



PACE

Modelling & Simulation of a rotary inverted pendulum

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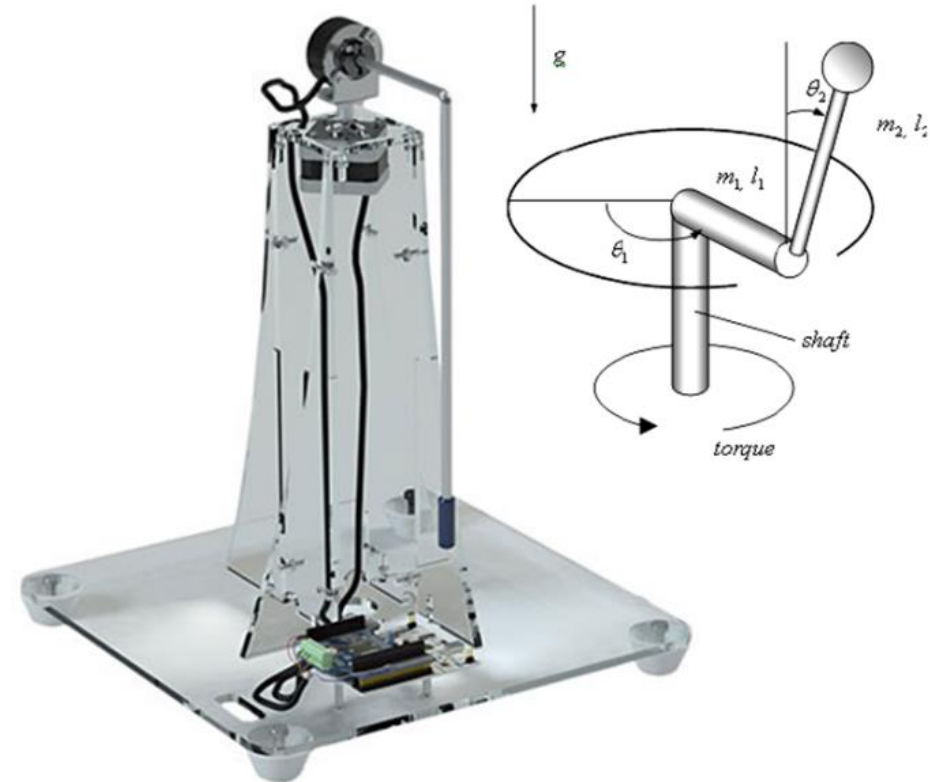
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Modelling

These slides outline the steps required to derive the non-linear equations of motion (aka the dynamic model) of a rotary pendulum with only one actuator.

Challenging topics involved:

- Vector calculus
- Rigid Body Kinematics & Dynamics
 - Lagrangian Mechanics
- Simulation of Ordinary Differential Equations (ODEs)



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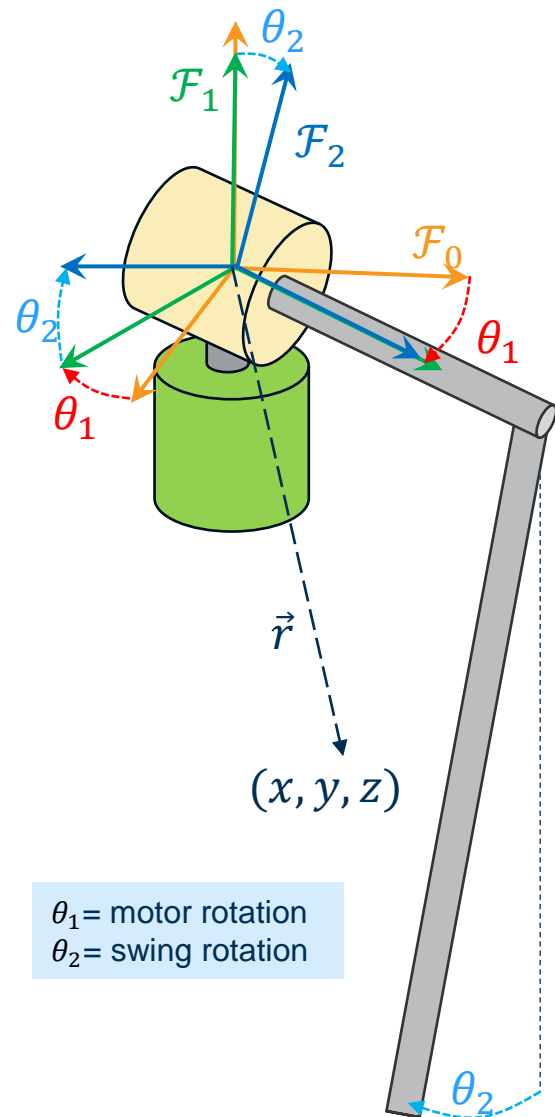
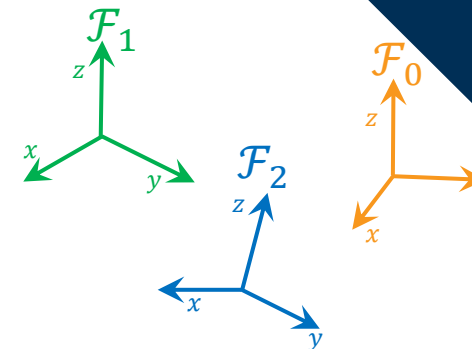
Important Kinematics

There are 3 frames of reference: \mathcal{F}_0 , \mathcal{F}_1 and \mathcal{F}_2 .

