

Fundamentals of Solid State Physics

Thermal Properties

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Thermal Properties

- Heat Capacity (Thermal Capacity) 热容
- Thermal Conductivity 热导
- Thermal Expansion 热膨胀
- ...

Thermal Properties

- Thermal properties are the combinations of properties of **lattice vibration (phonons)** and **free electrons**
- For insulators, there are no free electron. Thermal properties of **lattice vibration (phonons)** dominate.
- For metals,
thermal properties = **phonon part + free electron part**

Thermal capacity

$$C_V = C_{V,p} + C_{V,e}$$

Thermal conductivity

$$\kappa = \kappa_p + \kappa_e$$

Fundamentals of Solid State Physics

Thermal Properties - Phonons

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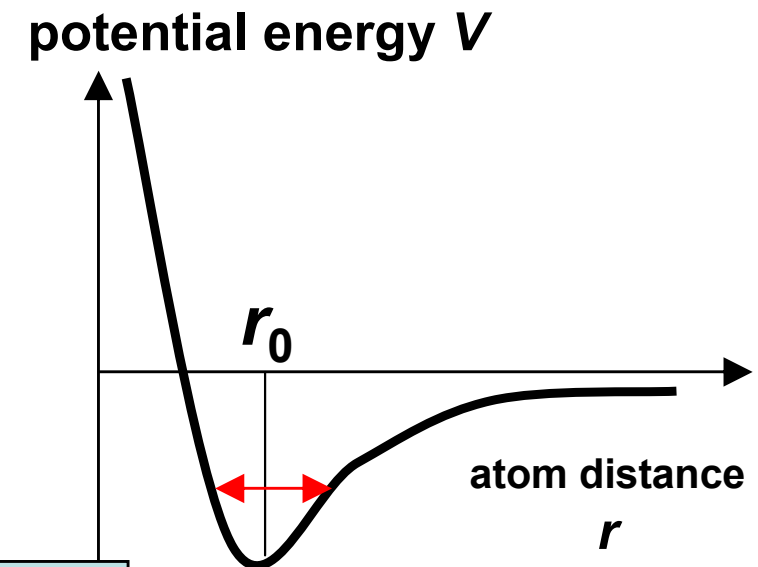
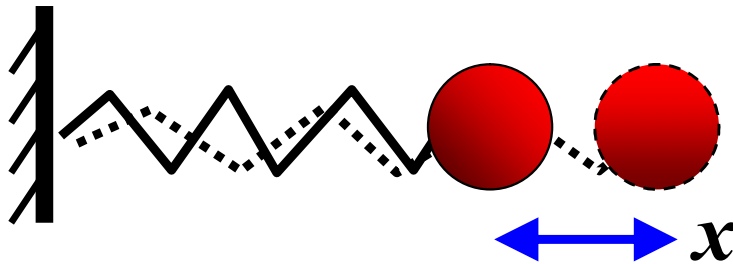


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Harmonic Oscillator: Classical Theory

- Vibration amplitude is continuous
- Energy is continuous, and temperature dependent



energy of one spring + one atom
= potential energy + kinetic energy

$$E = \frac{1}{2} Kx^2 + \frac{1}{2} mv^2 = \frac{1}{2} k_B T + \frac{1}{2} k_B T = k_B T$$

Thermal vibration
around r_0

Equipartition Theorem (能量均分定理)

Internal Energy 内能

- Total vibration energy of a crystal

- all the springs + all the atoms ($3NL$)

