

MEAM 520

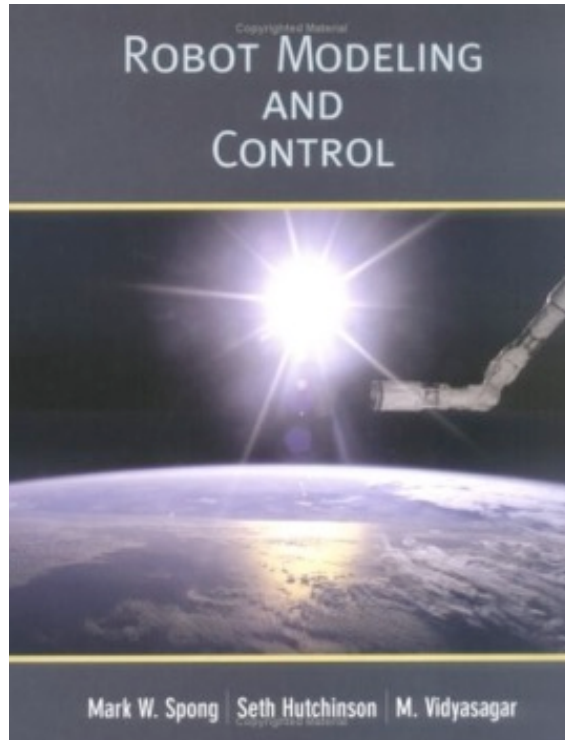
Lecture 20: Joint Space Dynamics

Cynthia Sung, Ph.D.

Mechanical Engineering & Applied Mechanics

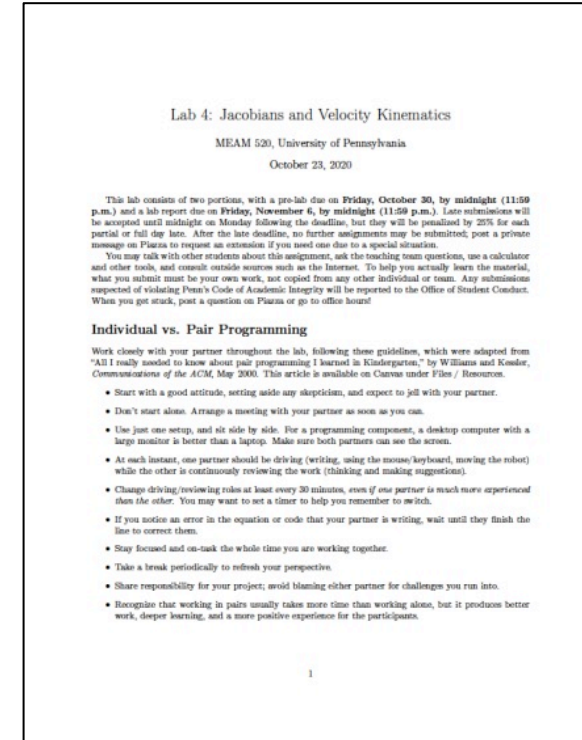
University of Pennsylvania

Today: Dynamics



Chapter 7: Dynamics

- Read 7.1-7.3



Lab 4 due 11/6

Lab 5 (last lab!) due 11/20

Previously: Manipulator Jacobian

$$J = \begin{bmatrix} J_v \\ J_\omega \end{bmatrix}$$

(6 x n) Jacobian
a.k.a. manipulator Jacobian
a.k.a. geometric Jacobian

(3 x n) linear velocity Jacobian

(3 x n) angular velocity Jacobian

$$J_v(\vec{q}) = \begin{bmatrix} \frac{\partial x}{\partial q_1} & \frac{\partial x}{\partial q_2} & \cdots & \frac{\partial x}{\partial q_n} \\ \frac{\partial y}{\partial q_1} & \frac{\partial y}{\partial q_2} & \cdots & \frac{\partial y}{\partial q_n} \\ \frac{\partial z}{\partial q_1} & \frac{\partial z}{\partial q_2} & \cdots & \frac{\partial z}{\partial q_n} \end{bmatrix}$$

forward velocity kinematics

$$\xi = J(q)\dot{q}$$

(6 x 1) body velocity

(6 x n) Jacobian

(n x 1) joint velocities

$$J_\omega = [\rho_1 \hat{\mathbf{z}} \quad \rho_2 \mathbf{R}_1^0 \hat{\mathbf{z}} \quad \rho_3 \mathbf{R}_2^0 \hat{\mathbf{z}} \quad \cdots \quad \rho_n \mathbf{R}_{n-1}^0 \hat{\mathbf{z}}]$$

$$\rho_i = \begin{cases} 0 & \text{for prismatic} \\ 1 & \text{for revolute} \end{cases}$$

inverse velocity kinematics

$$\dot{q} = J^{-1} \xi$$

Previously: Static Force/Torque Relationships

$$\begin{array}{ccccc}
 (n \times 1) & (n \times 6) & (6 \times 1) & & \\
 \vec{\tau} & = & J^T(\vec{q}) & \vec{F} & \\
 \uparrow & & \uparrow & & \uparrow \\
 \text{joint} & & & & \text{endpoint} \\
 \text{forces and} & & & & \text{forces and} \\
 \text{torques} & & & & \text{torques} \\
 & & \uparrow & & \\
 & & \text{Jacobian} & & \\
 & & \text{matrix} & & \\
 & & \text{transpose} & &
 \end{array}$$

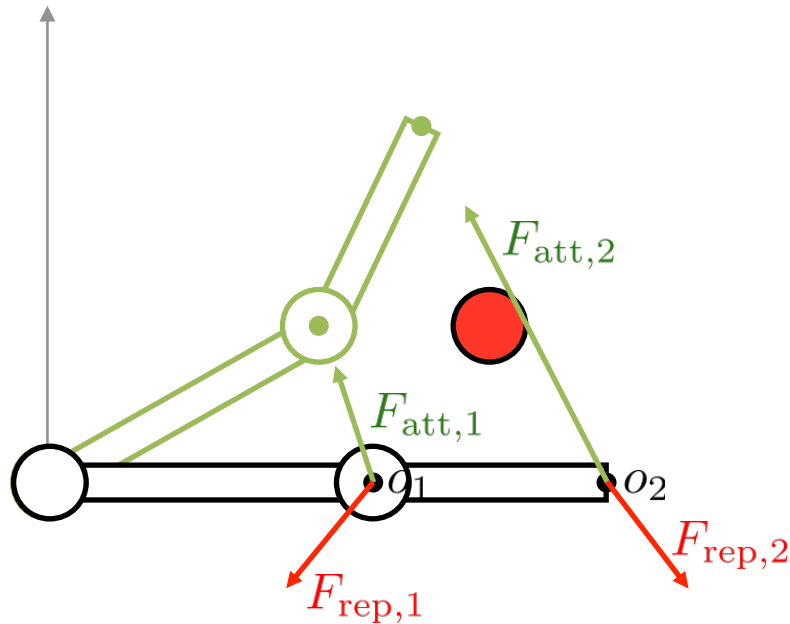
Simplest to think about for
a 3-DOF robot with all
revolute joints.
We want to output a force
at the tip.

$$\begin{array}{ccccc}
 (3 \times 1) & (3 \times 3) & (3 \times 1) & & \\
 \vec{\tau} & = & J^T(\vec{q}) & \vec{F} & \\
 \uparrow & & \uparrow & & \uparrow \\
 \text{joint} & & & & \text{endpoint} \\
 \text{torques} & & & & \text{forces} \\
 & & \uparrow & & \\
 & & \text{Jacobian} & & \\
 & & \text{matrix} & & \\
 & & \text{transpose} & &
 \end{array}$$

Derivation

$$\begin{aligned}
 \vec{\tau}^\top d\vec{q} &= \vec{F}^\top d\vec{x} \\
 d\vec{x} &= J_v d\vec{q} \\
 \vec{\tau}^\top d\vec{q} &= \vec{F}^\top J_v d\vec{q} \\
 \vec{\tau}^\top &= \vec{F}^\top J_v \\
 \vec{\tau} &= J_v^\top \vec{F}
 \end{aligned}$$

