P(x)]_{0}^{x} + c [Q(x)]_{0}^{x}$$

$$= -c [P(x) - P(0)] + c [Q(x) - Q(0)]$$

$$\int_{0}^{x} g(s) ds = c \int_{0}^{x} -P'(x) dx + c \int_{0}^{x} Q'(x) dx$$

$$\int_{0}^{x} g(s) ds = -c [P(x)]_{0}^{x} + c [Q(x)]_{0}^{x}$$

$$= -c [P(x) - P(0)] + c [Q(x) - Q(0)]$$

$$\frac{1}{c} \int_{0}^{x} g(s) ds = [(Q(x) - P(x)) - (Q(0) - P(0))]$$

$$\int_{0}^{x} g(s) ds = c \int_{0}^{x} -P'(x) dx + c \int_{0}^{x} Q'(x) dx$$

$$\int_{0}^{x} g(s) ds = -c [P(x)]_{0}^{x} + c [Q(x)]_{0}^{x}$$

$$= -c [P(x) - P(0)] + c [Q(x) - Q(0)]$$

$$\frac{1}{c} \int_{0}^{x} g(s) ds = [(Q(x) - P(x)) - (Q(0) - P(0))]$$

$$\implies \frac{1}{c} \int_{0}^{x} g(s) ds + (Q(0) - P(0)) = [Q(x) - P(x)]$$

$$\int_{0}^{x} g(s) ds = c \int_{0}^{x} -P'(x) dx + c \int_{0}^{x} Q'(x) dx$$

$$\int_{0}^{x} g(s) ds = -c [P(x)]_{0}^{x} + c [Q(x)]_{0}^{x}$$

$$= -c [P(x) - P(0)] + c [Q(x) - Q(0)]$$

$$\frac{1}{c} \int_{0}^{x} g(s) ds = [(Q(x) - P(x)) - (Q(0) - P(0))]$$

$$\implies \frac{1}{c} \int_{0}^{x} g(s) ds + (Q(0) - P(0)) = [Q(x) - P(x)]$$

$$\frac{1}{c} \int_0^x g(s) \, ds + (Q(0) - P(0)) = [Q(x) - P(x)]$$

Adding and subtracting the above two equations, we have

$$Q(x) = \frac{1}{2} \left[ f(x) + \frac{1}{c} \int_0^x g(s) \, ds + (Q(0) - P(0)) \right]$$

$$\frac{1}{c} \int_0^x g(s) \, ds + (Q(0) - P(0)) = [Q(x) - P(x)]$$

Adding and subtracting the above two equations, we have

$$Q(x) = \frac{1}{2} \left[ f(x) + \frac{1}{c} \int_0^x g(s) \, ds + (Q(0) - P(0)) \right]$$

$$P(x) = \frac{1}{2} \left[ f(x) - \frac{1}{c} \int_0^x g(s) \, ds - (Q(0) - P(0)) \right]$$

$$\frac{1}{c} \int_0^x g(s) \, ds + (Q(0) - P(0)) = [Q(x) - P(x)]$$

Adding and subtracting the above two equations, we have

$$Q(x) = \frac{1}{2} \left[ f(x) + \frac{1}{c} \int_0^x g(s) \, ds + (Q(0) - P(0)) \right]$$

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The solution to wave equation is given by

$$u(x,t) = P(x-ct) + Q(x+ct)$$

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Hence, 
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This is called D'Alembert's solution to the wave equation. BC's are not considered!

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$$u(x,t) = \frac{k}{2} \left( \sin \pi (x - ct) + \sin \pi (x + ct) \right)$$
  
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Use Plancherel's identity to find a relation between the solution and the initial data (in terms of  $L^2$ -norm) of  $u_{tt} = u_{xx}$  with  $u(x,0) = u_0$ ,  $u_t(x,0) = u_1 = 0$ .

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$$\begin{aligned} \hat{u}_{tt} + \omega^2 \hat{u} &= 0, \ \hat{u}(\omega, 0) = \hat{u}_0, \ \hat{u}_t(\omega, 0) = \hat{u}_1 \\ \Longrightarrow \hat{u}(\omega, t) &= A e^{i\omega t} + B e^{-i\omega t} \\ \hat{u}(\omega, 0) &= \hat{u}_0 \Longrightarrow A + B = \hat{u}_0 \\ \hat{u}_t(\omega, 0) &= \hat{u}_1 \Longrightarrow i\omega (A e^{i\omega t} - B e^{-i\omega t}) \bigg|_{t=0} = \hat{u}_1 \Longrightarrow A - B = -\frac{i}{\omega} \hat{u}_1 \end{aligned}$$

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$$\int_{-\infty}^{\infty} |u(x,t)|^2 dx = \int_{-\infty}^{\infty} |\hat{u}(\omega,t)|^2 d\omega$$

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$$\begin{split} &\int_{-\infty}^{\infty} |u(x,t)|^2 dx = \int_{-\infty}^{\infty} |\hat{u}(\omega,t)|^2 d\omega \\ &\leq \frac{1}{2} \int_{-\infty}^{\infty} |\hat{u}_0(\omega)|^2 (|e^{i\omega t}|^2 + |e^{i\omega t}|^2) d\omega \end{split}$$

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$$\int_{-\infty}^{\infty} |u(x,t)|^2 dx \leq \int_{-\infty}^{\infty} |u_0|^2 dx \quad \text{(Plancherel's identity)}$$