$$U = 3NLk_B T$$

N - # of primitive cells

L - # of atoms in a primitive cell

- Heat capacity (Specific heat) 比热容

- energy per unit of temperature

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V = 3NLk_B$$

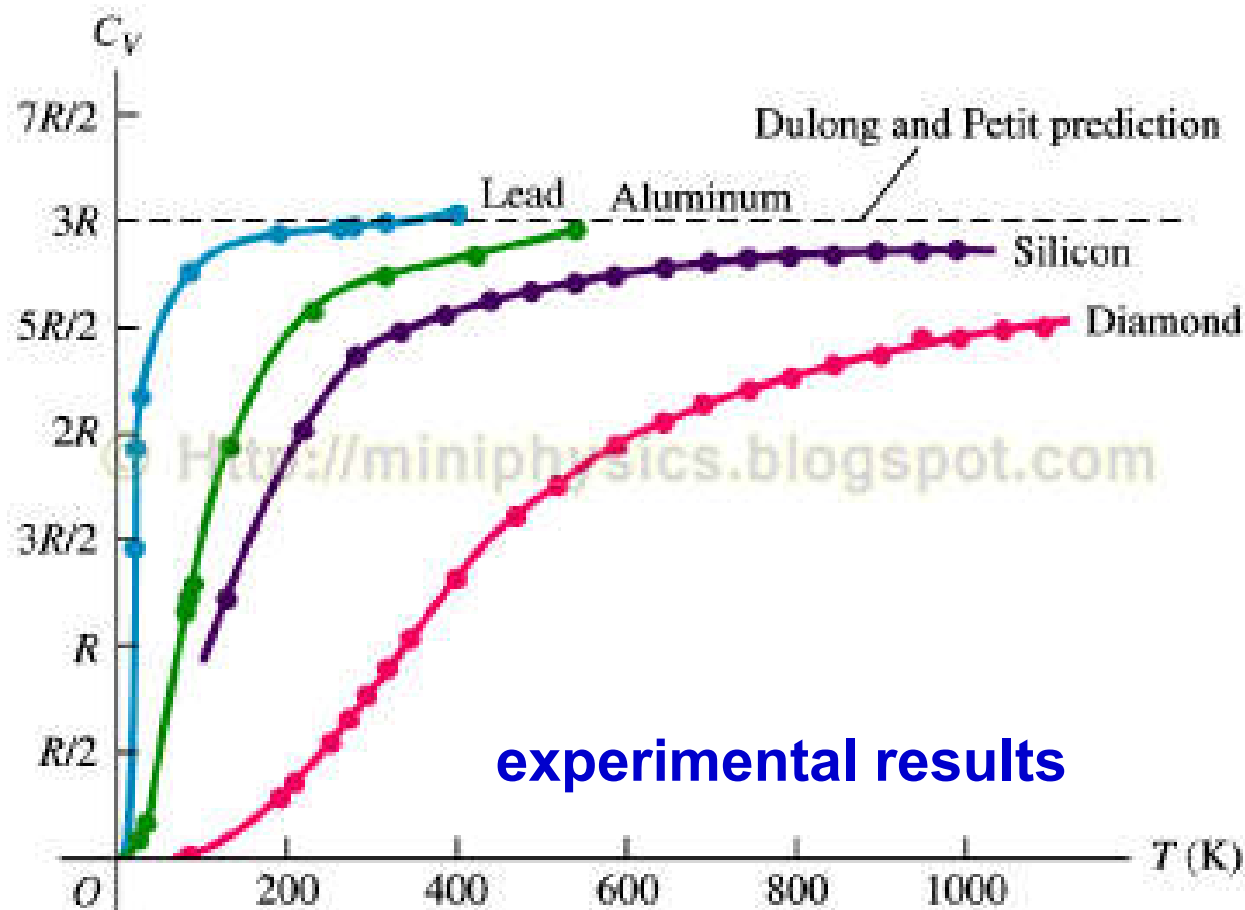
Dulong–Petit Law

In the system, every atom contributes an energy of $3k_B T$

Heat Capacity C_V

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V = 3NLk_B$$

Dulong–Petit Law

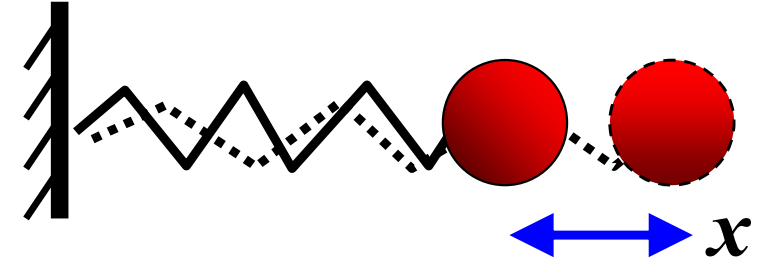


only valid at high temperature

experimental results

Harmonic Oscillator: Quantum Theory

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) + V(x) \cdot \psi(x) = E \psi(x)$$

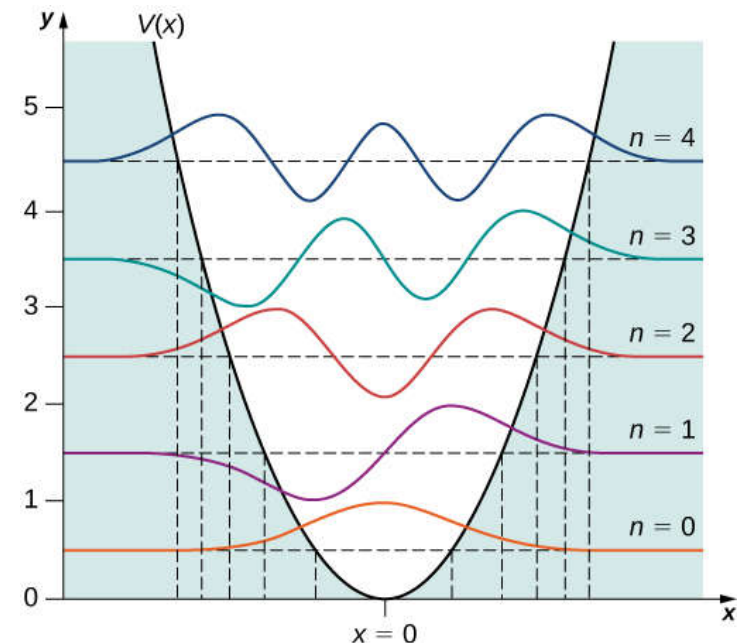


$$V(x) = \frac{1}{2} K x^2 = \frac{1}{2} m \omega^2 x^2$$

$$\omega = \sqrt{\frac{K}{m}}$$

→ $E_n = \left(\frac{1}{2} + n \right) \hbar \omega \quad n = 0, 1, 2, \dots$

Vibration energy is quantized



Harmonic Oscillator: Quantum Theory

■ Quantum theory

- Vibration energy is quantized
- At each ω state, the energy is the ground state energy ($\hbar\omega/2$) plus energy of n phonons ($\hbar\omega$)

$$E(\omega) = \left(\frac{1}{2} + n \right) \hbar\omega \quad n = 0, 1, 2, \dots$$

- average n in each state follows Bose-Einstein distribution

$$\bar{n} = \frac{1}{e^{\hbar\omega/k_B T} - 1}$$

- If there are N primitive cells, and L atoms in each cell, the total number of states is $3NL$

Internal Energy 内能

- Internal energy is **the ground state energy plus the energy of all the phonons**

$$U = U_0 + \sum_{i=1}^{3NL} \frac{\hbar \omega_i}{e^{\hbar \omega_i / k_B T} - 1}$$

$$= U_0 + \int_0^{\omega_{\max}} g(\omega) \frac{\hbar \omega}{e^{\hbar \omega / k_B T} - 1} d\omega$$

phonon energy

ground state

DOS

average phonon number
in each state

heat capacity

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V$$

Heat Capacity C_V

- At high temperature, $T \gg 0 \text{ K}$

$$\begin{aligned}
 U &= U_0 + \sum_{i=1}^{3NL} \frac{\hbar \omega_i}{e^{\hbar \omega_i / k_B T} - 1} \\
 &\approx U_0 + \sum_{i=1}^{3NL} \frac{\hbar \omega_i}{\hbar \omega_i / k_B T} \\
 &= U_0 + 3NLk_B T
 \end{aligned}$$

heat capacity

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V = 3NLk_B$$

constant

Dulong–Petit Law

Heat Capacity C_V

- At low and medium temperatures

For optical phonons:

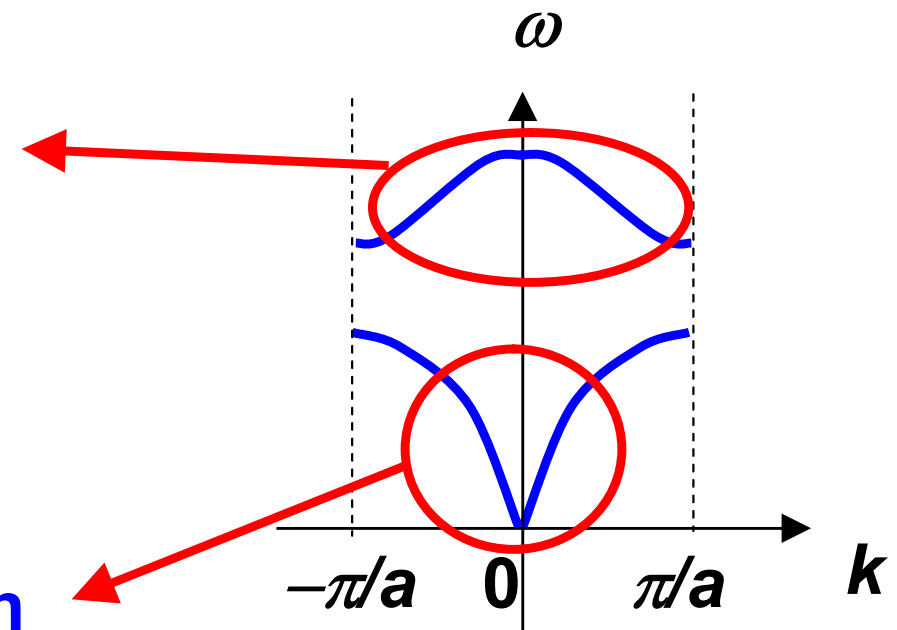
assume all the phonons
have the same ω_0

