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BURSTT: Bustling Universe Radio Survey Telescope in Taiwan

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Abstract

Fast Radio Bursts (FRBs) are bright millisecond-duration radio transients that appear about 1000 times per day, all-sky, for a fluence threshold 5 Jy ms at 600 MHz. The FRB radio-emission physics and the compact objects involved in these events are subjects of intense and active debate. To better constrain source models, the Bustling Universe Radio Survey Telescope in Taiwan (BURSTT) is optimized to discover and localize a large sample of rare, high-fluence, and nearby FRBs. This population is the most amenable to multi-messenger and multi-wavelength follow-up, which allows a deeper understanding of source mechanisms. BURSTT will provide horizon-to-horizon sky coverage with a half power field-of-view (FoV) of $\sim 10^4~\rm deg^2$, a 400 MHz effective bandwidth between 300 and 800 MHz, and subarcsecond localization, which is made possible using outrigger stations that are hundreds to thousands of km from the main array. Initially, BURSTT will employ 256 antennas. After tests of various antenna designs and optimizing the system's performance, we plan to expand to 2048 antennas. We estimate that BURSTT-256 will detect and localize $\sim 100~\rm bright$ ($\geqslant 100~\rm Jy~ms$) FRBs per year. Another advantage of BURSTT's large FoV and continuous operation will be its greatly enhanced monitoring of FRBs for repetition. The current lack of sensitive all-sky observations likely means that many repeating FRBs are currently cataloged as single-event FRBs.

Unified Astronomy Thesaurus concepts: Radio transient sources (2008); Astronomical instrumentation (799); Wide-field telescopes (1800); Very long baseline interferometry (1769)

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1. Introduction

Fast Radio Bursts (FRBs) are bright (\sim 1 Jy), millisecond flashes of radio light of uncertain astrophysical origin (Lorimer et al. 2007; Tendulkar et al. 2017; Macquart et al. 2020). With a fluence threshold above 5 Jy ms at 600 MHz, the all-sky occurrence rates of FRBs is \sim 1000 per day (e.g., CHIME/FRB Collaboration et al. 2018, 2021). Their spatial distribution is independent of the galactic latitude (Bhandari et al. 2018; Josephy et al. 2021).

The nature of FRBs (including their emission mechanism, central object type, and environment) is one of the most perplexing enigmas in astrophysics. Deepening the mystery, about 4% of FRB sources emit multiple bursts (CHIME/FRB Collaboration et al. 2021), the so-called "repeaters" (e.g., Spitler et al. 2016), while for most FRBs only one burst is observed (e.g., Petroff et al. 2019; CHIME/FRB Collaboration et al. 2021). A few FRBs have been reported with periodicities ranging from sub-seconds to several months (Chime/Frb Collaboration et al. 2020; Rajwade et al. 2020; The CHIME/FRB Collaboration et al. 2021). It is so far unclear whether repeaters and non-repeaters originate from astrophysically different populations, and whether there are multiple channels for FRB formation (Platts et al. 2019; Hashimoto et al. 2020).

About two dozen FRBs have been associated with a particular host galaxy (Chatterjee et al. 2017; Bhandari et al. 2020; Heintz et al. 2020; Bhandari et al. 2022), and a few repeaters have been pinpointed inside the host galaxy through Very-Long Baseline Interferometry (VLBI) (Marcote et al. 2017, 2020; Kirsten et al. 2022). Two repeating FRBs are associated with persistent radio sources (Marcote et al. 2017; Niu et al. 2021a), which show complicated polarization properties (Michilli et al. 2018; Anna-Thomas et al. 2022; Dai et al. 2022).

Only a small number of FRB's have been associated with host galaxies. Consequently, the distance to these sources has so far not been available, and therefore the luminosity function of FRBs is currently poorly constrained (Luo et al. 2018).

Although almost 15 years have passed since their discovery (Lorimer et al. 2007), there is no consensus about their origin despite an increasing number of recent detections (Platts et al. 2019; CHIME/FRB Collaboration et al. 2021). Beyond the central questions of the source composition and emission mechanism, it has been suggested that FRBs could be used to address key issues in cosmology and physics, including dark energy (e.g., Hashimoto et al. 2019; Liu et al. 2019), dark matter (e.g., Muñoz et al. 2016; Leung et al. 2022), testing of the general relativity (e.g., Wei et al. 2015; Hashimoto et al. 2021), and the missing baryon problem (e.g., Muñoz & Loeb 2018; Macquart et al. 2020).

There have been three major challenges in revealing the physical origins of FRBs. The first challenge is low detection probability. Existing telescopes have a limited field of view (FoV) but FRBs randomly appear in the sky, so the great majority of detectable FRBs are currently missed. The second challenge is their poor localization to their host, due to the insufficient spatial resolution of existing radio telescopes. The final challenge is that the high antenna gain of current telescopes means that the FRBs that are currently detected are often too distant to conduct multi-wavelength/multi-messenger observations to identify their progenitors, host galaxies, and simultaneous emission counterparts. All of these issues arise because existing radio telescopes are not tailored for FRBs.

The proposed Bustling Universe Radio Survey Telescope in Taiwan (BURSTT) will address these problems by detecting high-fluence (bright) FRBs in the nearby universe. However, these are rare because of the relatively small cosmic volume in the nearby universe, so we have maximized the FoV to increase the detection rate. BURSTT is a unique fisheye radio software telescope, with which we will observe at least 25 times more sky than any existing radio observatory, except the Survey for Transient Astronomical Radio Emission 2 (STARE2) (Bochenek et al. 2020a), which has a much lower sensitivity than BURSTT. By viewing so much more sky than other telescopes, and hence detecting a larger sample of bright and nearby sources, BURSTT will be unique in identifying multifrequency (i.e., optical and X-ray/gamma-ray) and multimessenger (gravitational-wave and neutrino) counterparts. In addition, BURSTT-2048 will add substantially to the the overall world-wide FRB detection rate.

In this paper, the scientific objectives of BURSTT are shown in Section 2. The potential technical aspects are illustrated in Section 3. Finally, the work is summarized in Section 4.

2. Science Objectives

2.1. The Extremely Wide Field of View of BURSTT

To achieve a complete census of FRBs, monitoring observations of the nearby universe with a very wide FoV are required (Bochenek et al. 2020a; Connor et al. 2021).

For any array of antennas, the FoV is inversely proportional to the effective collecting area of the array, so the wide field of BURSTT means that it will have a smaller collecting area than many previous telescopes. In estimating the detection rate (see below), the area and FOV factors partially compensate, and the resulting detection rate is not strongly dependent on the choice of FoV. However, for the wide-field (low collecting area) telescope, the detected sources are on average closer.

The redshift of the CHIME/FRB sample is in the range of 0.3–0.5 (Rafiei-Ravandi et al. 2021). The CHIME/FRB team is currently adding outriggers to CHIME, which will localize many FRB events in this redshift range (Cassanelli et al. 2022; Mena-Parra et al. 2022). For many of these, optical follow-up observations will be needed to get the redshift of the host

galaxy. In contrast, we anticipate that the median redshift of the BURSTT FRB sample will be in the range of 0.03–0.05. This means that many of the BURSTT host galaxies will be present in existing optical redshift survey catalogs (Cutri et al. 2003; Karachentsev et al. 2013; Bilicki et al. 2014, 2016; Dey et al. 2019), which will allow an immediate redshift determination to be made without the need for follow-up observations. This will allow for a rapid improvement in the measured FRB luminosity function (Shin et al. 2022).

