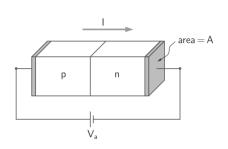
#### SEMICONDUCTOR DEVICES

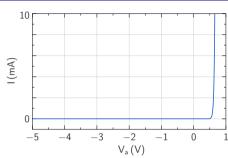
*p-n* Junctions: Part 2

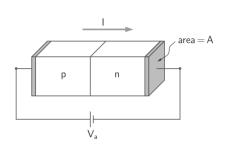


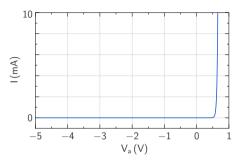
M.B.Patil
mbpatil@ee.iitb.ac.in
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Department of Electrical Engineering Indian Institute of Technology Bombay

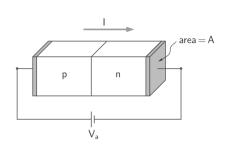


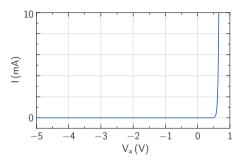




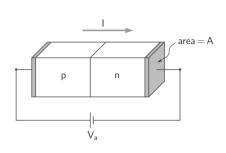


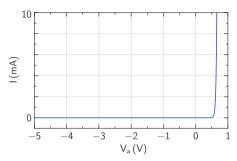
\* With  $V_a \approx 0.6\,\mathrm{V}$  a substantial current flows. When  $V_a$  is increased further, the current increases rapidly.



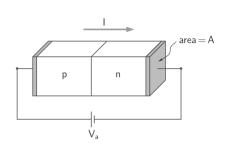


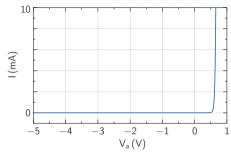
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- \* When a reverse bias (i.e.,  $V_a < 0$ ,V) is applied, the diode blocks conduction, i.e., the current is negligibly small.



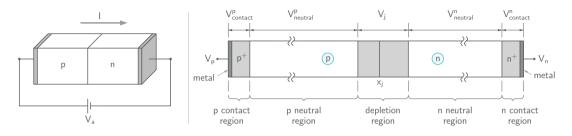


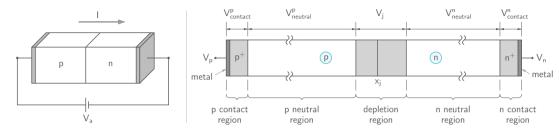
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  - We want to understand this "rectifying" behaviour.





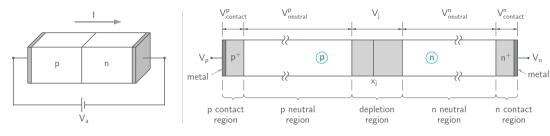
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- \* When a reverse bias (i.e.,  $V_a < 0$ ,V) is applied, the diode blocks conduction, i.e., the current is negligibly small.
  - We want to understand this "rectifying" behaviour.
- \* As we increase the forward bias, the current increases rapidly, and at some point, the device will get damaged because of overheating. For silicon diodes used in low-power applications, the forward voltge must be restricted to about 0.8 V.
  - (Note: Although we will show an applied forward/reverse bias with a battery, in practice, a battery is generally not connected directly across a diode.)



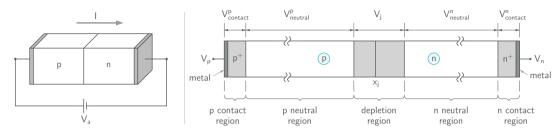


Consider a pn junction in equilibrium ( $V_a = 0 \text{ V}$ ).

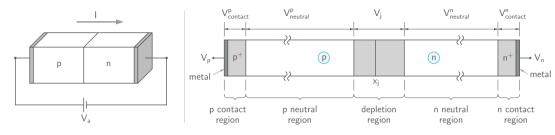
\*  $V_{\text{contact}}^{p}$  and  $V_{\text{contact}}^{n}$  are the voltage drops across the contact regions (which generally include heavily doped regions).



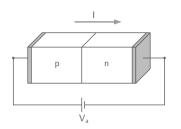
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- \*  $V_{\text{contact}}^p$  and  $V_{\text{contact}}^n$  remain constant irrespective of the applied voltage (to be discussed later).

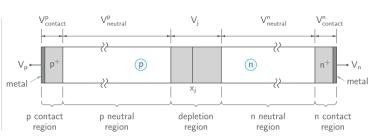


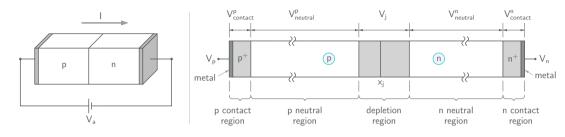
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- \*  $V_{\text{neutral}}^{p}$  and  $V_{\text{neutral}}^{n}$  are the voltage drops across the neutral p and n regions. In equilibrium, they are both zero.



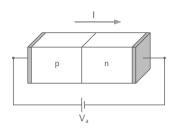
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- \*  $V_{\text{neutral}}^p$  and  $V_{\text{neutral}}^n$  are the voltage drops across the neutral p and n regions. In equilibrium, they are both zero.
- \* Even with current flow,  $V_{\text{neutral}}^{p}$  and  $V_{\text{neutral}}^{n}$  remain negligibly small since a very small electric field is sufficient to create the required  $J_{p}^{\text{drift}} = qp\mu_{p}\mathcal{E}$  or  $J_{n}^{\text{drift}} = qn\mu_{n}\mathcal{E}$  (note that p and n in these equations represent the *majority* carrier densities).

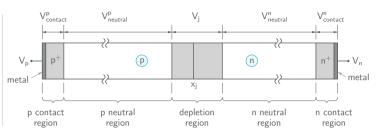






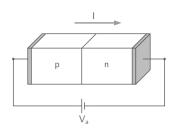
\*  $V_j$  is the voltage across the junction and is equal to  $V_{\rm bi}$  in equilibrium. We have shown, using Poisson's equation, that  $V_j \propto W^2$  where W is the depletion width.

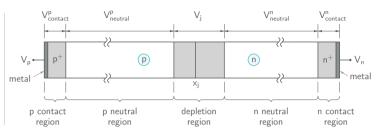




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- \* In equilibrium,  $V_p = V_n$ , and we get
  - (1):  $V_{
    m contact}^p V_{
    m bi} + V_{
    m contact}^n =$  0, taking voltage drop as positive.

(We assume that the signs of  $V_{\rm contact}^p$  and  $V_{\rm contact}^n$  are taken into account.)



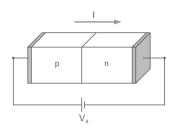


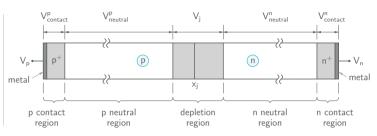
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(We assume that the signs of  $V_{\text{contact}}^p$  and  $V_{\text{contact}}^n$  are taken into account.)

\* When a bias is applied, we have

(2): 
$$V_{\text{contact}}^p - V_i + V_{\text{contact}}^n = V_p - V_n = V_a$$
.