Inertial
(globally
fixed) frame

Encoder-
fixed
frame

Pendulum-fixed
frame



θ_1 = motor rotation
 θ_2 = swing rotation

- Let's define some position vector \vec{r} . We can express this as a numerical (x, y, z) vector, but these values depend on the frame of reference that we choose to express the vector in.
- Let's define: " \vec{r} expressed in frame \mathcal{F}_0 " as $\mathbf{r}^{\mathcal{F}_0} \in \mathbb{R}^3$.
- For our pendulum system, it can be shown that:

$$\mathbf{r}^{\mathcal{F}_1} = \mathbf{R}_{10}(\theta_1) \mathbf{r}^{\mathcal{F}_0}$$

$$\mathbf{r}^{\mathcal{F}_2} = \mathbf{R}_{21}(\theta_2) \mathbf{r}^{\mathcal{F}_1}$$

$$\mathbf{r}^{\mathcal{F}_2} = \mathbf{R}_{21} \mathbf{R}_{10} \mathbf{r}^{\mathcal{F}_0}$$

$$\mathbf{R}_{20}(\theta_1, \theta_2) = \begin{bmatrix} \cos \theta_1 \cos \theta_2 & -\sin \theta_1 \cos \theta_2 & \sin \theta_2 \cos \theta_1 \\ \sin \theta_1 \cos \theta_2 & \cos \theta_1 \cos \theta_2 & \sin \theta_1 \sin \theta_2 \\ -\sin \theta_2 & 0 & \cos \theta_2 \end{bmatrix}$$

This is known as a $z \rightarrow y'$
(yaw-pitch) transformation

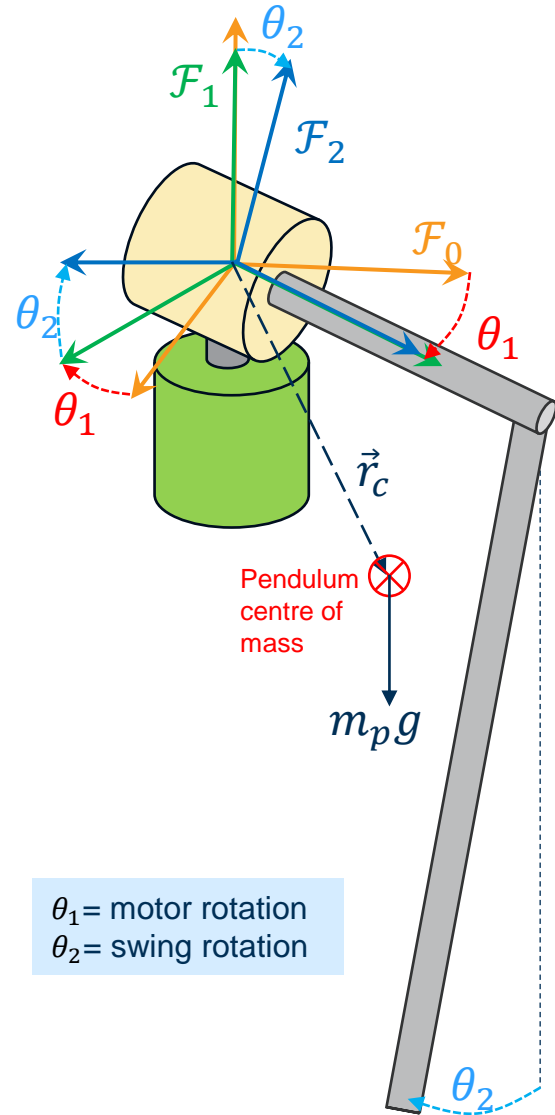
Rotation matrices \mathbf{R}

$$\mathbf{R}_{10} = \mathbf{R}_z(\theta_1) = \begin{bmatrix} \cos \theta_1 & -\sin \theta_1 & 0 \\ \sin \theta_1 & \cos \theta_1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$\mathbf{R}_{21} = \mathbf{R}_y(\theta_2) = \begin{bmatrix} \cos \theta_2 & 0 & \sin \theta_2 \\ 0 & 1 & 0 \\ -\sin \theta_2 & 0 & \cos \theta_2 \end{bmatrix}$$

\mathbf{R}_{20} describes the orientation of the pendulum body with respect to the inertial frame of reference.

Calculating the energy of the pendulum



Gravitational Potential Energy

$$V = m_p g \mathbf{r}_c \Big|_z$$

Vertical position of the centre of mass, expressed in the inertial frame.

$$\mathbf{r}_c = \mathbf{r}_c^{\mathcal{F}_0} = \mathbf{R}_{20}(\theta_1, \theta_2)^T \mathbf{r}_c^{\mathcal{F}_2}$$

Rotational Kinetic energy

$$T_{rot_p} = \frac{1}{2} \boldsymbol{\omega}_p^T \mathbf{I}_p \boldsymbol{\omega}_p$$

Pendulum inertia tensor,
calculated in the next slide

Angular velocity of pendulum,
expressed in the inertial frame. This
can be shown to be:

$$\boldsymbol{\omega}_p = \boldsymbol{\omega}_p^{\mathcal{F}_0} = \mathbf{R}_z(\theta_1)^T \begin{bmatrix} 0 \\ \dot{\theta}_2 \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 \end{bmatrix}$$

Linear Kinetic energy

$$T_{lin_p} = \frac{1}{2} m_p \dot{\mathbf{r}}_c^T \dot{\mathbf{r}}_c$$

Pendulum Mass

Linear velocity of pendulum centre of mass, expressed in the inertial frame.