To estimate the system equivalent flux density (SEFD) sensitivity of the BURSTT, we assume a conservative system temperature (Tsys) of $\sim \! 150$ K, which is achievable with commercially available LNAs, and an effective area for 256 antennas of $\sim \! 100$ m² at 600 MHz. The corresponding SEFD is $\sim \! 5000$ Jy. For an FRB with a duration of 1 ms across 400 MHz bandwidth with a fluence of 100 Jy ms, the corresponding signal-to-noise ratio (S/N) is $\sim \! 12$. In the future, we hope to achieve Tsys of 50 K, which will improve the sensitivity.

To estimate the event rate, we assume a Euclidean distribution to scale from a previously measured reference rate,

$$N(S) = N_0 \left(\frac{S}{S_0}\right)^{-1.5},\tag{1}$$

where N(S) is the event rate above fluence S and N_0 is the reference rate above reference fluence S_0 . We use the reference event rate reported by CHIME (CHIME/FRB Collaboration et al. 2021), 300 FRBs per day per sky, with $S_0 = 5$ Jy ms at 600 MHz, and two additional constraints: scattering time at 600 MHz below 10 ms and DM above 100 pc cm⁻³. The FoV of BURSTT is $\sim 10^4$ deg² or 0.24 sky. Thus, we estimate that BURSTT-256 can detect ~ 100 bright FRBs per year with the threshold of fluence higher than 100 Jy ms at 600 MHz. Note that the CHIME FRB rate is at the low end of published rate estimates, which makes our estimate conservative.

The properties of repeating FRBs are poorly constrained (e.g., Fonseca et al. 2020) because only a small number have been detected. Uncertain parameters include the event rate and its dependence on flux density, the possibility of non-Poisson distribution of events (Oppermann et al. 2018), and the host environment of these FRBs. With its all-sky coverage, BURSTT will likely uncover a larger fraction of repeaters, sharply localize the sources, and provide detailed data for each received pulse. This should substantially improve our understanding of repeating FRBs.

Figure 1 shows the FoV versus the SEFD and the effective area for the existing, planned, and future-concept FRB surveys, including CHIME (CHIME/FRB Collaboration et al. 2018), Australian Square Kilometre Array Pathfinder (ASKAP) (Macquart et al. 2010), the Canadian Hydrogen Observatory and Radio-transient Detector (CHORD) (Vanderlinde et al. 2019), STARE2 (Bochenek et al. 2020a), Galactic Radio

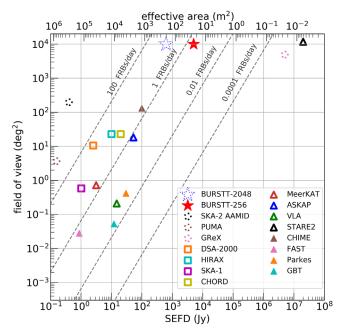


Figure 1. Comparison of BURSTT's FoV, effective collecting area, sensitivity (SEFD), and FRB detection rate (dashed lines) vs. existing (solid), planned (outline), and future-concept (dotted circle) observatories. Rates (dashed lines) were calibrated to CHIME, assuming Euclidean rates and 400 MHz bandwidth. Open triangles are sparse interferometers, which provide arcsecond localization and would require correlator upgrades to achieve these rates. The rate is a hypothetically upper limit, with the assumption of 24/7 FRB searches (which only CHIME/FRB does), as well as the optimal FRB searches with coherently beamforming for the interferometry (ASKAP, VLA). BURSTT is unique in the large FoV with enough sensitivity to detect a large sample of bright and nearby FRBs.

Explorer (GReX) (Connor et al. 2021), Deep Synoptic Array 2000 (DSA-2000), ¹⁹ Five-hundred-meter Aperture Spherical Telescope (FAST) (Jiang et al. 2019), Green Bank Telescope (GBT) (Connor et al. 2016), the Hydrogen Intensity and Real-time Analysis eXperiment (HIRAX) (Newburgh et al. 2016), MeerKAT (Rajwade et al. 2021), Parkes (Petroff et al. 2014), Packed Ultra-wideband Mapping Array (PUMA) (Slosar et al. 2019), Square Kilometre Array Phase 1 (SKA-1) (Braun et al. 2019), SKA-2 (Torchinsky et al. 2016), and the Very Large Array (VLA) (Law et al. 2018).

2.2. Immediate Localization by BURSTT

The critical step to reveal the origin of FRBs is to measure the accurate positions of FRBs (i.e., localization) to identify their progenitors. Out of more than 600 FRBs published to date (Petroff et al. 2016; CHIME/FRB Collaboration et al. 2021), there has been only one successful case of localization to a stellar progenitor, which identified the FRB as a Galactic magnetar, SGR 1935+2154 (Bochenek et al. 2020b; CHIME/FRB

¹⁹ https://www.deepsynoptic.org/overview

Collaboration et al. 2020; Kirsten et al. 2021), and clearly indicates the magnetar as the progenitor of this particular FRB.

Recently, three repeating FRB sources have been localized in nearby star-forming regions of galaxies (Bassa et al. 2017; Marcote et al. 2020; Piro et al. 2021). This may indicate young stellar populations as a possible origin of repeating FRBs. In contrast, another repeating FRB source was localized in the globular cluster in M81, a nearby galaxy, which suggests that FRB origins are to be found in old stellar populations (Kirsten et al. 2022). This apparent contradiction could be due to the small sample size. A statistical relation between localized FRBs and their host galaxies has been studied (Li & Zhang 2020), while more localized samples are necessary to understand the host environments. BURSTT will detect and localize ~ 100 bright FRB per year to further distinguish their local environment, and such statistics are important to understand the population (Leung et al. 2021; Cassanelli et al. 2022; Mena-Parra et al. 2022).

Very-Long Baseline Interferometry (VLBI) along with outrigger stations have been proposed to localize both non-repeating and repeating FRBs (Cassanelli et al. 2022). For the latter, the European VLBI Network (EVN) has localized several repeating FRBs to their hosts (Marcote et al. 2020; Kirsten et al. 2022). Along with the VLBI outrigger stations, BURSTT can localize FRBs to their hosts. For instance, an outrigger located a few hundred km away from the main array will provide $\sim 1''$ localization for the host galaxy. This localization will lead to the identification of FRB hosts and potentially the progenitors of the nearby sources, which will allow us to understand the nature of FRBs.

2.3. Long-term High Cadence Monitoring with BURSTT

The observational classification of FRBs as repeating and non-repeating has caused a major problem over the past decade: non-repeating FRBs can be contaminated by repeating FRBs because this classification is a purely observational definition (Ai et al. 2021; Chen et al. 2022). Due to limited observational time, a significant fraction of repeating bursts may be missed and thus such FRBs could be misclassified as non-repeating FRBs. However, long-term monitoring observations are extremely expensive and difficult using current and planned radio telescopes.

BURSTT will monitor a fixed large patch of the sky all the time (See Table 1). This is essential for non-stop monitoring of repeating and non-repeating FRBs. Repeating FRBs could originate from the repeating activities of progenitors, such as pulsars and magnetars, while non-repeating FRBs may be generated by catastrophic one-off events, such as compact merger systems (Platts et al. 2019). BURSTT will thus provide great statistics of source repetitions, as well as precise constraints on apparent non-repetition of one-off events.