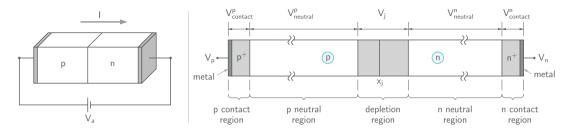




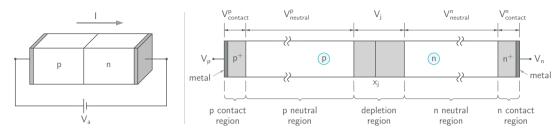
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(We assume that the signs of  $V_{\text{contact}}^p$  and  $V_{\text{contact}}^n$  are taken into account.)

- \* When a bias is applied, we have
  - (2):  $V_{\text{contact}}^p V_j + V_{\text{contact}}^n = V_p V_n = V_a$ .
- \* (1)-(2) gives  $-V_{bi} + V_j = -V_a$ , i.e.,  $V_j = V_{bi} V_a$

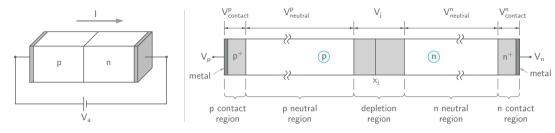


For an abrupt silicon pn junction, the built-in voltage is  $V_{\rm bi}=0.85\,\rm V$ . Let  $W_0$  and  $W_1$  denote the depletion widths for  $V_a=0\,\rm V$  and  $V_a=0.6\,\rm V$ , respectively. What is  $W_1/W_0$ ?



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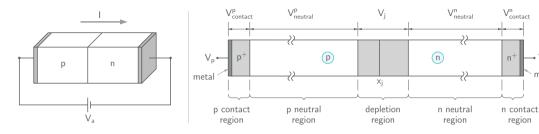
Solution:  $V_j \propto W^2 \rightarrow V_j = kW^2$ .



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Solution: 
$$V_j \propto W^2 \rightarrow V_j = kW^2$$
.

$$V_{\text{bi}} = kW_0^2 \quad \text{(for } V_a = 0 \text{ V)}$$



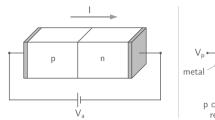
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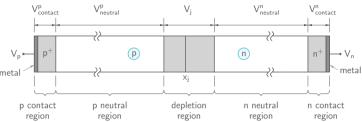
Solution: 
$$V_j \propto W^2 \rightarrow V_j = kW^2$$
.

$$V_{\text{bi}} = kW_0^2 \quad \text{(for } V_a = 0 \text{ V)}$$

$$V_{\rm bi} - 0.6 \, {\rm V} = kW_1^2$$
 (for  $V_a = 0.6 \, {\rm V}$ )

metal





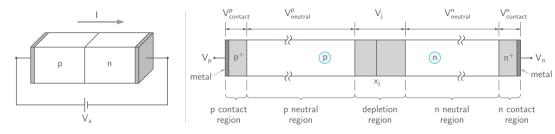
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Solution: 
$$V_j \propto W^2 \rightarrow V_j = kW^2$$
.

$$V_{\rm bi} = kW_0^2$$
 (for  $V_a = 0 \,\mathrm{V}$ )

$$V_{\rm bi} - 0.6 \, {\rm V} = kW_1^2 \quad \text{(for } V_a = 0.6 \, {\rm V}\text{)}$$

$$\rightarrow \frac{0.85-0.6}{0.85} = \left(\frac{W_1}{W_0}\right)^2 \rightarrow \frac{W_1}{W_0} = 0.54.$$



For an abrupt silicon pn junction, the built-in voltage is  $V_{\rm bi}=0.85\,{\rm V}$ . Let  $W_0$  and  $W_1$  denote the depletion widths for  $V_a=0\,{\rm V}$  and  $V_a=0.6\,{\rm V}$ , respectively. What is  $W_1/W_0$ ?

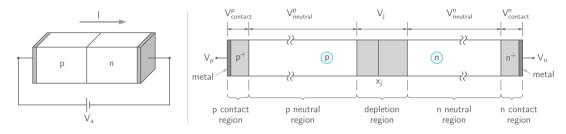
Solution: 
$$V_j \propto W^2 \rightarrow V_j = kW^2$$
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$$V_{\rm bi} = kW_0^2$$
 (for  $V_a = 0 \,\mathrm{V}$ )

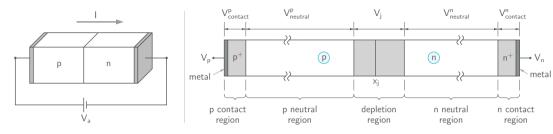
$$V_{\rm bi} - 0.6 \, {\rm V} = kW_1^2 \quad \text{(for } V_a = 0.6 \, {\rm V}\text{)}$$

$$\rightarrow \frac{0.85-0.6}{0.85} = \left(\frac{W_1}{W_0}\right)^2 \rightarrow \frac{W_1}{W_0} = 0.54.$$

Application of a forward bias of 0.6V causes the depletion regoin to shrink by a factor 0.54.

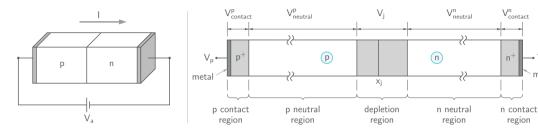


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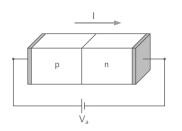


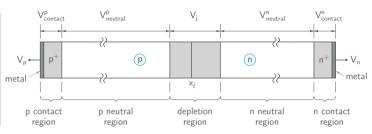
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Solution: 
$$V_j \propto W^2 \rightarrow V_j = kW^2$$
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$$V_{\text{bi}} = kW_0^2$$
 (for  $V_a = 0 \text{ V}$ )

metal



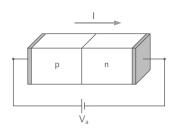


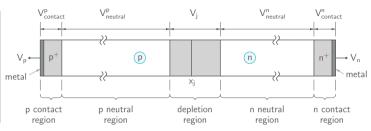
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Solution: 
$$V_j \propto W^2 \rightarrow V_j = kW^2$$
.

$$V_{\text{bi}} = kW_0^2 \quad \text{(for } V_a = 0 \text{ V)}$$

$$V_{\text{bi}} - (-2) V = kW_1^2$$
 (for  $V_a = -2 V$ )





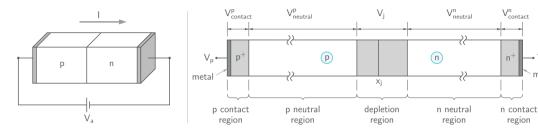
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Solution: 
$$V_j \propto W^2 \rightarrow V_j = kW^2$$
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$$V_{\text{bi}} = kW_0^2$$
 (for  $V_a = 0 \text{ V}$ )

$$V_{\text{bi}} - (-2) V = kW_1^2$$
 (for  $V_a = -2 V$ )

$$\rightarrow \frac{0.85 + 2}{0.85} = \left(\frac{W_1}{W_0}\right)^2 \rightarrow \frac{W_1}{W_0} = 1.83.$$



For an abrupt silicon pn junction, the built-in voltage is  $V_{\rm bi}=0.85\,{\rm V}$ . Let  $W_0$  and  $W_1$  denote the depletion widths for  $V_a=0\,{\rm V}$  and  $V_a=-2\,{\rm V}$  (i.e., a reverse bias  $V_R$  of  $2\,{\rm V}$ ), respectively. What is  $W_1/W_0$ ?