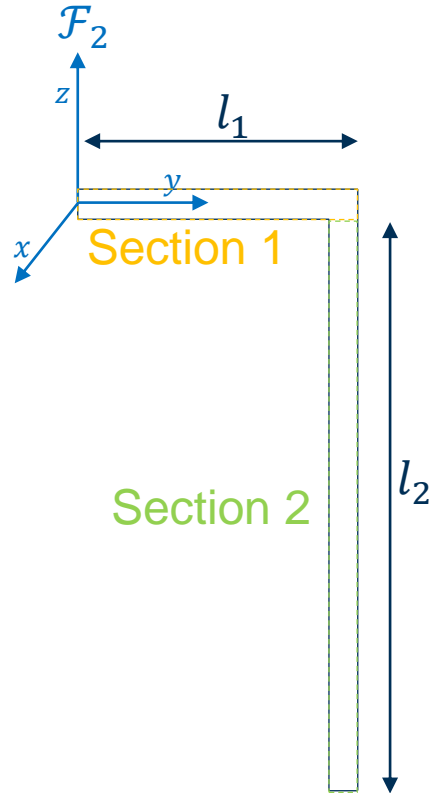
This can be shown to be:

$$\dot{\mathbf{r}}_c = \dot{\mathbf{r}}_c^{\mathcal{F}_0} = \dot{\mathbf{R}}_{20}(\theta_1, \theta_2)^T \mathbf{r}_c^{\mathcal{F}_2}$$

We can now define the systems 'Lagrangian': $L = T_{rot_p} + T_{lin_p} - V$

Aside: Approximating the pendulum Inertia Tensor I_p

I_p can be thought of as the 'rotational mass' of an object. This depends on the axis of rotation, which makes I_p a 3X3 matrix.



$$I_p^{\mathcal{F}_2} = I_1^{\mathcal{F}_2} + I_2^{\mathcal{F}_2}$$

Pendulum inertia tensor resolved in the moving (or 'body') frame \mathcal{F}_2

Section 1

$$I_1^{\mathcal{F}_2} = \begin{bmatrix} \frac{1}{4}m_1r^2 + \frac{1}{3}m_1l_1^2 & 0 & 0 \\ 0 & \frac{1}{2}m_1r^2 & 0 \\ 0 & 0 & \frac{1}{4}m_1r^2 + \frac{1}{3}ml_1^2 \end{bmatrix}$$

Section 2

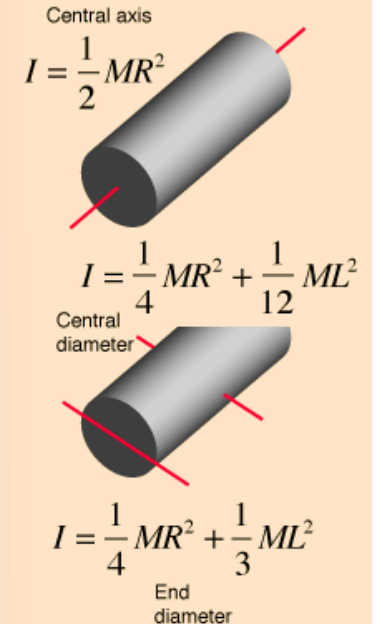
$$I_2^{\mathcal{F}_2} = \begin{bmatrix} \frac{1}{4}m_2r^2 + \frac{1}{12}m_2l_2^2 & 0 & 0 \\ 0 & \frac{1}{4}m_2r^2 + \frac{1}{12}m_2l_2^2 & 0 \\ 0 & 0 & \frac{1}{2}m_2r^2 \end{bmatrix} + m_2 \left(\begin{bmatrix} 0 \\ l_1 \\ -l_2 \\ 2 \end{bmatrix}^T \begin{bmatrix} 0 \\ l_1 \\ -l_2 \\ 2 \end{bmatrix} \mathbf{E}_3 - \begin{bmatrix} 0 \\ l_1 \\ -l_2 \\ 2 \end{bmatrix} \begin{bmatrix} 0 \\ l_1 \\ -l_2 \\ 2 \end{bmatrix}^T \right)$$

From the parallel axis theorem

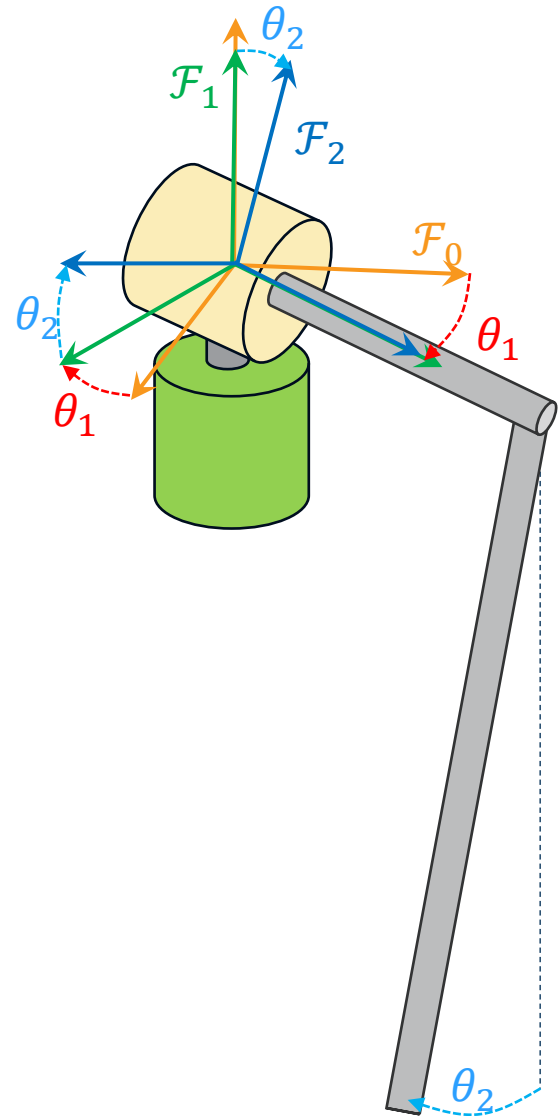
Assumption:
solid cylinder

(can improve this using measurements)

Moment of Inertia: Cylinder



Finding the state-space 'Equations of Motion'



Let's define our state coordinates, $q = \begin{bmatrix} \theta_1 \\ \theta_2 \end{bmatrix}$.
 $\rightarrow q$ completely describes the position of our pendulum

Euler-Lagrange Equations

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}} \right) - \frac{\partial L}{\partial q} = \boxed{F}$$

The 'external force' vector. If we ignore friction damping, then $F = \begin{bmatrix} u \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} u$, where u is the torque applied by the motor.