Previously: Potential Fields



$$F_{\text{att},i}(q) = -\zeta_i (o_i(q) - o_i(q_f))$$

when $\rho_i(q) \leq \rho_0$

$$F_{\text{rep},i}(q) = \eta_i \left(\frac{1}{\rho(o_i(q))} - \frac{1}{\rho_0} \right) \frac{1}{\rho^2(o_i(q))} \nabla \rho(o_i(q))$$

$$1. \quad q^0 \leftarrow q_s, i \leftarrow 0$$

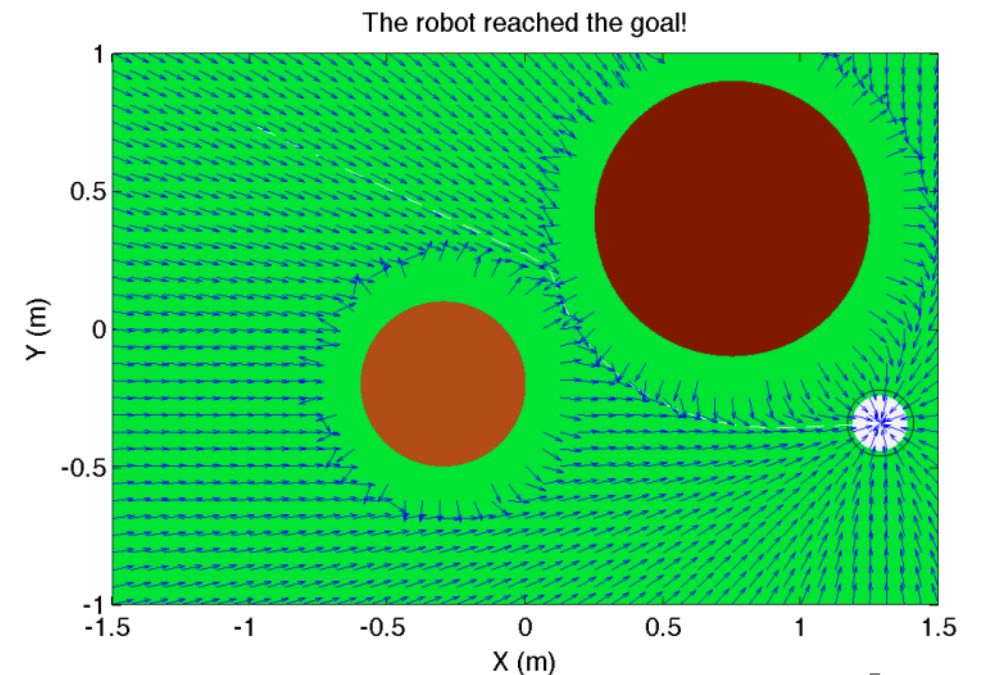
$$2. \quad \text{IF } \|q^i - q_f\| > \epsilon$$

$$q^{i+1} \leftarrow q^i + \alpha^i \frac{\tau(q^i)}{\|\tau(q^i)\|}$$

$$i \leftarrow i + 1$$

ELSE return $\langle q^0, q^1, \dots, q^i \rangle$

3. GO TO 2



Questions from you

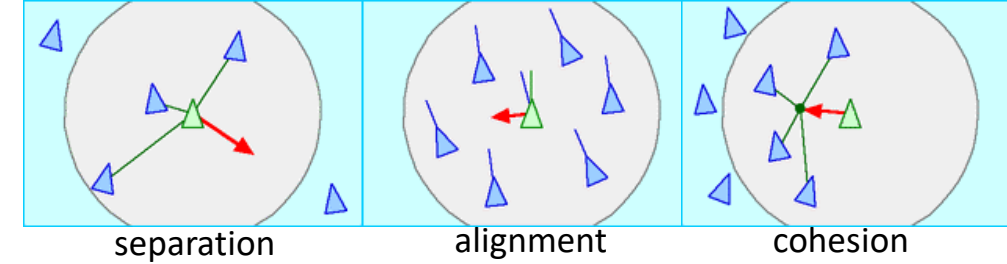
How do we know how to set the parameters?

The values of your parameters that will work depend highly on your environment and your robot.

The more complicated your system, the harder it will be to balance your parameters.

In general, increasing ζ will bring you to the goal faster, but may bring you closer to obstacles. Increasing η or ρ_0 will push you away more from obstacles, but if your environment is dense, it will create more local minima.

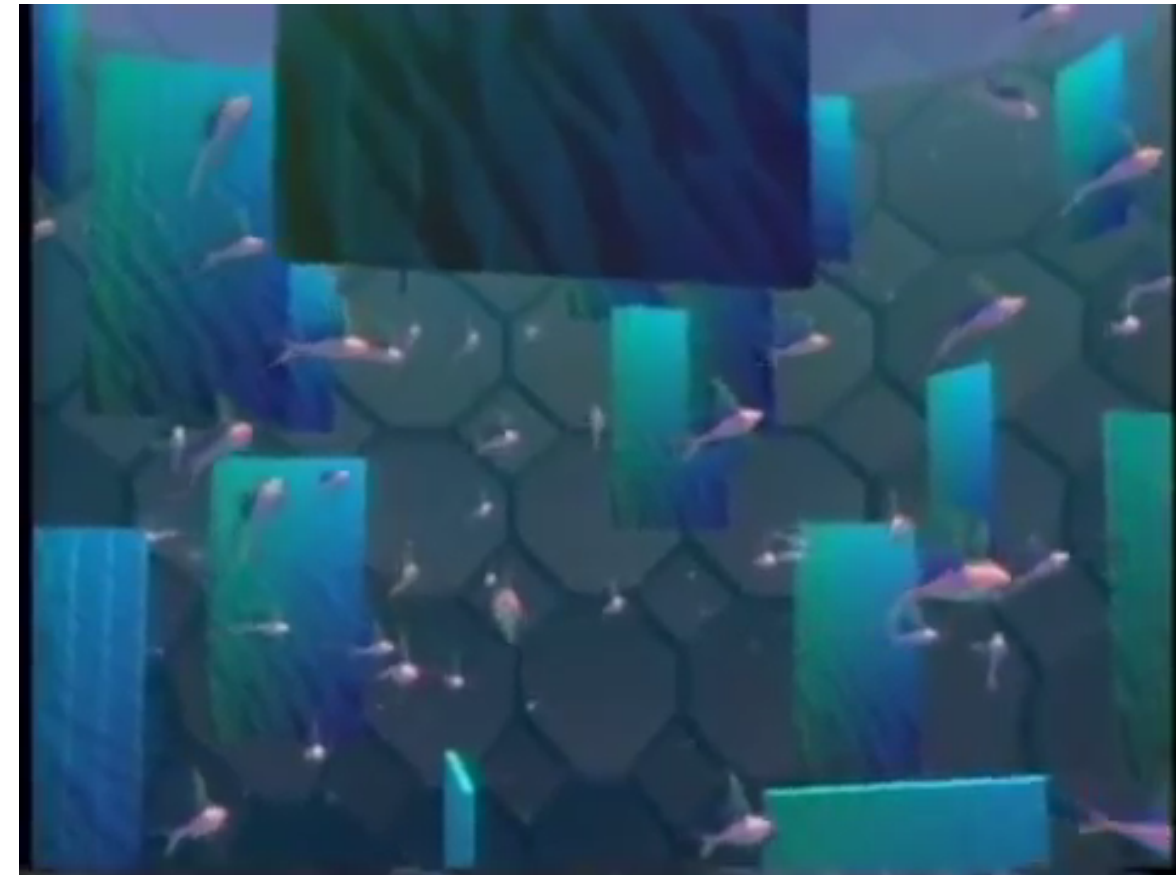
Questions from you



Why use potential fields if the robot gets stuck?

Potential fields are very strong in dynamic environments, where if there is a local minimum, it doesn't exist for very long.

Because you are applying “forces” on the robot, they also naturally extend to dynamic systems.