Solution: 
$$V_j \propto W^2 \rightarrow V_j = kW^2$$
.

$$V_{bi} = kW_0^2$$
 (for  $V_a = 0 \text{ V}$ )  
 $V_{bi} - (-2) \text{ V} = kW_1^2$  (for  $V_a = -2 \text{ V}$ )

$$\rightarrow \frac{0.85+2}{0.85} = \left(\frac{W_1}{W_0}\right)^2 \rightarrow \frac{W_1}{W_0} = 1.83.$$

Application of a reverse bias of 2V causes the depletion regoin to expand by a factor 1.83.

Forward bias Ec  $qV_{bi}$  $\Delta_{\mathsf{p}} \frac{1}{\mathsf{E}_{\mathsf{v}}} = \mathsf{E}_{\mathsf{v}}$ P qN<sub>d</sub> p n (n)  $-qN_a$  $V_{-\mathcal{E}_{\mathsf{m}}}$ 

ψ

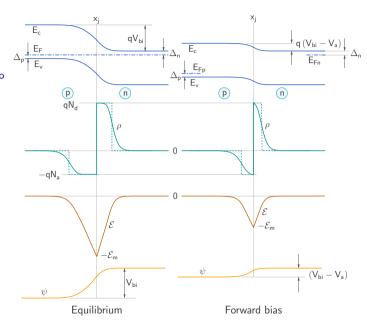
Forward bias

 $V_{bi}$ 

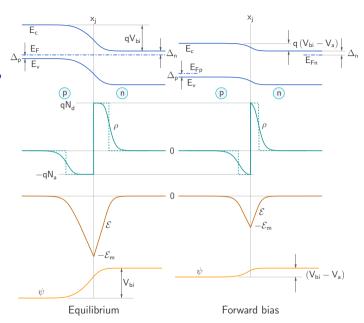
Equilibrium

 $(V_{bi} - V_a)$ 

\* The electrostatic conditions (viz.,  $\rho(x)$ ,  $\mathcal{E}(x)$ ,  $\psi(x)$ ) under forward bias are similar to the equilibrium situation except for a reduced junction voltage  $(V_{\rm bi}-V_a)$ .



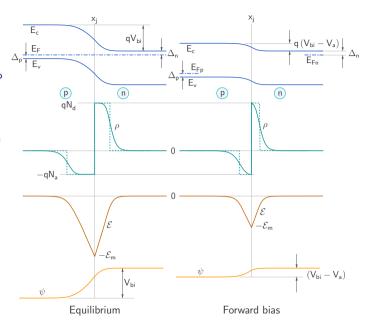
- \* The electrostatic conditions (viz.,  $\rho(x)$ ,  $\mathcal{E}(x)$ ,  $\psi(x)$ ) under forward bias are similar to the equilibrium situation except for a reduced junction voltage  $(V_{\rm bi}-V_a)$ .
- \* The depletion width, the maximum electric field, and the junction voltage drop decrease with forward bias.



- \* The electrostatic conditions (viz.,  $\rho(x)$ ,  $\mathcal{E}(x)$ ,  $\psi(x)$ ) under forward bias are similar to the equilibrium situation except for a reduced junction voltage ( $V_{\rm bi}-V_a$ ).
- \* The depletion width, the maximum electric field, and the junction voltage drop decrease with forward bias.
- \* Solving Poisson's equation using the depletion approximation, we get

$$W = \sqrt{rac{2\epsilon}{q}\left(rac{N_a + N_d}{N_a N_d}
ight)(V_{
m bi} - V_a)},$$

$$W_n = \frac{N_a}{N_a + N_d} W, \ W_p = \frac{N_d}{N_a + N_d} W.$$



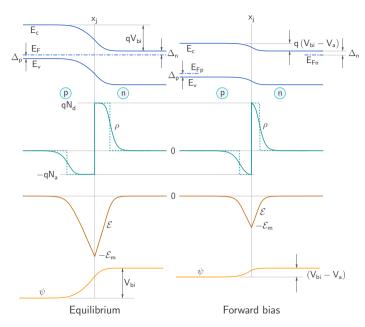
Forward bias Ec  $qV_{bi}$  $\Delta_{\mathsf{p}} \frac{1}{\mathsf{E}_{\mathsf{v}}} = \mathsf{E}_{\mathsf{v}}$ P qN<sub>d</sub> p n (n)  $-qN_a$  $V_{-\mathcal{E}_{\mathsf{m}}}$ ψ  $(V_{bi} - V_a)$ 

 $V_{bi}$ 

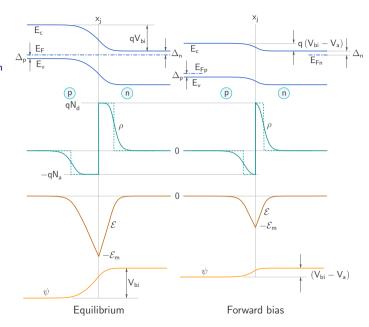
Forward bias

Equilibrium

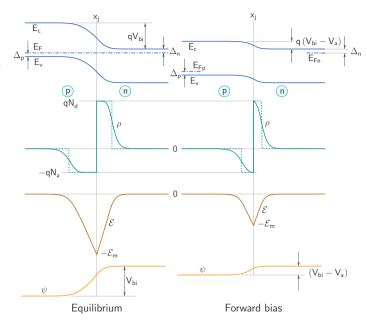
 Although the equilibrium condition is disturbed with an applied bias, the situation sufficiently far from the depletion region is hardly different.



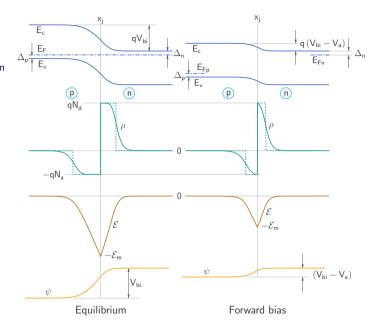
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- \* We can extend the Fermi level concept to describe carrier concentrations sufficiently far from the depletion region.



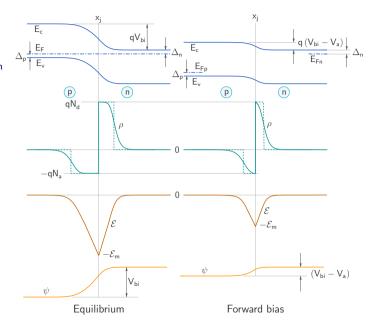
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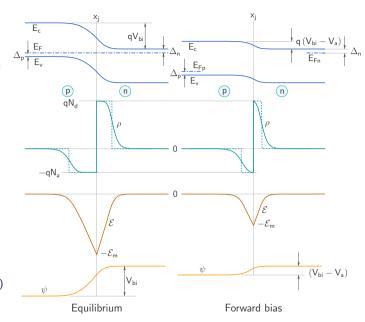
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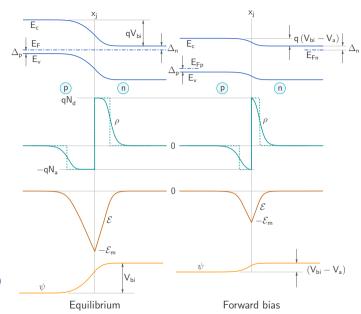
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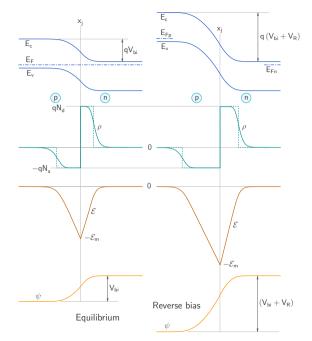