This equation basically creates the "F=ma" of the system: the equation that describes the behaviour of the pendulum. Also known as a "Dynamic Model".

By substituting L into this and expanding / rearranging, the result ends up looking something like this:

"Mass Matrix" $\boxed{M(q)}\ddot{q} + \boxed{N(q, \dot{q})} = F$



Vector describing
gravity & Coriolis
effects

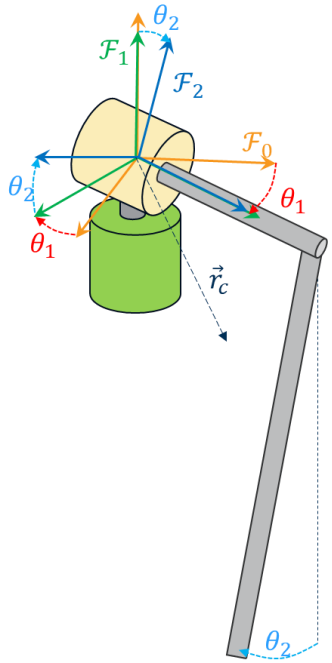
$$\dot{x} = \begin{bmatrix} \dot{q} \\ M^{-1}(F - N) \end{bmatrix} = \underbrace{\begin{bmatrix} \dot{q} \\ -M^{-1}N \end{bmatrix}}_{f(x)} + \underbrace{\begin{bmatrix} 0 \\ M^{-1} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \end{bmatrix}}_{g(x)} u$$

$$\boxed{\dot{x} = f(x) + g(x)u}$$

Often it is nice to express this in the equivalent 'first order ODE form', by defining a new set of variables

$$x = \begin{bmatrix} q \\ \dot{q} \end{bmatrix} = \begin{bmatrix} \theta_1 \\ \theta_2 \\ \dot{\theta}_1 \\ \dot{\theta}_2 \end{bmatrix}$$

Simulating the free-response system ($u=0$)



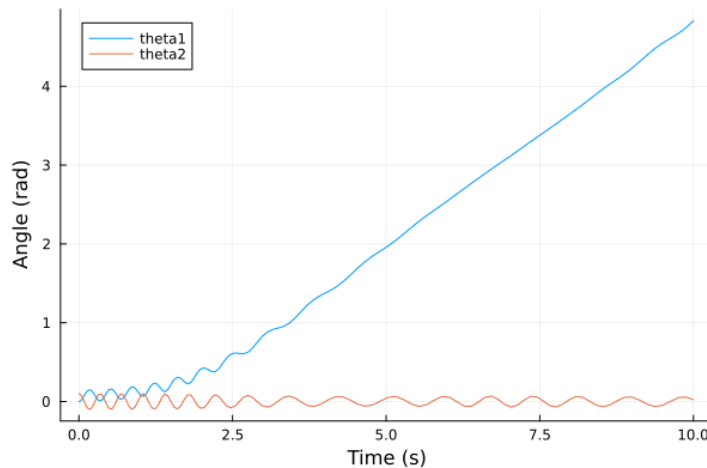
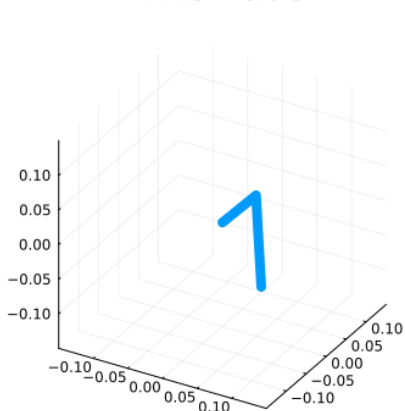
Time = 0.0 s

Given a starting state $x(t = 0) = x_0$, the model equations $\dot{x} = f(x)$ can be ‘solved’ (aka ‘simulated’) using a variety of techniques (aka ‘ODE solvers’).

- This can be done in Python, Matlab, C++, Julia and more.
 - I use ODE solvers in Julia’s DifferentialEquations.jl toolbox, this is just my personal preference.

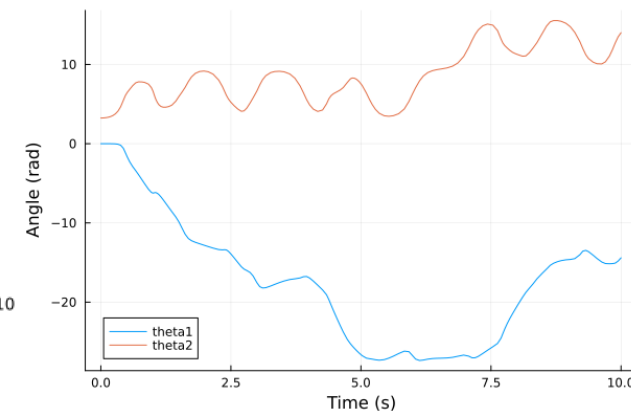
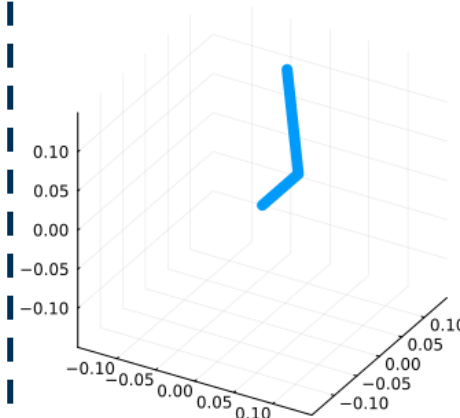
You can then create an animation of the model simulation.