Table 1
The Main Properties of the BURSTT

Quantity	Value		
Project	BURSTT-256	BURSTT-2048	
SEFD	\sim 5000 Jy	∼600 Jy	
Effective area	40-200 m ²	320-1600 m ²	
Number of antennas (main	256	2048	
station)			
(outrigger stations)	64		
Polarization	single		
E-W FoV	${\sim}100^{\circ}$		
N-S FoV	${\sim}100^{\circ}$		
Daily exposure time	24 hr (North pole)		
	\sim 10 hr (45°)		
	\sim 7 hr (Equator)		
Frequency range	300-800 MHz	TBD	
Bandwidth	400 MHz	$\geqslant 400 \text{ MHz}$	
Number of frequency	1024	TBD	
channels			
E-W baseline	~8000 km (Northern Taiwan to Hawaii)		
N-S baseline	~300 km (Northern to Southern Taiwan)		

BURSTT can monitor the northern hemisphere at least 7 hr per day (24 hr for North pole and ~7 hr for the equator). BURSTT's longer monitoring minimizes the chance it will miss any repeating FRBs, which will resolve the missing repeating FRB problem. Ravi (2019a) compared the event rate of nearby FRB samples and those of cataclysmic progenitor events. The event rate of nearby FRBs exceeds the cataclysmic progenitor events, which indicates that the FRB sample cannot be explained by cataclysmic events alone. Hence, Ravi (2019a) concluded that most observed cases of FRBs must originate from sources that emit several bursts over their lifetimes.

The CHIME/FRB Collaboration et al. (2021) reported 474 non-repeating FRBs and 62 repeating bursts from 18 repeaters, which implies the probability that a burst originates at a repeater is equal to 62/(62+474) = 11.5%. As exposure time increases (Ai et al. 2021), we expect that more repeating and bright bursts will be detected. For instance, Herrmann (2021) reported a high-fluence FRB (334 Jy ms) from the FRB 20201124A repeating source with observation times of 90 hr. As a result of the small number of detections, the properties of repeating FRBs are poorly constrained (e.g., Fonseca et al. 2020). This uncertainty includes the event rate and its dependence on flux density, the possibility of non-Poisson distribution of events (Oppermann et al. 2018), and the host environment of the FRB. With its all-sky, continuous sensitivity to high-fluence events, BURSTT will explore a unique parameter space for repeaters. Since BURSTT will detect ~100 high-fluence FRBs from the nearby universe, the non-detections of repeats will shed some light on the underlying population using statistical methods.

Finding repeaters requires either high sensitivity or long exposure time. To detect more repeaters, we take the advantage

of BURSTT's long exposure time (given by the large FoV) as a compensation to the sensitivity issue. We expect that BURSTT will detect high-fluence bursts from repeaters, and other telescopes with high sensitivity can do follow-up observations to detect more fainter bursts from the sources. For instance, the bright repeaters can then be readily followed-up by larger apertures, including the Five-hundred-meter Aperture Spherical Telescope (FAST), Giant Meterwave Radio Telescope (GMRT), and the Very Large Array (VLA). Historically, repetition rates increase rapidly with sensitivity, and thus we expect repeaters detected by the BUTSTT to have the highest rates in follow-up campaigns.

With the high cadence, BURSTT may detect bright pulses from pulsars and rotating radio transients (Good et al. 2021), and further study their connections to FRBs (Bij et al. 2021; Majid et al. 2021; Thulasiram & Lin 2021).

2.4. An FRB Telescope Dedicated to the Nearby Universe

Our plan to reveal the origin of FRBs is to explore the nearby universe because this maximizes the chance of detecting possible multi-wavelength and multi-messenger counterparts of FRBs, including FRB progenitors. For example, follow-up was key to the association of a repeating FRB with the Galactic magnetar SGR 1935+2154 (Bochenek et al. 2020b; CHIME/FRB Collaboration et al. 2020; Kirsten et al. 2021). This progenitor was identified because it is located nearby. Simultaneous X-ray emissions from SGR 1935+2154 were also detected with X-ray telescopes (e.g., Tavani et al. 2021). This observed association gives hope that FRBs in nearby galaxies may also have detectable X-ray counterparts. (Li et al. 2021). BURSTT will monitor a large fraction of nearby galaxies.

An important wavelength range to examine for counterparts is the optical. No optical-FRB coincidence has been found (MAGIC Collaboration et al. 2018; Tominaga et al. 2018) to date. In addition, improved constraints would be very valuable because the expected flux density of the optical counterpart strongly depends on the physical mechanism of the FRB; for example, the magnetosphere model predicts optical counterparts, whereas the maser model predicts a negligible flux density in optical (e.g., Yang et al. 2019; Yalinewich & Pen 2022).

BURSTT is the ideal compliment to the current generation of multi-messenger detectors, such as the Laser Interferometer Gravitational-Wave Observatory (LIGO) (Abbott et al. 2009), IceCube (Aartsen et al. 2017), and the Rubin Observatory's Legacy Survey of Space and Time (LSST) (LSST Science Collaboration et al. 2009), which (like BURSTT) monitor the nearby universe with near full-sky coverage. It will be instructive to search for FRBs associated with gravitational waves detected with the LIGO-Virgo-KAGRA Gravitational Wave Detector Network in near future (The LIGO Scientific Collaboration et al. 2022). Some FRB models predict an association with

Gravitational-Wave (Wei et al. 2018) or neutrino counterparts (Metzger et al. 2020). If FRBs originate from neutron star mergers, which are known to generate gravitational waves (Abbott et al. 2017) and potentially neutrinos (Fang & Metzger 2017; Kimura et al. 2018), then the expected time window of radio emission is on the millisecond timescale comparable to that of FRBs (e.g., Yamasaki et al. 2018). So far, no FRB has been found in association with gravitational-wave sources (e.g., Abbott et al. 2016; The LIGO Scientific Collaboration et al. 2022). If no significant GW-FRB associations can be detected in future despite the greatly improved detection rate of BURSTT, then the neutron star merger scenario for the origin of FRBs can be strongly constrained. Scenarios involving cosmic-ray acceleration (e.g., Li et al. 2014; Metzger et al. 2020) can produce neutrinos that can be detected by Large Volume Neutrino detectors, such as IceCube, KM3NeT (Adrián-Martínez et al. 2016), and Baikal-GVD (Avrorin et al. 2022). Although no neutrinos from FRBs have been detected so far, the probability of finding neutrino associations significantly improves with the total number of FRB detections (Aartsen et al. 2018, 2020). These observations would bring clarity to the role of FRBs within the non-thermal universe and constrain the still unknown acceleration mechanisms of the highest energy cosmic rays.

The diffuse Galactic foreground is known to be pervasive, bright, and dynamic, thus making it challenging to be observed properly. Only a few surveys have a large enough coverage to observe the low-frequency radio sky below 1 GHz (e.g., Haslam et al. 1981, 1982; Guzmán et al. 2011), including MSSS (Heald et al. 2015), GLEAM (Wayth et al. 2015), and TGSS (Intema et al. 2017). It is therefore complementary to generate low-frequency high resolution sky maps using (for example) the Long Wavelength Array (LWA) (e.g., Eastwood et al. 2018) and the Engineering Development Array 2 (EDA2) (e.g., Kriele et al. 2022). With the large FoV of BURSTT below 1 GHz, it has the potential to improve the removal of the Galactic foreground contamination in many Cosmic Microwave Background radiation and Epoch of Reionization works (see e.g., Ichiki 2014; Spinelli et al. 2021).

3. BURSTT Design

In this section, we describe the initial designs of the BURSTT system with 256 antennas, which serves as a test-bed for BURSTT-2048. Our technology goal for BURSTT-256 is to optimize the system, and compare component and algorithm designs.

3.1. Overview of the BURSTT Instrument and the Telescope Site

BURSTT-256 will consist of a combination one 256 antenna main station, and smaller outriggers at other sites in Taiwan and Hawaii.

We have surveyed several sites, and have permission to deploy the main station at the Fushan Botanical gardens²⁰ in Northern Taiwan. Our survey indicates acceptable radio frequency interference (RFI) conditions at that site. We are continuing to survey other sites, sheltered from RFI, for the outrigger stations in Taiwan and surrounding islands, as well as in Hawaii. The outriggers will each employ 64 antennas, allowing VLBI source localization (Leung et al. 2021; Cassanelli et al. 2022; Mena-Parra et al. 2022).