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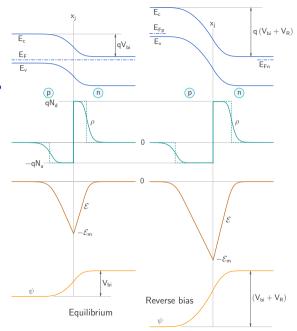


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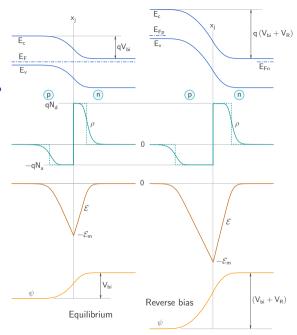




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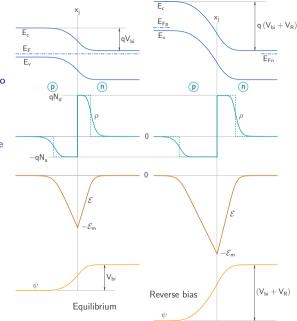
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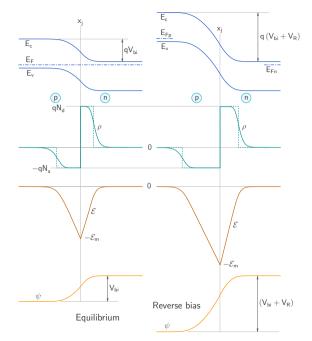


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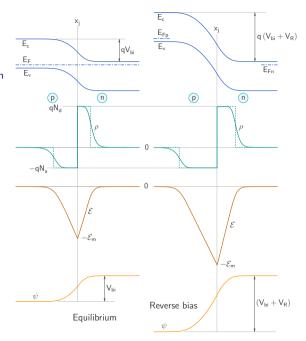
$$W = \sqrt{rac{2\epsilon}{q} \left(rac{ extsf{N}_{ extsf{a}} + extsf{N}_{ extsf{d}}}{ extsf{N}_{ extsf{a}} extsf{N}_{ extsf{d}}}
ight) (V_{ extsf{bi}} + V_{ extsf{R}})},$$

$$W_n = \frac{N_a}{N_a + N_d} \ W, \ \ W_\rho = \frac{N_d}{N_a + N_d} \ W.$$

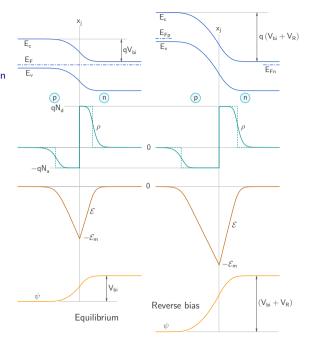




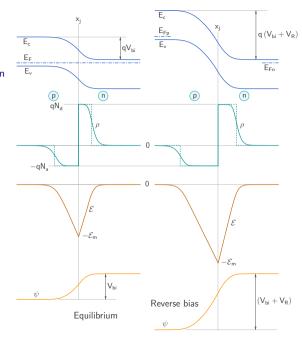
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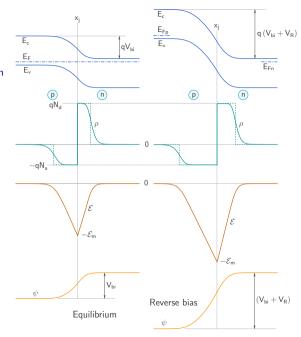
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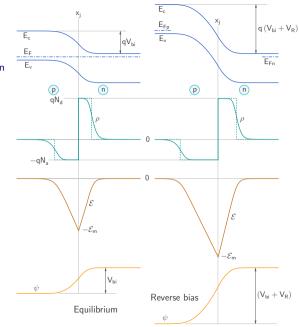
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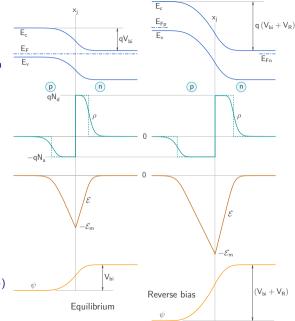
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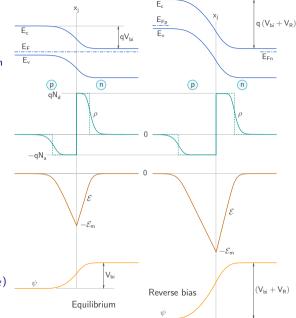


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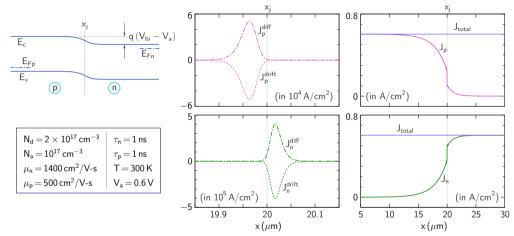
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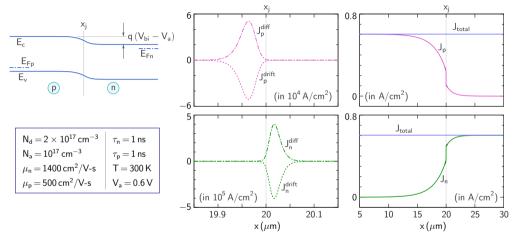
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### Current densities in forward bias



Near the junction,

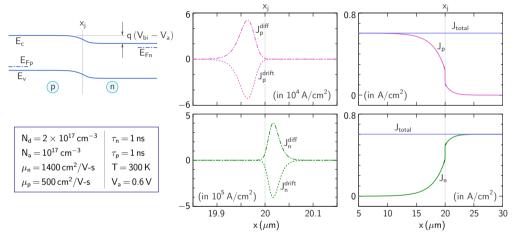
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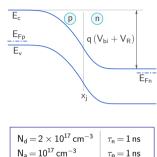
\* Although the equilibrium condition is disturbed, we still have  $J_p^{\text{diff}} \approx -J_p^{\text{drift}}$ , and  $J_n^{\text{diff}} \approx -J_n^{\text{drift}}$ .

#### Current densities in forward bias



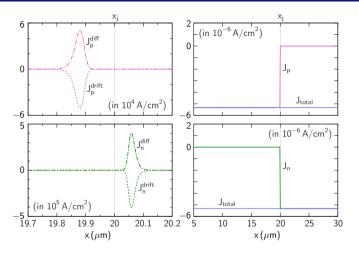
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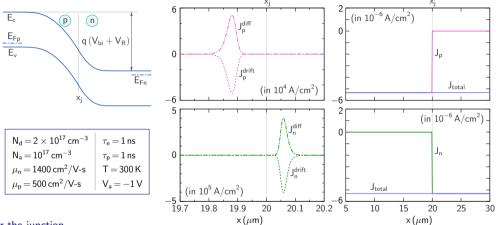
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- \* The net current densities  $J_n$  and  $J_p$  are much smaller than the drift and diffusion components.



