$$I = qA\left(\frac{D_{p}p_{n0}}{L_{p}} + \frac{D_{n}n_{p0}}{L_{n}}\right)\left(e^{V_{a}/V_{T}} - 1\right)$$

$$\begin{split} I &= qA \left( \frac{D_p p_{n0}}{L_p} + \frac{D_n n_{p0}}{L_n} \right) \left( e^{V_a/V_T} - 1 \right) \\ &= qA \left( \frac{D_p n_i^2}{N_d \sqrt{D_p \tau_p}} + \frac{D_n n_i^2}{N_a \sqrt{D_n \tau_n}} \right) \left( e^{V_a/V_T} - 1 \right) \end{split}$$

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$$\begin{split} I &= qA \left( \frac{D_\rho p_{n0}}{L_\rho} + \frac{D_n n_{\rho0}}{L_n} \right) \left( e^{V_a/V_T} - 1 \right) \\ &= qA \left( \frac{D_\rho n_i^2}{N_d \sqrt{D_\rho \tau_\rho}} + \frac{D_n n_i^2}{N_a \sqrt{D_n \tau_n}} \right) \left( e^{V_a/V_T} - 1 \right) \\ &= qA \left( \frac{n_i^2}{N_d} \sqrt{\frac{D_\rho}{\tau_\rho}} + \frac{n_i^2}{N_a} \sqrt{\frac{D_n}{\tau_n}} \right) \left( e^{V_a/V_T} - 1 \right). \end{split}$$

Different materials (T = 300 K):

$$\begin{split} I &= qA \left( \frac{D_p p_{n0}}{L_p} + \frac{D_n n_{p0}}{L_n} \right) \left( e^{V_a/V_T} - 1 \right) \\ &= qA \left( \frac{D_p n_i^2}{N_d \sqrt{D_p \tau_p}} + \frac{D_n n_i^2}{N_a \sqrt{D_n \tau_n}} \right) \left( e^{V_a/V_T} - 1 \right) \\ &= qA \left( \frac{n_i^2}{N_d} \sqrt{\frac{D_p}{\tau_p}} + \frac{n_i^2}{N_a} \sqrt{\frac{D_n}{\tau_n}} \right) \left( e^{V_a/V_T} - 1 \right). \end{split}$$

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Semiconductor	$N_c$ (cm <sup>-3</sup> )	$N_{\nu}$ (cm <sup>-3</sup> )	$E_g$ (eV)	$n_i$ (cm <sup>-3</sup> )
Ge	$1.04\times10^{19}$	$6.0  imes 10^{18}$	0.664	$2.33\times10^{13}$
Si	$2.8  imes 10^{19}$	$1.04  imes 10^{19}$	1.12	$1.02\times10^{10}$
GaAs	$4.7  imes 10^{17}$	$7.0  imes 10^{18}$	1.424	$2.1  imes 10^6$

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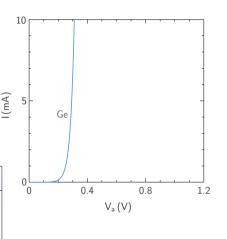
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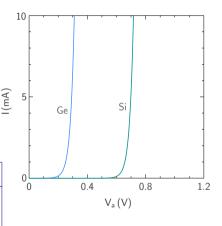
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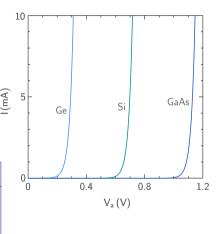
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 $I = I_s \left( e^{V_a/V_T} - 1 \right)$ , where

– 
$$1)$$
, where

The temperature dependence of  $I_s$  comes

mainly from  $n_i(T)$ .

$$p^2 \sqrt{D}$$
  $p^2$ 

 $I_s = qA\left(rac{n_i^2}{N_d}\sqrt{rac{D_p}{ au_p}} + rac{n_i^2}{N_a}\sqrt{rac{D_n}{ au_n}}
ight).$ 

$$\sqrt{D_{-}}$$
  $n_{-}^2$ 























 $I = I_s (e^{V_a/V_T} - 1)$ , where

$$\frac{1}{n^2}$$

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 $n_i(T) = \sqrt{N_c(T)N_v(T)} \exp\left(-\frac{E_g(T)}{2kT}\right),$ 

where  $N_c \propto T^{3/2}$ ,  $N_V \propto T^{3/2}$ .

