

Lattice gauge theory and the discretization of Dirac operators

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Based on a joint work with

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Physicists

Mathematicians

Then (FFM05)

• For $0 < \alpha \ll 1$, we have

the index of a twisted Dirac op D on $\mathbb{T}^d = (\mathbb{R}/\mathbb{Z})^d$

the index of the Wilson-Dirac op D_w on $\widehat{\mathbb{T}}^d = (\alpha \mathbb{Z}/\mathbb{Z})^d$

• The same holds for Clifford index, family index,
and equivariant index.

My motivation comes from

• Seiberg-Witten theory

• $D_{\text{FF}} = PL$

in dimension 4.

$$\underbrace{\dots}_{\text{a}} = \widehat{\mathbb{T}}^d$$

§1 Index, spectral flow, and K-theory

For simply. d : even

$$\mathbb{T}^d := (\mathbb{R}/\mathbb{Z})^d$$

E Hermitian

\mathbb{T}^d \mathbb{Z}_2 -graded Clifford module (odd)

$$\hookrightarrow \exists \gamma: E \rightarrow E$$

$$\exists c_j: E \rightarrow E \quad (j=1, \dots, d)$$

$$\text{sc. } \gamma = \gamma^*$$

$$c_j = -c_j^*$$

$$\gamma^2 = \text{id}$$

$$c_j^2 = -\text{id}$$

$$c_i c_j + c_j c_i = 0 \quad \text{if } i \neq j$$

$$\gamma c_i + c_i \gamma = 0$$

$\exists \gamma: E \rightarrow E$

$\exists c_j: E \rightarrow E \quad (j=1, \dots, d)$

$$\text{sc} \quad \gamma = \gamma^* \quad c_j = -c_j^*$$

$$\gamma^2 = \text{id} \quad c_j^2 = -\text{id}$$

$$c_i c_j + c_j c_i = 0 \quad \text{if } i \neq j$$

$$\forall c_i + c_j \gamma = 0$$

Fix a connection ∇ on E .

Dirac op $D: \Gamma(E) \rightarrow \Gamma(E)$

$$D \bar{\Phi} := \sum_{j=1}^d c_j \nabla_j \bar{\Phi}$$

Ex $d=2$

$$\gamma = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad c_1 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad c_2 = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} = 2(\overline{0})^\Gamma$$

$$\rightsquigarrow D = c_1 \frac{\partial}{\partial x} + c_2 \frac{\partial}{\partial y} = \begin{pmatrix} 0 & -\underbrace{\left(\frac{\partial}{\partial x} + i \frac{\partial}{\partial y} \right)}_{\mathcal{D}} \\ \underbrace{\left(\frac{\partial}{\partial x} + i \frac{\partial}{\partial y} \right)}_{\mathcal{D}^\Gamma} & 0 \end{pmatrix}$$

Dirac op $D : \Gamma(\epsilon) \rightarrow \Gamma(\epsilon)$

$$D\bar{\Psi} := \sum_{j=1}^d c_j \nabla_j \bar{\Psi}$$

D is self-adjoint and odd

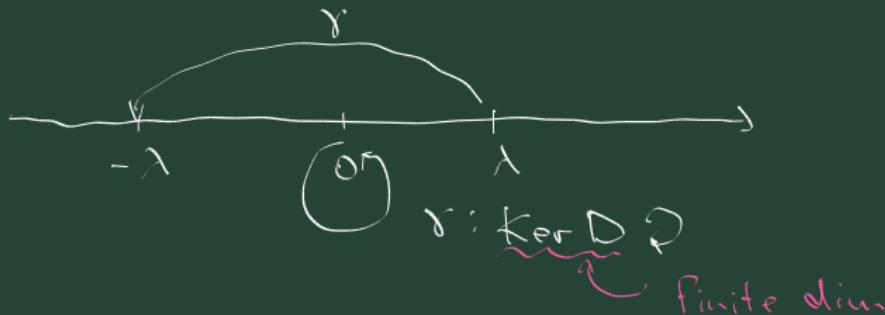
$$D = D^*$$

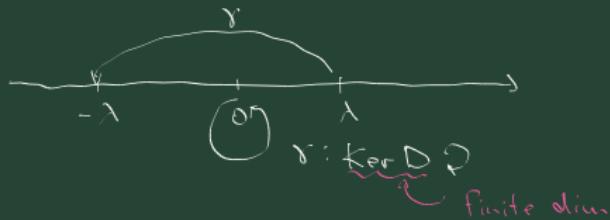
$$\underbrace{DD + \gamma D}_{s} = 0$$

If $D\bar{\Psi} = \lambda \bar{\Psi}$,

$$D(\gamma \bar{\Psi}) = -\gamma D\bar{\Psi} = -\gamma(\lambda \bar{\Psi}) = (\gamma \lambda)(\bar{\Psi})$$