$$\text{Example 1: } x_0 = \begin{bmatrix} \theta_{10} \\ \theta_{20} \\ \dot{\theta}_{10} \\ \dot{\theta}_{20} \end{bmatrix} = \begin{bmatrix} 0 \\ 0.1 \\ 0 \\ 0 \end{bmatrix}, \quad u = 0$$



$$\text{Example 2: } x_0 = \begin{bmatrix} 0 \\ \pi + 0.1 \\ 0 \\ 0 \end{bmatrix}, \quad u = 0$$

Time = 0.0 s

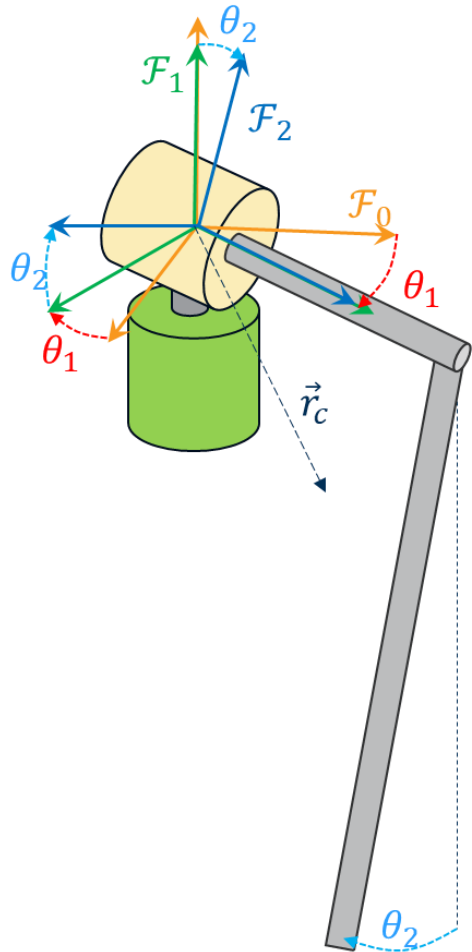


Very chaotic!
(no damping is modelled)

Adding damping to the model

In reality, there will be *viscous damping* at the joints.

- This is a friction torque proportional to the joint velocity:



$$\tau_{fr_i} \approx -\boxed{d_{fr_i}} \dot{\theta}_i$$

A damping constant of proportionality for joint i

In vector form:

$$\tau_{fr} \approx -\boxed{D} \dot{\theta}$$

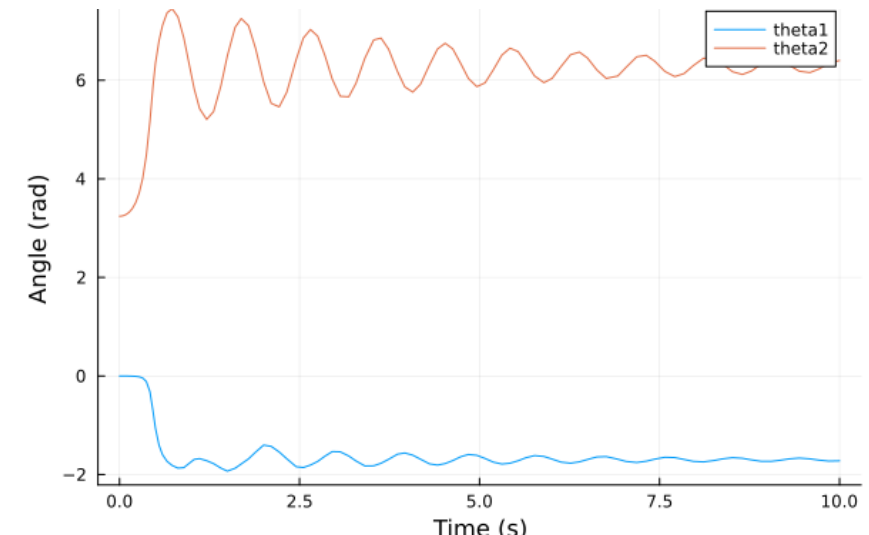
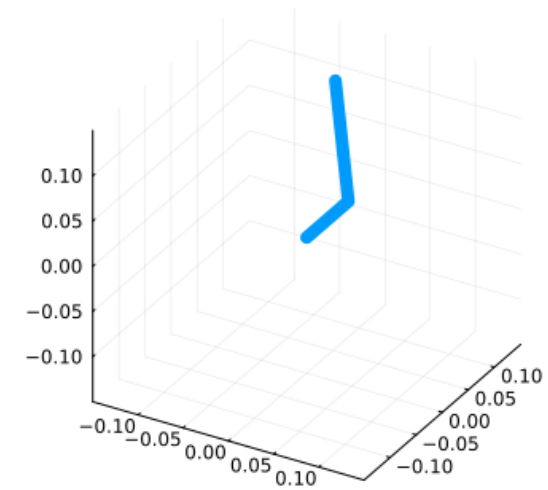
A diagonal Damping matrix

The model now becomes:

$$M(q)\ddot{q} + N(q, \dot{q}) + D\dot{\theta} = F$$

$$\dot{x} = \underbrace{\begin{bmatrix} \dot{q} \\ M^{-1}(F - N + D\dot{\theta}) \end{bmatrix}}_{f(x,u)}$$