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$$I_{s} = qA \left( \frac{n_{i}^{2}}{N_{d}} \sqrt{\frac{D_{p}}{\tau_{p}}} + \frac{n_{i}^{2}}{N_{a}} \sqrt{\frac{D_{n}}{\tau_{n}}} \right).$$



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 $n_i(T) = \sqrt{N_c(T)N_v(T)} \exp\left(-\frac{E_g(T)}{2kT}\right),$ 

with  $E_g(0) = 1.17 \text{ eV}$ ,  $\alpha = 4.73 \times 10^{-4}$ ,  $\beta = 636$ .

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Temperature dependence
$$I = I_s (e^{V_a/V_T} - 1), \text{ where}$$

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 $I_s = qA\left(\frac{n_i^2}{N_d}\sqrt{\frac{D_p}{\tau_D}} + \frac{n_i^2}{N_a}\sqrt{\frac{D_n}{\tau_D}}\right).$ 1.08 150 200 250 300 350 400 T(K)

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 $I = I_s (e^{V_a/V_T} - 1)$ , where

$$E_{\mathcal{B}}(T)$$
 for silicon is given by  $E_{\mathcal{B}}(T) = E_{\mathcal{B}}(0) - \frac{\Delta T}{T+\beta}$  (eV), with  $E_{\mathcal{B}}(0) = 1.17$  eV,  $\alpha = 4.73 \times 10^{-4}$ ,  $\beta = 636$ .

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$$\frac{n_i^2}{N_i} \sqrt{\frac{D_n}{D_n}}$$

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300 350 400 T(K)

$$s(T) = \sqrt{N_c(T)N_V(T)} \epsilon$$

$$n_i(T) = \sqrt{N_c(T)N_v(T)} \exp\left(-\frac{E_g(T)}{2kT}\right),$$
  
where  $N_c \propto T^{3/2}$ ,  $N_v \propto T^{3/2}$ .

$$E_g(T)$$
 for silicon is given by  $E_g(T) = E_g(0) - \frac{\alpha T^2}{T + \beta}$  (eV),

As 
$$T$$
 increases,  $E_g/2kT$  decreases, and  $n_i$  increases substantially because of the exponential function  $\rightarrow I$  increases.

1.16

the exponential function  $\rightarrow I$  increases.

Temperature dependence  $10^{14}$ 1.16  $I = I_s (e^{V_a/V_T} - 1)$ , where  $10^{10}$  $n_i (cm^{-3})$  $I_s = qA\left(\frac{n_i^2}{N_A}\sqrt{\frac{D_p}{T_p}} + \frac{n_i^2}{N_a}\sqrt{\frac{D_n}{T_p}}\right).$  $10^{6}$ The temperature dependence of  $I_s$  comes 1.08 mainly from  $n_i(T)$ . 250 300 350 400 200 250 300 150 200 150 350 400

T(K)

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The temperature dependence of  $I_s$  comes mainly from  $n_i(T)$ .

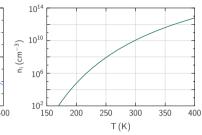
$$n_i(T) = \sqrt{N_c(T)N_v(T)} \exp\left(-\frac{E_g(T)}{2kT}\right),$$

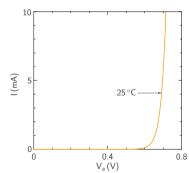
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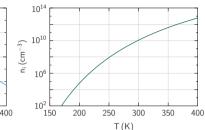
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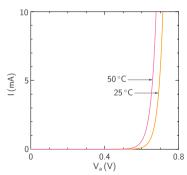
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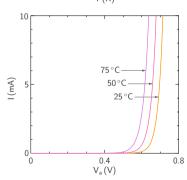
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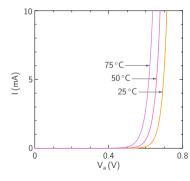
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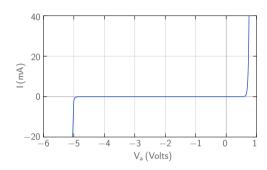
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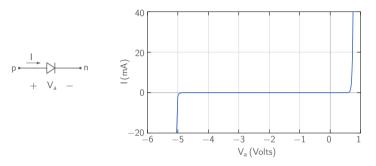
As T increases,  $E_g/2kT$  decreases, and  $n_i$  increases substantially because of the exponential function  $\to I$  increases.

