Read each question carefully. Avoid simple mistakes. (Use the back of the page if necessary). You must show your work in order to get credits or partial credits.

1. [10pts] For each of the following equations, state the order and whether it is nonlinear, linear inhomogeneous, or linear homogeneous; provide reasons.

(a)
$$u_{tt} - c^2 u_{xx} + u^3 = 0$$

(b)
$$u_t - ku_{xx} + 6u_{xt} = 0$$

$$(c) u_{xx} + u_{yy} + 2xyu_{xy} = 0$$

(d)
$$u_{tt} - u_{xx} = \sin u$$

(e)
$$u_t - u_{xxx} + 6uu_x = 0$$

(f)
$$\frac{du}{dt} = G(u)$$

- 2. [15pts] Solve the first order $2u_t + 3u_x = 0$ with auxiliary condition $u = \sin x$ when t = 0. (Hint: Express the solution as, according to characteristic curve method, u(t,x) = f(bt ax))
- 3. [15pts] Solve $u_{tt} = c^2 u_{xx}$, $u(x,0) = e^x$, $u_t(x,0) = \sin x$.
- 4. [10pts] Find the solution of the Laplace equation $u_{xx} + u_{yy} = 0$ for $x, y \in [0, 1]$. Use separation of variables.
- 5. [10pts] Solve the wave equation $u_{tt} u_{xx} = 0$ for $x \in [0, \ell], t \in \mathbb{R}$.
- 6. [10pts] Use change of variables t' = at + bx, x' = bt ax to solve

$$au_t + bu_x = f(x, t), (t, x) \in R \times \mathbb{R}$$
 (1)

7. [20pts] The solution to $u_t - ku_{xx} = 0$, u(0, x) = f(x) is given by

$$u(t,x) = e^{-kt\Delta} f(x) = \frac{1}{(4k\pi t)^{n/2}} \int_{\mathbb{R}^n} e^{-\frac{|x-y|^2}{4kt}} f(y) dy := \int_{\mathbb{R}^n} p_t(x,y) f(y) dy,$$

where $p_t(x, y)$ denote the fundamental solution or heat kernel. Show that

$$a) \int_{\mathbb{R}^n} e^{-|x|^2} dx = \pi^{n/2}$$

$$b) \int p_t(x,y) dx = 1 \qquad \forall t > 0$$

$$c) \text{ Given any } \delta > 0, \lim_{t \to 0} \int_{|x-y| \ge \delta} p_t(x,y) dy = 0$$

$$d) \text{ If } f \in L^p, 1 \le p \le \infty, \text{ then } u(t,x) \to f(x) \text{ as } t \to 0.$$

8. [20pts] a) Consider the Schrödinger equation

$$i\partial_t u + \Delta u + \varepsilon |u|^p u = 0$$

Show that the conserved quantities (constants of motion) are

$$Q(t) = \int |u|^2 dx$$

$$E(t) = \frac{1}{2} \int (|\nabla_x u|^2 + \frac{\varepsilon}{p+1} |u|^{p+1}) dx$$

(Hint: Use integration by parts to prove dE(t)/dt = 0, thus E(t) must be a constant)

b) How about wave equation

$$\partial_{tt}u - u_{xx} + \varepsilon |u|^p u = 0$$

Is
$$E(t) = \frac{1}{2} \int_{\mathbb{R}^n} (u_t^2 + |\nabla u|^2 + \frac{\varepsilon}{p+1} |u|^{p+1}) dx$$
 conserved in time?

9. [15pts] Solve the heat equation with convection

$$u_t - ku_{xx} + Vu_x = 0,$$
 $-\infty < x < \infty$
 $u(x,0) = \phi(x)$

where V is a constant. (Hint: Sub y = x - Vt)

10. [10pts] Solve $u_{tt} = 9u_{xx}$ in $x \in (0, \frac{\pi}{2})$, $u(x, 0) = \cos x$, $u_t(x, 0) = 0$, $u_x(0, t) = 0$, $u(\frac{\pi}{2}, t) = 0$.

