

Effective potential

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$$V(\rho) = \frac{1}{2} \lambda \rho^2 + v \rho - c \sqrt{2\rho}$$

$$\rho = \frac{1}{2} \sigma^2$$

$$E_q = \sqrt{q^2 + m_f^2}$$

$$+ \frac{2}{3} N_f N_c \int \frac{d^3 q}{(2\pi)^3} \frac{q^2}{E_q} \left[ 1 - n_F(E_q, \mu) - n_F(E_q, -\mu) \right]$$

Vacuum part :

$$\int \frac{d^3 q}{(2\pi)^3} \frac{q^2}{(q^2 + m_f^2)^{\frac{1}{2}}} = \int \frac{d^3 q}{(2\pi)^3} \left[ \frac{q^2 + m^2}{(q^2 + m^2)^{\frac{1}{2}}} - \frac{m^2}{(q^2 + m^2)^{\frac{1}{2}}} \right]$$

$$= \int \frac{d^3 q}{(2\pi)^3} \frac{1}{(q^2 + m^2)^{-\frac{1}{2}}} - m^2 \int \frac{d^3 q}{(2\pi)^3} \frac{1}{(q^2 + m^2)^{\frac{1}{2}}} \quad \textcircled{1} \quad \textcircled{2}$$

For dimensional regularization:

$$\int \frac{d^D p}{(2\pi)^D} \frac{1}{(p^2 + L)^\alpha} = \frac{\Gamma(a - \frac{D}{2}) L^{\frac{D}{2} - a}}{(4\pi)^{\frac{D}{2}} \Gamma(a)}$$

So  $\textcircled{1}$ :  $\int \frac{d^3 q}{(2\pi)^3} \frac{1}{(q^2 + m^2)^{-\frac{1}{2}}}$

$$\Rightarrow \int \frac{d^D q}{(2\pi)^D} \frac{1}{(q^2 + m^2)^{-\frac{1}{2}}}$$

$$= \mu^{2\varepsilon} \int \frac{d^D q}{(2\pi)^D} \frac{1}{(q^2 + m^2)^{-\frac{1}{2}}}$$

$$= \mu^{2\varepsilon} \frac{\Gamma(\frac{1}{2} - \frac{D}{2}) (m^2)^{\frac{D}{2} + \frac{1}{2}}}{(4\pi)^{\frac{D}{2}} \Gamma(-\frac{1}{2})}$$

set  $\varepsilon = \frac{3-D}{2}$

then  $D = 3 - 2\varepsilon$

and  $\alpha = -\frac{1}{2}$

$$= \mu^{2\varepsilon} \frac{\Gamma(\varepsilon-2) (m^2)^{2-\varepsilon}}{(4\pi)^{\frac{3}{2}-\varepsilon} \Gamma(-\frac{1}{2})}$$

Expand around  $\varepsilon=0$  (2)

$$= -\frac{m^4}{32\pi^2} \left\{ \frac{1}{\varepsilon} - \frac{1}{2} [-3 + 2\gamma_E + 4 \log(\frac{m}{2\sqrt{\pi}\mu})] \right\}$$

②:  $m^2 \int \frac{d^3 q}{(2\pi)^3} \frac{1}{(q^2 + m^2)^{\frac{1}{2}}}$

$a = \frac{1}{2}$

$$= m^2 \mu^{2\varepsilon} \int \frac{d^D q}{(2\pi)^D} \frac{1}{(q^2 + m^2)^{\frac{1}{2}}}$$

$$= m^2 \mu^{2\varepsilon} \frac{\Gamma(\frac{1}{2} - \frac{D}{2}) (m^2)^{\frac{D}{2} - \frac{1}{2}}}{(4\pi)^{\frac{D}{2}} \Gamma(\frac{1}{2})}$$

$$= m^2 \mu^{2\varepsilon} \frac{\Gamma(\varepsilon-1) (m^2)^{1-\varepsilon}}{(4\pi)^{\frac{3}{2}-\varepsilon} \Gamma(\frac{1}{2})}$$

Expand around  $\varepsilon=0$

$$= \frac{m^4}{8\pi^2} \left\{ -\frac{1}{\varepsilon} + [-1 + \gamma_E + 2 \log(\frac{m}{2\sqrt{\pi}\mu})] \right\}$$

So ①-②:

$$= -\frac{m^4}{32\pi^2} \left\{ \frac{1}{\varepsilon} - \frac{1}{2} [-3 + 2\gamma_E + 4 \log(\frac{m}{2\sqrt{\pi}\mu})] \right\}$$

$$- \frac{m^4}{8\pi^2} \left\{ -\frac{1}{\varepsilon} + [-1 + \gamma_E + 2 \log(\frac{m}{2\sqrt{\pi}\mu})] \right\}$$

$$= -\frac{m^4}{32\pi^2} \left\{ \frac{1}{\varepsilon} + \frac{3}{2} - \gamma_E - 2 \log(\frac{m}{2\sqrt{\pi}\mu}) \right\}$$

$$- \frac{m^4}{32\pi^2} \left\{ -\frac{4}{\varepsilon} - 4 + 4\gamma_E + 8 \log(\frac{m}{2\sqrt{\pi}\mu}) \right\}$$