For silicon, the  $\emph{I-V}$  curve shifts by about  $-2\,\mathrm{mV/^\circ C}$  as the temperature is increased.

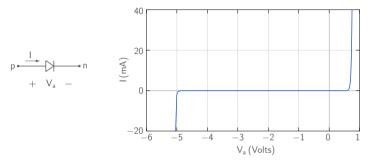




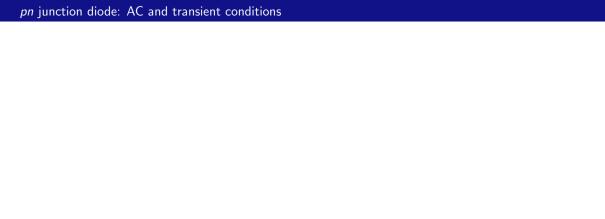




\* A real diode cannot withstand indefinitely large reverse voltages and "breaks down" at some point as  $V_R$  is increased.



- \* A real diode cannot withstand indefinitely large reverse voltages and "breaks down" at some point as  $V_B$  is increased.
- \* Reverse breakdown can be due to
  - impact ionisation (avalanche breakdown)
  - tunneling (Zener breakdown)

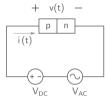


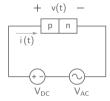
#### pn junction diode: AC and transient conditions

\* We have looked at the DC behaviour of a pn junction diode so far. We now want to consider  $V_a$  (the applied voltage) varying with time.

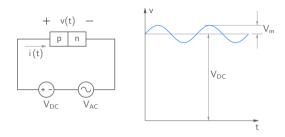
#### pn junction diode: AC and transient conditions

- \* We have looked at the DC behaviour of a pn junction diode so far. We now want to consider  $V_a$  (the applied voltage) varying with time.
- \* Two situations are of interest:
  - \* Small-signal behaviour (AC): With  $V_a(t) = V_{DC} + V_m \sin \omega t$ , how does the current vary with time when  $V_m$  is "small?"
  - \* Large-signal behaviour: The variation in the applied voltage is not small. In particular, we are interested in the turn-off and turn-on transients.

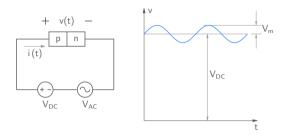




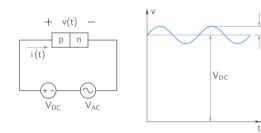
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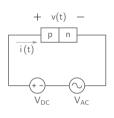


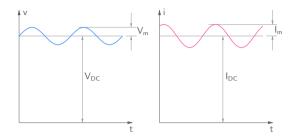
- \* Let  $v(t) = V_{DC} + V_m \sin \omega t$ .
- \* If  $V_m$  is "small," the current is also sinusoidal, i.e.,  $i(t) = I_{DC} + I_m \sin(\omega t + \phi)$ .



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 $I_{DC}$ 





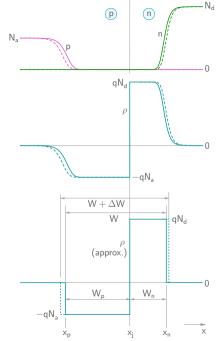
- \* Let  $v(t) = V_{DC} + V_m \sin \omega t$ .
- \* If  $V_m$  is "small," the current is also sinusoidal, i.e.,  $i(t) = I_{DC} + I_m \sin(\omega t + \phi)$ .
- \* In small-signal analysis, we are interested in the relationship between the sinusoidal parts of the current and voltage, in particular, the ratio of the current and voltage phasors,  $I_m \angle \phi / V_m \angle 0$ .

\* A pn junction diode conducts negligibly small current with a DC reverse bias.

