MA 207 - Differential Equations-II

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Start with

by John von Neumann

"The sciences do not try to explain, they hardly even try to interpret, they mainly make models. By a model is meant a mathematical construct which, with the addition of certain verbal interpretations, describes observed phenomena. The justification of such a mathematical construct is solely and precisely that it is expected to work—that is, correctly to describe phenomena from a reasonably wide area. Furthermore, it must satisfy certain esthetic criteria—that is, in relation to how much it describes, it must be rather simple."



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The ultimate goal of a meteorologist is to set up differential equations of the movements of the air and to obtain, as their integral, the general atmospheric circulation, and as particular integrals the cyclones, anticyclones, tornados, and thunderstorms.

Andria Maurinic - DURTESTATS COM

Lecture 9



Outline of the lecture

- Elliptic PDE: Laplace Equation, Poisson Equation
- Uniqueness of Solution
- Heat Conduction Equation Separation of variables
- Non-homogeneous boundary data
- Non-homogeneous equation

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- Mixed BC : Let $\partial \Omega = \partial \bar{\Omega}_D \cup \partial \bar{\Omega}_N$, $\partial \Omega_D \cap \partial \Omega_N = \phi$.

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$$u = g_1 \text{ on } \partial\Omega_D, \ \frac{\partial u}{\partial \nu} = g_2 \text{ on } \partial\Omega_N.$$



Pierre-Simon, marquis de Laplace (23 March 1749 – 5 March 1827)

"French truely polymath who made important to contributions to engineering, mathematics, statistics, physics, astronomy, and philosophy. Called French Newton His five-volumes on Celestial Mechanics (1799-1825) were classic . Formulated Laplace's equation, and pioneered the Laplace transform The Laplacian differential operator, widely used in mathematics, is also named after him. He restated and developed the nebular hypothesis of the origin of the Solar System and He was one of the first scientists to postulate the existence of black holes and the notion of gravitational collapse.



Baron Simeon Denis Poisson (21 June 1781 – 25 April 1840)

"French mathematician, engineer, and physicist who made many scientific advances Known for his work on definite integrals, electromagnetic theory, and probability. In probability: Poisson Distribution is named after him.

Worked on Celestial Mechanics (extended works of Laplace and Lagrange)
300 papers published



oisson.



Digression-Gauss Divergence Theorem

Theorem

Let Ω be bounded domain whose boundary $\partial \Omega$ is piece-wise smooth and orientable.

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Theorem

Let Ω be bounded domain whose boundary $\partial \Omega$ is piece-wise smooth and orientable. Let \bar{F} be a vector function which is continuously differentiable. Then,

$$\underbrace{\int_{\Omega} \nabla \cdot \vec{F} \, dx}_{\sum_{i=1}^{n} \int_{\Omega} \frac{\partial F_{i}}{\partial x_{i}} \, dx} = \underbrace{\int_{\partial \Omega} \vec{F} \cdot \nu \, ds}_{\sum_{i=1}^{n} \int_{\partial \Omega} F_{i} \nu_{i} \, ds}.$$

$$\int_{\Omega} \Delta u v \, dx = \int_{\partial \Omega} v \frac{\partial u}{\partial \nu} \, ds - \int_{\Omega} \nabla u \cdot \nabla v \, dx$$

$$\int_{\Omega} \Delta u v \ dx = \int_{\partial \Omega} v \frac{\partial u}{\partial \nu} \ ds - \int_{\Omega} \nabla u \cdot \nabla v \ dx$$

Proof:

$$\int_{\Omega} \Delta u v \ dx = \int_{\Omega} \nabla \cdot (\nabla u) v \ dx$$

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Proof:

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has a unique solution.

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Now w = 0 on $\partial \Omega$.

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Exercise: 1. Verify for other BC's.

2. Tut. Sheet 4 - Qns. 6, 7(i), 8

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Claim: $v(x,t) \equiv 0$.



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Integrate from 0 to t to obtain

$$\int_{\Omega} v(x,t)^2 dx + 2 \int_0^t \int_{\Omega} |\nabla v(x,t)|^2 dx = 0$$

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 $\implies u_1 = u_2$; hence the solution is unique.