The Einstein Model
(爱因斯坦模型)

For acoustic phonons:

assume linear ω - k relation

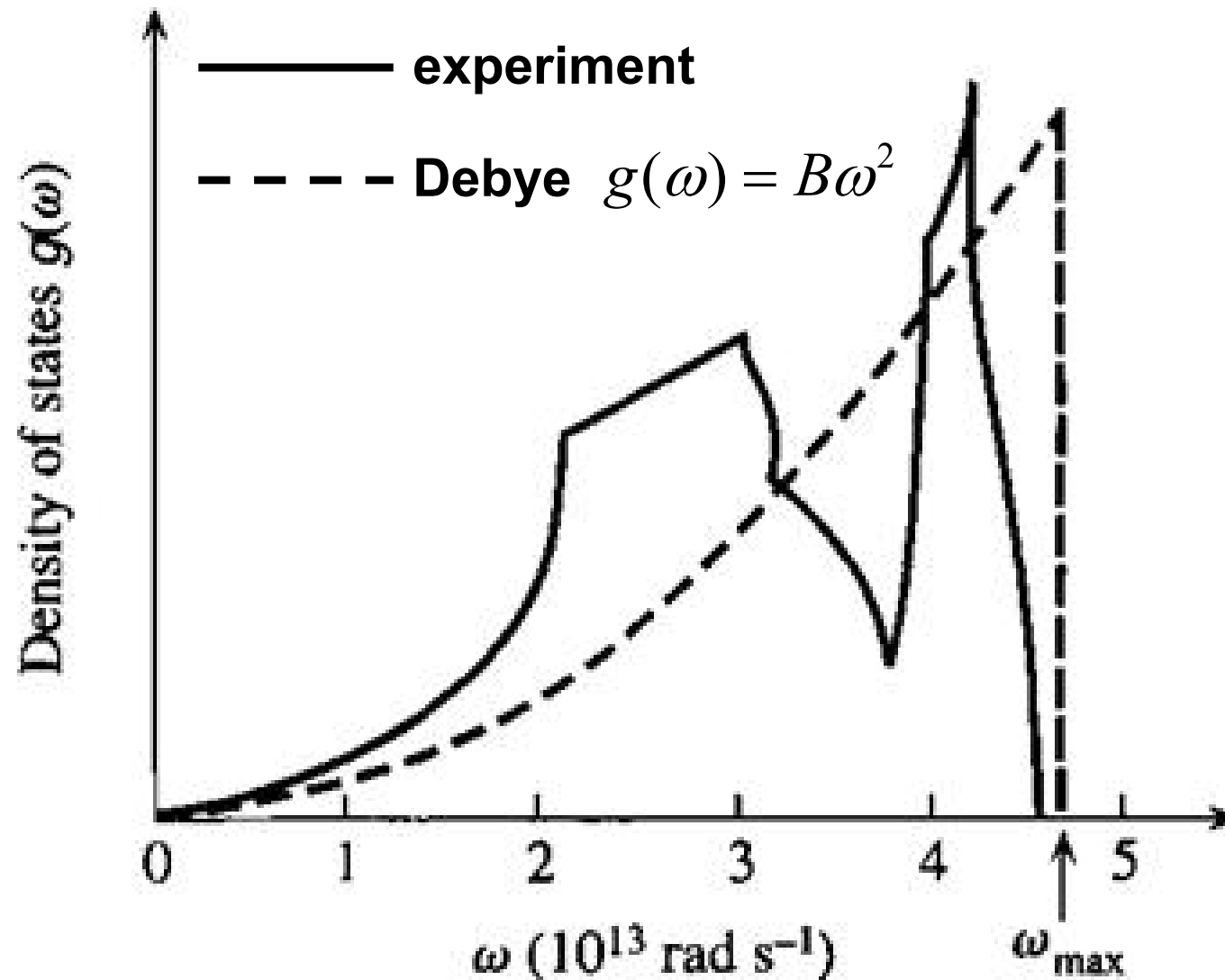
The Debye Model (德拜模型)



$$g(\omega) = B\omega^2$$

Phonon Density of States $g(\omega)$

DOS for copper



Heat Capacity C_V

- At very low temperature, $T \rightarrow 0$ K
- Most phonons have $\omega \rightarrow 0$, just focus on the acoustic branch

The Debye Model

$$U = U_0 + \int_0^{\omega_{\max}} g(\omega) \frac{\hbar \omega}{e^{\hbar \omega / k_B T} - 1} d\omega$$

$$\approx U_0 + \int_0^{\omega_{\max}} B \omega^2 \frac{\hbar \omega}{e^{\hbar \omega / k_B T} - 1} d\omega$$

heat capacity

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V \approx \frac{12\pi^4}{5} N k_B \left(\frac{T}{\theta_D} \right)^3 \propto T^3$$

**Debye
 T^3 Law**

$$\theta_D = \frac{\hbar \omega_D}{k_B} = \frac{\hbar v_g}{k_B} \left(\frac{6\pi^2 N}{V} \right)^{1/3}$$

Debye Temperature

Heat Capacity C_V

- At very low temperature, $T \rightarrow 0$ K
- Most phonons have $\omega \rightarrow 0$, just focus on the acoustic branch

$$\theta_D = \frac{\hbar \omega_D}{k_B} = \frac{\hbar v_g}{k_B} \left(\frac{6\pi^2 N}{V} \right)^{1/3}$$

Debye Temperature is around room temperature for most materials

	θ_D (K)
C (diamond)	2230
Si	645
Al	428
Cu	343

Heat Capacity C_V

- At median temperature, assume all the phonons in an optical branch have frequency ω_0