Boids, Stanley & Stella Breaking the Ice, 1987
<https://www.youtube.com/watch?v=3bTqWsvQyzE>

Dynamics

Kinematics: motion of the robot without consideration of the forces/torques producing motion

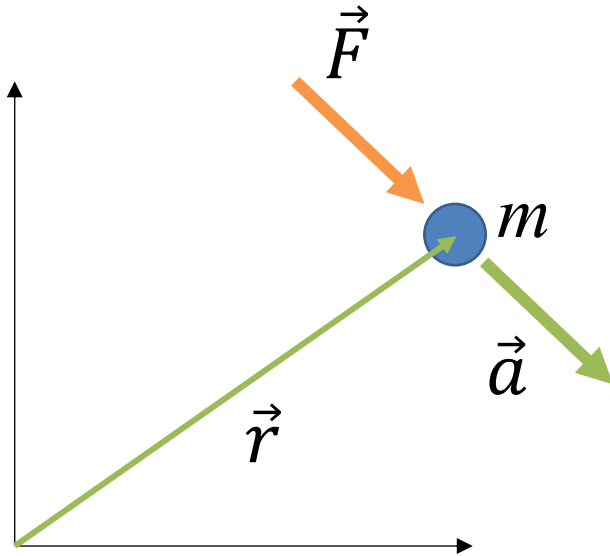
Dynamics: Relationship between forces and motion

Particle Dynamics



Hibbeler Ch. 13.1-13.2

Beer Ch. 12.1



Position: $\vec{r}(t)$

Velocity: $\vec{v}(t) = \frac{d\vec{r}}{dt}$

Acceleration: $\vec{a}(t) = \frac{d\vec{v}}{dt} = \frac{d^2\vec{r}}{dt^2}$

Newton's Second Law: $\vec{F}(t) = m\vec{a}(t)$

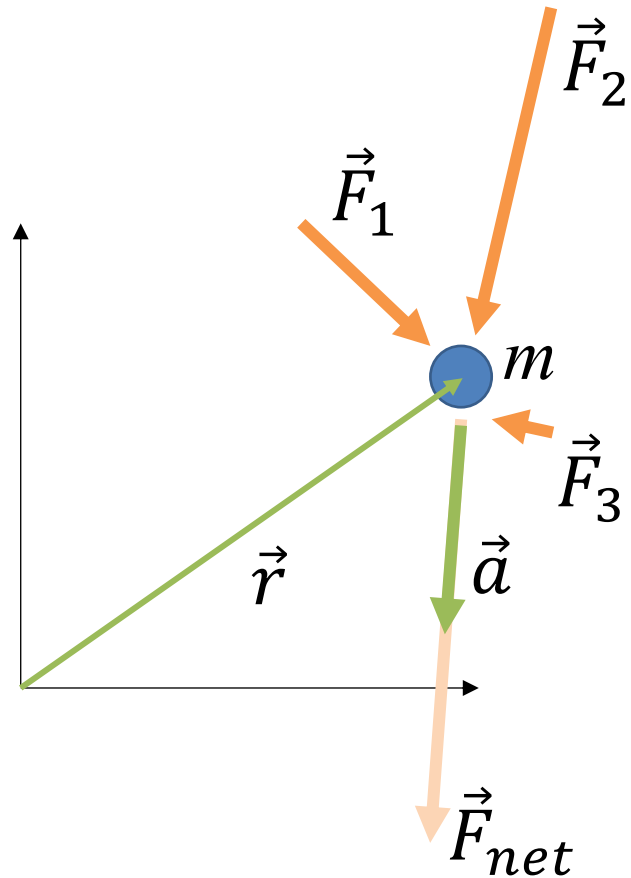
$$\vec{F}(t) = m \frac{d\vec{v}}{dt}$$

$$\vec{F}(t) = \frac{d\vec{p}}{dt} \text{ linear momentum}$$

Particle Dynamics



Hibbeler Ch. 14
Beer Ch. 13.1-13.2



$$\sum_i \vec{F}_i(t) = m\vec{a}(t)$$

$$\text{Kinetic Energy: } K = \frac{1}{2} m \vec{v}^T \vec{v}$$

$$\text{Work: } W = \int \vec{F}_{net} \cdot d\vec{r}$$

$$W = \int \vec{F}_C \cdot d\vec{r} + \int \vec{F}_{NC} \cdot d\vec{r}$$

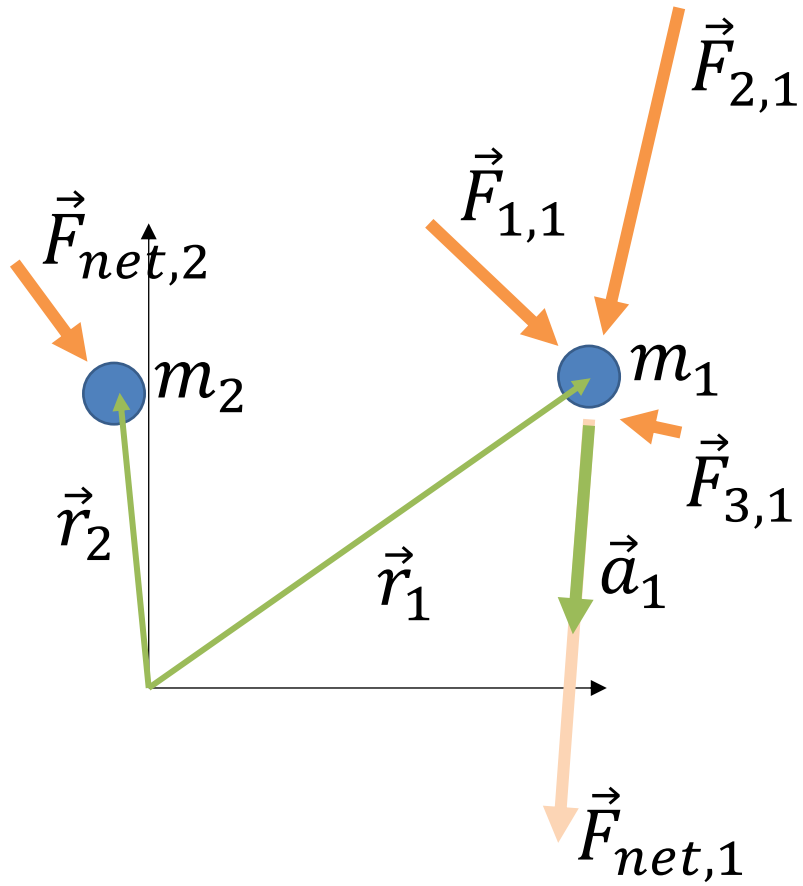
$$\text{Potential Energy: } P = - \int \vec{F}_C \cdot d\vec{r}$$

Multiple Particles



Hibbeler Ch. 13.3, 14.3

Beer Ch. 14



Particle j : $\vec{F}_{net,j}(t) = m_j \vec{a}_j(t)$

Kinetic Energy: $K = \sum_j \frac{1}{2} m_j \vec{v}_j^T \vec{v}_j$

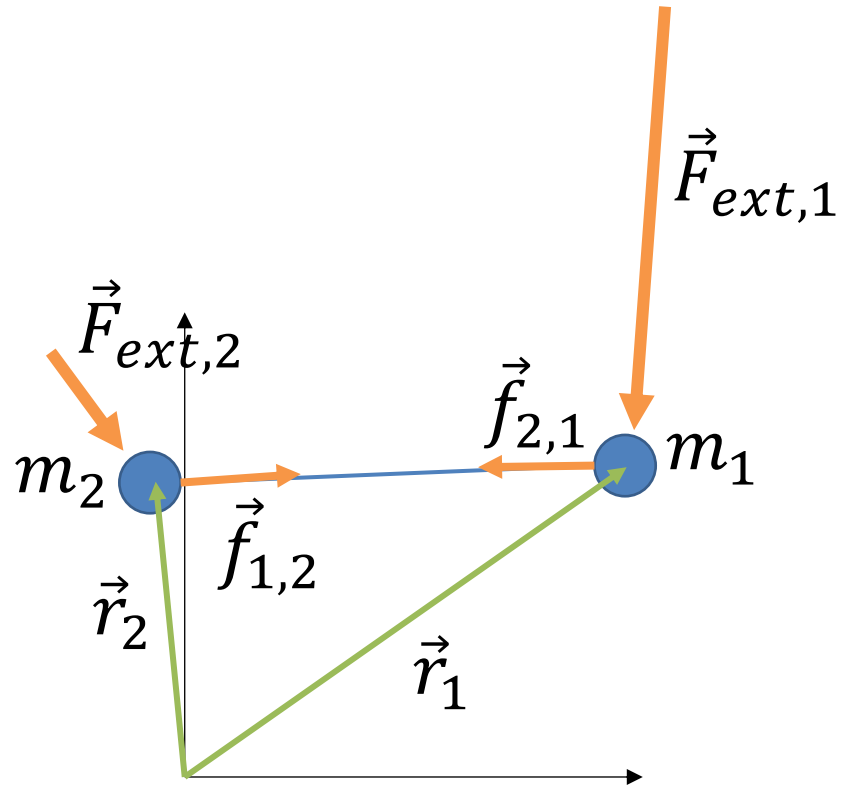
Potential Energy: $P = - \sum_j \int \vec{F}_{C,j} \cdot d\vec{r}_j$