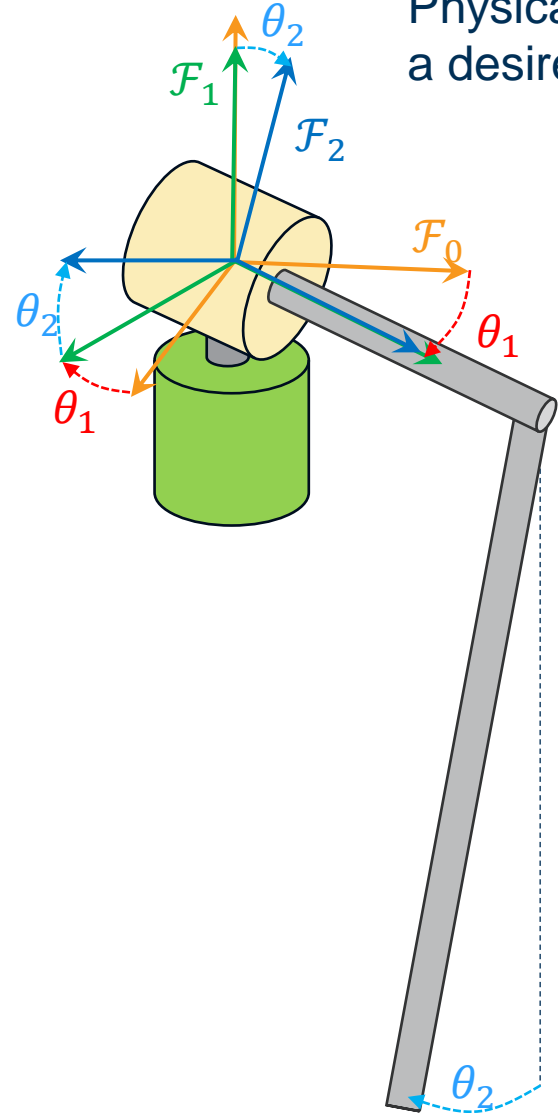
Time = 0.0 s



Stepper motor stabilization control

What if our control input is $u = \ddot{\theta}_1$, instead of the motor torque?

Let's assume we can precisely control the acceleration of the stepper motor $\ddot{\theta}_1$. Physically, this means the motor imparts a torque on the system that ensures $\ddot{\theta}_1$ follows a desired acceleration, $u(t)$. We can consider this as a *constraint*!



-> Lets express $\ddot{\theta}_1 = \underbrace{[1 \ 0]}_A \ddot{q}$

Dynamics equations

$$M\ddot{q} + N + A^T \tau_1 = 0 \quad (1)$$

rearrange

$$\ddot{q} = M^{-1}(-N - A^T \tau_1) = 0$$

Constraint equations

$$A\ddot{q} - u(t) = 0$$

Sub into

$$AM^{-1}(-N - A^T \tau_1) - u(t) = 0$$

rearrange

$$\tau_1 = -(AM^{-1}A^T)^{-1}(AM^{-1}N + u)$$

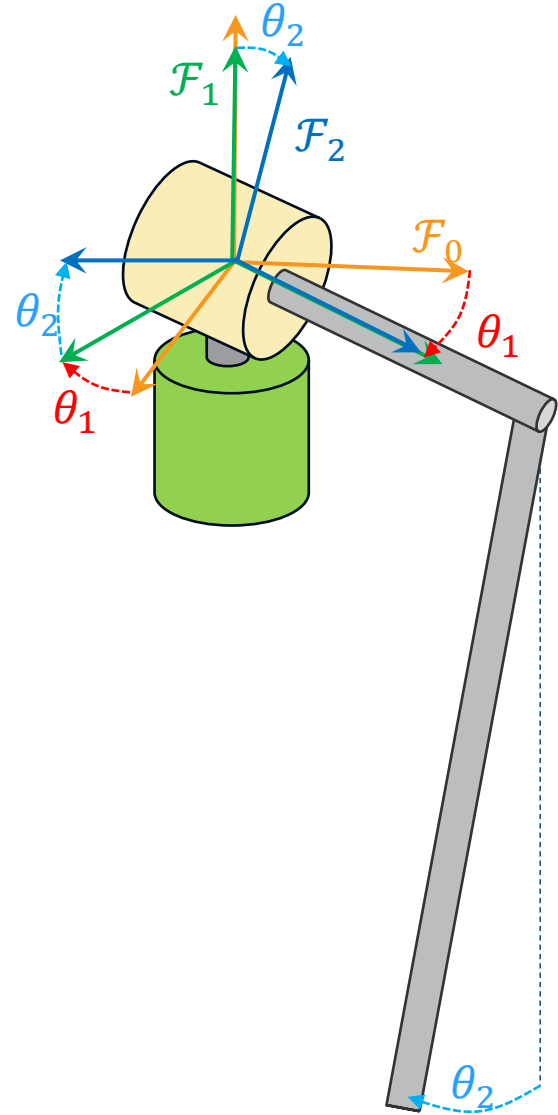
Sub into (1)

$$M\ddot{q} + N - A^T(AM^{-1}A^T)^{-1}(AM^{-1}N + u) = 0$$

$$M\ddot{q} + N_a = B_a u$$

$$\begin{cases} N_a = N - A^T(AM^{-1}A^T)^{-1}AM^{-1}N \\ B_a = A^T(AM^{-1}A^T)^{-1} \end{cases}$$

What if our control input is $u = \ddot{\theta}_1$, instead of the motor torque?