$N_d = 2 \times 10^{17}  \text{cm}^{-3}$	$ au_{n} = 1ns$
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$\mu_{\mathrm{n}}$ $=$ 1400 cm $^2/\mathrm{V}$ -s	T = 300  K
$\mu_{ m p}$ $=$ $500{ m cm}^2/{ m V}$ -s	$V_a = -1 V$

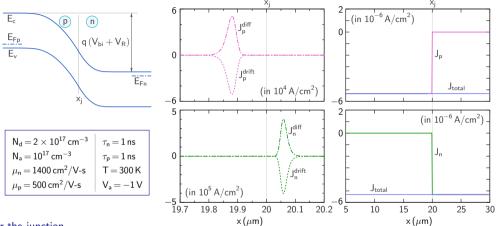
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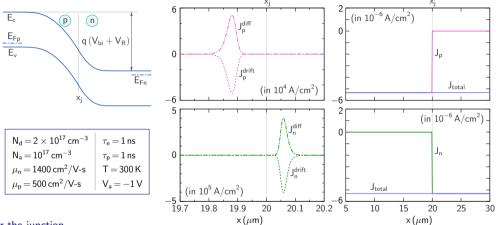
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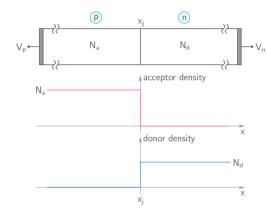


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- \* Note that  $J_{\text{total}}$  in reverse bias is negligibly small compared to the forward bias case (0.7 A/cm<sup>2</sup> for  $V_a = 0.6 \,\text{V}$ ). For all practical purposes, we can say that the current is zero for reverse bias.

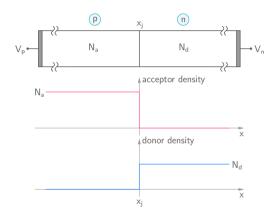
M. B. Patil, IIT Bombay

Definitions:



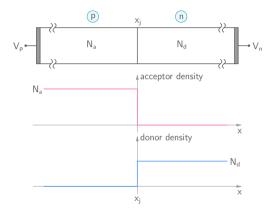
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 $p_{p0}$ : equilibrium hole density in the neutral p-region



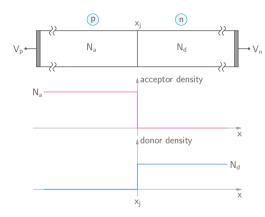
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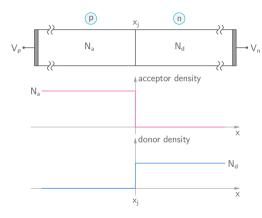
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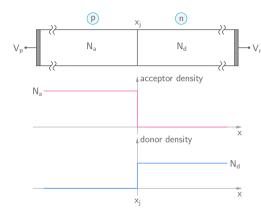
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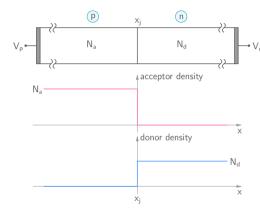
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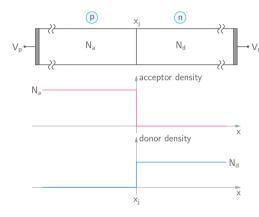
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Example:  $N_a = 5 \times 10^{16} \, \mathrm{cm}^{-3}$ ,  $N_d = 10^{18} \, \mathrm{cm}^{-3}$  ( $T = 300 \, \mathrm{K}$ ).



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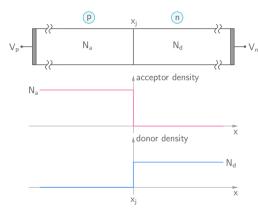
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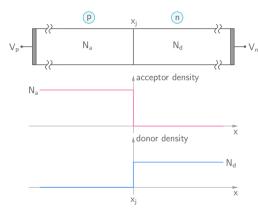
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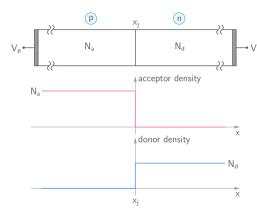
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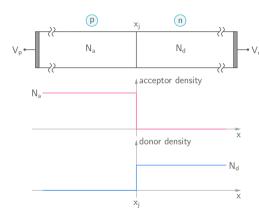
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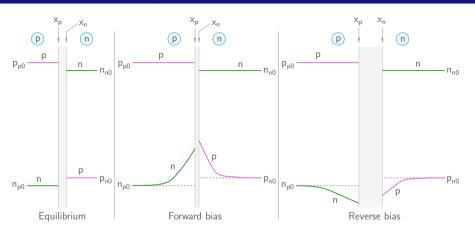
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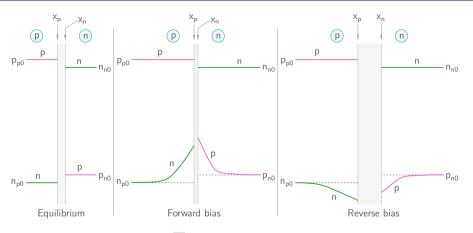
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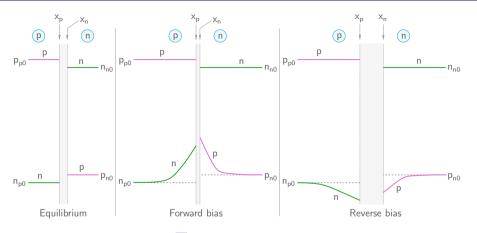
$$p_{n0} = \frac{n_i^2}{n_{n0}} = \frac{(1.5 \times 10^{10})^2}{10^{18}} = 2.25 \times 10^2 \,\mathrm{cm}^{-3}.$$



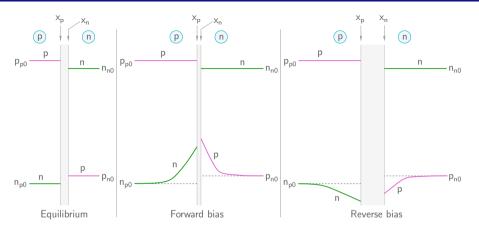




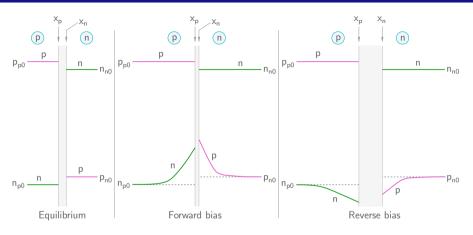
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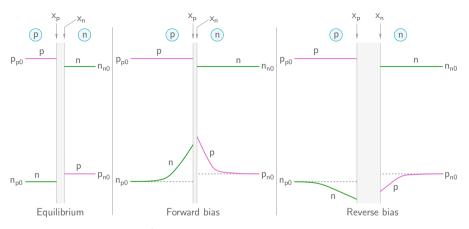
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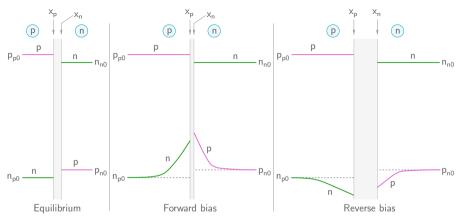
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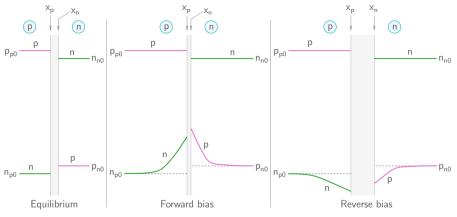
$$J_p^{
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m drift}$$



$$J_p^{
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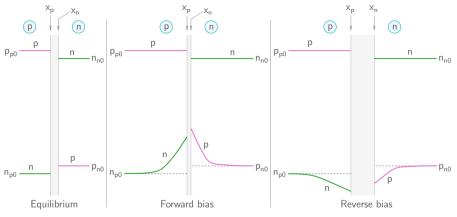


$$\label{eq:Jpdiff} \mathcal{J}_{p}^{\rm diff} \approx -\mathcal{J}_{p}^{\rm drift} \ \, \rightarrow q\,\mu_{p}\,p\,\mathcal{E} = qD_{p}\frac{dp}{dx}, \ \, {\rm i.e.,} \ \, \mathcal{E} = -\frac{d\psi}{dx} = \frac{D_{p}}{\mu_{p}}\frac{1}{p}\frac{dp}{dx}.$$