The Einstein Model

$$U = U_0 + N \frac{\hbar \omega_0}{e^{\hbar \omega_0 / k_B T} - 1}$$

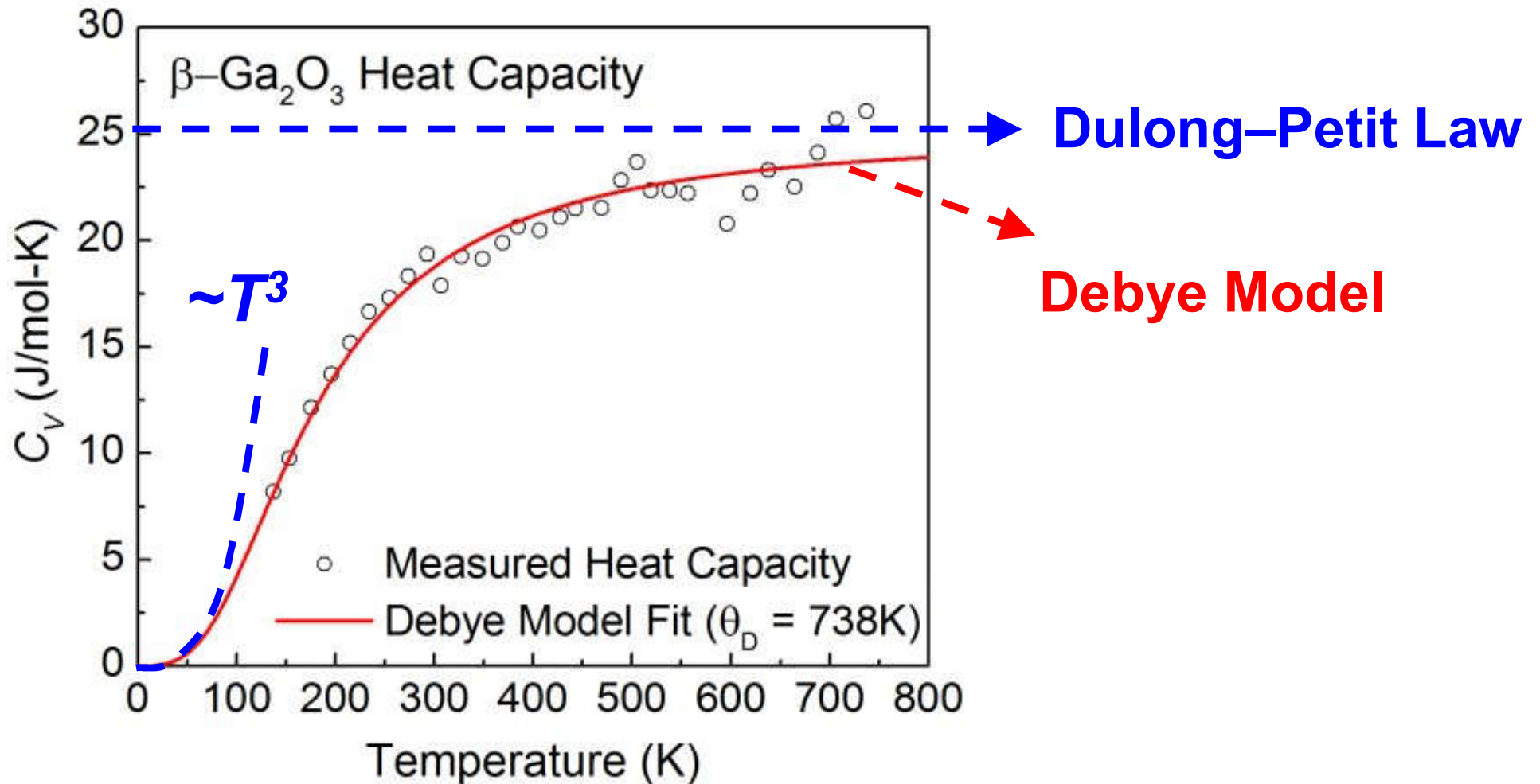
**heat capacity
(of one branch)**

$$\begin{aligned} C_V &= \left(\frac{\partial U}{\partial T} \right)_V = N k_B \left(\frac{\hbar \omega_0}{k_B T} \right)^2 \frac{e^{\hbar \omega_0 / k_B T}}{(e^{\hbar \omega_0 / k_B T} - 1)^2} \\ &= N k_B \left(\frac{\theta_E}{T} \right)^2 \frac{e^{\theta_E / T}}{(e^{\theta_E / T} - 1)^2} \end{aligned}$$

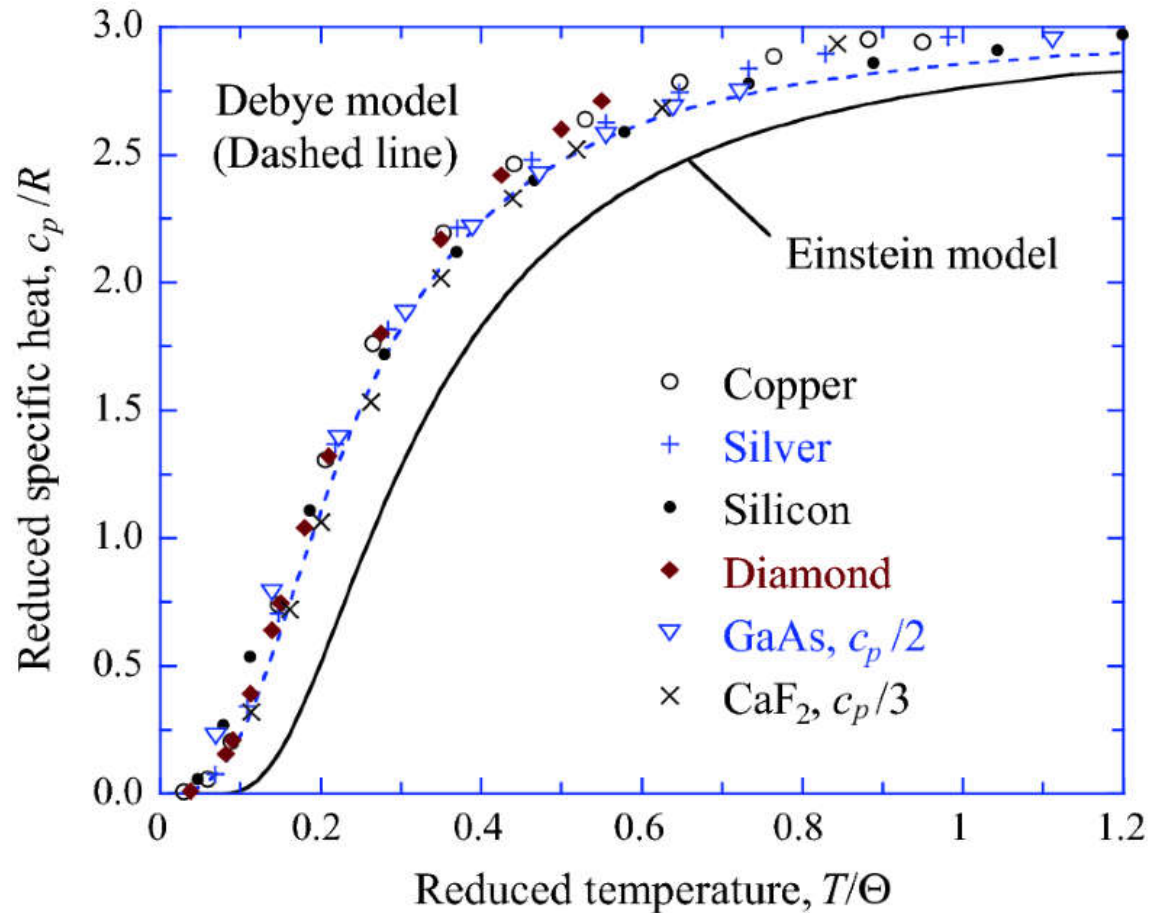
$$\theta_E = \frac{\hbar \omega_0}{k_B}$$