We will process an effective bandwidth of 400 MHz selected within the 300–800 MHz analog range through direct digital polyphase filterbanks (PFB). The CHIME/FRB team has demonstrated that processing a 400 MHz bandwidth is practical (CHIME/FRB Collaboration et al. 2018). we would like to explore the FRB rate at frequencies below 400 MHz, which is at the bottom of the CHIME band. Our design offers frequency agility. Should a shift in band improve the rate, then BURSTT will be able to adapt through a simple software change.

Due to its modular nature, BURSTT is flexible enough to be expanded by deploying more antennas or outrigger stations. We plan to expand the main station from 256 antennas (BURSTT-256) to 2048 antennas (BURSTT-2048), and more outrigger stations could be deployed with longer baselines, which will increase the localization precision. The main properties of the BURSTT are summarized in Table 1.

For the BURSTT-256 project, the main station will search for bursts with fluence higher than 100 Jy ms, which is mentioned in Section 2.1. Assuming the duration of 1 ms and the bandwidth of 400 MHz, the corresponding S/N at the main station is 12. Once the main station detects a candidate burst, the data in the ring buffer in the outrigger will be copied out for the offline analysis. We expect a burst with a fluence of 100 Jy ms would yield an S/N at least of six in the cross-correlation analysis for the localization purpose.

We plan to use single-polarization antennas in BURSTT. These are simpler to design and build than dual-polarization antennas, and optimization of the antenna-LNA match is more effective when only one polarization is involved. We plan to start with all antennas in one polarization direction within the main-station array because this enhances sensitivity for linearly polarized sources, such as FRBs. Later, we have the option to rotate some antennas to the orthogonal polarization if we determine that dual-polarization data is essential to test FRB source models.

Even in the initial configuration, we plan to provide some polarized properties of all detected FRBs, such as rotation measure and polarization position angle swings (Masui et al. 2015), using two sub-stations with orthogonal polarization at one of the outrigger sites.

The back-end includes correlator hardware that will process the radio signals and accompanying software including fast Fourier transform (FFT) techniques, and establish the capability to locate FRB events. Much of the development effort of the BURSST program lies in the development and testing of real-time software. BURSTT will develop an upgraded version of the real-time processing digitizer and FRB search engines employed for CHIME (The CHIME Collaboration et al. 2022).

An overview of the front- and back-ends of the main station is given in Figure 2. The front-end receives the signal from the sky through the bandpass (BP) filter and the noise amplifiers. The back-end digitizes the signals, performs the frequency-domain channelization, forms beams on the sky, and searches for FRBs. The hardware at outrigger stations will be similar to the main station, except that no beamforming or FRB search will be performed. Rather the outriggers just record the arriving signal streams. This is discussed further in Sections 3.2 and 3.3.

3.2. Front-End Design and Development

BURSTT will use Log-Periodic Dipole Array (LPDA) antennas. The LPDA antennas, which provide the required gain (7-9 dBi) with a simple structure, have been characterized and measured for use in numerous scientific and industrial applications. Millions of LPDAs have been deployed over the decades for use as home television antennas. Because LPDA antennas have a structure that can withstand strong winds, they are suitable for the environment in Taiwan, which experiences typhoons. Figure 3 shows some photographs of various prototype antennas that are currently in development. The first, which is shown in the left-hand panel, is a standard design that has been characterized well but it is expensive to manufacture. In addition, the long elements for low frequencies would be weak against strong winds. The second, which is shown in the middle panel, is the planar LPDA, which is a cost-effective design that simplifies the manufacturing process. However, because the elements have a long band-like structure, they can easily vibrate in the wind. The third design, which is shown in the right-hand panel, is the planar LPDA attached to a quadrangular pyramid structure that is composed of insulator plates, which is costeffective and raises the structure resonance frequency of the elements so that they can withstand strong winds. However, the structure can be heated by the Sun, which can degrade the amplifier's noise temperature. We will continue development and will test several prototypes in the field.

Received signals pass through the front-end electronics (FEE) before arriving at the digitizer. We will use a sampling rate of 1600 MHz, so our passband lies in the first Nyquist zone. This means that no mixers or local oscillators are needed in the RF system. To minimize noise temperature due to transmission line losses, the first-stage low-noise amplifier (LNA) is closely integrated into the feed point of the LDPA (as shown in Figure 4). The external second-stage FEE module consists of two equalizers cascaded with second-stage amplifiers (LNAs) and filter modules to eliminate RFI at low

²⁰ https://fushan.tfri.gov.tw/en/index.php

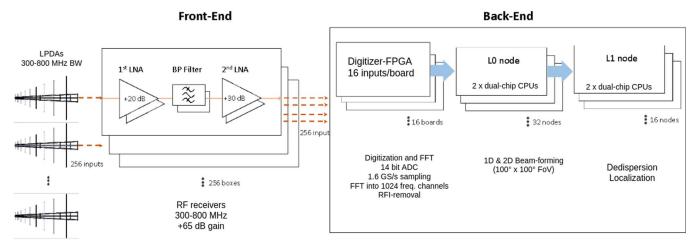


Figure 2. The system diagram of the BURSTT main station.

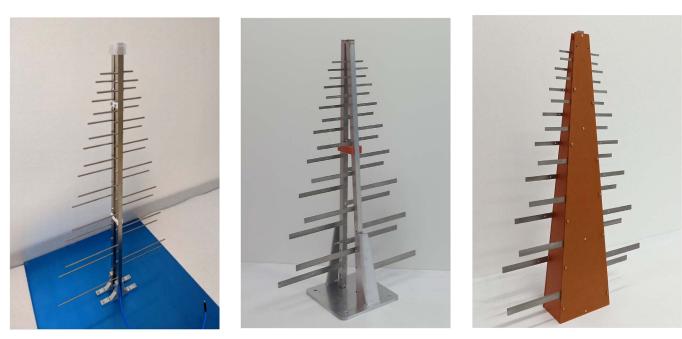


Figure 3. Photographs of the prototype LPDA antennas in development. Standard LPDA (left-hand panel), planar LPDA (middle panel), and planar LPDA on a quadrangular pyramid structure (right-hand panel).

frequencies, and to prevent aliased detection of signals and noise above 800 MHz. Figure 5 shows the measured forward gain and the noise temperature (20–30 K) of a first-stage LNA that we are developing. Comparing this to typical sky brightness temperature values, which fall with frequency from $\sim\!50$ K at 300 MHz to $\sim\!6$ K at 800 MHz, we see that the further effort to reduce LNA noise temperature can improve sensitivity, particularly at the high end of our frequency range. The required system gain is 65 dB, which is obtained by combining the first-stage LNA gain (25–30 dB), the second-stage amplifier assembly (35–40 dB), and slope-equalizer

losses (-5 dB). We will continue to optimize the LNA design and plan to install an in-house developed low-noise amplifier in the 2048 antenna array.

The antenna array of BURSTT-256 station has a footprint of 32×32 m, and the distance to the data acquisition (DAQ) hut where the digitizer is located is up to 50 m, which requires cable lengths of ${\sim}60$ m (Figure 6). When using the most popular and reliable coaxial cable, LMR-400, the transfer losses are -4 dB and -48 dB for 200 MHz and 900 MHz, respectively, which can be corrected in the RF slope equalizer or in the correlator after channelization.