Particles with Constraints



Hibbeler Ch. 13.3

Beer Ch. 14



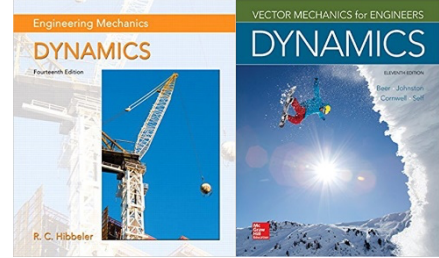
Internal forces: $\vec{f}_{i,j}(t) = -\vec{f}_{j,i}(t)$

$$\sum_j \sum_i \vec{f}_{i,j} = 0$$

$$\vec{F}_{net,sys} = \sum_j \left(\vec{F}_{ext,j} + \cancel{\sum_i \vec{f}_{i,j}} \right)$$

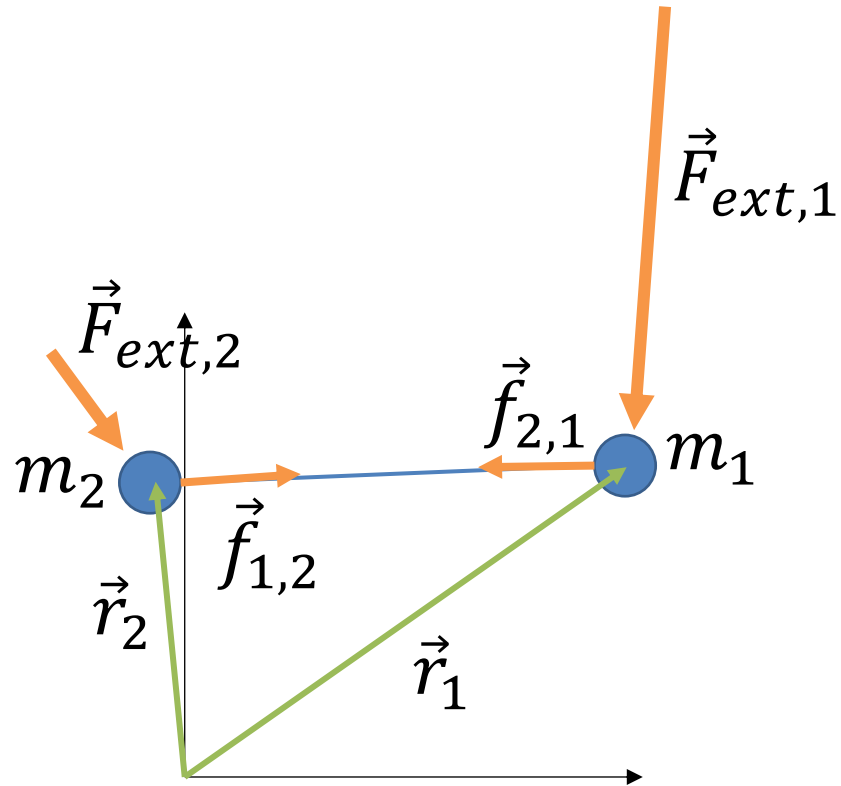
$$\vec{F}_{net,sys} = \sum_j \vec{F}_{ext,j}$$

Particles with Constraints



Hibbeler Ch. 13.3

Beer Ch. 14



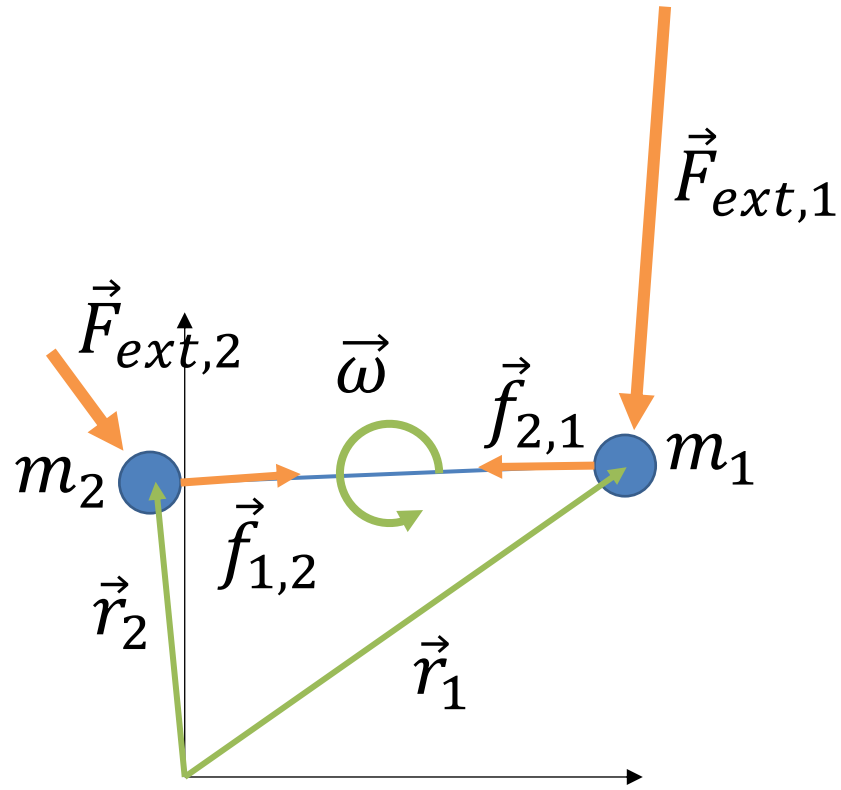
$$\vec{F}_{net,sys} = \sum_j \vec{F}_{ext,j} = \underbrace{\sum_j m_j \vec{a}_j}_{\text{Newton's 2nd Law}}$$

$$\vec{F}_{net,sys} = m_{tot} \sum_j \frac{m_j \vec{a}_j}{m_{tot}}$$

$$= m_{tot} \sum_j \frac{m_j \left(\frac{d^2 \vec{r}_j}{dt^2} \right)}{m_{tot}}$$

$$= m_{tot} \frac{d^2}{dt^2} \boxed{\frac{\sum_j m_j \vec{r}_j}{m_{tot}}} \quad \text{Center of Mass}$$

Particles with Constraints

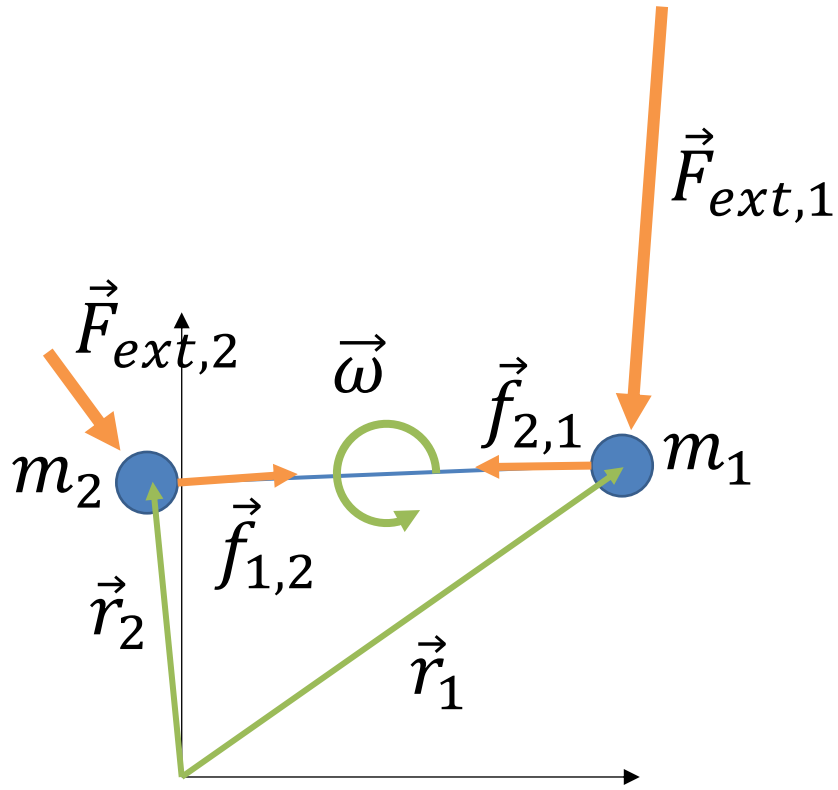


Inertia tensor: $[I]_{3 \times 3}$

Euler Equation: $\sum \vec{\tau}_{COM} = [I]_{COM} \vec{\alpha}$

Euler Equation: $\sum \vec{\tau}_p = [I]_p \vec{\alpha} + \vec{r}_{p/COM} \times m_{tot} \vec{a}_p$