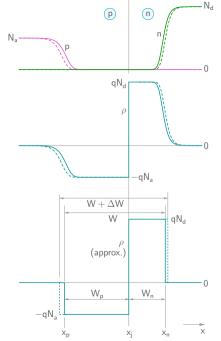
- \* A pn junction diode conducts negligibly small current with a DC reverse bias.
- \* With a time-varying applied reverse bias, it can conduct an appreciable current.

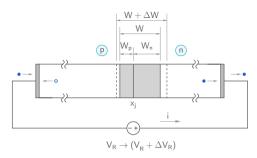
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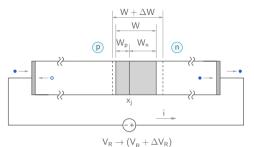
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- \* This change is made possible by removal of majority carriers.

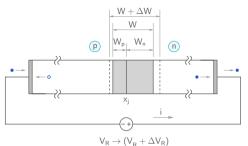






\* Movement of majority carriers is relatively fast, and the time scale involved is  $\sim \tau = \frac{\epsilon_s}{q\mu_n n}$  for electrons.

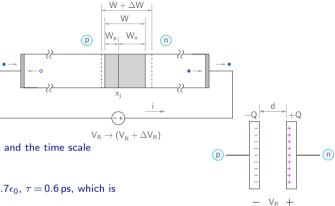
For  $n=10^{16}\,\mathrm{cm^{-3}}$ ,  $\mu_n=1000\,\mathrm{cm^2/V}$ -s,  $\epsilon_s=11.7\epsilon_0$ ,  $\tau=0.6\,\mathrm{ps}$ , which is negligibly small for all practical purposes.



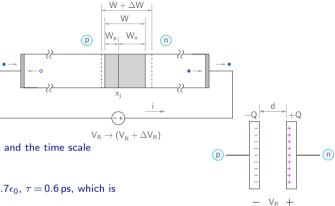
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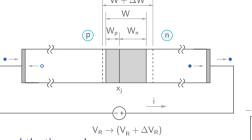
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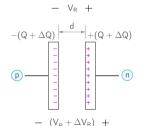
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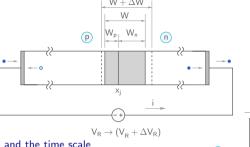
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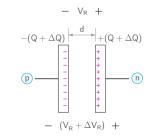


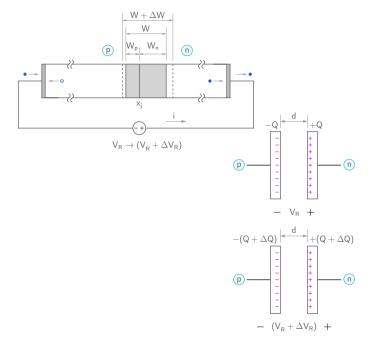
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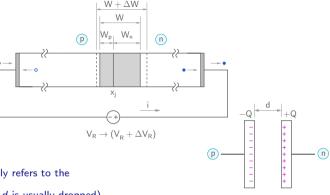
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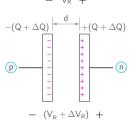
Note: For simplicity, we have not shown  $V_{\rm bi}$  in the figure; the drop across the junction is actually  $V_{\rm bi}+V_{\rm R}$ , as seen before.



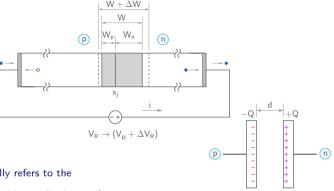




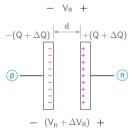
\* In semiconductor devices, "capacitance" generally refers to the differential capacitance  $C_d = \frac{dQ}{dV}$  (the subscript d is usually dropped).



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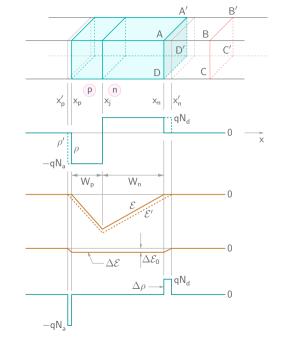


- \* In semiconductor devices, "capacitance" generally refers to the differential capacitance  $C_d=\frac{dQ}{dV}$  (the subscript d is usually dropped).
- \* For a reverse-biased pn junction,  $C = \frac{\Delta Q}{\Delta V_P}$ .