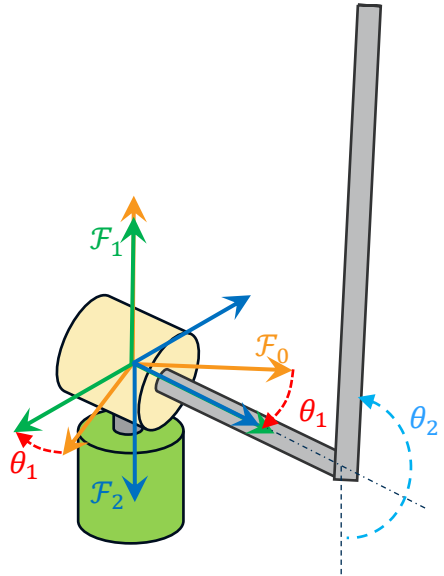


$$M\ddot{q} + N_a = B_a u \quad \left\{ \begin{array}{l} N_a = N - A^T(AM^{-1}A^T)^{-1}AM^{-1}N \\ B_a = A^T(AM^{-1}A^T)^{-1} \end{array} \right.$$

First order form

$$\dot{x} = \begin{bmatrix} \dot{q} \\ \ddot{q} \end{bmatrix} = \underbrace{\begin{bmatrix} \dot{q} \\ -M^{-1}N_a \end{bmatrix}}_{f(x)} + \underbrace{\begin{bmatrix} 0 \\ M^{-1}B_a \end{bmatrix}}_{g(x)} u$$

What should $u(t)$ be if we want to stabilise the pendulum?

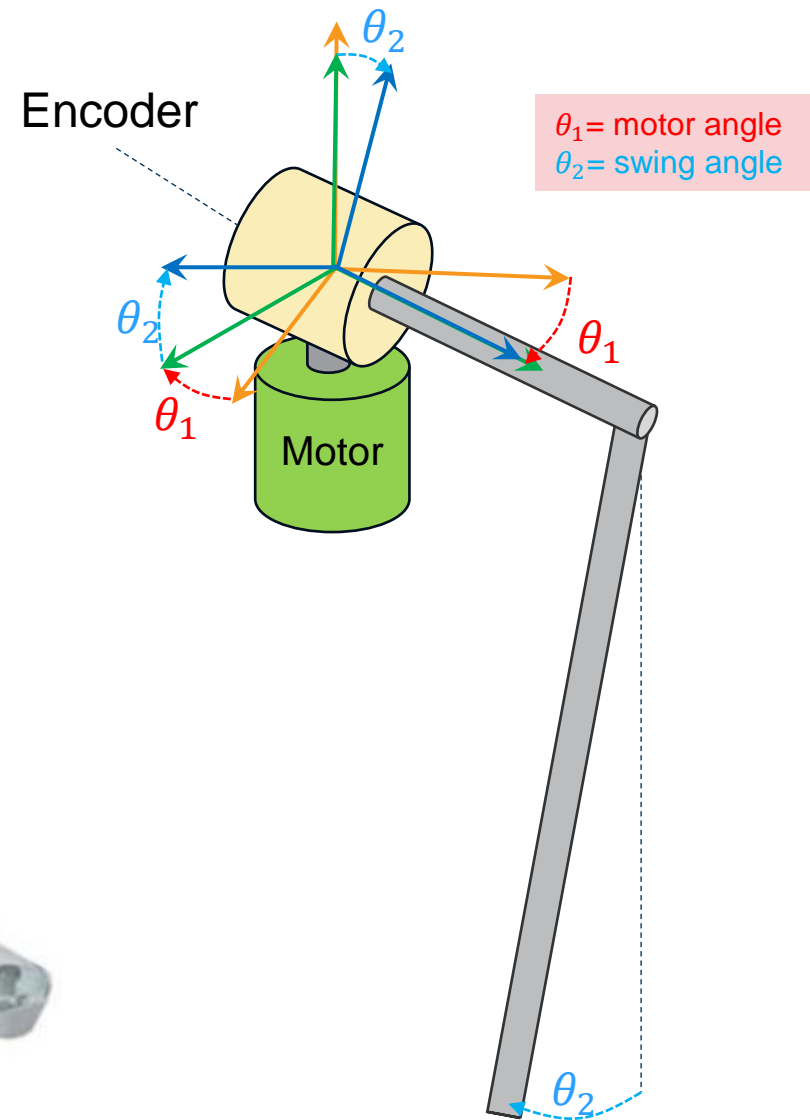


$$\dot{x} = f(x) + Bu$$

$$\Delta \dot{x} = \begin{bmatrix} \dot{\theta}_1 \\ \theta_2 \\ \dot{\theta}_1 \\ \dot{\theta}_2 \end{bmatrix} \approx \left. \frac{\partial f(x)}{\partial x} \right|_{x_{equil}} \Delta x + Bu$$

$$x_{equil} = \begin{bmatrix} \theta_1 \\ \theta_2 \\ \dot{\theta}_1 \\ \dot{\theta}_2 \end{bmatrix}_{equil} = \begin{bmatrix} \theta_1 \\ k\pi \\ 0 \\ 0 \end{bmatrix}, \text{ where } k = \text{any integer}$$