Particles with Constraints



Position: $\vec{r}(t)$

Velocity: $\vec{v}(t) = \frac{d\vec{r}}{dt}$

Acceleration: $\vec{a}(t) = \frac{d\vec{v}}{dt} = \frac{d^2\vec{r}}{dt^2}$

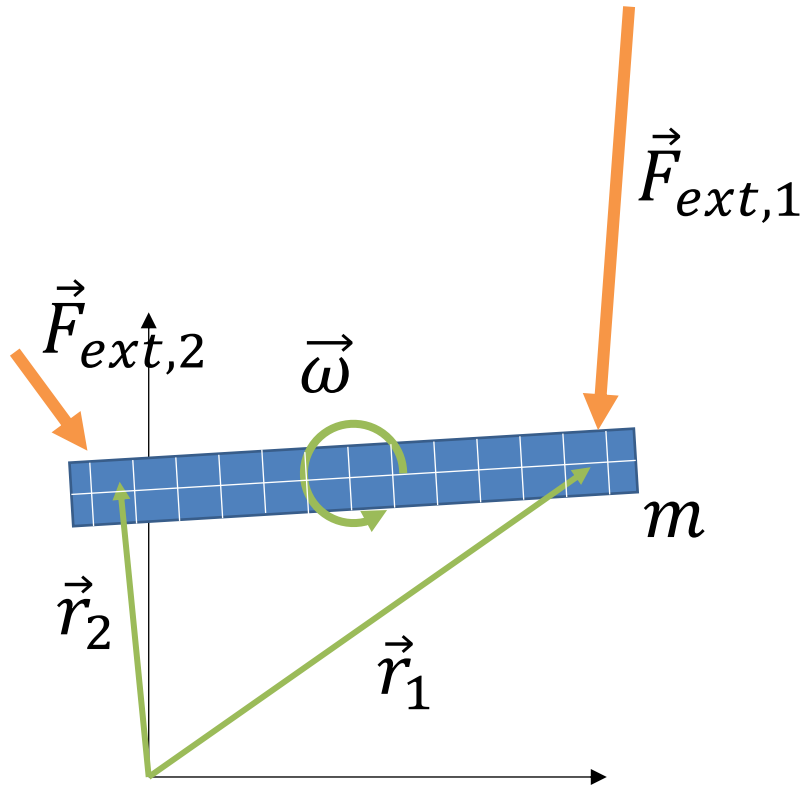
Kinematic Constraints

$$\vec{r}_2 = \vec{r}_1 + \vec{r}_{2/1}$$

$$\vec{v}_2 = \frac{d\vec{r}_2}{dt} = \vec{v}_1 + \vec{\omega} \times \vec{r}_{2/1}$$

$$\vec{a}_2 = \frac{d\vec{v}_2}{dt} = \vec{a}_1 + \vec{\alpha} \times \vec{r}_{2/1} + \vec{\omega} \times (\vec{\omega} \times \vec{r}_{2/1})$$

Rigid Bodies



$$\vec{r}_{COM} = \frac{1}{m} \int \vec{r} dm$$

$$I = \begin{bmatrix} I_{xx} & I_{xy} & I_{xz} \\ I_{yx} & I_{yy} & I_{yz} \\ I_{zx} & I_{zy} & I_{zz} \end{bmatrix}$$

*I depends on
your frame!*

$$I_{xx} = \iiint (y^2 + z^2) dm$$

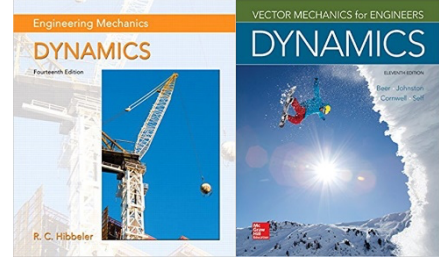
$$I_{xy} = I_{yx} = - \iiint xy dm$$

$$I_{yy} = \iiint (x^2 + z^2) dm$$

$$I_{xz} = I_{zx} = - \iiint xz dm$$

$$I_{zz} = \iiint (x^2 + y^2) dm$$

$$I_{yz} = I_{zy} = - \iiint yz dm$$



Hibbeler Ch. 21.1

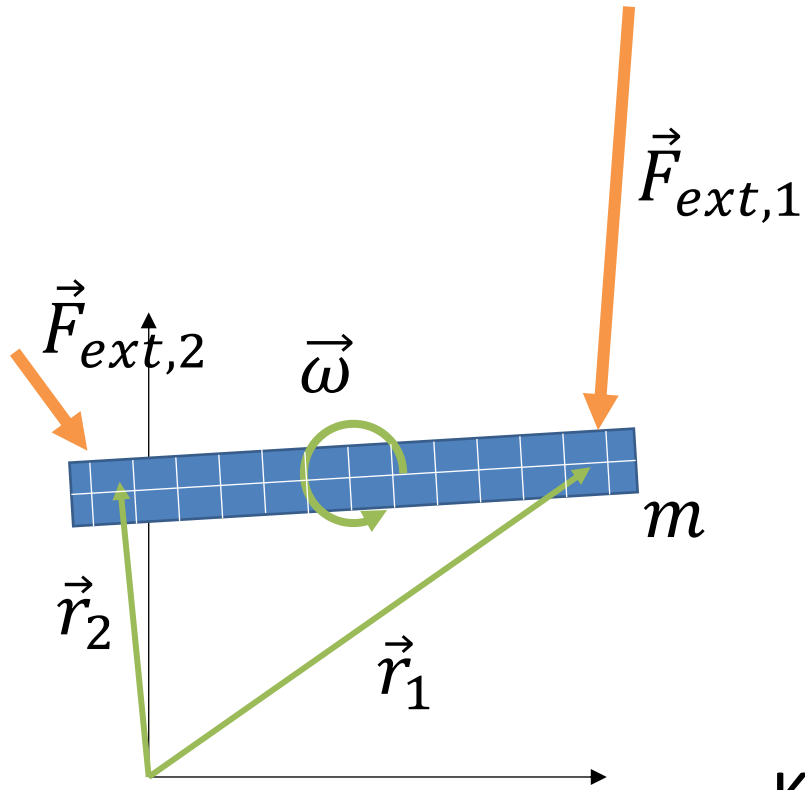
Beer Ch. 18.2

Rigid Bodies



Hibbeler Ch. 21.3

Beer Ch. 18.1

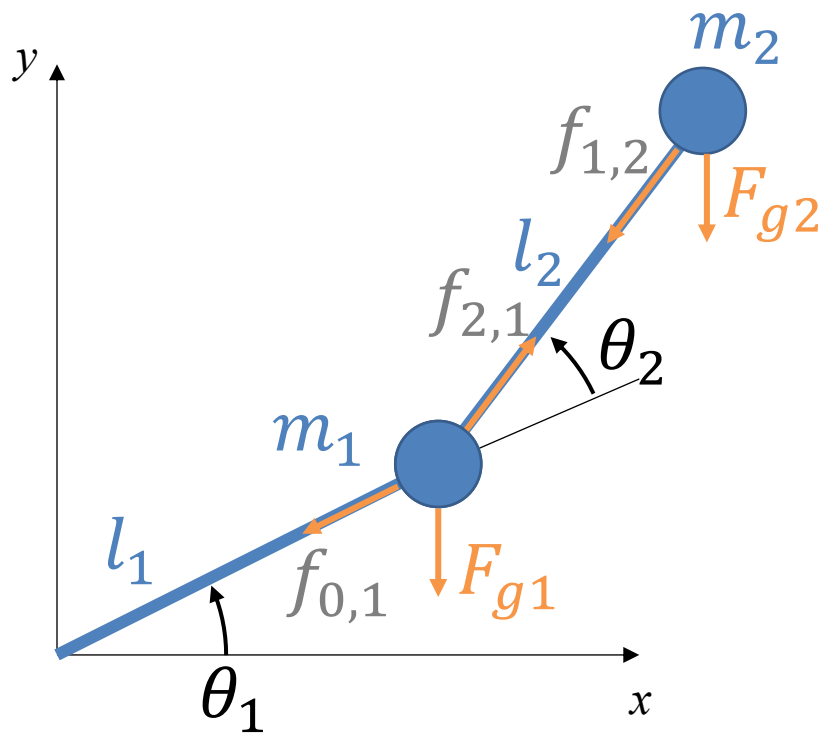


$$\text{Kinetic Energy: } K = \frac{1}{2} m \vec{v}_{COM}^T \vec{v}_{COM} + \frac{1}{2} \vec{\omega}^T I_{COM} \vec{\omega}$$

$$\text{Gravitational Potential Energy: } P_g = m z_{COM}$$

Example: RR Manipulator w/ mass concentrated at ends of links

AKA the double pendulum



Newton:

$$m_1 \vec{a}_1 = \vec{F}_{g1} + \vec{f}_{0,1} + \vec{f}_{2,1}$$

$$m_2 \vec{a}_2 = \vec{F}_{g2} + \vec{f}_{1,2}$$

9 unknowns
9 equations

Constraints:

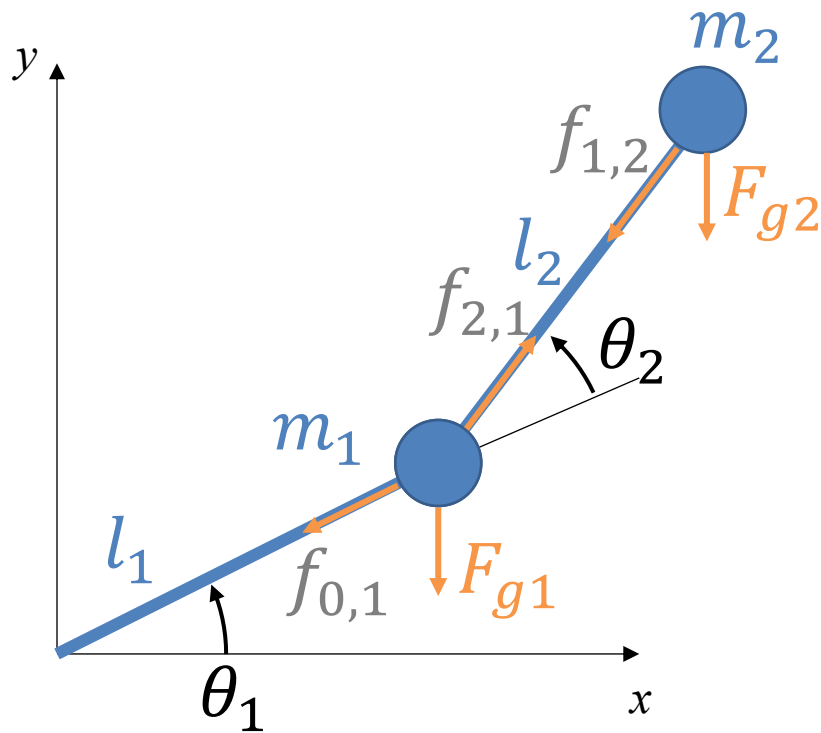
$$\vec{a}_1 = \vec{\alpha}_1 \times \vec{r}_1 - \omega_1^2 \vec{r}_1$$

$$\vec{a}_2 = \vec{a}_1 + \vec{\alpha}_2 \times \vec{r}_{2/1} - \omega_2^2 \vec{r}_{2/1}$$

$$f_{1,2} = f_{2,1}$$

Example: RR Manipulator w/ mass concentrated at ends of links

AKA the double pendulum



EOM (sub $\alpha = \ddot{\theta}$, $\omega = \dot{\theta}$):

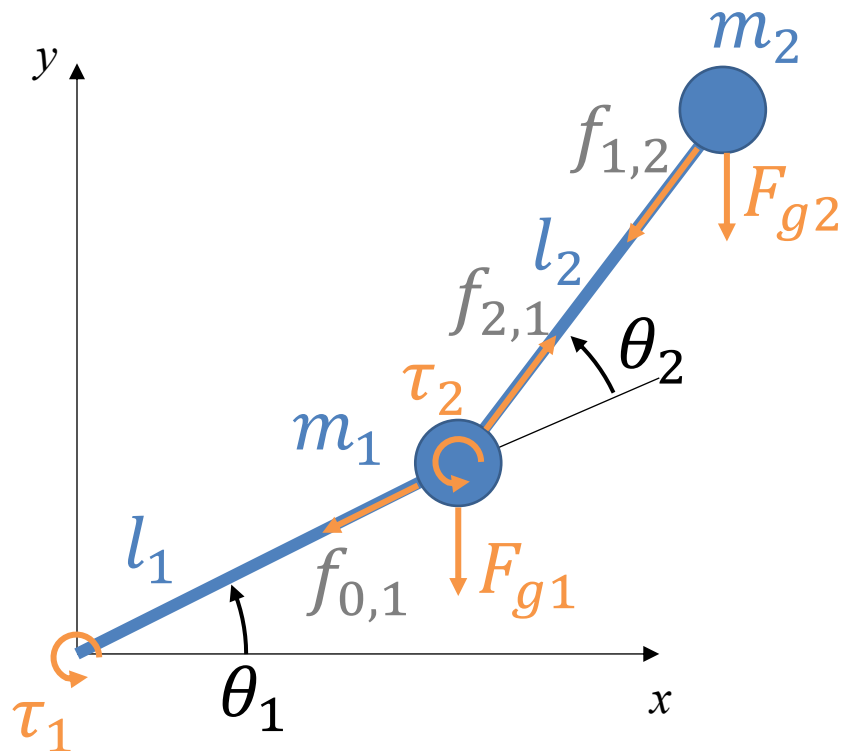
$$\begin{aligned} & [m_1 l_1^2 + m_2 (l_1^2 + l_2^2 + 2l_1 l_2 c_2)] \ddot{\theta}_1 \\ & + [m_2 (l_2^2 + l_1 l_2 c_2)] \ddot{\theta}_2 - m_2 l_1 l_2 s_2 (2\dot{\theta}_1 \dot{\theta}_2 + \dot{\theta}_2^2) \\ & + m_1 g l_1 c_1 + m_2 g (l_1 c_1 + l_2 c_2) = 0 \end{aligned}$$

$$\begin{aligned} & [m_2 (l_2^2 + l_1 l_2 c_2)] \ddot{\theta}_1 + m_2 l_2 \ddot{\theta}_2 + m_2 l_1 l_2 s_2 \dot{\theta}_1^2 \\ & + m_2 g l_2 c_2 = 0 \end{aligned}$$

Example: RR Manipulator w/ mass concentrated at ends of links

AKA the double pendulum

EOM (sub $\alpha = \ddot{\theta}$, $\omega = \dot{\theta}$):



coefficients of \ddot{q}_i depend only on q

$$[m_1 l_1^2 + m_2(l_1^2 + l_2^2 + 2l_1 l_2 c_2)] \ddot{\theta}_1$$

$$+ [m_2(l_2^2 + l_1 l_2 c_2)] \ddot{\theta}_2 - m_2 l_1 l_2 s_2 (2\dot{\theta}_1 \dot{\theta}_2 + \dot{\theta}_2^2)$$

$$+ m_1 g l_1 c_1 + m_2 g (l_1 c_1 + l_2 c_2) = \tau_1$$

$$[m_2(l_2^2 + l_1 l_2 c_2)] \ddot{\theta}_1 + m_2 l_2 \ddot{\theta}_2 + m_2 l_1 l_2 s_2 \dot{\theta}_1^2$$

$$+ m_2 g l_2 c_{12} = \tau_2$$

centrifugal and Coriolis terms
depend on q and \dot{q}

gravitational terms depend only on q

The Manipulator Equation

We can write this as a matrix equation

$$\tau = \underline{D(q)}\ddot{q} + C(q, \dot{q})\dot{q} + g(q)$$

SHV uses a bit of strange notation.