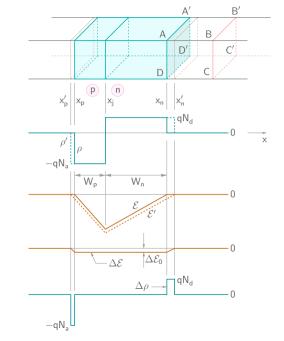


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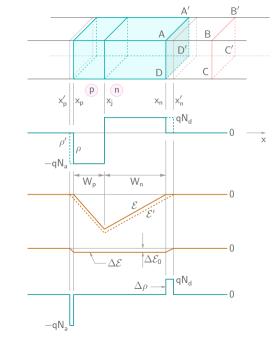
$$V_{\rm bi} + V_R = -\int_{x_p}^{x_n} \mathcal{E}(x) dx$$



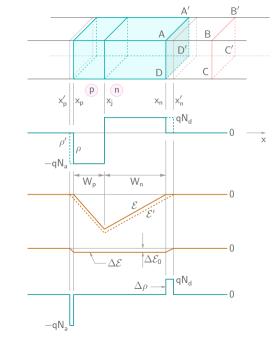
$$\begin{aligned} V_{\mathsf{b}\mathsf{i}} + V_R &= -\int_{\mathsf{x}_p}^{\mathsf{x}_n} \mathcal{E}(\mathsf{x}) d\mathsf{x} \\ V_{\mathsf{b}\mathsf{i}} + V_R + \Delta V_R &= -\int_{\mathsf{x}_p'}^{\mathsf{x}_n'} \mathcal{E}'(\mathsf{x}) d\mathsf{x} \end{aligned}$$



$$\begin{aligned} V_{\text{bi}} + V_R &= -\int_{x_p}^{x_n} \mathcal{E}(x) dx \\ V_{\text{bi}} + V_R + \Delta V_R &= -\int_{x_p'}^{x_n'} \mathcal{E}'(x) dx \\ &\to \Delta V_R = -\int_{x_-'}^{x_n'} (\mathcal{E}'(x) - \mathcal{E}(x)) dx = -\int_{x_-'}^{x_n'} \Delta \mathcal{E}(x) dx \end{aligned}$$



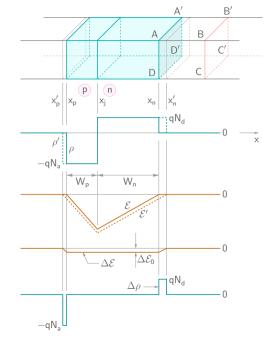
$$\begin{split} V_{\text{bi}} + V_R &= -\int_{x_p}^{x_n} \mathcal{E}(x) dx \\ V_{\text{bi}} + V_R + \Delta V_R &= -\int_{x_p'}^{x_n'} \mathcal{E}'(x) dx \\ &\to \Delta V_R = -\int_{x_p'}^{x_n'} (\mathcal{E}'(x) - \mathcal{E}(x)) dx = -\int_{x_p'}^{x_n'} \Delta \mathcal{E}(x) dx \\ &= \Delta \mathcal{E}_0 W \text{ as } \Delta V_R \to 0 \text{ V}. \end{split}$$



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 $\Delta Q$ , the total charge in the Gaussian box between AA'D'D and BB'C'C, is given by

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 $=\Delta \mathcal{E}_0 W$  as  $\Delta V_R \rightarrow 0 V$ .

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$$V_{bi} + V_R + \Delta V_R = -\int_{x_p'}^{x_n'} \mathcal{E}'(x) dx$$

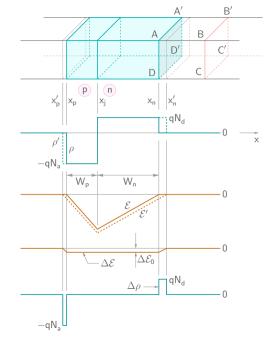
$$\to \Delta V_R = -\int_{x_p'}^{x_n'} (\mathcal{E}'(x) - \mathcal{E}(x)) dx = -\int_{x_p'}^{x_n'} \Delta \mathcal{E}(x) dx$$

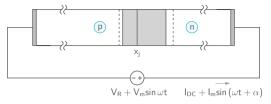
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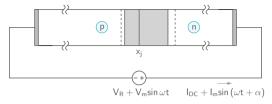
$$\rightarrow C_J = \left. \frac{\Delta Q}{\Delta V_R} \right|_{\Delta V_R \rightarrow 0} = \frac{A \epsilon_s}{W}.$$

 $C_J$  is called the "junction capacitance" or "depletion layer capacitance."