11. [Bonus 10pts] Solve the inhomogeneous diffusion equation on the halfline with Dirichlet boundary condition

$$u_t - ku_{xx} = f(x,t)$$
 $0 < x, t < \infty$
 $u(x,0) = \phi(x), u(0,t) = 0$

using the method of reflection.

Solutions 2. The general solution for the transport equation is given by u(x,t)=f(bt-ax); see Lecture notes or Section 1.2 in Strauss. Sub the initial condition with t=0 into this, we get $f(-ax)=\sin x$. Therefore $f(x)=-\sin(x/a)$.

3. The general solution for the wave equation on the line $(-\infty, \infty)$ is given by u(x,t) = f(x+ct) + g(x-ct). Substituting the initial conditions into this expression yields

$$\begin{cases} f(x) + g(x) = e^x \\ cf'(x) - cg'(x) = \sin x \end{cases}$$

Now taking derivative on the first equation we obtain $f'(x) + g'(x) = e^x$; this, together with the second equation will lead to a solution of f' and g'. From there antiderivative will recover f and g.

4. Write u(x,y) = X(x)Y(y) and sub into $\Delta u = 0$. We then have

$$\frac{X''(x)}{X(x)} = -\frac{Y''(y)}{Y(y)} = -\lambda$$

where $\lambda \geq 0$ is a constant (or eigenvalue of d^2/dx^2).

- 5. Write u(x, y) = T(t)X(x),
- 6. This is a transport equation with inhomogeneous term f. A detailed solution was given in the handout Monday Feb 22.
 - 7. a)

$$\int_{\mathbb{R}^n} e^{-|x|^2} dx = \int_{-\infty}^{\infty} \cdots \int_{-\infty}^{\infty} e^{-x_1^2} \cdots e^{-x_n^2} dx_1 \cdots dx_n$$
$$= \int_{-\infty}^{\infty} e^{-x_1^2} dx_1 \cdots \int_{-\infty}^{\infty} e^{-x_n^2} dx_n$$

There are a few ways to evaluate $\int_{-\infty}^{\infty} e^{-x^2} dx$. For instance using polar coordinate (r, θ) in 2 dimensions for $x = (x_1, x_2)$.

$$\int_{\mathbb{R}^2} e^{-|x|^2} dx_1 dx_2 = \int_0^{2\pi} d\theta \int_0^{\infty} e^{-r^2} r dr$$
$$= 2\pi \int_0^{\infty} e^{-r^2} r dr = -\pi \int_0^{\infty} d(e^{-r^2}) = \pi$$

Meanwhile,

$$\int_{\mathbb{R}^2} e^{-|x|^2} dx_1 dx_2 = \int_{\mathbb{R}^2} e^{-x_1^2} e^{-x_2^2} dx_1 dx_2$$
$$= \int_{\mathbb{R}} e^{-x_1^2} dx_1 \int_{\mathbb{R}} e^{-x_2^2} dx_2$$

- which shows $\int_{-\infty}^{\infty} e^{-x^2} dx = \sqrt{\pi}$ b) $\int p_t(x,y) dx = 1$ follows from a) by scaling or change of variables.
- c) For fixed $\delta > 0$, a change of variables $u = (y x)/\sqrt{4kt}$ gives

$$\int_{|x-y| \ge \delta} \frac{1}{(4k\pi t)^{n/2}} e^{-\frac{|x-y|^2}{4kt}} dy = \int_{|u| \ge \delta/\sqrt{4kt}} e^{-|u|^2} du \to 0 \quad as \ t \to 0$$

d) Let $f \in L^p$, $1 \leq p \leq \infty$. We may assume $f \in C_b(\mathbb{R}^n)$, that is, f is bounded continuous function. Then

$$u(t,x) - f(x) = \int_{|x-y| < \delta} p_t(x,y) [f(y) - f(x)] dy + \int_{|x-y| \ge \delta} p_t(x,y) [f(y) - f(x)] dy := I_1 + I_2.$$