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$$= -\frac{m^4}{32\pi^2} \left\{ -\frac{3}{\epsilon} - \frac{1}{2} + 3\gamma_E + 6 \log\left(\frac{m}{2\sqrt{s}\mu}\right) \right\}$$
$$= -\frac{3m^4}{16\pi^2} \log\left(\frac{m}{\mu}\right) + \text{divergent} + \text{constant}$$

Here we use  $\overline{MS}$

$$= -\frac{3m^4}{16\pi^2} \log\left(\frac{m}{\mu}\right)$$

So vacuum part of  $V(\phi)$  is

$$V_{\text{vac}}(\phi) = \frac{2}{3} N_c N_f (-1) \frac{3m^4}{16\pi^2} \log\left(\frac{m}{\mu}\right)$$
$$= -\frac{N_c N_f}{8\pi^2} m^4 \log\left(\frac{m}{\mu}\right)$$

this result is same to  
arXiv:1005.3166

Pion wave function renormalization at  $p_0=0$  and  $\vec{p}=0$ .

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$$\begin{aligned}\Sigma &= T \sum_n \int \frac{d^3 q}{(2\pi)^3} \text{Tr} \left[ i \gamma_5 h T_i \bar{G}_2(q) i \gamma_5 h T_j \bar{G}_2(q-p) \right] \\ &= -\frac{1}{2} \delta_{ij} h^2 \oint \bar{G}_2(q) \bar{G}_2(q-p) \text{Tr} \left[ \gamma_5 (-i \not{q} + m) \gamma_5 (-i \not{q-p} + m) \right] \\ &= -\frac{1}{2} \delta_{ij} h^2 \cdot 4N_c \oint \bar{G}_2(q) \bar{G}_2(q-p) [q(q-p) + m^2] \\ &= -2N_c h^2 \delta_{ij} \oint \bar{G}_2(q) \bar{G}_2(q-p) [q_0(q_0-p_0) + \vec{q}(\vec{q}-\vec{p}) + m^2]\end{aligned}$$

Set  $p_0=0$  then,

$$\begin{aligned}&= -2N_c h^2 \delta_{ij} \oint \bar{G}_2(q) \bar{G}_2(q-p) [q_0^2 + \vec{q}(\vec{q}-\vec{p}) + m^2] \\ &= -2N_c h^2 \delta_{ij} \oint \bar{G}_2(q) \bar{G}_2(q-p) [q_0^2 + \vec{q}^2 + m^2 - \vec{q} \cdot \vec{p}] \\ &= -2N_c h^2 \delta_{ij} \oint \bar{G}_2(q) \bar{G}_2(q-p) [\bar{G}_2^{-1}(q) - \vec{q} \cdot \vec{p}] \\ &= -2N_c h^2 \delta_{ij} \oint [\bar{G}_2(q-p) - \bar{G}_2(q) \bar{G}_2(q-p) \vec{q} \cdot \vec{p}] \quad \textcircled{1}\end{aligned}$$

Here we can expand  $\bar{G}_2(q-p)$  around  $\vec{p}=0$ .

$$\begin{aligned}\bar{G}_2(q-p) &= \frac{1}{q_0^2 + x' + m^2} \quad \text{Here } x = \vec{q}^2 \\ &\quad x' = (\vec{q}-\vec{p})^2 \\ &= \bar{G}_2(q) - \bar{G}_2^2(q) (x' - x) + \bar{G}_2^3(q) (x' - x)^2 \\ &= \bar{G}_2(q) - \bar{G}_2^2(q) (\vec{q}^2 + \vec{p}^2 - 2\vec{p} \cdot \vec{q} - \vec{q}^2)\end{aligned}$$

$$\begin{aligned}
& + \bar{G}_2^3(q) (\vec{q}^2 + \vec{p}^2 - 2\vec{p} \cdot \vec{q} - \vec{q}^2)^2 \\
& = \bar{G}_2(q) - \bar{G}_2^2(q) (-2\vec{p} \cdot \vec{q} + \vec{p}^2) + \bar{G}_2^3(q) (4(\vec{p} \cdot \vec{q})^2 - 4\vec{p}^2(\vec{p} \cdot \vec{q}) + \vec{p}^4)
\end{aligned}$$