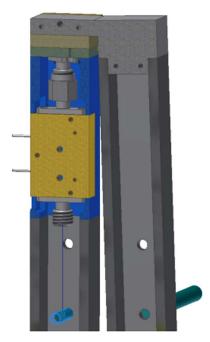


Figure 4. First-stage LNA (yellow) to be integrated into the LPDA: The feed point of the LDPA is directly connected to the input port of the first-stage LNA, the amplified signal is then transmitted to the second-stage FEE module by coaxial cable inside the support rod of the LPDA.

3.2.1. The Beam Pattern

Figure 7 shows simulated far field radiation patterns of the LPDA antenna calculated using the Ansys high-frequency structure simulator (HFSS) software. With the HFSS software, we utilize the finite element method to compute the EM field distribution of components. The 3 dB beam widths (FWHM) of the prototype LPDA antenna are around 80° on the H plane and 60° on the E plane from 300 to 800 MHz (Table 2).

3.3. Back-End Design and Development

When the BURSTT main station detects a candidate FRB, it sends a SAVE message to the outriggers, which then transfer recently buffered data from RAM to disk. The outriggers and main station transmit the data to a VLBI processing facility for localization processing. This process synthesizes an instrument with a nominal VLBI resolution $\lambda/(2D)$ of a 7800 km \times 270 km baselines with 13 mas and 370 mas resolution in E-W and N-S, respectively. This allows not only unique identification of each FRB's host galaxy but also ~ 100 parsec-scale localization of the FRB within that galaxy.

3.3.1. Digitizer and Engines

We employ the Xilinx ZCU216 field programmable gate array (FPGA) boards as the platform of our digitizer and channelizer. The ZCU216 is equipped with a radio frequency

system on chip (RFSoC) FPGA that can process 16 RF inputs. RFSoC systems include an embedded analog to digital converter (ADC) inside the FPGA chip, which simplifies wiring and saves physical space. This is the highest bandwidth RFSoC that is currently available. Figure 8 summarizes the functional block diagram of the RFSoC F-engine. We will be clocking the ADC at sampling rate 1600MHz for a Nyquist bandwidth of 800 MHz. A polyphase filter bank will channelize the data, forming 2048 bands of width ∼390 kHz, and we will select 1024 of these for further analysis, choosing RFI-clear bands. We then re-quantized to resolution 4+4i bits to reduce the data rate. Gathering all of the 16 input frequencydomain data streams, a 100G Ethernet system will transmit the 100G UDP packets to the server for further data processing. The ADCs have 14 bit resolution, which is substantially higher than the 8 bits sampling that has been commonly used to date.

3.3.2. Preprocessing and RFI Mitigation

We discuss the steps of preprocessing in this section, including the RFI-removal techniques, beamforming, and frequency channelization.

For BURSTT-256, the antennas will be organized into 16 columns, each with 16 antennas. The 16 antennas in a N-S column will be connected to one FPGA F-engine. It is thus efficient to perform 1D beamforming within the FPGA from Ng et al. (2017). Each beam from the 16 F-engines is then similarly combined in the CPU array to form the E-W direction beams, This process occurs independently for each frequency channel. This approach is similar in spirit to the 2D-FFT telescope (Tegmark & Zaldarriaga 2009) although with an implementation that is tailored to our F-engine design. For BURSTT-2048, a second layer of 2D spatial FFT will be carried out in the CPU array to synthesize all full-resolution beams.

The bandpass calibration and the system SEFD will be derived from observations of the Sun during the day. A pulsed broadband noise source will be used to monitor the gain variation during the night. Phase calibration will be achieved by using the many radio sources that are continuously available within the large FoV.

RFI Surveys in the Fushan candidate site shows that about 60% of the 300–800 MHz frequency range is free of noticeable RFI, while an additional 15% is affected by weak RFI. The FPGA F-engine offers a possibility of mitigating the weak RFI contamination using the using a spatial filtering technique applied to time streams from the antennas. (Kocz et al. 2010). The basic idea is that for a spatially fixed RFI source, the cross-correlation between any two antennas will record a distinct phase relation. The covariance matrix of all the antennas contains all of the phase information that one can measure with the array. Any source, not just the RFI sources, will contribute to the covariance matrix. However, for signals that are above

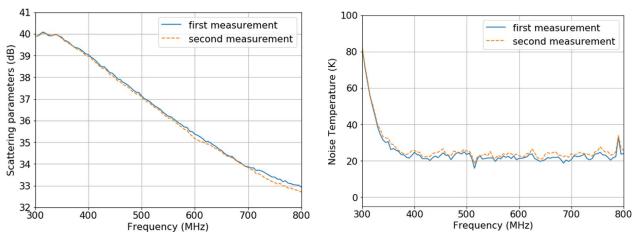


Figure 5. Measured performance of the two samples of the first-stage LNA, (left-hand panel) measured forward gain with the output port connected with an equalizer, an additional equalizer will be installed in the final system to provide +-0.5 dB gain flatness over 350–800 MHz, (right-hand panel) measured noise temperature.

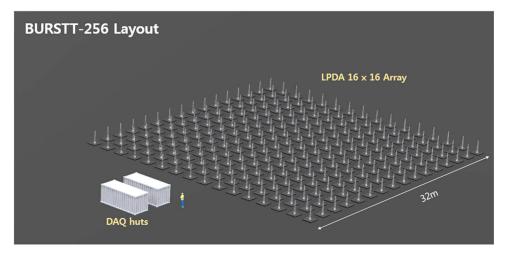


Figure 6. BURSTT 256 antenna array station layout.

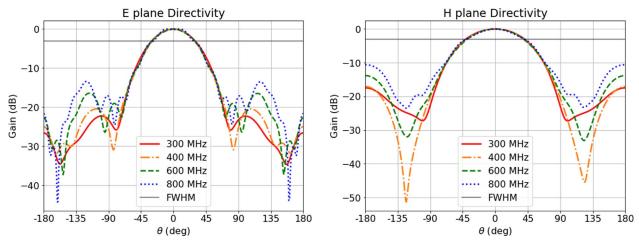


Figure 7. Normalized directivity of the prototype antenna simulated using HFSS, (left-hand panel) the E plane directivity and (right-hand panel) H plane directivity.

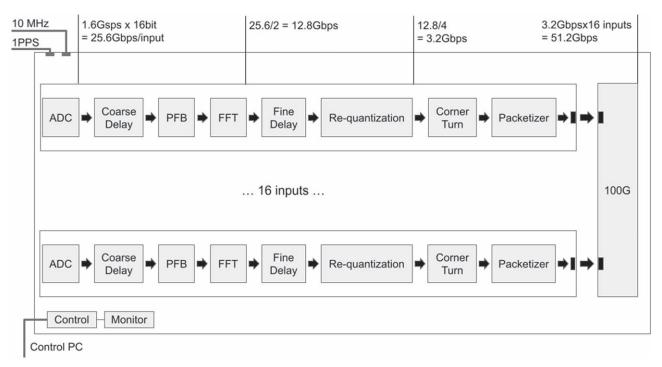


Figure 8. A functional block diagram of the BURSTT F-engine. Each FPGA-based RFSoC board will process 16 inputs with an initial bandwidth of 0–800 MHz. After channelization, a choice of 400 MHz bandwidth is retained. The re-quantization block further reduces the data rate by a factor of 4. Each board will send out a total of 51.2 Gbps of data over the 100G Ethernet interface to a GPU processing node.