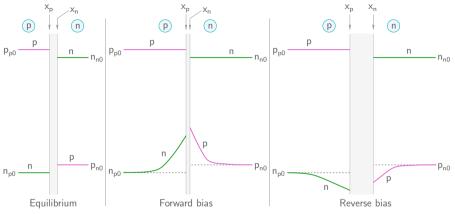
$$J_{\rho}^{\rm diff} \approx -J_{\rho}^{\rm drift} \ \, \rightarrow q \, \mu_{\rho} \, p \, \mathcal{E} = q D_{\rho} \frac{d\rho}{dx}, \; {\rm i.e.,} \ \, \mathcal{E} = -\frac{d\psi}{dx} = \frac{D_{\rho}}{\mu_{\rho}} \frac{1}{\rho} \frac{d\rho}{dx}. \label{eq:definition}$$

$$rac{D}{\mu} = rac{kT}{q} 
ightarrow \int d\psi = -V_T \, \int rac{1}{p} \, dp$$



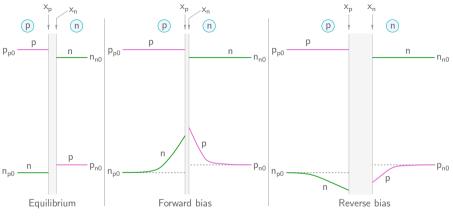
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ho(x_2)}{
ho(x_1)}$$

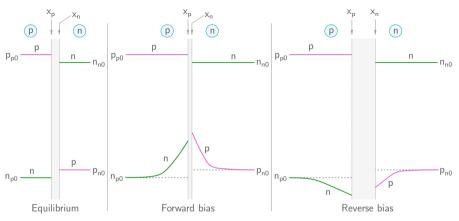


$$J_{p}^{\rm diff} \approx -J_{p}^{\rm drift} \ \, \rightarrow q \, \mu_{p} \, p \, \mathcal{E} = q D_{p} \frac{dp}{dx}, \ \, {\rm i.e.,} \ \, \mathcal{E} = -\frac{d\psi}{dx} = \frac{D_{p}}{\mu_{p}} \frac{1}{p} \frac{dp}{dx}.$$

$$\frac{D}{\mu} = \frac{kT}{q} \rightarrow \int d\psi = -V_T \int \frac{1}{p} dp \rightarrow \psi \Big|_{x_1}^{x_2} = -V_T \log \frac{p(x_2)}{p(x_1)} \rightarrow \frac{p(x_n)}{p(x_p)} = \exp \left(\frac{\psi(x_p) - \psi(x_n)}{V_T}\right).$$

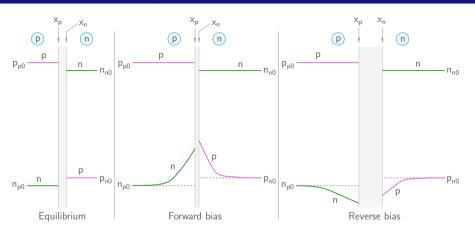


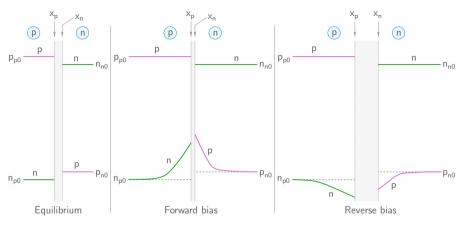
$$J_p^{ ext{diff}} pprox - J_p^{ ext{drift}} 
ightarrow rac{p(x_n)}{p(x_p)} = \expigg(rac{\psi(x_p) - \psi(x_n)}{V_T}igg).$$



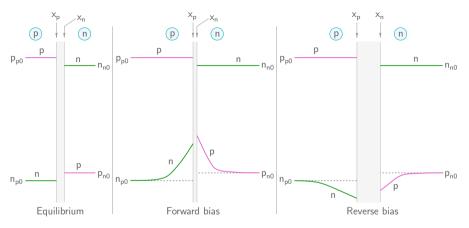
$$J_p^{\text{diff}} \approx -J_p^{\text{drift}} \to \frac{p(x_n)}{p(x_p)} = \exp\bigg(\frac{\psi(x_p) - \psi(x_n)}{V_T}\bigg).$$

$$J_n^{\mathrm{diff}} pprox - J_n^{\mathrm{drift}} 
ightarrow rac{n(x_n)}{n(x_p)} = \exp\left(rac{\psi(x_n) - \psi(x_p)}{V_T}
ight).$$

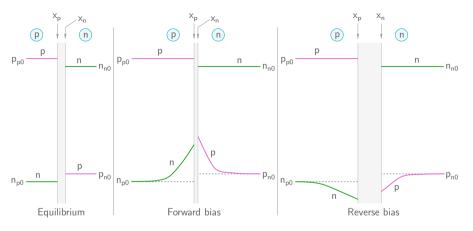




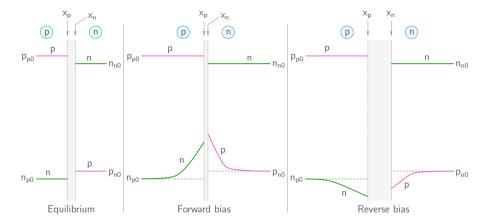
\* When a bias is applied, the minority carrier concentrations in the neutral regions can change substantially.

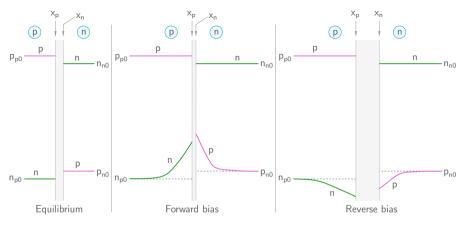


- \* When a bias is applied, the minority carrier concentrations in the neutral regions can change substantially.
- \* There is a corresponding change in the majority carrier concentrations as well, and it serves to keep these regions charge-neutral.