$\rightsquigarrow \text{Spec}(D)$ is symmetric around the origin





Def

$$\text{ind}(D) := \text{fr}(\gamma|_{\ker D})$$

$$= \dim \{\tilde{x} \mid D\tilde{x} = 0 \text{ and } \gamma\tilde{x} = \tilde{y}\}$$

$$- \dim \{\tilde{x} \mid D\tilde{x} = 0 \text{ and } \gamma\tilde{x} = -\tilde{y}\}$$

$$\left\{ \begin{array}{l} \text{Def} \\ \text{ind}(D) := \text{tr}(\gamma | \ker D) \\ = \dim \{\tilde{x} \mid D\tilde{x} = 0 \text{ and } \gamma \tilde{x} = \tilde{x}\} \\ - \dim \{\tilde{x} \mid D\tilde{x} = 0 \text{ and } \gamma \tilde{x} = -\tilde{x}\} \end{array} \right.$$

Next we consider a family

$$D + w\gamma$$

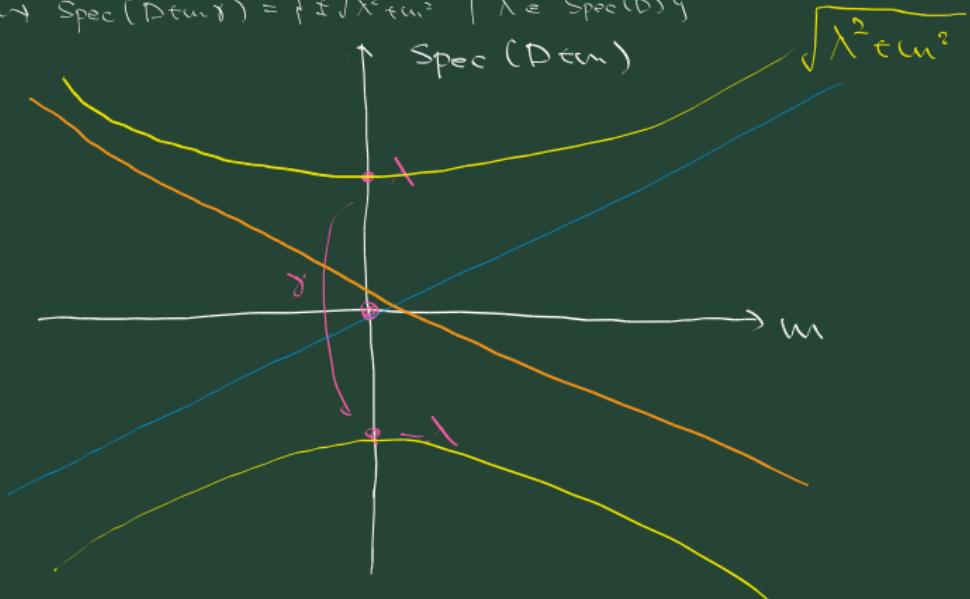
for $w \in [-1, 1]$

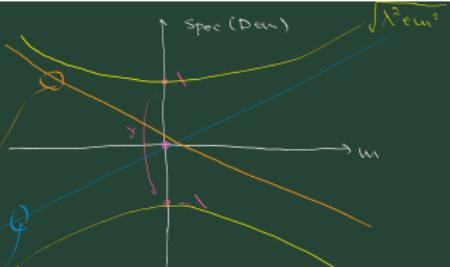
$$\begin{aligned} (D + w\gamma)^2 &= D^2 + w(D\gamma + \gamma D) + w^2 \gamma^2 \\ &= D^2 + w^2 \end{aligned}$$

$$\rightsquigarrow \text{Spec}(D + w\gamma) = \left\{ \pm \sqrt{\lambda^2 + w^2} \quad \mid \lambda \in \text{Spec}(D) \right\}$$

$$\begin{aligned}
 (\text{D} + m\gamma)^2 &= D^2 + m(D\gamma + \gamma D) + m^2\gamma^2 \\
 &= D^2 + m^2
 \end{aligned}$$

$$\rightsquigarrow \text{Spec}(\text{D} + m\gamma) = \{ \pm \sqrt{\lambda^2 + m^2} \mid \lambda \in \text{Spec}(D) \}$$





Def

$$\text{ind}(D) := \text{fr}(\gamma |_{\ker D})$$

$$= \dim \{\tilde{s} \mid D\tilde{s} = 0 \text{ and } \gamma \tilde{s} = \tilde{s}\}$$