Table 2
The FWHM of the Beam Pattern

	300 MHz	400 MHz	600 MHz	800 MHz
E plane	58°0	56°0	58°0	54°.0
H plane	80°.0	78°.0	77°.0	80°.0

the thermal noise, one can identify them by solving the eigenvalues and eigenvectors of the covariance matrix. This procedure is performed independently on each spectral channel. The antennas' voltage data can be projected onto the eigenvectors to form the eigenmodes. The eigenmodes with the strongest variances (i.e., the largest eigenvalues) are assumed to be from the RFI source. It is therefore possible to remove the RFI contribution by nulling or zeroing out the RFI eigenmodes and deprojecting them back to the antenna data space. The antenna data after the RFI removal are passed onto the beamforming stage.

Figure 9 gives an example of this technique applied to the TV station signal that was recorded in Fushan Botanical Garden (the main-station site) with a four antenna test system. Three TV stations, each with a 6 MHz bandwidth, are within the recorded 40 MHz bandwidth. We note here that the eigenmode decomposition works better when the variance of the antennas are comparable. For this example, the four

antennas have been normalized to the mean of the four antennas at each spectral channel. It is shown that for the TV signal centered on 581 MHz, the signal is reduced by about 20 dB after nulling out the strongest eigenmode. The TV signal centered on 569 MHz is only reduced by a similar amount after nulling two of the strongest eigenmodes. At Fushan, we may be receiving the TV signals from multiple towers or multi-path propagation may be an important factor. As the project expands the testing to more sites and with more antennas, the RFI-removal technique will be refined.

The performance of the RFI mitigation improves as the number of antennas is increased. Using the entire array in principle gives the best RFI nulling, but at substantial increase in the cost of the computation. To keep this system simple, we plan to update the RFI nulling solutions just once several hours or once a day. For the 16 antennas connected to one F-engine, applying the nulling solution is accomplished simply by multiplying the 16 antenna time streams for each frequency channel by a 16×16 complex matrix, similar to the beamforming operation. In fact, the two operations can be combined. Therefore, the RFI mitigation can be applied for each column of antennas at no additional real-time computational cost. An additional test will determine if 16 antenna RFI mitigation is sufficient or if more antennas should be combined to achieve the desired improvement.

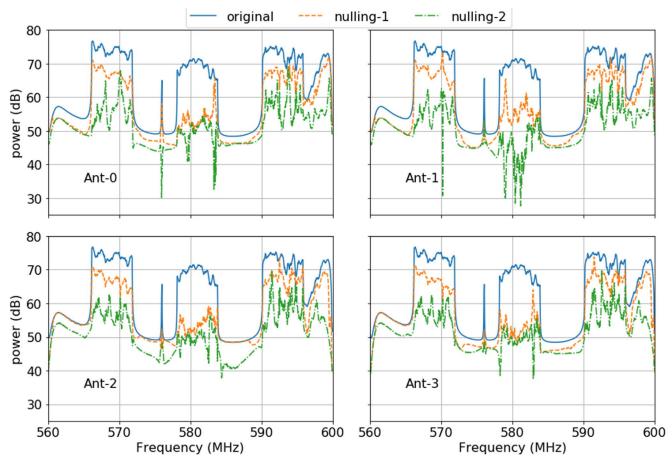


Figure 9. The RFI-removal testing at the Fushan site. The four panels correspond to the four antennas. The blue-solid line shows the originally received power as a function of observing frequency. The orange-dotted line shows the received power after nulling of the strongest eigenmode. The green-dashed line corresponds to the received power after nulling of strongest two eigenmodes.

3.3.3. Dedispersion, Search, and Follow-up

Searching for FRBs with unknown dispersion measure (DM) presents a challenge. This search has high computational cost and for BURSTT the search must be carried out real-time (Petroff et al. 2019) because the raw data rate is too high for disk storage. Search algorithms such as the Fast Dispersion Measure Transform (FDMT) (Zackay & Ofek 2017) or the tree-dedispersion algorithm (Taylor 1974; Masui et al. 2015; CHIME/FRB Collaboration et al. 2018) convert the intensity versus frequency and time into intensity versus DM and time data, along with the S/N estimates. Peaks in this data with S/N higher than 10 are initial FRB candidates. These dedipsersion algorithms are well tested: the tree-dedispersion algorithm in real-time has been used by CHIME/FRB (CHIME/FRB Collaboration et al. 2018), the upgraded Molonglo Observatory Synthesis Telescope (UTMOST) (Farah et al. 2019), the Commensal Radio Astronomy FAST Survey (CRAFTS) (Niu et al. 2021b), and Deep Synoptic Array 10 (DSA-10)

(Ravi et al. 2019), and the FDMT algorithm in real-time has been implemented in ASKAP (Bannister et al. 2017).

Since the DM for the majority of FRBs is less than 1000 pc $\rm cm^{-3}$ (CHIME/Pulsar Collaboration et al. 2021) and the BURSTT will be searching for nearby FRBs, we will initially set the BURSTT DM range to 0–1000 pc $\rm cm^{-3}$.

In addition to transmitting automated SAVE messages to out outriggers, BURSTT will also send out real-time alert of the FRB events with an initial localization shortly after the detection to the community (Petroff et al. 2017; CHIME/FRB Collaboration et al. 2018) to allow rapid or retrospective multi-frequency and multi-messenger follow-up observations.

4. Discussion and Summary

Over 600 FRBs have been published since the first report in 2007 (Lorimer et al. 2007; Petroff et al. 2019; CHIME/FRB Collaboration et al. 2021), while the complex phenomena cannot be totally explained by theories (Pen 2018; Platts et al. 2019). The presence (or lack) of counterparts in other bands

will provide critical constraints on both the emission mechanism and progenitor systems of FRBs (Popov & Pshirkov 2016; Cunningham et al. 2019). Candidate models (e.g., compact object mergers and young magnetar flares) yield very different predictions for multi-wavelength counterparts (Popov & Postnov 2013; Totani 2013; Pen 2018), yet progress has been slow due to a lack of precise and rapid FRB localizations to date.

BURSTT with a large FoV of $\sim 10^4~\rm deg^2$ will monitor the whole visible sky to detect and localise $\sim 100~\rm bright$ and nearby FRBs per year. First, the large FoV yields a high cadence monitoring of FRBs, which is crucial to statistically understand the repeater and the apparent one-off events. Second, the detection and localization of nearby FRBs offers the best chance for counterpart identification, which allows for timely follow-up observations at other wavelengths, including X-ray, infrared, and so on. A large sample of $\sim 100~\rm bright$ FRBs per year will bring unprecedented new evidence about the nature of FRBs.

Understanding the physical processes through nearby and bright FRBs opens the possibility of further cosmological applications. Moreover, determining FRB host galaxies, and hence distances, has the potential to enable large scale structure studies and unique determinations of the content of the intergalactic medium (IGM). This includes constraints on the location of the "missing baryons" in the universe, whether in the IGM property, or in the halos of galaxies—the circumgalactic medium (e.g., McQuinn 2014; Muñoz & Loeb 2018; Ravi 2019b). Some studies suggest that the structures of galaxy halos can be uniquely constrained by FRBs (Prochaska & Zheng 2019; Ocker et al. 2021), for comparison with models of galaxy evolution. Faraday rotation of polarized FRBs could constrain the origin and evolution of cosmic magnetism (Akahori et al. 2016). Uncertainties in FRBs' dispersion measures have been proposed to constrain Einstein's weak equivalence principle (Hashimoto et al. 2021).

The BURSTT design is scalable. Initially, 256 antennas will be deployed, which will help us to understand and improve the system. Because of its modular design, the telescope can be expanded simply by adding more main-station antennas, outriggers and correlation processors, with a corresponding increase in collecting area, localization precision, and sensitivity. In the BURSTT design, we have eliminated the dishes, cylindrical reflectors, or RF-summed phased-arrays that are used in other radio telescopes. Therefore, the BURSTT design concentrates cost in the real-time signal-processing hardware. Most of this hardware is commercial-off-the-shelf equipment, and is used for a wide range of general computing applications. Computing costs have fallen with Moore's law for decades and they continue to fall. BURSTT is therefore a technology pathfinder in an area that can be expected to grow, both for scientific and financial reason. If computing costs continue to fall, then there is no reason to stop the expansion of BURSTT

at 2048 elements because further increases in collecting area allow an increase the volume and the redshift range of the survey.