Einstein Temperature

Heat Capacity C_V - Example



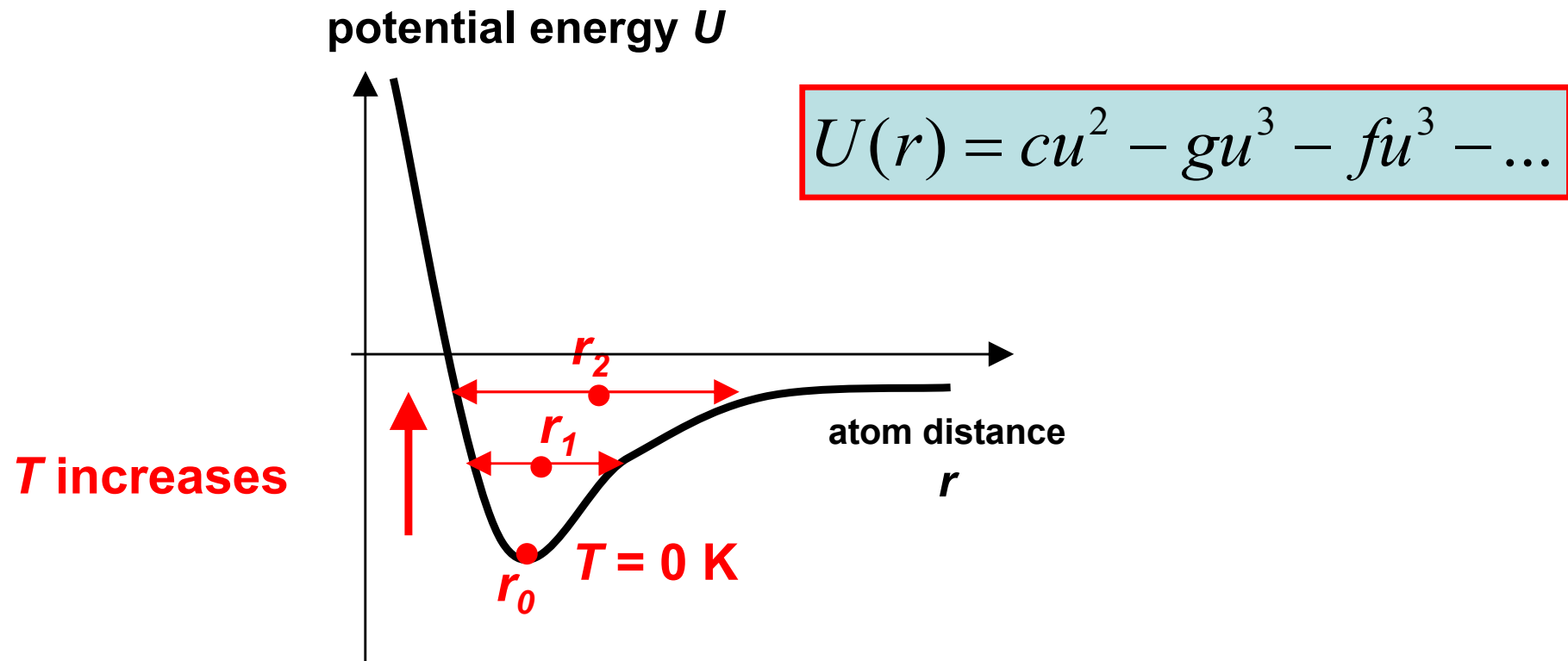
Heat Capacity C_V - Example



**The Debye Model
matches better
with experimental
results**

Thermal Expansion 热膨胀

- Thermal expansion originates from the *anharmonic* nature of the potential
- Vibration increases with temperature

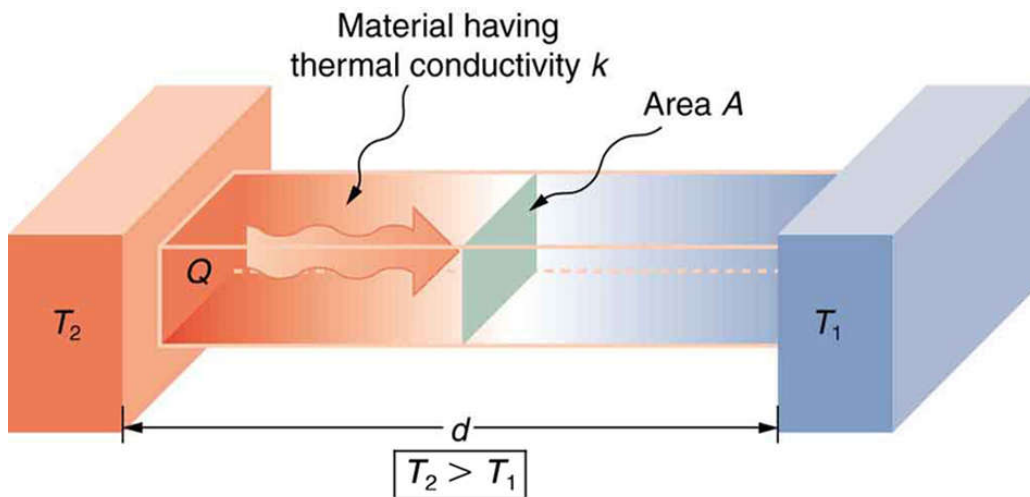


Thermal Conduction 热导

■ Fourier's Law

- heat flux is proportional to the temperature gradient

$$Q = -\kappa \cdot \frac{dT}{dx}$$



<https://www.khanacademy.org>

Q - heat flux (W/m^2)

κ - thermal conductivity (W/m/K)

T - temperature (K)

Thermal Conductivity 热导率

- Thermal conductivity κ

$$\kappa = \frac{1}{3} C_V v_g l = \frac{1}{3} C_V v_g^2 \tau_p$$

Ashcroft & Mermin, p20

C_V - thermal capacity

v_g - sound speed

l - phonon mean free path

τ_p - phonon relaxation time

l and τ_p is dependent on crystal structure, defects, impurities, ...