Given $\varepsilon > 0$, there exists $\delta > 0$ such that $|f(y) - f(x)| < \varepsilon$ whenever $|y-x|<\delta$. Since

$$|I_1| \le \varepsilon \int_{|x-y| < \delta} p_t(x,y) dy \le \varepsilon \int_{\mathbb{R}^n} p_t(x,y) dy = \varepsilon$$
and $|I_2| \le 2C \int_{|x-y| \ge \delta} p_t(x,y) dy \to 0 \text{ as } t \to 0,$

we see immediately that $u(t,x) - f(x) \to 0$ as $t \to 0$.

- 8. a) Show Q'(t) = 0 using the Schrödinger equation and integration by
- b) Similar method.

9. Sub y = x - Vt and we have $u(x,t) = u(y + Vt,t) := \tilde{u}(y,t)$, where x = x(y,t) = y + Vt. By chain rule,

$$\begin{split} \frac{\partial}{\partial t} \tilde{u}(y,t) &= \frac{\partial}{\partial t} [u(y+Vt,t)] = \frac{\partial}{\partial x} u(x,t) \cdot V + \frac{\partial}{\partial t} u(x,t) \\ \frac{\partial}{\partial y} \tilde{u}(y,t) &= \frac{\partial}{\partial y} [u(y+Vt,t)] = \frac{\partial}{\partial x} u(x,t) \cdot 1. \end{split}$$

Sub the above into the $u_t - ku_{xx} + Vu_x = 0$, $u(x,0) = \phi(x)$. We find that $\tilde{u}_t - k\tilde{u}_{yy} = 0$, $\tilde{u}(y,0) = \phi(y)$. Thus $\tilde{u}(y,t) = \int p_t(y,z)\phi(z)dz$ which gives that

$$u(x,t) = \int p_t(x - Vt, z)\phi(z)dz$$
$$= \frac{1}{(4k\pi t)^{n/2}} \int_{\mathbb{R}} e^{-\frac{|x - Vt - z|^2}{4kt}} \phi(z)dz.$$

10. By reflection method as described in Section 3.3, extend u(x,t) as odd function on $(-\frac{\pi}{2}, \frac{\pi}{2})$, and then periodically on \mathbb{R} , which we call v(x,t). Also do the same for ϕ and ψ . Then v satisfies the wave equation on \mathbb{R} with initial condition $v(x,0) = \tilde{\phi}$, $v_t(x,0) = \tilde{\psi}$, with a little abuse of notation we will continue to call them ϕ , psi. Thus the solution is given by

$$v(x,t) = \frac{1}{2}(\phi(x+ct) + \phi(x-ct)) + \int_{x-ct}^{x+ct} \psi(s)ds$$

Note that the boundary condition is automatically satisfied because of the odd symmetries on **both** $x=0, x=\pi/2$. As in the text, the strip $(0,\pi/2) \times \mathbb{R}$ divided into diamond regions. Inside each of these diamond, the wave propagates through a chain of reflections against the boundaries, $\phi(x\pm ct)$ has different sign depending on the number of reflections, similarly for $\pm \psi$. After some simple calculations, the express of v(x,t) has a domain of dependence resulting from those reflections.

11. The inhomogeneous problem on the line has the solution

$$v(x,t) = e^{t\Delta}\phi + \int_0^t e^{(t-s)\Delta}f(x,s)ds$$
$$= \int p_t(x,y)\phi(y)dy + \int_0^t ds \int p_{t-s}(x,y)f(y,s)dy$$

Make odd extension of u to v with $\phi(-x) = -\phi(x)$, then restrict $u = v|_{\mathbb{R} \times \mathbb{R}_+}$ to obtain the solution.