Finally, in addition to deciphering the nature of the mysterious FRBs, BURSTT will open up a vast new discovery space for non-FRB surveys due to its unique all-sky collecting area. BURSTT can contribute to the study of self-triggering and reconstruction of cosmic-ray induced extensive air showers (Schröder 2017), as well as other extreme energy phenomena, such as Ultra-High Energy Cosmic Rays (Ackermann et al. 2022), and Ultra-High Energy neutrinos (Coleman et al. 2022). Moreover, BURSTT could also potentially contribute to LIGO counterparts (The LIGO Scientific Collaboration et al. 2022), Extreme Scattering Events (ESEs) (Kerr et al. 2018), 21 cm absorbers (Margalit & Loeb 2016), pulsar searches (CHIME/Pulsar Collaboration et al. 2021), interplanetary scintillation (Hewish et al. 1964), and so on. Consequently, BURSTT will be a promising frontier of radio astronomy in the foreseeable future.

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References

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Aartsen, M. G., Ackermann, M., Adams, J., et al. 2017, JInst, 12, P03012 Aartsen, M. G., Ackermann, M., Adams, J., et al. 2018, ApJ, 857, 117 Aartsen, M. G., Ackermann, M., Adams, J., et al. 2020, ApJ, 890, 111 Abbott, B. P., Abbott, R., Adhikari, R., et al. 2009, RPPh, 72, 076901 Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2016, PhRvD, 93, 122003
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Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017, ApJL, 848, L12
Ackermann, M., Agarwalla, S. K., Alvarez-Muñiz, J., et al. 2022, arXiv:2203.
Adrián-Martínez, S., Ageron, M., Aharonian, F., et al. 2016, JPhG, 43, 084001
Ai, S., Gao, H., & Zhang, B. 2021, ApJL, 906, L5
Akahori, T., Ryu, D., & Gaensler, B. M. 2016, ApJ, 824, 105
Anna-Thomas, R., Connor, L., Burke-Spolaor, S., et al. 2022, arXiv:2202.
Avrorin, A. V., Avrorin, A. D., Ayinutdinov, V. M., et al. 2022, JETP,
  134, 399
Bannister, K. W., Shannon, R. M., Macquart, J. P., et al. 2017, ApJL, 841, L12
Bassa, C. G., Tendulkar, S. P., Adams, E. A. K., et al. 2017, ApJL, 843, L8
Bhandari, S., Heintz, K. E., Aggarwal, K., et al. 2022, AJ, 163, 69
Bhandari, S., Keane, E. F., Barr, E. D., et al. 2018, MNRAS, 475, 1427
Bhandari, S., Sadler, E. M., Prochaska, J. X., et al. 2020, ApJL, 895, L37
Bij, A., Lin, H.-H., Li, D., et al. 2021, ApJ, 920, 38
Bilicki, M., Jarrett, T. H., Peacock, J. A., Cluver, M. E., & Steward, L. 2014,
     pJS, 210, 9
Bilicki, M., Peacock, J. A., Jarrett, T. H., et al. 2016, ApJS, 225, 5
Bochenek, C. D., McKenna, D. L., Belov, K. V., et al. 2020a, PASP, 132,
  034202
Bochenek, C. D., Ravi, V., Belov, K. V., et al. 2020b, Natur, 587, 59
Braun, R., Bonaldi, A., Bourke, T., Keane, E., & Wagg, J. 2019, arXiv:1912.
   12699
Cassanelli, T., Leung, C., Rahman, M., et al. 2022, AJ, 163, 65
Chatterjee, S., Law, C. J., Wharton, R. S., et al. 2017, Natur, 541, 58
Chen, B. H., Hashimoto, T., Goto, T., et al. 2022, MNRAS, 509, 1227
CHIME/FRB Collaboration, Amiri, M., Andersen, B. C., et al. 2020, Natur,
CHIME/FRB Collaboration, Amiri, M., Andersen, B. C., et al. 2021, ApJS,
  257, 59
CHIME/FRB Collaboration, Amiri, M., Bandura, K., et al. 2018, ApJ, 863, 48
CHIME/FRB Collaboration, Andersen, B. C., Bandura, K. M., et al. 2020,
CHIME/Pulsar Collaboration, Amiri, M., Bandura, K. M., et al. 2021, ApJS,
Coleman, A., Eser, J., Mayotte, E., et al. 2022, arXiv:2205.05845
Connor, L., Lin, H.-H., Masui, K., et al. 2016, MNRAS, 460, 1054
Connor, L., Shila, K. A., Kulkarni, S. R., et al. 2021, PASP, 133, 075001
Cunningham, V., Cenko, S. B., Burns, E., et al. 2019, ApJ, 879, 40
Cutri, R. M., Skrutskie, M. F., van Dyk, S., et al. 2003, 2MASS All Sky
  Catalog of point sources (Pasadena, CA: IPAC)
Dai, S., Feng, Y., Yang, Y. P., et al. 2022, arXiv:2203.08151
Dey, A., Schlegel, D. J., Lang, D., et al. 2019, AJ, 157, 168
Eastwood, M. W., Anderson, M. M., Monroe, R. M., et al. 2018, AJ, 156, 32
Fang, K., & Metzger, B. D. 2017, ApJ, 849, 153
Farah, W., Flynn, C., Bailes, M., et al. 2019, MNRAS, 488, 2989
Fonseca, E., Andersen, B. C., Bhardwaj, M., et al. 2020, ApJL, 891, L6
Good, D. C., Andersen, B. C., Chawla, P., et al. 2021, ApJ, 922, 43
Guzmán, A. E., May, J., Alvarez, H., & Maeda, K. 2011, A&A, 525, A138
Hashimoto, T., Goto, T., On, A. Y. L., et al. 2020, MNRAS, 498, 3927
Hashimoto, T., Goto, T., Santos, D. J. D., et al. 2021, PhRvD, 104, 124026
Hashimoto, T., Goto, T., Wang, T.-W., et al. 2019, MNRAS, 488, 1908
Haslam, C. G. T., Klein, U., Salter, C. J., et al. 1981, A&A, 100, 209
Haslam, C. G. T., Salter, C. J., Stoffel, H., & Wilson, W. E. 1982, A&AS, 47, 1
Heald, G. H., Pizzo, R. F., Orrú, E., et al. 2015, A&A, 582, A123
Heintz, K. E., Prochaska, J. X., Simha, S., et al. 2020, ApJ, 903, 152
Herrmann, W. 2021, ATel, 14556, 1
Hewish, A., Scott, P. F., & Wills, D. 1964, Natur, 203, 1214
Ichiki, K. 2014, PTEP, 2014, 06B109
Intema, H. T., Jagannathan, P., Mooley, K. P., & Frail, D. A. 2017, A&A,
  598, A78
Jiang, P., Yue, Y., Gan, H., et al. 2019, SCPMA, 62, 959502
Josephy, A., Chawla, P., Curtin, A. P., et al. 2021, ApJ, 923, 2
Karachentsev, I. D., Makarov, D. I., & Kaisina, E. I. 2013, AJ, 145, 101
Kerr, M., Coles, W. A., Ward, C. A., et al. 2018, MNRAS, 474, 4637
Kimura, S. S., Murase, K., Bartos, I., et al. 2018, PhRvD, 98, 043020
Kirsten, F., Marcote, B., Nimmo, K., et al. 2022, Natur, 602, 585
Kirsten, F., Snelders, M. P., Jenkins, M., et al. 2021, NatAs, 5, 414
Kocz, J., Briggs, F. H., & Reynolds, J. 2010, AJ, 140, 2086
```