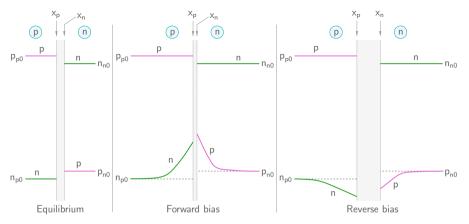


- \* When a bias is applied, the minority carrier concentrations in the neutral regions can change substantially.
- \* There is a corresponding change in the majority carrier concentrations as well, and it serves to keep these regions charge-neutral.
- \* Low-level injection:  $\Delta n \approx \Delta p \ll p_{p0}$  in the neutral p-region  $\to p(x) \approx p_{p0}$  for  $x \le x_p$  $\Delta p \approx \Delta n \ll n_{n0}$  in the neutral n-region  $\to n(x) \approx n_{n0}$  for  $x \ge x_n$



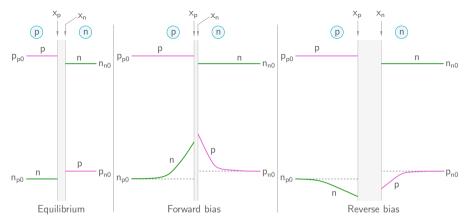


\* 
$$\frac{p(x_n)}{p(x_p)} = \exp\left(\frac{\psi(x_p) - \psi(x_n)}{V_T}\right)$$
,  $\frac{n(x_n)}{n(x_p)} = \exp\left(\frac{\psi(x_n) - \psi(x_p)}{V_T}\right)$ . Also,  $\psi(x_n) - \psi(x_p) = V_j$ .



$$* \frac{p(x_n)}{p(x_p)} = \exp\left(\frac{\psi(x_p) - \psi(x_n)}{V_T}\right), \quad \frac{n(x_n)}{n(x_p)} = \exp\left(\frac{\psi(x_n) - \psi(x_p)}{V_T}\right). \text{ Also, } \psi(x_n) - \psi(x_p) = V_j.$$

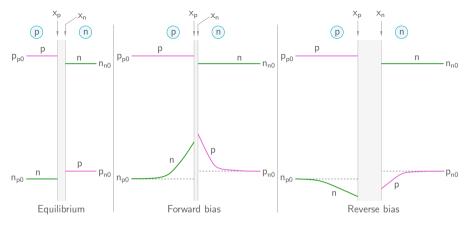
\* Low-level injection:  $p(x) \approx p_{p0}$  for  $x \le x_p$ , and  $n(x) \approx n_{p0}$  for  $x \ge x_p$ .



$$* \frac{p(x_n)}{p(x_p)} = \exp\left(\frac{\psi(x_p) - \psi(x_n)}{V_T}\right), \quad \frac{n(x_n)}{n(x_p)} = \exp\left(\frac{\psi(x_n) - \psi(x_p)}{V_T}\right). \text{ Also, } \psi(x_n) - \psi(x_p) = V_j.$$

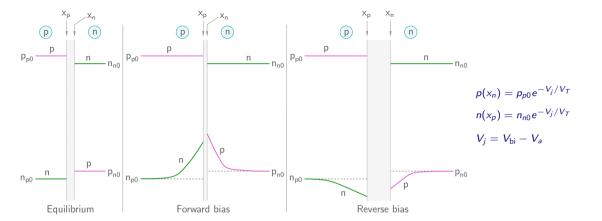
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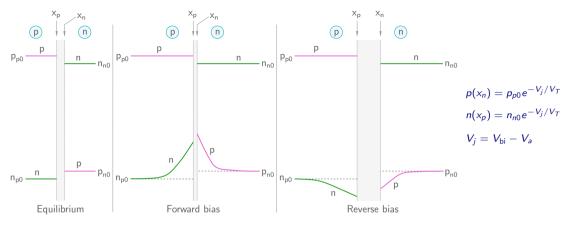
$$ightarrow rac{p(x_n)}{p_{p0}} = e^{-V_j/V_T} 
ightarrow p(x_n) = p_{p0}e^{-V_j/V_T}$$



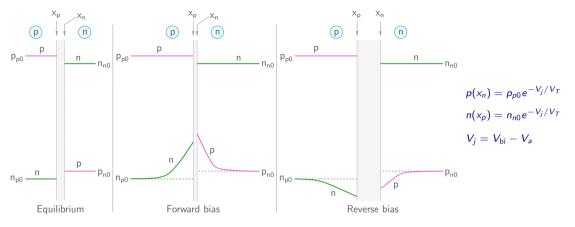
$$* \frac{p(x_n)}{p(x_p)} = \exp\left(\frac{\psi(x_p) - \psi(x_n)}{V_T}\right), \quad \frac{n(x_n)}{n(x_p)} = \exp\left(\frac{\psi(x_n) - \psi(x_p)}{V_T}\right). \text{ Also, } \psi(x_n) - \psi(x_p) = V_j.$$

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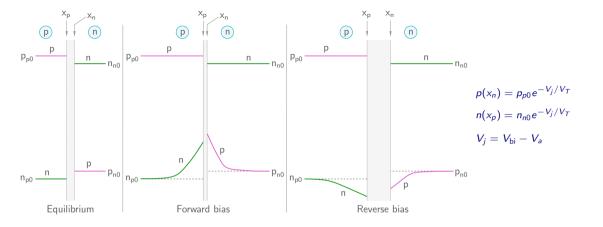




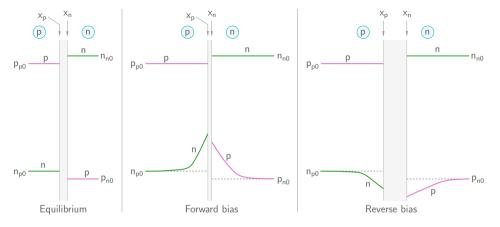
Equilibrium: 
$$p(x_n) = p_{n0} = p_{p0} \exp\left(\frac{-V_{\text{bi}}}{V_T}\right)$$
,  $n(x_p) = n_{p0} = n_{n0} \exp\left(\frac{-V_{\text{bi}}}{V_T}\right)$ .



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With bias:  $p(x_n) = p_{p0} \exp\left(\frac{-V_{\text{bi}} + V_a}{V_T}\right)$ ,  $n(x_p) = n_{n0} \exp\left(\frac{-V_{\text{bi}} + V_a}{V_T}\right)$ .

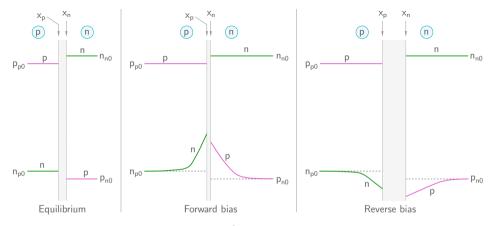


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 $p(x_n) = p_{n0} \exp\left(\frac{V_a}{V_T}\right)$ ,  $n(x_p) = n_{p0} \exp\left(\frac{V_a}{V_T}\right)$ .

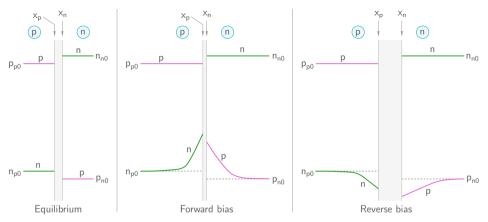


Example: Consider an abrupt, uniformly doped silicon pn junction at  $T=300\,\mathrm{K}$ , with  $N_a=5\times10^{16}\,\mathrm{cm}^{-3}$  and  $N_d=10^{18}\,\mathrm{cm}^{-3}$ . Compute the depletion width and the minority carrier densities at the depletion region edges ( $x_p$  and  $x_n$ ) for an applied bias of  $+0.3\,\mathrm{V}$ ,  $+0.6\,\mathrm{V}$ ,  $-1\,\mathrm{V}$ ,  $-5\,\mathrm{V}$ .