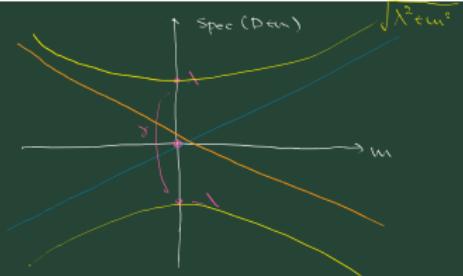
$$- \dim \{\tilde{s} \mid D\tilde{s} = 0 \text{ and } \gamma \tilde{s} = -\tilde{s}\}$$

If $D\tilde{s} = 0$ and $\gamma \tilde{s} = \pm \tilde{s}$,

$$(D + m\delta)\tilde{s} = \pm m\tilde{s}.$$

multiplicity = $\dim \{\tilde{s} \mid D\tilde{s} = 0 \text{ and } \gamma \tilde{s} = +\tilde{s}\}$

multiplicity = $\dim \{\tilde{s} \mid D\tilde{s} = 0 \text{ and } \gamma \tilde{s} = -\tilde{s}\}$



$$\boxed{\text{multiplicity} - \text{multiplicity}} = \text{ind}(D)$$

!!

$Sf(\{D_t\}_{t \in [0,1]}, \gamma_{m \in \{0,1\}})$
spectral flow.

$$\begin{array}{c}
 \text{Rmk} \quad K^0_{cpt} \underset{\Psi}{\cong} K^1(E_{\{0,1\}}, \{ \pi(q) \}) \\
 \text{ind}(D) \qquad \qquad \qquad \downarrow \\
 \qquad \qquad \qquad Sf(\{D_t\}_{t \in [0,1]}, \gamma_{m \in \{0,1\}})
 \end{array}$$

\S Dirac op on the lattice

Fix $N \in \{1, 2, \dots\}$, $a := \frac{1}{N}$

$$\hat{\mathbb{H}}^d := (\mathbb{Z}/N\mathbb{Z})^d$$

$$\begin{matrix} E \\ \downarrow \\ \hat{\mathbb{H}}^d \end{matrix}$$

$$e_j = \begin{pmatrix} 0 \\ \vdots \\ 1 \\ \vdots \\ 0 \end{pmatrix}_{\in \mathbb{J}}$$

Using the parallel transport of ∇ on E ,

$$\left\{ \begin{array}{l} \nabla_j^a \bar{\Phi}(x) := \frac{\bar{\Phi}(x + ae_j) - \bar{\Phi}(x)}{a} \quad \text{for } x \in \hat{\mathbb{H}}^d \\ \nabla_j^b \bar{\Phi}(x) := \frac{\bar{\Phi}(x) - \bar{\Phi}(x - ae_j)}{a} \end{array} \right.$$

$$\left\{ \begin{array}{l} \hat{\nabla}_j^f \bar{\Phi}(x) := \frac{\bar{\Phi}(x + a\epsilon_j) - \bar{\Phi}(x)}{a} \quad \text{for } x \in \mathbb{T}^d \\ \hat{\nabla}_j^h \bar{\Phi}(x) := \frac{\bar{\Phi}(x) - \bar{\Phi}(x - a\epsilon_j)}{a} \end{array} \right. \quad \dim(\Gamma(\hat{\epsilon})) < \infty$$

$$\rightsquigarrow (\hat{\nabla}_j^f)^* = - (\hat{\nabla}_j^h)$$

$$\gamma : \mathbb{E}^+ \xrightarrow{\sim} \mathbb{E}^-$$

$\gamma = +1 \quad \gamma = -1$

Lattice covariant derivative | Waine Dirac \leftrightarrow

$$\hat{\nabla}_j := \frac{\hat{\nabla}_j^f + \hat{\nabla}_j^h}{2}$$

$$\hat{D} := \sum c_j \hat{\nabla}_j$$

$$\rightsquigarrow (\hat{\nabla}_j)^* = - \hat{\nabla}_j$$

$$\begin{cases} \rightsquigarrow \begin{cases} \hat{D}^* = \hat{D} \\ \hat{D}\delta + \delta\hat{D} = 0 \end{cases} \\ \rightsquigarrow \boxed{\text{ind}(\hat{D}) = 0} \end{cases}$$

⑨ Wilson-Dirac operator

The Wilson term.

$$W := \frac{a}{z} \gamma \sum_{j=1}^d (\hat{\nabla}_j^f)^* (\hat{\nabla}_j^f)$$

Rank No continuous counterpart.