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Exercise: Tut. Sheet 4- Qns. 7 (ii), (iii), (iv), 10

Exercise: Show that the problem

$$\begin{aligned} u_{tt} - \Delta u &= f \text{ in } \Omega, \ t \in (0, T] \\ BC: \quad & \frac{\partial u}{\partial \nu} = g \text{ on } \partial \Omega, \ t \in [0, T] \\ IC: \quad & u(x, 0) = u_0(x), \ u_t(x, 0) = u_1(x), \ x \in \Omega \end{aligned}$$

has a unique solution.

Exercise: Tut. Sheet 4 - Qn. 7(iv), 9

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- The rod is sufficiently thin so that the heat is distributed equally over the cross-section at time t.

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- The surface of the rod is insulated and hence there is no heat loss through the boundary.

The temperature distribution of the rod is given by the solution of the initial-boundary value problem :

$$u_t = ku_{xx}, \quad 0 < x < I, \quad t > 0$$
 (k > 0 is the thermal conductivity)

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BC: u(0,t) = 0, \ u(I,t) = 0
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u_t = ku_{xx}, \quad 0 < x < I, \quad t > 0 (k > 0 is the thermal conductivity)

BC : u(0,t) = 0, \quad u(I,t) = 0

IC : u(x,0) = \phi(x) \quad 0 \le x \le I.
```

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BOUNDED DOMAIN, HOMOGENEOUS BCs.

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That is, X(x) is the solution of the eigenvalue problem defined by:

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The cases $\lambda = 0$ and $\lambda < 0$ give only trivial solutions to the EVP.

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Hence the eigenvalues are $\lambda = (\frac{n\pi}{I})^2$ and eigenfunctions are

$$\lambda = (\frac{n\pi}{l})^2$$
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$$X_n(x) = B_n \sin(\frac{n\pi x}{I}).$$

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Integrating, we obtain In $T(t) = -k\lambda t + \ln A' \Longrightarrow T(t) = A'e^{-k\lambda t}$. For each λ_n , $T_n(t) = A'_n e^{-k\lambda_n t} = A'_n e^{-k(\frac{n\pi}{l})^2 t}$.

$$T'(t) + \lambda kT(t) = 0$$

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Integrating, we obtain
$$\begin{split} & \text{In } T(t) = -k\lambda t + \text{In } A' \Longrightarrow T(t) = A'e^{-k\lambda t}. \\ & \text{For each } \lambda_n, \ T_n(t) = A'_n e^{-k\lambda_n t} = A'_n e^{-k(\frac{n\pi}{l})^2 t}. \\ & \text{Let } u_n(x,t) = X_n(x)T_n(t). \ \text{Now, } \sum_{n=1}^N u_n(x,t) \text{ satisfies the heat equation.} \end{split}$$

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Let
$$u_n(x,t) = X_n(x)T_n(t)$$
. Now, $\sum_{n=1}^{\infty} u_n(x,t)$ satisfies the heat

equation.

Under suitable assumptions on the initial data (which will be

specified later), $u(x,t) = \sum_{n=1}^{\infty} u_n(x,t)$ also satisfies the heat equation.

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 ASSUMPTION (A1)

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$$u(x,0) = \phi(x) \Longrightarrow \sum_{n=1}^{\infty} A_n \sin(\frac{n\pi x}{I}) = \phi(x)$$