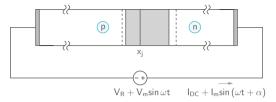
For an abrupt, uniformly doped silicon pn junction, with  $N_a=10^{17}\,\mathrm{cm}^{-3}$  and  $N_d=2\times10^{16}\,\mathrm{cm}^{-3}$ , and area  $=0.01\,\mathrm{cm}^2$ , calculate the capacitance (i.e., the differential capacitance) for an applied reverse bias of  $V_R=2\,\mathrm{V}$  (  $T=300\,\mathrm{K}$ ).



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Solution: The built-in voltage is

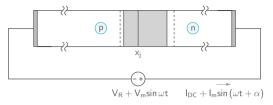
$$V_{\rm bi} = rac{kT}{q} \, \log \left( rac{N_a N_d}{n_i^2} 
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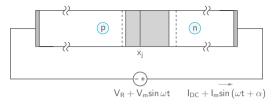


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Capacitance 
$$C_J = \frac{A \, \epsilon_s}{W} \ = \frac{0.01 \, \mathrm{cm}^2 \times 11.7 \times 8.85 \times 10^{-14} \, \mathrm{F/cm}}{0.464 \times 10^{-4} \, \mathrm{cm}}$$



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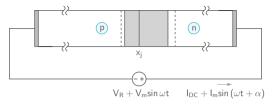
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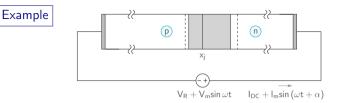
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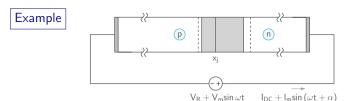
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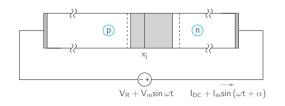
$$= 0.223 \, \mathrm{nF}.$$





$$C_J = rac{A\epsilon_s}{W} = A\epsilon_s \sqrt{rac{qN_a}{2\epsilon_s(V_{
m bi}-V_a)}}.$$

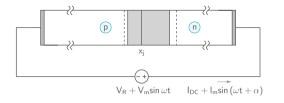




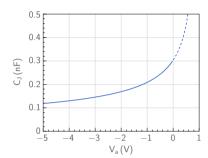
$$C_J = rac{A\epsilon_s}{W} = A\epsilon_s \sqrt{rac{qN_a}{2\epsilon_s(V_{\mathrm{bi}} - V_a)}}.$$

$$\frac{1}{C_I^2} = \frac{1}{(A\epsilon_s)^2} \frac{2\epsilon_s(V_{\text{bi}} - V_a)}{qN_a} = \frac{2}{qN_a\epsilon_s A^2} (V_{\text{bi}} - V_a).$$

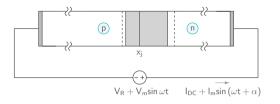




$$\begin{aligned} C_J &= \frac{A\epsilon_s}{W} = A\epsilon_s \sqrt{\frac{qN_a}{2\epsilon_s(V_{bi} - V_a)}}.\\ \frac{1}{C_J^2} &= \frac{1}{(A\epsilon_s)^2} \frac{2\epsilon_s(V_{bi} - V_a)}{qN_a} = \frac{2}{qN_a\epsilon_s A^2} (V_{bi} - V_a). \end{aligned}$$

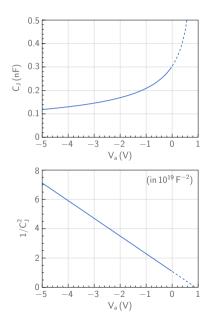




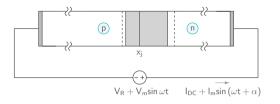


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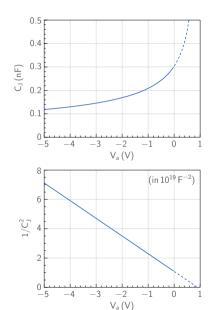