$$\rightsquigarrow W^* = W$$

$$- W\gamma = \gamma W \quad \leftarrow W \text{ is } \underline{\text{not}} \text{ odd}$$

The Wilson term

$$W := \frac{a}{z} \gamma \sum_{j=1}^d (\hat{\nabla}_j^f)^* (\hat{\nabla}_j^f)$$

$$\hat{\nabla}_j^f \bar{\Phi}(x) := \frac{\Phi(x + ae_j) - \bar{\Phi}(x)}{a}$$

$$\hat{\nabla}_j := \frac{\hat{\nabla}_j^f + \hat{\nabla}_j^b}{z}$$

$$\hat{\nabla}_j^b \bar{\Phi}(x) := \frac{\bar{\Phi}(x) - \bar{\Phi}(x - ae_j)}{a}$$

naive Dirac cp

$$\hat{D} := \sum c_j \hat{\nabla}_j$$

$$(\hat{D}_w)^* = D_w$$

Def Wilson-Dirac cp

} D_w is not odd.

$$D_w := \hat{D} + W$$

$$\begin{cases} \text{Def} \quad \text{Wilson-Divac op} \\ D_w := D + W \end{cases} \quad \left. \begin{array}{l} (D_w)^* = D_w \\ D_w \text{ is } \underline{\text{not}} \text{ odd} \end{array} \right\}$$

Thus, it is not so obvious to calculate the index of D_w
 \rightsquigarrow spectral flow

$$\begin{cases} \text{Def} \\ \text{For } 0 < \alpha \ll 1, \\ \text{ind}(D_w) := \text{sf}\left(\{D_w + m\delta\}_{m \in \mathbb{Z}_{\geq 0}}\right) \end{cases}$$

Lemma (FFMDSCE)

$$\begin{cases} \exists \alpha_0 = \frac{1}{N_0} \quad \forall \alpha \leq \frac{1}{N} \in (0, \alpha_0) \\ \text{s.t. } D_w + m\delta \text{ is } \underline{\text{invertible}} \text{ at } m = \pm 1 \text{ and } m = - \end{cases}$$

§ Main theorem

$$\boxed{\begin{array}{l} E \\ \downarrow \\ \mathbb{F}^d \\ D := \sum c_j \mathbb{F}_j \\ \text{ind}(D) \end{array}}$$

$$a \ll 1$$



$$\boxed{\begin{array}{l} \hat{E} = E |_{\mathbb{F}^d} \\ \downarrow \\ \hat{\mathbb{F}}^d = (\mathbb{F}_j |_{\mathbb{F}^d}) \\ \hat{D} := \sum c_j \hat{\mathbb{F}}_j \\ W := \frac{a}{z} \gamma \sum (\hat{\mathbb{F}}_j)^* (\hat{\mathbb{F}}_j) \\ D_W := \hat{D} + W \\ \text{ind}(D_W) := \text{sf}(\{D_W + a\gamma\}) \end{array}}$$

$$\begin{array}{ccc} K^0(pt) & \xrightarrow{\quad \cong \quad} & k^1(E(\cdot, \cdot), \{I(\cdot)\}) \\ \downarrow & & \downarrow \\ \text{ind}(D) & = & \text{sf}(D + a\gamma) \end{array}$$

$$\text{ind}(D_W) := \text{sf}(D_W + a\gamma)$$

$$\boxed{\begin{array}{l} E \\ \downarrow \\ T^d \\ D := \sum c_j T^d_j \\ \text{ind}(D) \end{array}}$$

$$a < 1 \quad \rightsquigarrow$$

$$\boxed{\begin{array}{l} \hat{E} = E |_{T^d} \\ \downarrow \\ \hat{T}^d = (a^{\infty}, \dots) \\ \hat{B} := \sum c_j \hat{T}^d_j \\ W := \frac{a}{2} r \sum (\hat{T}^d_j)^* (\hat{T}^d_j) \\ D_w := \hat{B} + W \end{array}}$$

$$K^0(\rho) \xrightarrow{\cong} K^1(B, \{A_i\}_{i=1}^k)$$

$$\downarrow$$

$$\text{ind}(D) = \text{sf}(D + w)$$

$$\text{ind}(D_w) := \text{sf}(D_w + w)$$

Then (FFM $\circ\Sigma^*$)

$$\exists \alpha_0 = \frac{1}{n_0} \quad \forall \alpha = \frac{1}{n} \in (0, \alpha_0)$$

$$\text{ind}(D) = \text{ind}(D_w)$$

The same holds for Clifford, ind., family inners, and equivariant rules.