

## SAMURAI Analysis Technique

The analysis uses an incremental form of the variational cost function that avoids the inversion of the background error covariance matrix by using a control variable  $\hat{\mathbf{x}}$ , similar to the forms in Barker et al. (2004, Eq. 2) and Gao et al. (2004, Eq. 7):

$$J(\hat{\mathbf{x}}) = \frac{1}{2}\hat{\mathbf{x}}^T\hat{\mathbf{x}} + \frac{1}{2}(\mathbf{H}\mathbf{C}\hat{\mathbf{x}} - d)^T\mathbf{R}^{-1}(\mathbf{H}\mathbf{C}\hat{\mathbf{x}} - d) \quad (\text{B1})$$

where  $\mathbf{H}$  denotes the linearized observation operator,  $\mathbf{C}$  denotes the square root of the background error covariance matrix,  $\mathbf{R}$  denotes the observation error covariance matrix, and  $d \equiv \mathbf{y} - h(\mathbf{x}_b)$  denotes the difference between the observations  $\mathbf{y}$  and the nonlinear observation operator applied to the background state estimate  $h(\mathbf{x}_b)$ . In the current study,  $h$  and  $\mathbf{H}$  are equivalent. The cost function is minimized using a conjugate gradient algorithm (Polak 1971; Press et al. 2002) to find the atmospheric state where the gradient with respect to  $\hat{\mathbf{x}}$  is 0:

$$\nabla J(\hat{\mathbf{x}}) = (\mathbf{I} + \mathbf{C}^T\mathbf{H}^T\mathbf{R}^{-1}\mathbf{H}\mathbf{C})\hat{\mathbf{x}} - \mathbf{C}^T\mathbf{H}^T\mathbf{R}^{-1}d \quad (\text{B2})$$

We can express the transform from the control variable to an analysis increment as an operator sequence  $\delta\mathbf{x} = \mathbf{C}\hat{\mathbf{x}} = \mathbf{SDF}\hat{\mathbf{x}}$ . The  $\mathbf{SDF}$  matrix transforms represent the cubic B-spline transform, standard deviation of the background error, and recursive filter operators, respectively. For brevity, only the one-dimensional transform for the radial direction is illustrated here. An additional transformation in the  $z$  direction follows. The spline transform

755 **S** is given by:

$$\begin{aligned}
\mathbf{S} &= (\mathbf{P} + \mathbf{Q})^{-1} \\
\mathbf{P} &= [p_{mm'}]^T, \quad p_{mm'} = \int_D \phi_m(r) \phi_{m'}(r) dr \\
\mathbf{Q} &= [q_{mm'}]^T, \quad q_{mm'} = \int_D \epsilon_q(r) \phi_m'''(r) \phi_{m'}'''(r) dr \\
\phi_m(r) &= \Phi\left(\frac{r - r_m}{\Delta r}\right), \quad \text{for } m \in M \text{ and } r \in D
\end{aligned}
\tag{B3}$$

757 where  $\Phi$  is the cubic B-spline given by:

$$\begin{aligned}
&\frac{1}{6}(1 - |\xi|)^3 - \frac{2}{3}(1 - |\xi|)^3 \quad \text{if } 1 \geq |\xi| \geq 0 \\
&\frac{1}{6}(2 - |\xi|)^3 \quad \text{if } 2 \geq |\xi| \geq 1 \\
&0 \quad \text{if } |\xi| \geq 2
\end{aligned}
\tag{B4}$$

759 Note that  $\mathbf{P}$  becomes the identity matrix for orthogonal basis functions, but is required  
760 for computing the cubic B-Spline coefficients. Here  $\mathbf{Q}$  is a third derivative constraint to  
761 reduce Gibb's oscillations in the spline transform (Ooyama 2002), with  $\epsilon_q$  a specified effective  
762 spatial filter cutoff length. The resulting increment or analysis can then be evaluated at any  
763 point in the physical domain through the inner product of the basis functions and the spline  
764 coefficients. For example, for the radial wind

$$u(r) = \boldsymbol{\phi}^T \mathbf{x} + \tilde{\boldsymbol{\phi}}^T \tilde{\mathbf{x}}
\tag{B5}$$

766 where the second term on the right hand side represents imposed boundary conditions  
767 (Ooyama 2002).

768 For the current study, the background error correlations were assumed to be Gaussian  
769 and isotropic, and were calculated using an efficient recursive filter operator that replicates  
770 the effects of this correlation (Purser et al. 2003). The operator combination  $\mathbf{D}\mathbf{F}$  is the  
771 application of the background error covariance matrix, where  $\mathbf{D}$  is the standard deviation  
772 of the background errors and  $\mathbf{F}$  is the recursive filter. The standard deviation of the back-  
773 ground errors were purposefully set very high given no prior knowledge of background state

774 other than the objective analysis from BM08. A large background error standard deviation  
 775 has the detrimental side effect of making the spline analysis unconstrained in data-poor re-  
 776 gions. The recursive filter length scale acts as both an effective distance for the influence of  
 777 the observations and as a spatial filter. A large length scale helps spread the information  
 778 provided by an observation across data gaps, but also removes fine-scale detail from the  
 779 analysis. Sensitivity tests indicated that a  $6\Delta$  background error length scale (relative to the  
 780 grid scale) was a good compromise between smoothing and data density constraints, and was  
 781 adequate for the current study.

782 The control variable state vector is constructed in axisymmetric cylindrical space and is  
 783 given by  $\mathbf{x} = \{\rho r v, \psi, S', q', \rho'_a\}^T$ , where  $\rho$  denotes the moist density including water vapor,  
 784  $\psi$  denotes the transverse streamfunction,  $S$  denotes the moist static energy ( $c_p T + Lq +$   
 785  $gz$ ),  $\rho_a$  denotes the dry air density, primes represent departures from a static, background  
 786 reference state, and the remaining symbols are as defined in Section 2a. The mean tropical  
 787 sounding from Jordan (1958) was used as the reference state in this study. Radial and vertical  
 788 momentum were recovered from the streamfunction using  $\rho u = -\partial\psi/\partial z$  and  $\rho w = \partial\psi/\partial r$ .  
 789 The vertical and radial resolution of the analyses used in the study were 100 m and 1 km,  
 790 respectively.