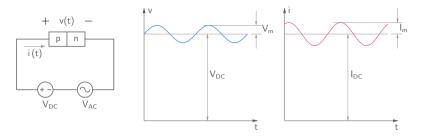
Solution: The junction capacitance is given by

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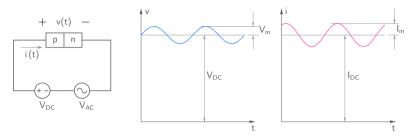
$$\frac{1}{C_I^2} = \frac{1}{(A\epsilon_s)^2} \frac{2\epsilon_s (V_{\rm bi} - V_a)}{qN_a} = \frac{2}{qN_a\epsilon_s A^2} (V_{\rm bi} - V_a).$$

 $ightarrow 1/C_J^2$  versus  $V_a$ : Slope gives  $N_a$ ; x-intercept gives  $V_{
m bi}$ .



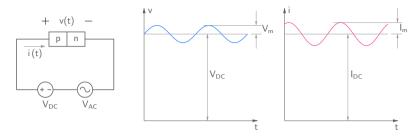


 $\mbox{Small signal} \rightarrow \mbox{With a sinusoidal input, the output (voltage or current) should} \\ \mbox{also be sinusoidal, i.e., it should not be distorted.}$ 



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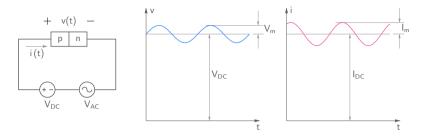
For a reverse-biased 
$$pn$$
 junction,  $C_J = \frac{dQ}{dV_a} = \frac{A\epsilon_s}{W(V_a)} = \frac{K}{\sqrt{V_{bi} - V_a}}$ , with  $K = A\epsilon_s \sqrt{\frac{qN_aN_d}{2\epsilon_s(N_a + N_d)}}$ .



Small signal  $\to$  With a sinusoidal input, the output (voltage or current) should also be sinusoidal, i.e., it should not be distorted.

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With 
$$V_a(t) = -(V_R + V_m \sin \omega t)$$
,  $i(t) = \frac{dQ}{dt} = \frac{dQ}{dV_a} \frac{dV_a}{dt} = C_J(V_a) \times (-\omega V_m \cos \omega t)$ .

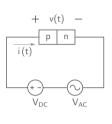


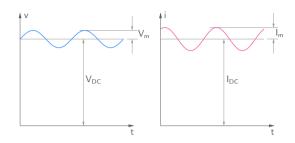
Small signal  $\rightarrow$  With a sinusoidal input, the output (voltage or current) should also be sinusoidal, i.e., it should not be distorted.

For a reverse-biased 
$$pn$$
 junction,  $C_J = \frac{dQ}{dV_a} = \frac{A\epsilon_s}{W(V_a)} = \frac{K}{\sqrt{V_{bi} - V_a}}$ , with  $K = A\epsilon_s \sqrt{\frac{qN_aN_d}{2\epsilon_s(N_a + N_d)}}$ .

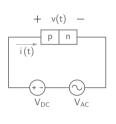
With 
$$V_a(t) = -(V_R + V_m \sin \omega t)$$
,  $i(t) = \frac{dQ}{dt} = \frac{dQ}{dV_a} \frac{dV_a}{dt} = C_J(V_a) \times (-\omega V_m \cos \omega t)$ .

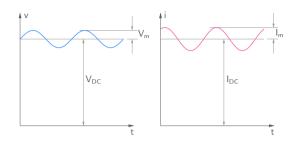
 $\rightarrow i(t)$  is sinusoidal if  $C_J$  can be treated as a constant.





$$egin{aligned} & v_{a}(t) = -(V_R + V_m \sin \omega t) 
ightarrow -(V_R + V_m) < v_a < -(V_R - V_m). \ & C_J^{\min} = rac{K}{\sqrt{V_{\mathrm{bi}} + V_R + V_m}}. \end{aligned}$$