Most people call this matrix H or M .

where

$D(q)$ is the $n \times n$ mass matrix (inertia terms)

$C(q, \dot{q})$ is the $n \times n$ matrix of centrifugal (square of joint velocities) and Coriolis (product of two different joint velocities) terms

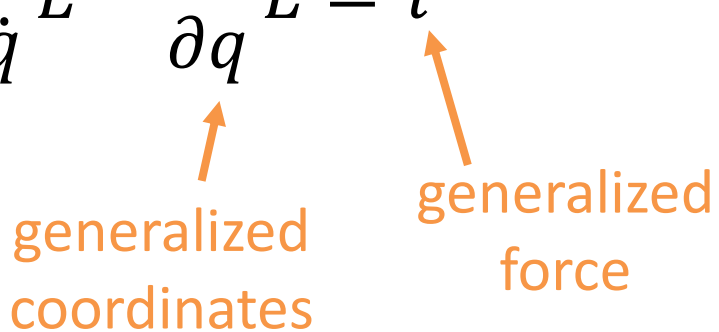
$g(q)$ is a $n \times 1$ vector of gravitational terms

Another Method: Euler-Lagrange Equation

Derivation SHV 7.1.3

Lagrangian: $L = K - P$

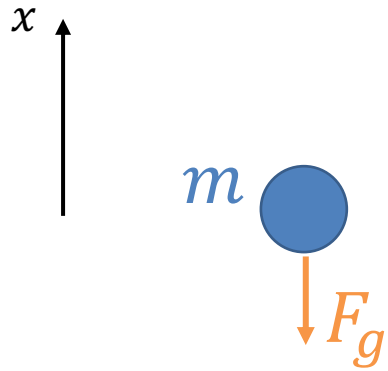
EOM: $\frac{d}{dt} \frac{\partial}{\partial \dot{q}} L - \frac{\partial}{\partial q} L = \tau$



generalized
coordinates

generalized
force

Example: Particle under Gravity



Kinetic energy: $K = \frac{1}{2} m \dot{x}^2$

Potential energy: $P = mgx$

Lagrangian: $L = K - P = \frac{1}{2} m \dot{x}^2 - mgx$

$$\frac{\partial}{\partial x} L = -mg \quad \frac{\partial}{\partial \dot{x}} L = m\dot{x}$$

EOM: $\frac{d}{dt} \frac{\partial}{\partial \dot{q}} L - \frac{\partial}{\partial q} L = \tau$

$$m\ddot{x} + mg = \tau$$

Euler-Lagrange Equation

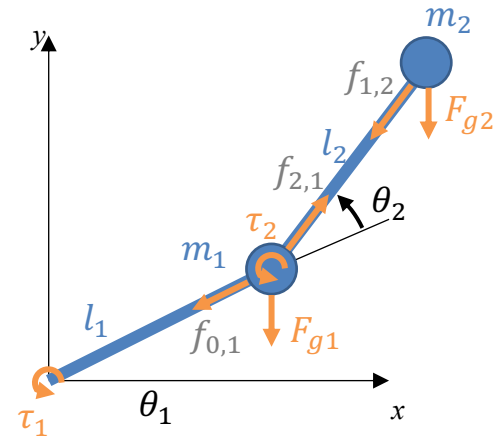
Kinetic Energy K

$$\text{Link 1: } \begin{bmatrix} x_1 \\ y_1 \end{bmatrix} = \begin{bmatrix} l_1 c_1 \\ l_1 s_1 \end{bmatrix} \Rightarrow \begin{bmatrix} \dot{x}_1 \\ \dot{y}_1 \end{bmatrix} = \begin{bmatrix} -l_1 s_1 \dot{\theta}_1 \\ l_1 c_1 \dot{\theta}_1 \end{bmatrix}$$

$$\text{Link 2: } \begin{bmatrix} x_2 \\ y_2 \end{bmatrix} = \begin{bmatrix} l_1 c_1 + l_2 c_{12} \\ l_1 s_1 + l_2 s_{12} \end{bmatrix} \Rightarrow \begin{bmatrix} \dot{x}_2 \\ \dot{y}_2 \end{bmatrix} = \begin{bmatrix} (-l_1 s_1 - l_2 s_{12}) \dot{\theta}_1 - l_2 s_{12} \dot{\theta}_2 \\ (l_1 c_1 + l_2 c_{12}) \dot{\theta}_1 + l_2 c_{12} \dot{\theta}_2 \end{bmatrix}$$

$$K = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 v_2^2$$

$$= \frac{1}{2} m_1 l_1^2 \dot{\theta}_1^2 + \frac{1}{2} m_2 [(l_1^2 + 2l_1 l_2 c_2 + l_2^2) \dot{\theta}_1^2 + 2(l_2^2 + l_1 l_2 c_2) \dot{\theta}_1 \dot{\theta}_2 + l_2^2 \dot{\theta}_2^2]$$



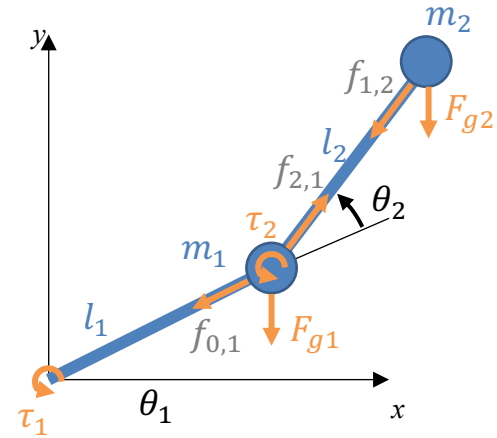
Example: RR manipulator

Potential Energy P

$$\begin{bmatrix} x_1 \\ y_1 \end{bmatrix} = \begin{bmatrix} l_1 c_1 \\ l_1 s_1 \end{bmatrix}$$

$$\begin{bmatrix} x_2 \\ y_2 \end{bmatrix} = \begin{bmatrix} l_1 c_1 + l_2 c_{12} \\ l_1 s_1 + l_2 s_{12} \end{bmatrix}$$

$$\begin{aligned} P &= m_1 g y_1 + m_2 g y_2 \\ &= m_1 g l_1 s_1 + m_2 g (l_1 s_1 + l_2 s_{12}) \end{aligned}$$



Example: RR manipulator

Equation of Motion

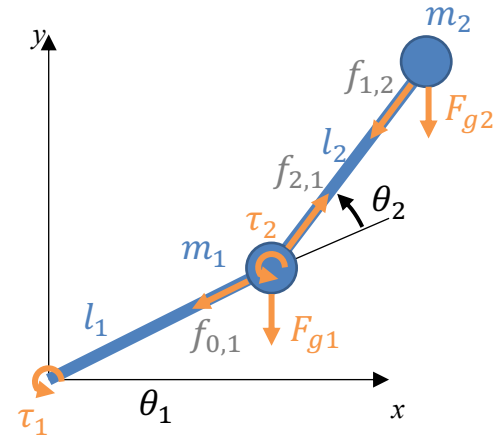
$$L = K - P \quad \tau = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q}$$

$$K = \frac{1}{2} m_1 l_1^2 \dot{\theta}_1^2 + \frac{1}{2} m_2 [(l_1^2 + 2l_1 l_2 c_2 + l_2^2) \dot{\theta}_1^2 + 2(l_2^2 + l_1 l_2 c_2) \dot{\theta}_1 \dot{\theta}_2 + l_2^2 \dot{\theta}_2^2]$$

$$P = m_1 g l_1 s_1 + m_2 g (l_1 s_1 + l_2 s_{12})$$

$$\frac{\partial L}{\partial \dot{q}} = \frac{\partial K}{\partial \dot{q}} = \begin{bmatrix} m_1 l_1^2 \dot{\theta}_1 + m_2 (l_1^2 + 2l_1 l_2 c_2 + l_2^2) \dot{\theta}_1 + m_2 (l_2^2 + l_1 l_2 c_2) \dot{\theta}_2 \\ m_2 (l_2^2 + l_1 l_2 c_2) \dot{\theta}_1 + m_2 l_2^2 \dot{\theta}_2 \end{bmatrix}$$

$$\frac{\partial L}{\partial q} = \begin{bmatrix} -m_1 g l_1 c_1 - m_2 g l_1 c_1 - m_2 g l_2 c_{12} \\ -m_2 l_1 l_2 s_2 \dot{\theta}_1^2 - m_2 l_1 l_2 s_2 \dot{\theta}_1 \dot{\theta}_2 - m_2 g l_2 c_{12} \end{bmatrix}$$