 $(n_i = 1.5 \times 10^{10} \, \text{cm}^{-3} \text{ for silicon at } T = 300 \, \text{K.})$ 

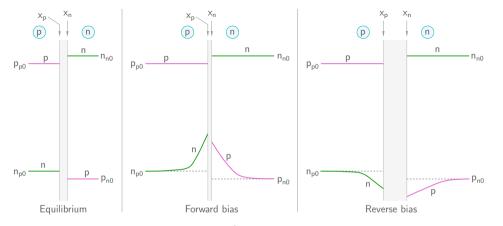


Solution: 
$$p_{p0} \approx N_a = 5 \times 10^{16} \, \mathrm{cm}^{-3} \rightarrow n_{p0} = \frac{n_i^2}{p_{p0}} = \frac{(1.5 \times 10^{10})^2}{5 \times 10^{16}} = 4.5 \times 10^3 \, \mathrm{cm}^{-3}.$$



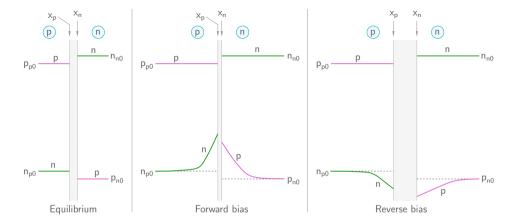
Solution: 
$$p_{\rho 0} \approx N_a = 5 \times 10^{16} \,\mathrm{cm}^{-3} \rightarrow n_{\rho 0} = \frac{n_i^2}{p_{\rho 0}} = \frac{(1.5 \times 10^{10})^2}{5 \times 10^{16}} = 4.5 \times 10^3 \,\mathrm{cm}^{-3}.$$

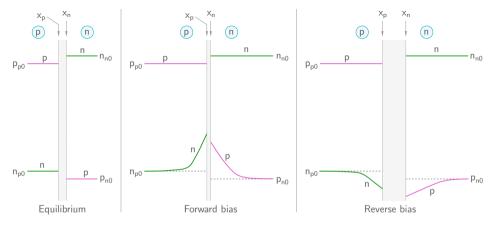
$$n_{n 0} \approx N_a = 1 \times 10^{18} \,\mathrm{cm}^{-3} \rightarrow p_{n 0} = \frac{n_i^2}{n_{n 0}} = \frac{(1.5 \times 10^{10})^2}{1 \times 10^{18}} = 2.25 \times 10^2 \,\mathrm{cm}^{-3}.$$



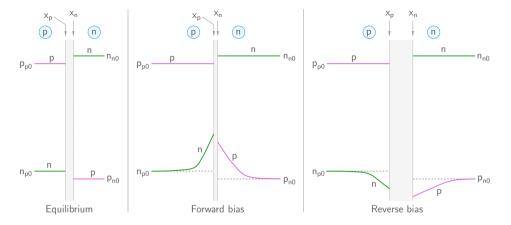
Solution: 
$$p_{p0} \approx N_a = 5 \times 10^{16} \,\mathrm{cm}^{-3} \rightarrow n_{p0} = \frac{n_i^2}{p_{p0}} = \frac{(1.5 \times 10^{10})^2}{5 \times 10^{16}} = 4.5 \times 10^3 \,\mathrm{cm}^{-3}.$$
  $n_{n0} \approx N_a = 1 \times 10^{18} \,\mathrm{cm}^{-3} \rightarrow p_{n0} = \frac{n_i^2}{n_{n0}} = \frac{(1.5 \times 10^{10})^2}{1 \times 10^{18}} = 2.25 \times 10^2 \,\mathrm{cm}^{-3}.$ 

$$V_{\text{bi}} = V_T \log \left( \frac{N_a N_d}{n_c^2} \right) = (0.0259 \,\text{V}) \log \left( \frac{5 \times 10^{16} \times 10^{18}}{(1.5 \times 10^{10})^2} \right) = 0.86 \,\text{V}.$$

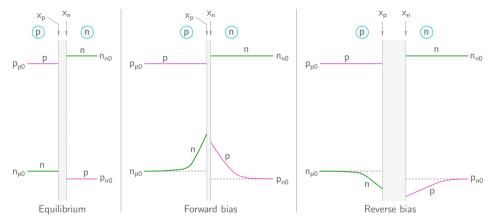




$$V_a = 0.3 \, \mathrm{V}$$
:  $W = \sqrt{rac{2\epsilon}{q} \, rac{N_a + N_d}{N_a N_d} \, (V_{\mathrm{bi}} - V_a)} = 0.12 \, \mathrm{\mu m}$ .



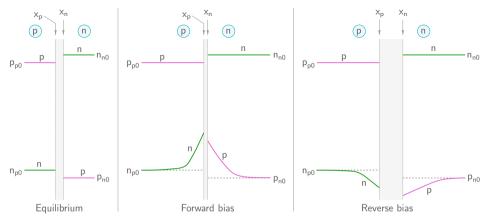
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  $n(x_p) = n_{p0} \, \exp\left(\frac{V_a}{V_T}\right) = 4.5 \times 10^3 \times \exp\left(\frac{0.3}{0.0259}\right) = 4.83 \times 10^8 \,\mathrm{cm}^{-3}$ .



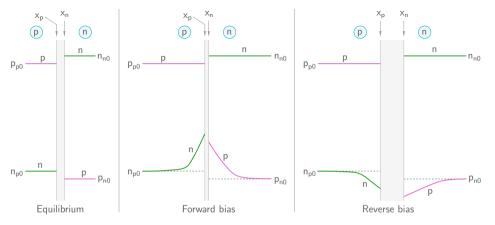
$$V_a = 0.3 \, \text{V}$$
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$$p(x_n) = p_{n0} \exp\left(\frac{V_a}{V_T}\right) = 2.25 \times 10^2 \times \exp\left(\frac{0.3}{0.0259}\right) = 2.41 \times 10^7 \,\mathrm{cm}^{-3}.$$

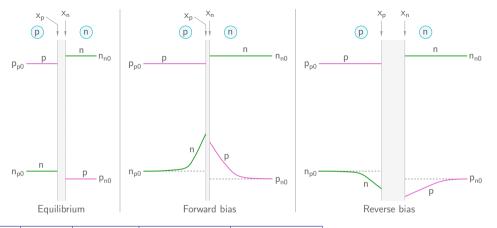


V <sub>a</sub> (V)	<i>W</i> (μm)	$\mathcal{E}_m$ (kV/cm)	$n(x_p)$ (cm <sup>-3</sup> )	$p(x_n)$ (cm <sup>-3</sup> )
0.6	0.08	61.3	$5.18 \times 10^{13}$	$2.59 \times 10^{12}$
0.3	0.12	90.4	$4.83 \times 10^{8}$	$2.41 \times 10^{7}$
0.0	0.15	112.2	$4.50 \times 10^{3}$	$2.25 \times 10^{2}$
-1.0	0.22	165.3	$7.68\times10^{-14}\approx0$	$3.84\times10^{-15}\approx0$
-5.0	0.40	293.6	≈ 0	≈ 0



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 With forward bias, the minority carrier concentrations can increase by several orders of magnitude.



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-5.0	0.40	293.6	≈ 0	$\approx 0$

- With forward bias, the minority carrier concentrations can increase by several orders of magnitude.
- \* With reverse bias, the minority carrier concentrations become very small and can be replaced with zero for all practical purposes.