

考虑一维双原子链(1D diatomic chain):

- 1) 计算长波极限 $(q \rightarrow 0)$ 下声学支和光学支格波的色散关系 $\omega \sim q$;
- 2) 分析第一布里渊区边界处的振动特点;
- 3) 当m = M时, 画出第一布里渊区内的色散关系 $\omega \sim q$, 并与一维单原子链 (1D monoatomic chain) 的情形进行比较.

提交时间: 3月19日之前

提交方式:手写(写明姓名学号)并拍照,通过本班课代表统一提交



考虑一维双原子链(1D diatomic chain):

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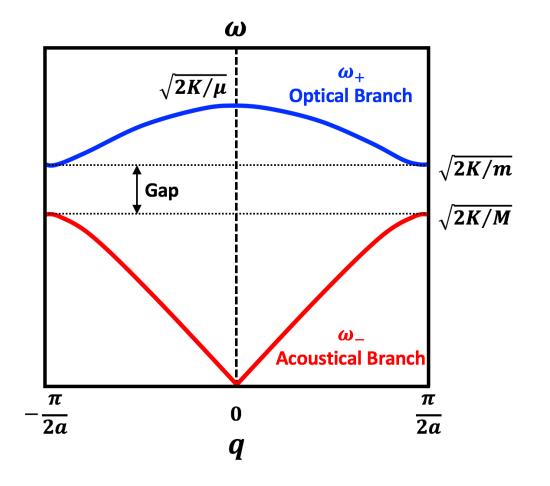
$$\omega^{2} = \begin{cases} \omega_{+}^{2} = \frac{K}{\mu} \left[1 + \sqrt{1 - \frac{4\mu^{2}}{mM}} \sin^{2}(aq) \right] & q \to 0 \\ \omega^{2} = \begin{cases} \omega_{+}^{2} = \frac{K}{\mu} \left[1 - \sqrt{1 - \frac{4\mu^{2}}{mM}} \sin^{2}(aq) \right] & q \to 0 \end{cases} & \omega_{+}^{2} \to 0 \end{cases}$$

注意对比一维单原子链:
$$\omega = 2\sqrt{\frac{K}{m}} \left| \sin\left(\frac{1}{2}aq\right) \right| \xrightarrow{q \to 0} \omega \to a\sqrt{\frac{K}{m}} |q|$$



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2) 分析第一布里渊区边界处的振动特点;

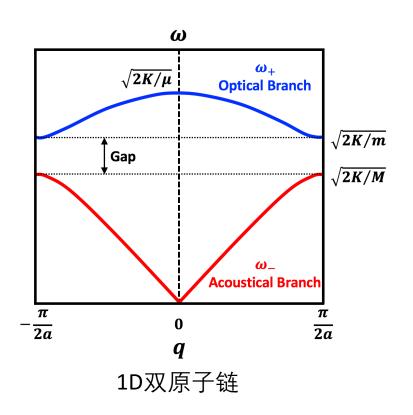
$$q = \pm \frac{\pi}{2a} \qquad \omega_+^2 \to \frac{2K}{m} \qquad \omega_-^2 \to \frac{2K}{M}$$

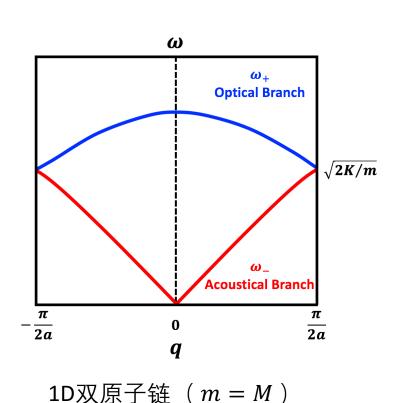


考虑一维双原子链(1D diatomic chain):

3) 当m = M时,画出第一布里渊区内的色散关系 $\omega \sim q$,并与一维单原子链

(1D monoatomic chain)的情形进行比较.





 $\omega_{\max} = 2\sqrt{\frac{K}{m}}$ 1D单原子链

Chapter 3.2: 课后作业



考虑一维量子谐振子
$$\hat{H} = \frac{1}{2m}\hat{p}^2 + \frac{1}{2}m\omega^2\hat{x}^2$$
,

证明:

1)
$$[\hat{a}, \hat{a}^+] = 1$$

2)
$$[\hat{a}^{+}\hat{a}, \hat{a}^{+}] = \hat{a}^{+}$$

3)
$$\left[\hat{a}^{+}\hat{a},\hat{a}\right] = -\hat{a}$$

4)
$$\widehat{H} = \hbar\omega \left(\widehat{a}^{+}\widehat{a} + \frac{1}{2}\right)$$

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证明要点: 利用
$$[\hat{x}, \hat{p}] = i\hbar$$

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$$\hat{a} = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{x} + \frac{i}{m\omega} \hat{p} \right) \qquad \hat{a}^{+} = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{x} - \frac{i}{m\omega} \hat{p} \right)$$