Q: Which material has the highest thermal conductivity?

Fundamentals of Solid State Physics

Thermal Properties of Free Electrons

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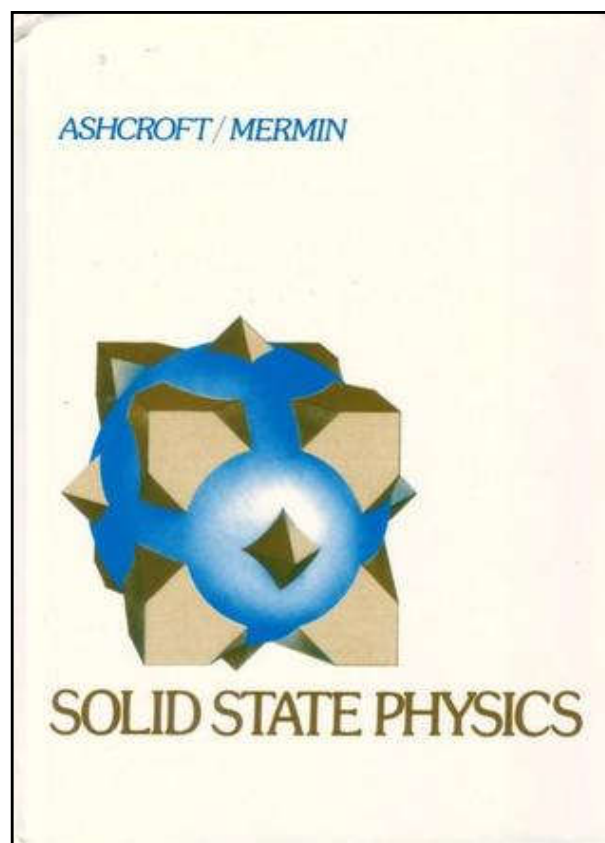


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Review

- Lecture 3.1, Sommerfeld Model
- Ashcroft & Mermin, Chapter 2



Density of States (DOS) 态密度

$$g(E) = \frac{dn}{dE}$$

DOS - number of energy states/levels per unit energy in $[E, E+dE]$, per unit volume

$$k = (3\pi^2 n)^{1/3}$$

$$E = \frac{\hbar^2 k^2}{2m_e}$$



$$n = \frac{1}{3\pi^2} \left(\frac{2m_e}{\hbar^2} \right)^{3/2} E^{3/2}$$

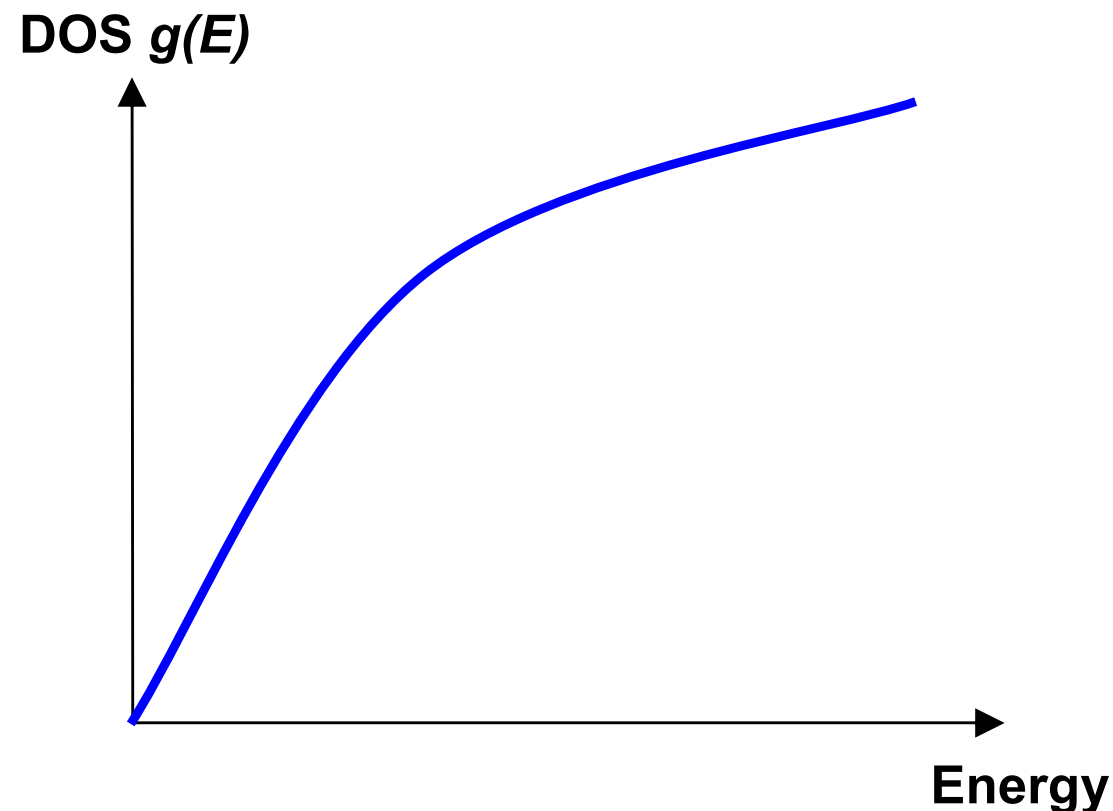
n - free electron density



$$g(E) = \frac{dn}{dE} = \frac{1}{2\pi^2} \left(\frac{2m_e}{\hbar^2} \right)^{3/2} E^{1/2}$$

Density of States (DOS) 态密度

$$g(E) = \frac{dn}{dE} = \frac{1}{2\pi^2} \left(\frac{2m_e}{\hbar^2} \right)^{3/2} E^{1/2}$$