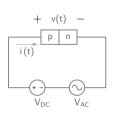


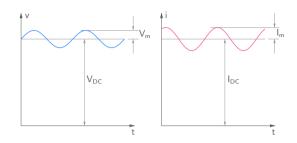
$$v_a(t) = -(V_R + V_m \sin \omega t) \rightarrow -(V_R + V_m) < v_a < -(V_R - V_m).$$

$$C_J^{\min} = rac{\mathcal{K}}{\sqrt{V_{\mathrm{bi}} + V_R + V_m}}, \quad C_J^{\max} = rac{\mathcal{K}}{\sqrt{V_{\mathrm{bi}} + V_R - V_m}}.$$

Consider one of these two extreme values.

$$C_J^{\text{max}} = \frac{\mathcal{K}}{\sqrt{V_{\text{bi}} + V_R - V_m}} = \frac{\mathcal{K}}{\sqrt{V_{\text{bi}} + V_R}} \times \frac{1}{\sqrt{1 - \frac{V_m}{V_{\text{bi}} + V_R}}} \approx \frac{\mathcal{K}}{\sqrt{V_{\text{bi}} + V_R}} \left(1 + \frac{1}{2} \frac{V_m}{V_{\text{bi}} + V_R}\right).$$





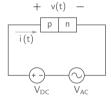
$$v_a(t) = -(V_R + V_m \sin \omega t) \to -(V_R + V_m) < v_a < -(V_R - V_m).$$

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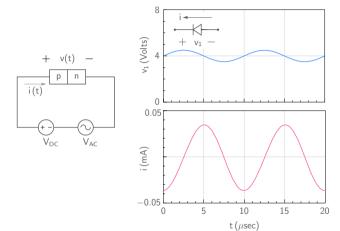
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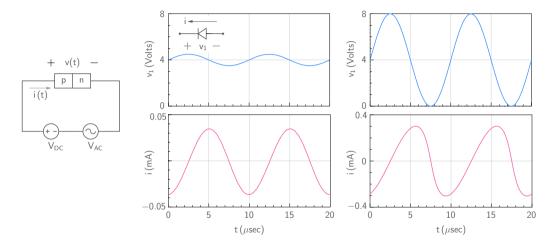
If  $\frac{V_m}{2(V_{i,j} + V_D)} \ll 1$ ,  $C_J$  can be treated as a constant.



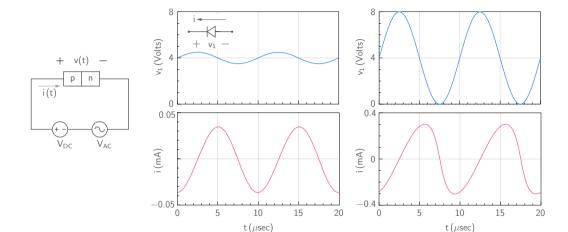
\* Small-signal condition:  $\frac{V_m}{2\left(V_{\mathrm{bi}}+V_R\right)}\ll 1.$ 



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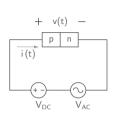


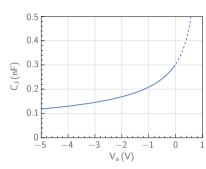
\* Small-signal condition:  $\frac{V_m}{2\left(V_{\mathrm{bi}}+V_R\right)}\ll 1.$ 

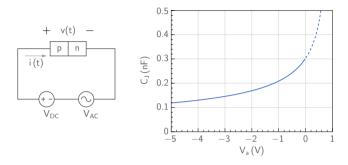


\* Small-signal condition: 
$$\frac{V_m}{2\left(V_{\rm bi}+V_R\right)}\ll 1.$$

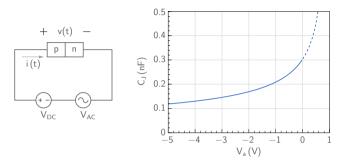
\* If the small-signal condition is not satisfied, i(t) shows distortion.



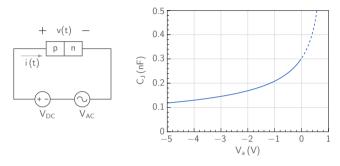




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- \* In these devices, the doping density profiles are designed so as to get a large capacitance change for a small change in reverse bias.