Example: RR manipulator

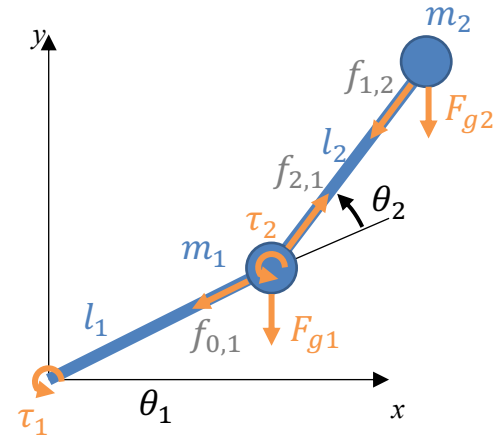
Equation of Motion

$$L = K - P \quad \tau = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q}$$

$$\frac{\partial L}{\partial \dot{q}} = \begin{bmatrix} m_1 l_1^2 \dot{\theta}_1 + m_2 (l_1^2 + 2l_1 l_2 c_2 + l_2^2) \dot{\theta}_1 + m_2 (l_2^2 + l_1 l_2 c_2) \dot{\theta}_2 \\ m_2 (l_2^2 + l_1 l_2 c_2) \dot{\theta}_1 + m_2 l_2^2 \dot{\theta}_2 \end{bmatrix}$$

$$\frac{\partial L}{\partial q} = \begin{bmatrix} -m_1 g l_1 c_1 - m_2 g l_1 c_1 - m_2 g l_2 c_{12} \\ -m_2 l_1 l_2 s_2 \dot{\theta}_1^2 - m_2 l_1 l_2 s_2 \dot{\theta}_1 \dot{\theta}_2 - m_2 g l_2 c_{12} \end{bmatrix}$$

$$\tau = \begin{bmatrix} m_1 l_1^2 \ddot{\theta}_1 + m_2 (l_1^2 + 2l_1 l_2 c_2 + l_2^2) \ddot{\theta}_1 - 2m_2 l_1 l_2 s_2 \dot{\theta}_1 \dot{\theta}_2 + m_2 (l_2^2 + l_1 l_2 c_2) \ddot{\theta}_2 \\ -m_2 l_1 l_2 s_2 \dot{\theta}_2^2 + m_1 g l_1 c_1 + m_2 g l_2 c_{12} \\ m_2 (l_2^2 + l_1 l_2 c_2) \ddot{\theta}_2 + m_2 l_2^2 \ddot{\theta}_2 + m_2 l_1 l_2 s_2 \dot{\theta}_1^2 + m_2 g l_2 c_{12} \end{bmatrix}$$



$$\tau = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q}$$

Observations

Kinetic Energy $K = \frac{1}{2} m_1 \vec{v}_1^\top \vec{v}_1 + \frac{1}{2} m_2 \vec{v}_2^\top \vec{v}_2$

$$K = \frac{1}{2} m_1 (J_{v1} \dot{q})^\top (J_{v1} \dot{q}) + \frac{1}{2} m_2 (J_{v2} \dot{q})^\top (J_{v2} \dot{q})$$

Linear velocity Jacobian: $v_i = J_{vi} \dot{q}$

$$K = \frac{1}{2} m_1 \dot{q}^\top J_{v1}^\top J_{v1} \dot{q} + \frac{1}{2} m_2 \dot{q}^\top J_{v2}^\top J_{v2} \dot{q}$$

$(AB)^\top = B^\top A^\top$

$$K = \frac{1}{2} \dot{q}^\top \underbrace{(m_1 J_{v1}^\top J_{v1} + m_2 J_{v2}^\top J_{v2})}_{\text{Function of } q} \dot{q}$$

$\Rightarrow \frac{\partial}{\partial \dot{q}} (\quad) = 0$

$$\frac{\partial}{\partial \dot{q}} K = \frac{1}{2} [(m_1 J_{v1}^\top J_{v1} + m_2 J_{v2}^\top J_{v2}) \dot{q}]^\top + \frac{1}{2} \dot{q}^\top (m_1 J_{v1}^\top J_{v1} + m_2 J_{v2}^\top J_{v2})$$

$$= \dot{q}^\top \boxed{(m_1 J_{v1}^\top J_{v1} + m_2 J_{v2}^\top J_{v2})}$$

Inertia Matrix D $\left\{ \begin{array}{l} \text{symmetric} \\ \text{positive definite} \end{array} \right.$

$$\tau = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q}$$

Observations

$$\text{Lagrangian: } L = K - P = \frac{1}{2} \dot{q}^\top D \dot{q} - P$$

all terms contain \dot{q} depends only on q

$$\text{Manipulator equation: } \tau = D(q)\ddot{q} + C(q, \dot{q})\dot{q} + g(q)$$

computed using D only

$$\frac{\partial}{\partial q} P$$

N -link manipulator w/ mass concentrated at ends of links $\tau = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q}$

Inertia:

$$N = 2: D = m_1 J_{v1}^\top J_{v1} + m_2 J_{v2}^\top J_{v2}$$

$$\text{general case: } D = \sum_{i=1}^N m_i J_{vi}^\top J_{vi}$$

Gravity:

$$N = 2: P = m_1 g l_1 s_1 + m_2 g (l_1 s_1 + l_2 s_{12})$$

$$\text{general case: } P = \sum_{i=1}^N m_i \vec{g} \cdot \vec{r}_i$$

$$g(q) = \frac{\partial}{\partial q} P$$

$$\tau = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q}$$

What about **C**?

$$L = \frac{1}{2} \dot{q}^\top D \dot{q} - P = \frac{1}{2} \sum_{i,j} d_{ij} \dot{q}_i \dot{q}_j - P$$

$$\frac{\partial}{\partial q_k} L = \frac{1}{2} \sum_{i,j} \frac{\partial}{\partial q_k} d_{ij} \dot{q}_i \dot{q}_j - \frac{\partial}{\partial q_k} P$$

gravitational terms – ignore from here on

$$\frac{\partial}{\partial \dot{q}_k} L = \sum_j d_{kj} \dot{q}_j$$

$$\frac{d}{dt} \frac{\partial}{\partial \dot{q}_k} L = \sum_j d_{kj} \ddot{q}_j + \sum_j \left[\frac{d}{dt} d_{kj} \right] \dot{q}_j$$

$$(C\dot{q})_k = \sum_{i,j} \left(\frac{\partial d_{kj}}{\partial q_i} - \frac{1}{2} \frac{\partial d_{ij}}{\partial q_k} \right) \dot{q}_i \dot{q}_j$$

$$= \sum_j d_{kj} \ddot{q}_j + \sum_{i,j} \frac{\partial d_{kj}}{\partial q_i} \dot{q}_i \dot{q}_j$$

inertia terms – ignore from here on

$$(C\dot{q})_k = \sum_{i,j} \frac{1}{2} \left(\frac{\partial d_{kj}}{\partial q_i} + \frac{\partial d_{ki}}{\partial q_j} - \frac{\partial d_{ij}}{\partial q_k} \right) \dot{q}_i \dot{q}_j$$

Christoffel symbols

Manipulator Equation

$$D(q)\ddot{q} + C(q, \dot{q})\dot{q} + g(q)$$

$$D = \sum_{i=1}^N m_i J_{vi}^T J_{vi}$$

$$g = \frac{\partial}{\partial q} \sum_{i=1}^N m_i \vec{g} \cdot \vec{r}_i$$

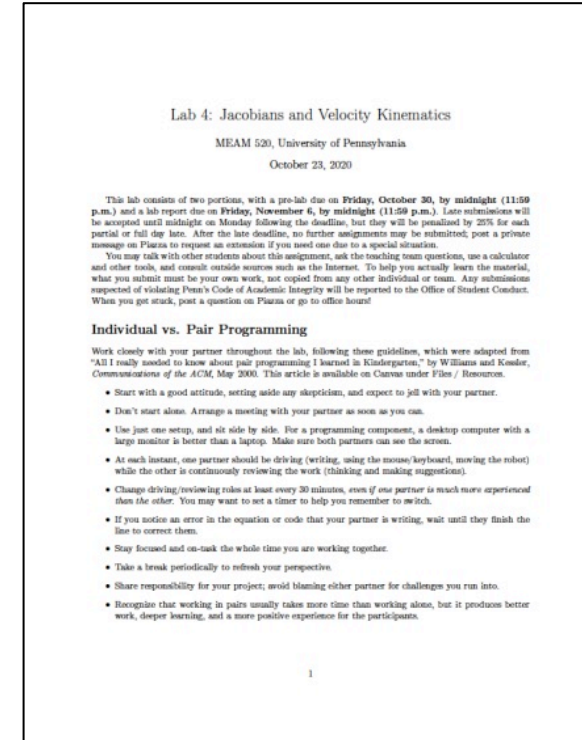
$$(C\dot{q})_k = \sum_{i,j} \frac{1}{2} \left(\frac{\partial d_{kj}}{\partial q_i} + \frac{\partial d_{ki}}{\partial q_j} - \frac{\partial d_{ij}}{\partial q_k} \right) \dot{q}_i \dot{q}_j \quad \text{or} \quad c_{kj} = \sum_i \frac{1}{2} \left(\frac{\partial d_{kj}}{\partial q_i} + \frac{\partial d_{ki}}{\partial q_j} - \frac{\partial d_{ij}}{\partial q_k} \right) \dot{q}_i$$

Next time: Non-point masses and alternative methods



Chapter 7: Dynamics

- Read 7.4-7.7



Lab 4 due 11/6

Lab 5 (last lab!) due 11/20