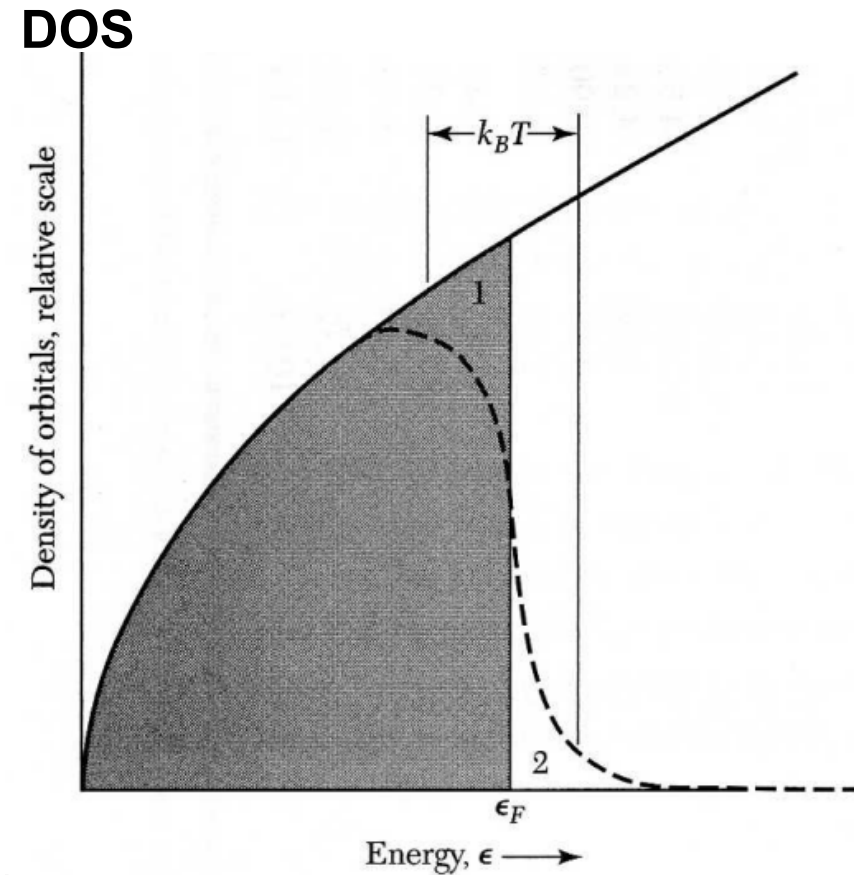
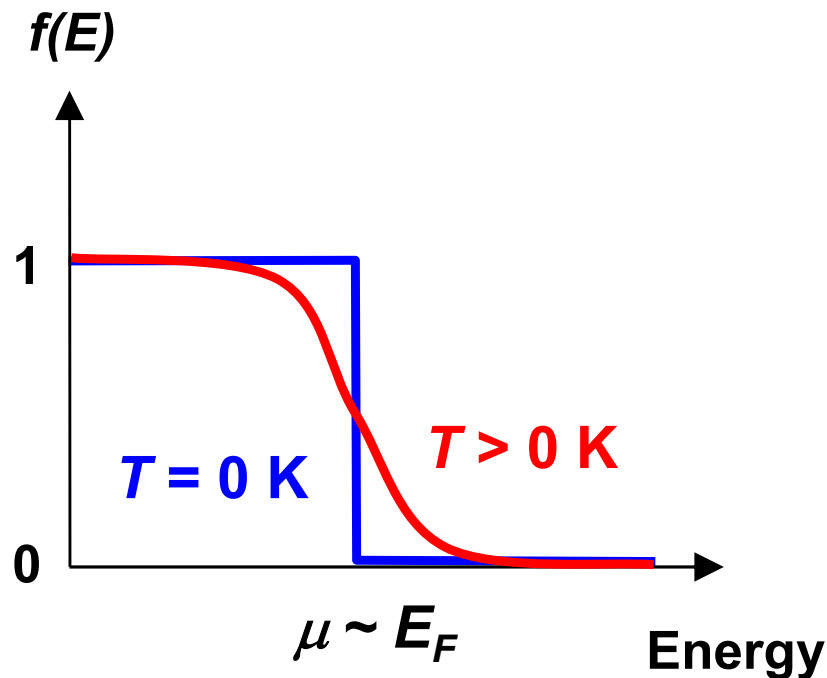
3D Free Electrons

Density of Electrons

Density of electrons = DOS * probability

$$f(E)g(E)$$

$$f(E) = \frac{1}{e^{(E-\mu)/k_B T} + 1}$$



When $T > 0 \text{ K}$, some electrons are excited to higher states (from 1 to 2)

Internal Energy 内能

- Internal energy is **the energy of all the free electrons**

$$U = \int_0^{+\infty} g(E) f(E) E dE$$

electron energy

DOS

average electron number
in each state
(Fermi-Dirac Distribution)

The Sommerfeld Model

Internal Energy 内能

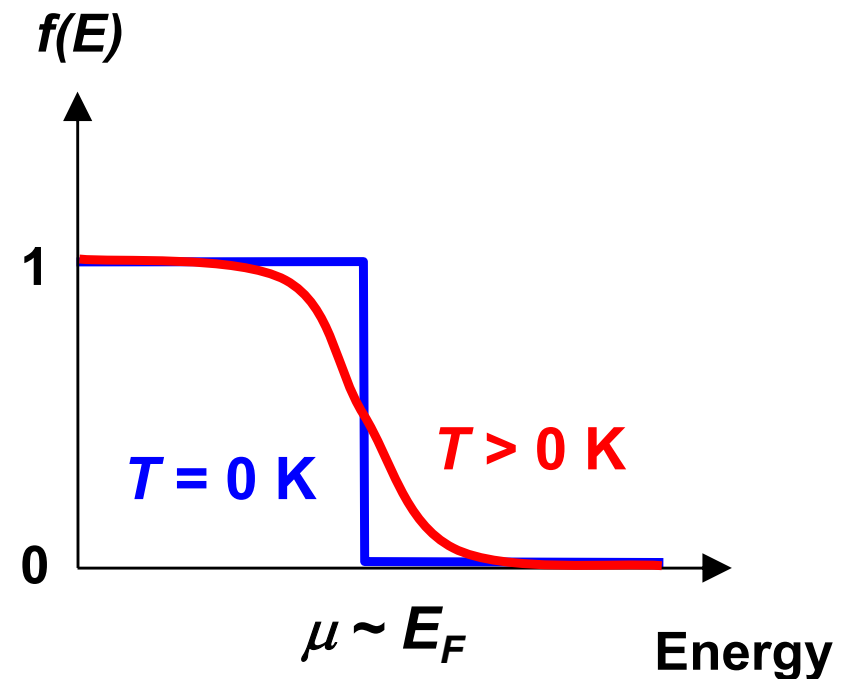
- When $T = 0$ K

$$U_0 = \int_0^{E_F} g(E) \cdot E dE$$

$$= \frac{3}{5} E_F$$

Homework 4.4

$$f(E) = \frac{1}{e^{(E-\mu)/k_B T} + 1}$$



Heat Capacity C_V

- When $T > 0$ K

$$U = U_0 + \frac{\pi^2}{6} (k_B T)^2 g(E_F)$$

heat capacity

$$C_{V,e} = \left(\frac{\partial U}{\partial T} \right)_V = \frac{\pi^2}{2} \frac{T}{T_F} n k_B \propto T$$