 ASSUMPTION (A2) -

Fourier sine series of $\phi(x)$ converges to $\phi(x)$.

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Hence,
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Hence, $A_n = \frac{2}{l} \int_0^l \phi(x) \sin(\frac{n\pi x}{l}) dx$. The formal solution of the

heat equation with Dirichlet BC & IC is given by

$$u(x,t) = \sum_{n=1}^{\infty} \left(\frac{2}{l} \int_{0}^{l} \phi(x) \sin(\frac{n\pi x}{l}) dx \right) \sin(\frac{n\pi x}{l}) e^{-k(\frac{n\pi}{l})^{2}t}.$$



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Introduce a new function:

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$$U(x,t) = \frac{1}{I} \left[(I-x)g(t) + xh(t) \right].$$

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Let $v(x,t) = u(x,t) - U(x,t)$.

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$$v_t - kv_{xx} = (u_t - ku_{xx}) - (U_t - kU_{xx})$$

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Therefore,

$$v_{t} - kv_{xx} = (u_{t} - ku_{xx}) - (U_{t} - kU_{xx})$$

$$= 0 - \left[\frac{1}{I}(I - x)g'(t) + xh'(t)\right] \qquad (U_{xx} = 0)$$



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$$v(I,t) = u(I,t) - U(I,t) = 0$$

Therefore, v(x, t) is a solution of the non-homogeneous heat equation on [0, l] with homogeneous boundary data,

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$$v(0,t) = u(0,t) - U(0,t)$$

$$= g(t) - g(t) = 0.$$

$$v(I,t) = u(I,t) - U(I,t) = 0$$

Therefore, v(x, t) is a solution of the non-homogeneous heat equation on [0, l] with homogeneous boundary data,

$$v_{t} - kv_{xx} = -\left[\frac{1}{I}(I - x)g'(t) + xh'(t)\right]$$

$$v(x, 0) = \phi(x) - \left[\frac{1}{I}(I - x)g(0) + xh(0)\right]$$

$$v(0, t) = v(I, t) = 0.$$

Hence,
$$v_t - kv_{xx} = -\left|\frac{1}{I}(I-x)g'(t) + xh'(t)\right|$$

$$v(x,0) = u(x,0) - U(x,0)$$

$$= \phi(x) - \left[\frac{1}{l}(l-x)g(0) + xh(0)\right].$$

$$v(0,t) = u(0,t) - U(0,t)$$

$$= g(t) - g(t) = 0.$$

$$v(l,t) = u(l,t) - U(l,t) = 0$$

Therefore, v(x, t) is a solution of the non-homogeneous heat equation on [0, I] with homogeneous boundary data,

$$v_{t} - kv_{xx} = -\left[\frac{1}{I}(I - x)g'(t) + xh'(t)\right]$$

$$v(x, 0) = \phi(x) - \left[\frac{1}{I}(I - x)g(0) + xh(0)\right]$$

$$v(0, t) = v(I, t) = 0.$$

$$u(x, t) = 0$$

u(x,t) = v(x,t) + U(x,t)

$$u_t - ku_{xx} = h(x, t), \quad 0 < x < I, \ t > 0$$

$$u_t - ku_{xx} = h(x, t), \quad 0 < x < l, \ t > 0$$

BC: $u(0, t) = 0, \ u(l, t) = 0$

$$u_t - ku_{xx} = h(x, t), \quad 0 < x < I, \ t > 0$$

 $BC: \ u(0, t) = 0, \ u(I, t) = 0$
 $IC: \ u(x, 0) = \phi(x) \quad 0 \le x \le I.$

 $h(x, t) \neq 0$ - (non-homogeneous) Consider the following eigenvalue problem:

$$y'' + \lambda y = 0, \quad 0 < x < I,$$

$$u_t - ku_{xx} = h(x, t), \quad 0 < x < l, \ t > 0$$

 $BC : u(0, t) = 0, \ u(l, t) = 0$
 $IC : u(x, 0) = \phi(x) \quad 0 \le x \le l.$

 $h(x, t) \neq 0$ - (non-homogeneous) Consider the following eigenvalue problem:

$$y'' + \lambda y = 0, \ 0 < x < I,$$

 $y(0) = y(I) = 0$

$$u_t - ku_{xx} = h(x, t), \quad 0 < x < I, \ t > 0$$

 $BC: \ u(0, t) = 0, \ u(I, t) = 0$
 $IC: \ u(x, 0) = \phi(x) \quad 0 \le x \le I.$

 $h(x, t) \neq 0$ - (non-homogeneous) Consider the following eigenvalue problem:

$$y'' + \lambda y = 0, \ 0 < x < I,$$

 $y(0) = y(I) = 0$

We know that $\lambda_n = (\frac{n\pi}{l})^2$ are the eigenvalues and

$$y_n = \sqrt{\frac{2}{I}} \sin(\frac{n\pi}{I}x)$$
 are the corresponding eigenfunctions.

$$u_t - ku_{xx} = h(x, t), \quad 0 < x < I, \ t > 0$$

 $BC: \ u(0, t) = 0, \ u(I, t) = 0$
 $IC: \ u(x, 0) = \phi(x) \quad 0 \le x \le I.$

 $h(x, t) \neq 0$ - (non-homogeneous) Consider the following eigenvalue problem:

$$y'' + \lambda y = 0, \ 0 < x < I,$$

 $y(0) = y(I) = 0$

We know that $\lambda_n = (\frac{n\pi}{I})^2$ are the eigenvalues and

$$y_n = \sqrt{\frac{2}{l}} \sin(\frac{n\pi}{l}x)$$
 are the corresponding eigenfunctions.

Expand
$$h(x, t) = \sum_{n=1}^{\infty} \gamma_n(t) y_n(x)$$



$$u_t - ku_{xx} = h(x, t), \quad 0 < x < I, \ t > 0$$

 $BC: \ u(0, t) = 0, \ u(I, t) = 0$
 $IC: \ u(x, 0) = \phi(x) \quad 0 \le x \le I.$

 $h(x, t) \neq 0$ - (non-homogeneous) Consider the following eigenvalue problem:

$$y'' + \lambda y = 0, \ 0 < x < I,$$

 $y(0) = y(I) = 0$

We know that $\lambda_n = (\frac{n\pi}{I})^2$ are the eigenvalues and

$$y_n = \sqrt{\frac{2}{I}} \sin(\frac{n\pi}{I}x)$$
 are the corresponding eigenfunctions.

Expand
$$h(x,t) = \sum_{n=1}^{\infty} \gamma_n(t) y_n(x)$$
 and $u(x,t) = \sum_{n=1}^{\infty} \alpha_n(t) y_n(x)$.

$$u_t - ku_{xx} = h(x, t), \quad 0 < x < I, \ t > 0$$

 $BC: \ u(0, t) = 0, \ u(I, t) = 0$
 $IC: \ u(x, 0) = \phi(x) \quad 0 \le x \le I.$

 $h(x, t) \neq 0$ - (non-homogeneous) Consider the following eigenvalue problem:

$$y'' + \lambda y = 0, \ 0 < x < I,$$

 $y(0) = y(I) = 0$

We know that $\lambda_n = (\frac{n\pi}{I})^2$ are the eigenvalues and

$$y_n = \sqrt{\frac{2}{I}} \sin(\frac{n\pi}{I}x)$$
 are the corresponding eigenfunctions.

Expand
$$h(x,t) = \sum_{n=1}^{\infty} \gamma_n(t) y_n(x)$$
 and $u(x,t) = \sum_{n=1}^{\infty} \alpha_n(t) y_n(x)$.

$$\sum_{n=1}^{\infty} \frac{d}{dt} \alpha_n(t) y_n(x) - k \sum_{n=1}^{\infty} \alpha_n(t) y_n''(x) = \sum_{n=1}^{\infty} \gamma_n(t) y_n(x).$$

$$\sum_{n=1}^{\infty} \frac{d}{dt} \alpha_n(t) y_n(x) - k \sum_{n=1}^{\infty} \alpha_n(t) y_n''(x) = \sum_{n=1}^{\infty} \gamma_n(t) y_n(x).$$

Now
$$y_n''(x) = -\lambda_n y_n(x)$$
.

$$\sum_{n=1}^{\infty} \frac{d}{dt} \alpha_n(t) y_n(x) - k \sum_{n=1}^{\infty} \alpha_n(t) y_n''(x) = \sum_{n=1}^{\infty} \gamma_n(t) y_n(x).$$

Now
$$y_n''(x) = -\lambda_n y_n(x)$$
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On substitution, we obtain

Plug these expressions in the PDE to obtain

$$\sum_{n=1}^{\infty} \frac{d}{dt} \alpha_n(t) y_n(x) - k \sum_{n=1}^{\infty} \alpha_n(t) y_n''(x) = \sum_{n=1}^{\infty} \gamma_n(t) y_n(x).$$

Now $y_n''(x) = -\lambda_n y_n(x)$.

On substitution, we obtain