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Swing-up

Trajectory Optimisation

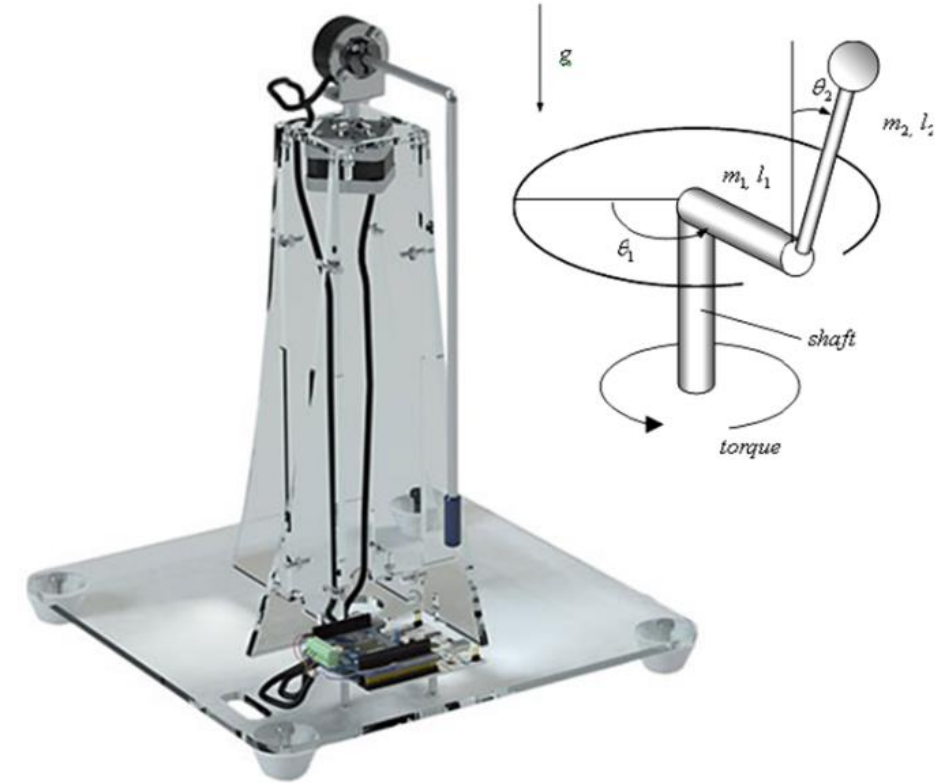
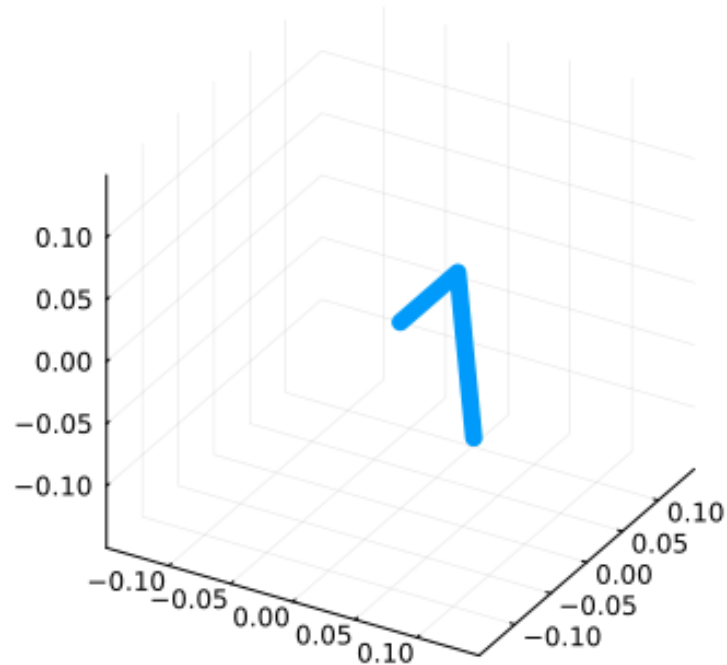
We want to 'swing up' the pendulum.

How can we do this intelligently? Can we consider:

- Motor Torque / Velocity / acceleration limits
- Spatial constraints

One approach is trajectory optimisation

Time = 0.0 s



How?

System Identification

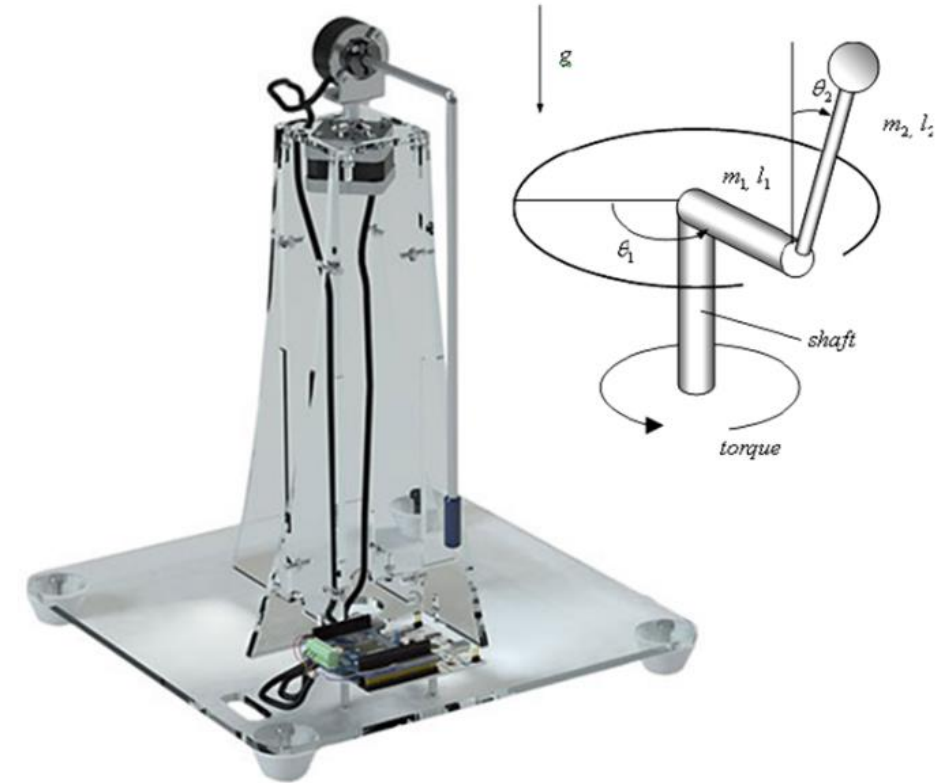
Why System Identification?

Our analytical model may be poor due to making bad estimates of:

$$m_1, m_2, d_{fr_2}, \mathbf{r}_c^{\mathcal{F}_2}$$

Can we improve our model parameters using observed data? YES

$$\mathbf{p} = m_1, \\ m_2, d_{fr_2}, \mathbf{r}_c^{\mathcal{F}_2}$$



How?

Next steps

- Improve model parameters & verify
- Implement on real set-up