T_F - Fermi temperature ($\sim 10^4$ K)

Only a few electrons around E_F contribute to $C_{V,e}$.
 At room temperature, for free electrons $C_{V,e} \ll Nk_B$
 much smaller than C_V from phonons

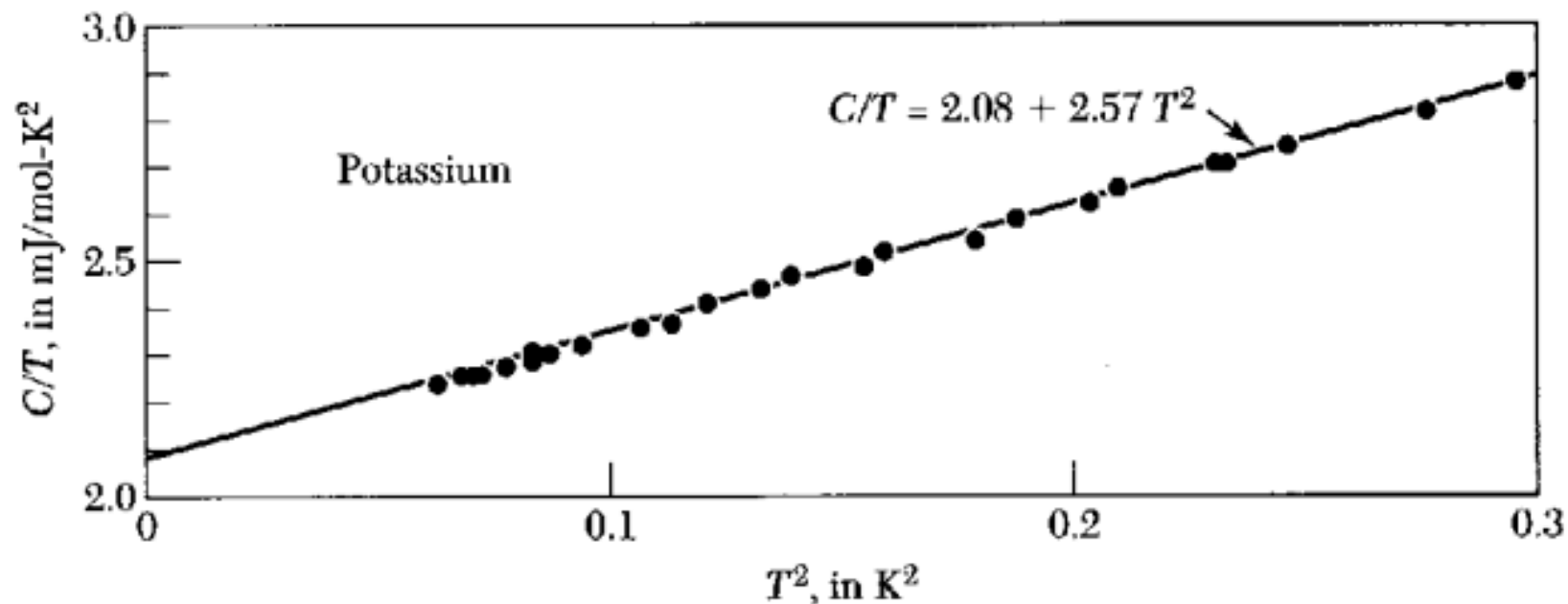
Heat Capacity C_V

- For metals at very low temperature $T \sim 0$ K
 - Thermal properties = phonon part + free electron part

$$C_V = C_{V,p} + C_{V,e} = AT^3 + \gamma T$$



$$C_V / T = AT^2 + \gamma$$



Thermal Conductivity 热导率

- Thermal conductivity κ for free electrons

$$\kappa_e = \frac{1}{3} C_V v_F l = \frac{1}{3} C_V v_F^2 \tau_e$$

Ashcroft & Mermin, p20

C_V - thermal capacity

v_F - Fermi velocity

l - electron mean free path

τ_e - electron relaxation time

Thermal Conductivity 热导率

- Thermal conductivity κ for metals

$$\kappa = \kappa_p + \kappa_e = \frac{1}{3} C_{V,p} v_g^2 \tau_p + \frac{1}{3} C_{V,e} v_F^2 \tau_e$$

For conductive metals like Cu or Ag, $v_F \gg v_g$
electron part dominates

Thermal Conductivity 热导率

- Relationship of thermal conductivity κ and electron conductivity σ for certain metals

□ Lorentz number L

Homework 8.7

$$L = \frac{\kappa}{\sigma T} = \frac{\pi^2}{3} \left(\frac{k_B}{e} \right)^2 = 2.44 \times 10^{-8} \text{ W} \cdot \Omega / \text{K}^2$$

Agree with experimental results
(Wiedemann and Franz law in 1853)

metals at 300 K	Cu	Ag	Au	Al	Fe	Pb
$L (10^{-8} \text{ W} \cdot \Omega / \text{K}^2)$	2.30	2.31	2.35	2.23	2.47	2.45

Q: Which metal has the highest thermal conductivity?

Thermal Conductivity 热导率

	κ (W/m/K)
C (diamond)	2000
Cu	400
C (graphite)	~200
Si	130
glass	1
paper	0.05

Q: High thermal conductivities of diamond and graphite have different origins. Why?

Summary

- Thermal properties are the combinations of properties of **lattice vibration (phonons)** and **free electrons**
- For insulators, there are no free electron. Thermal properties of **lattice vibration (phonons)** dominate.
- For metals,
thermal properties = **phonon part + free electron part**

Thermal capacity

$$C_V = C_{V,p} + C_{V,e}$$

Thermal conductivity

$$\kappa = \kappa_p + \kappa_e$$

Thank you for your attention