$$\sum_{n=1}^{\infty} \left(\frac{d}{dt} \alpha_n(t) + k \lambda_n \alpha_n(t) - \gamma_n(t) \right) y_n(t) = 0$$

Plug these expressions in the PDE to obtain

$$\sum_{n=1}^{\infty} \frac{d}{dt} \alpha_n(t) y_n(x) - k \sum_{n=1}^{\infty} \alpha_n(t) y_n''(x) = \sum_{n=1}^{\infty} \gamma_n(t) y_n(x).$$

Now $y_n''(x) = -\lambda_n y_n(x)$.

On substitution, we obtain

$$\sum_{n=1}^{\infty} \left(\frac{d}{dt} \alpha_n(t) + k \lambda_n \alpha_n(t) - \gamma_n(t) \right) y_n(t) = 0$$

Taking innerproduct with $y_m(t)$ and using orthonormality of eigen functions, we obtain

Plug these expressions in the PDE to obtain

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Now $y_n''(x) = -\lambda_n y_n(x)$.

On substitution, we obtain

$$\sum_{n=1}^{\infty} \left(\frac{d}{dt} \alpha_n(t) + k \lambda_n \alpha_n(t) - \gamma_n(t) \right) y_n(t) = 0$$

Taking innerproduct with $y_m(t)$ and using orthonormality of eigen functions, we obtain

$$\frac{d}{dt}\alpha_n(t) + k\lambda_n\alpha_n(t) = \gamma_n(t), \quad n = 1, 2, 3, \cdots,$$

Plug these expressions in the PDE to obtain

$$\sum_{n=1}^{\infty} \frac{d}{dt} \alpha_n(t) y_n(x) - k \sum_{n=1}^{\infty} \alpha_n(t) y_n''(x) = \sum_{n=1}^{\infty} \gamma_n(t) y_n(x).$$

Now $y_n''(x) = -\lambda_n y_n(x)$.

On substitution, we obtain

$$\sum_{n=1}^{\infty} \left(\frac{d}{dt} \alpha_n(t) + k \lambda_n \alpha_n(t) - \gamma_n(t) \right) y_n(t) = 0$$

Taking innerproduct with $y_m(t)$ and using orthonormality of eigen functions, we obtain

$$\frac{d}{dt}\alpha_n(t) + k\lambda_n\alpha_n(t) = \gamma_n(t), \quad n = 1, 2, 3, \cdots,$$

$$\implies \alpha_n(t) = \alpha_n(0)e^{-\lambda_n t} + \int_0^t e^{-\lambda_n(t-s)}\gamma_n(s) ds$$

Plug these expressions in the PDE to obtain

$$\sum_{n=1}^{\infty} \frac{d}{dt} \alpha_n(t) y_n(x) - k \sum_{n=1}^{\infty} \alpha_n(t) y_n''(x) = \sum_{n=1}^{\infty} \gamma_n(t) y_n(x).$$

Now $y_n''(x) = -\lambda_n y_n(x)$.

On substitution, we obtain

$$\sum_{n=1}^{\infty} \left(\frac{d}{dt} \alpha_n(t) + k \lambda_n \alpha_n(t) - \gamma_n(t) \right) y_n(t) = 0$$

Taking innerproduct with $y_m(t)$ and using orthonormality of eigen functions, we obtain

$$\frac{d}{dt}\alpha_n(t) + k\lambda_n\alpha_n(t) = \gamma_n(t), \quad n = 1, 2, 3, \cdots,$$

$$\implies \alpha_n(t) = \alpha_n(0)e^{-\lambda_n t} + \int_0^t e^{-\lambda_n(t-s)}\gamma_n(s) ds$$

$$u(x,0) = \phi(x) \Longrightarrow \phi(x) = \sum_{n=1}^{\infty} \alpha_n(0) y_n(x)$$

where
$$\alpha_n(0) = \int_0^1 \phi(x) y_n(x) dx$$
.

$$u(x,0) = \phi(x) \Longrightarrow \phi(x) = \sum_{n=1}^{\infty} \alpha_n(0) y_n(x)$$

where
$$\alpha_n(0) = \int_0^1 \phi(x) y_n(x) dx$$
.

Hence
$$u(x,t) = \sum_{n=1}^{\infty} \alpha_n(t) y_n(x)$$
 with

$$u(x,0) = \phi(x) \Longrightarrow \phi(x) = \sum_{n=1}^{\infty} \alpha_n(0) y_n(x)$$

where
$$\alpha_n(0) = \int_0^I \phi(x) y_n(x) dx$$
.

Hence
$$u(x,t) = \sum_{n=1}^{\infty} \alpha_n(t) y_n(x)$$
 with

$$\implies \alpha_n(t) = \alpha_n(0)e^{-\lambda_n t} + \int_0^t e^{-\lambda_n(t-s)}\gamma_n(s) ds.$$
 (CHECK!)

Exercise: Tut. Sheet 5 - Qns. 7, 8, 9 (i)