证明: 1)
$$[\hat{a}, \hat{a}^+] = 1$$

$$[\hat{a}, \hat{a}^+] = \hat{a} \ \hat{a}^+ - \hat{a}^+ \hat{a}$$

$$=\sqrt{\frac{m\omega}{2\hbar}}\Big(\hat{x}+\frac{i}{m\omega}\hat{p}\Big)\sqrt{\frac{m\omega}{2\hbar}}\Big(\hat{x}-\frac{i}{m\omega}\hat{p}\Big)-\sqrt{\frac{m\omega}{2\hbar}}\Big(\hat{x}-\frac{i}{m\omega}\hat{p}\Big)\sqrt{\frac{m\omega}{2\hbar}}\Big(\hat{x}+\frac{i}{m\omega}\hat{p}\Big)$$

$$= -\frac{i}{\hbar}(\hat{x}\hat{p} - \hat{p}\hat{x})$$

$$= 1$$



证明要点: 利用[
$$\hat{a}$$
, \hat{a}^+] = 1 即 \hat{a} \hat{a}^+ — \hat{a}^+ \hat{a} = 1

证明: 2)
$$[\hat{a}^+\hat{a}, \hat{a}^+] = \hat{a}^+$$

$$[\hat{a}^{+}\hat{a}, \hat{a}^{+}] = \hat{a}^{+}\hat{a}\hat{a}^{+} - \hat{a}^{+}\hat{a}^{+}\hat{a} = \hat{a}^{+}(1 + \hat{a}^{+}\hat{a}) - \hat{a}^{+}\hat{a}^{+}\hat{a} = \hat{a}^{+}$$

证明: 3)
$$[\hat{a}^+\hat{a},\hat{a}]=-\hat{a}$$

$$[\hat{a}^{\dagger}\hat{a},\hat{a}] = \hat{a}^{\dagger}\hat{a}\hat{a} - \hat{a}\hat{a}^{\dagger}\hat{a} = (\hat{a}\hat{a}^{\dagger} - 1)\hat{a} - \hat{a}\hat{a}^{\dagger}\hat{a} = -\hat{a}$$



证明要点: 利用
$$\hat{x} = \sqrt{\frac{\hbar}{2} \frac{1}{m\omega}} (\hat{a}^+ + \hat{a})$$
 $\hat{p} = i \sqrt{\frac{\hbar}{2} m\omega} (\hat{a}^+ - \hat{a})$ $[\hat{a}, \hat{a}^+] = 1$

证明: 4)
$$\widehat{H} = \hbar\omega\left(\widehat{a}^{\dagger}\widehat{a} + \frac{1}{2}\right)$$

$$\widehat{H} = \frac{1}{2m}\widehat{p}^2 + \frac{1}{2}m\omega^2\widehat{x}^2 = -\frac{\hbar\omega}{4}(\widehat{a}^+ - \widehat{a})(\widehat{a}^+ - \widehat{a}) + \frac{\hbar\omega}{4}(\widehat{a}^+ + \widehat{a})(\widehat{a}^+ + \widehat{a})$$

$$=\frac{\hbar\omega}{2}(\hat{a}^{+}\hat{a}+\hat{a}\hat{a}^{+})$$

$$=\frac{\hbar\omega}{2}(\hat{a}^{+}\hat{a}+\hat{a}^{+}\hat{a}+1)$$

$$= \hbar\omega \left(\hat{a}^{+}\hat{a} + \frac{1}{2}\right)$$

Chapter 3.3: 课后作业



设晶体中每个振动模的零点振动能为 $\frac{1}{2}\hbar\omega$,使用德拜模型和爱因斯坦模型分别求晶体的零点振动能。

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解题要点: 充分理解态密度 $g(\omega)$ 的物理意思,并利用态密度将"求和"化为"积分"

解:

$$E = \sum_{j=1}^{3N} \frac{1}{2} \hbar \omega_{j} = \int_{0}^{\omega} \frac{1}{2} \hbar \omega g(\omega) d\omega$$
 设N为原胞个数,每个原胞1个原子

$$g(\omega) = \frac{3V}{2\pi^2 \overline{c}^3} \omega^2$$

德拜模型:
$$g(\omega) = \frac{3V}{2\pi^2 \overline{c}^3} \omega^2 \qquad \omega_{\rm D} = \overline{c} \left[6\pi^2 \left(\frac{N}{V} \right) \right]^{1/3}$$

$$E_{\text{Debye}} = \int_0^{\omega_{\text{D}}} \frac{1}{2} \hbar \omega g(\omega) d\omega = \frac{9}{8} N \hbar \omega_{\text{D}}$$



解题要点: 充分理解态密度 $g(\omega)$ 的物理意思,并利用态密度将"求和"化为"积分"

解:

$$E = \sum_{i=1}^{3N} \frac{1}{2} \hbar \omega_{j} = \int_{0}^{\omega} \frac{1}{2} \hbar \omega g(\omega) d\omega$$

设N为原胞个数,每个原胞1个原子

爱因斯坦模型:
$$g(\omega) = 3N\delta(\omega - \omega_{\rm E})$$

$$\int_{-\infty}^{+\infty} f(x)\delta(x - x_0) = f(x_0)$$

$$E_{\text{Einstein}} = \int_{0}^{+\infty} \frac{1}{2} \hbar \omega g(\omega) d\omega = \frac{3}{2} N \hbar \omega_{\text{E}}$$