```
Kriele, M. A., Wayth, R. B., Bentum, M. J., Juswardy, B., & Trott, C. M. 2022,
  PASA, 39, e017
Law, C. J., Bower, G. C., Burke-Spolaor, S., et al. 2018, ApJS, 236, 8
Leung, C., Kader, Z., Masui, K. W., et al. 2022, Phys. Rev. D, 106, 043017
Leung, C., Mena-Parra, J., Masui, K., et al. 2021, AJ, 161, 81
Li, C. K., Lin, L., Xiong, S. L., et al. 2021, NatAs, 5, 378
Li, X., Zhou, B., He, H.-N., Fan, Y.-Z., & Wei, D.-M. 2014, ApJ, 797, 33
Li, Y., & Zhang, B. 2020, ApJL, 899, L6
Liu, B., Li, Z., Gao, H., & Zhu, Z.-H. 2019, PhRvD, 99, 123517
Lorimer, D. R., Bailes, M., McLaughlin, M. A., Narkevic, D. J., &
   Crawford, F. 2007, Sci, 318, 777
LSST Science Collaboration, Abell, P. A., Allison, J., et al. 2009,
  arXiv:0912.0201
Luo, R., Lee, K., Lorimer, D. R., & Zhang, B. 2018, MNRAS, 481, 2320
Macquart, J.-P., Bailes, M., Bhat, N. D. R., et al. 2010, PASA, 27, 272
Macquart, J. P., Prochaska, J. X., McQuinn, M., et al. 2020, Natur, 581, 391
MAGIC Collaboration, Acciari, V. A., Ansoldi, S., et al. 2018, MNRAS,
   481, 2479
Majid, W. A., Pearlman, A. B., Prince, T. A., et al. 2021, ApJL, 919, L6
Marcote, B., Nimmo, K., Hessels, J. W. T., et al. 2020, Natur, 577, 190
Marcote, B., Paragi, Z., Hessels, J. W. T., et al. 2017, ApJL, 834, L8
Margalit, B., & Loeb, A. 2016, MNRAS, 460, L25
Masui, K., Lin, H.-H., Sievers, J., et al. 2015, Natur, 528, 523
McQuinn, M. 2014, ApJL, 780, L33
Mena-Parra, J., Leung, C., Cary, S., et al. 2022, AJ, 163, 48
Metzger, B. D., Fang, K., & Margalit, B. 2020, ApJL, 902, L22
Michilli, D., Seymour, A., Hessels, J. W. T., et al. 2018, Natur, 553, 182
Muñoz, J. B., Kovetz, E. D., Dai, L., & Kamionkowski, M. 2016, PhRvL, 117,
Muñoz, J. B., & Loeb, A. 2018, PhRvD, 98, 103518
Newburgh, L. B., Bandura, K., Bucher, M. A., et al. 2016, Proc. SPIE, 9906,
Ng, C., Vanderlinde, K., Paradise, A., et al. 2017, in XXXII International
   Union of Radio Science General Assembly & Scientific Symposium (URSI
   GASS) (New York: IEEE), 20174
Niu, C. H., Aggarwal, K., Li, D., et al. 2021a, Natur, 606, 873
Niu, C.-H., Li, D., Luo, R., et al. 2021b, ApJL, 909, L8
```

Ocker, S. K., Cordes, J. M., & Chatterjee, S. 2021, ApJ, 911, 102

Petroff, E., Barr, E. D., Jameson, A., et al. 2016, PASA, 33, e045

Petroff, E., Hessels, J. W. T., & Lorimer, D. R. 2019, A&ARv, 27, 4

Pen, U.-L. 2018, NatAs, 2, 842

Oppermann, N., Yu, H.-R., & Pen, U.-L. 2018, MNRAS, 475, 5109

```
Petroff, E., Houben, L., Bannister, K., et al. 2017, arXiv:1710.08155
Petroff, E., van Straten, W., Johnston, S., et al. 2014, ApJL, 789, L26
Piro, L., Bruni, G., Troja, E., et al. 2021, A&A, 656, L15
Platts, E., Weltman, A., Walters, A., et al. 2019, PhR, 821, 1
Popov, S. B., & Postnov, K. A. 2013, arXiv:1307.4924
Popov, S. B., & Pshirkov, M. S. 2016, MNRAS, 462, L16
Prochaska, J. X., & Zheng, Y. 2019, MNRAS, 485, 648
Rafiei-Ravandi, M., Smith, K. M., Li, D., et al. 2021, ApJ, 922, 42
Rajwade, K., Stappers, B., Williams, C., et al. 2021, arXiv:2103.08410
Rajwade, K. M., Mickaliger, M. B., Stappers, B. W., et al. 2020, MNRAS,
  495, 3551
Ravi, V. 2019a, NatAs, 3, 928
Ravi, V. 2019b, ApJ, 872, 88
Ravi, V., Catha, M., D'Addario, L., et al. 2019, Natur, 572, 352
Schröder, F. G. 2017, PrPNP, 93, 1
Shin, K., Masui, K. W., Bhardwaj, M., et al. 2022, arXiv:2207.14316
Slosar, A., Ahmed, Z., Alonso, D., et al. 2019, BAAS, 51, 53
Spinelli, M., Bernardi, G., Garsden, H., et al. 2021, MNRAS, 505, 1575
Spitler, L. G., Scholz, P., Hessels, J. W. T., et al. 2016, Natur, 531, 202
Tavani, M., Casentini, C., Ursi, A., et al. 2021, NatAs, 5, 401
Taylor, J. H. 1974, A&AS, 15, 367
Tegmark, M., & Zaldarriaga, M. 2009, PhRvD, 79, 083530
Tendulkar, S. P., Bassa, C. G., Cordes, J. M., et al. 2017, ApJL, 834, L7
The CHIME Collaboration, Amiri, M., Bandura, K., et al. 2022, ApJS, 261, 29
The CHIME/FRB Collaboration, Andersen, B. C., Bandura, K., et al. 2021,
    Natur, 607, 256
The LIGO Scientific Collaboration, the Virgo Collaboration, the KAGRA
  Collaboration, et al. 2022, arXiv:2203.12038
Thulasiram, P., & Lin, H.-H. 2021, MNRAS, 508, 1947
Tominaga, N., Niino, Y., Totani, T., et al. 2018, PASJ, 70, 103
Torchinsky, S. A., Broderick, J. W., Gunst, A., Faulkner, A. J., &
  van Cappellen, W. 2016, arXiv:1610.00683
Totani, T. 2013, PASJ, 65, L12
Vanderlinde, K., Liu, A., Gaensler, B., et al. 2019, Canadian Long Range Plan
  for Astronomy and Astrophysics White Papers, 2020, 28
Wayth, R. B., Lenc, E., Bell, M. E., et al. 2015, PASA, 32, e025
Wei, J.-J., Gao, H., Wu, X.-F., & Mészáros, P. 2015, PhRvL, 115, 261101
Wei, J.-J., Wu, X.-F., & Gao, H. 2018, ApJL, 860, L7
Yalinewich, A., & Pen, U.-L. 2022, MNRAS, 515, 5682
Yamasaki, S., Totani, T., & Kiuchi, K. 2018, PASJ, 70, 39
Yang, Y.-P., Zhang, B., & Wei, J.-Y. 2019, ApJ, 878, 89
Zackay, B., & Ofek, E. O. 2017, ApJ, 835, 11
```