then

$$\begin{aligned}
① & = -2N_c h^2 \delta_{ij} \oint \left\{ \bar{G}_2(q) - \bar{G}_2^2(q) (-2\vec{p} \cdot \vec{q} + \vec{p}^2) + \bar{G}_2^3(q) (4(\vec{p} \cdot \vec{q})^2 - 4\vec{p}^2(\vec{p} \cdot \vec{q}) + \vec{p}^4) \right. \\
& \quad \left. - \bar{G}_2^2(q) \left[ \bar{G}_2(q) - \bar{G}_2^2(q) (-2\vec{p} \cdot \vec{q} + \vec{p}^2) + \bar{G}_2^3(q) (4(\vec{p} \cdot \vec{q})^2 - 4\vec{p}^2(\vec{p} \cdot \vec{q}) + \vec{p}^4) \right] \vec{q} \cdot \vec{p} \right\} \\
& = -2N_c h^2 \delta_{ij} \oint \left\{ \bar{G}_2(q) - \bar{G}_2^2(q) (-2\vec{p} \cdot \vec{q} + \vec{p}^2) + \bar{G}_2^3(q) (4(\vec{p} \cdot \vec{q})^2 - 4\vec{p}^2(\vec{p} \cdot \vec{q}) + \vec{p}^4) \right. \\
& \quad \left. - \bar{G}_2^2(q) \vec{q} \cdot \vec{p} + \bar{G}_2^3(q) (-2(\vec{p} \cdot \vec{q})^2 + \vec{p}^2 \vec{q} \cdot \vec{p}) - \bar{G}_2^4(q) (4(\vec{p} \cdot \vec{q})^3 - 4\vec{p}^2(\vec{p} \cdot \vec{q})^2 + \vec{p}^4 \vec{q} \cdot \vec{p}) \right\}
\end{aligned}$$

Keep  $\vec{p}^2$  term

$$\begin{aligned}
& = -2N_c h^2 \delta_{ij} \oint \left\{ -\vec{p}^2 \bar{G}_2^2(q) + 4(\vec{p} \cdot \vec{q})^2 \bar{G}_2^3(q) - 2(\vec{p} \cdot \vec{q})^2 \bar{G}_2^3(q) \right\} \\
& = -2N_c h^2 \delta_{ij} \oint \left\{ -\vec{p}^2 \bar{G}_2^2(q) + 2(\vec{p} \cdot \vec{q})^2 \bar{G}_2^3(q) \right\} \\
& = -2N_c h^2 \delta_{ij} |\vec{p}|^2 \oint \left[ -\bar{G}_2^2(q) + 2|\vec{q}|^2 \cos^2 \theta \bar{G}_2^3(q) \right]
\end{aligned}$$

$$= -2N_c h^2 \delta_{ij} |\vec{p}|^2 \frac{1}{n} \int \frac{d^3 q}{(2\pi)^3} \left[ -\bar{G}_2^2(q) + 2|\vec{q}|^2 \cos^2 \theta \bar{G}_2^3(q) \right]$$

extract  $\vec{q}^2$  from propagator

$$= -2N_c h^2 \delta_{ij} \frac{\vec{p}^2}{\vec{q}^3} \frac{1}{\vec{q}} \frac{1}{n} \int \frac{d^3 q}{(2\pi)^3} \left[ -\bar{G}_2^2(q) + 2\cos^2 \theta \bar{G}_2^3(q) \right]$$

Here  $\tilde{G}_2(\tilde{q})$  is the dimensionless propagator

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$$\tilde{G}_2(\tilde{q}) = \frac{1}{\frac{q^2}{2} + 1 + \frac{m^2}{q^2}} = \frac{1}{q_0^2 + 1 + \tilde{m}^2}$$

Then perform the Matsubara summation:

$$= -2 N_c h^2 \delta_{ij} \frac{\vec{P}^2}{q^3} \int \frac{d^3 q}{(2\pi)^3} \left[ -F_2(q, m^2) + 2 \cos^2 \theta F_3(q, m^2) \right]$$

Then

$$\Sigma_\pi^s(0) = -2 N_c h^2 \frac{1}{q^3} \int \frac{d^3 q}{(2\pi)^3} (-F_2 + 2 \cos^2 \theta F_3)$$

Vacuum part of threshold functions:

$$F_2 = \frac{q^3}{4(q^2 + m^2)^{\frac{3}{2}}} \quad F_3 = \frac{3 q^5}{16(q^2 + m^2)^{\frac{5}{2}}}$$

Regularization of  $F_2$ :

$$\begin{aligned} & \int \frac{d^3 q}{(2\pi)^3} \frac{1}{q^3} \frac{q^3}{4(q^2 + m^2)^{\frac{3}{2}}} = \frac{1}{4} \int \frac{d^3 q}{(2\pi)^3} \frac{1}{(q^2 + m^2)^{\frac{3}{2}}} \\ &= \frac{1}{4} \mu^{2\varepsilon} \int \frac{d^D q}{(2\pi)^D} \frac{1}{(q^2 + m^2)^{\frac{3}{2}}} \\ &= \frac{1}{4} \mu^{2\varepsilon} \frac{\Gamma(\frac{3}{2} - \frac{D}{2}) (m^2)^{\frac{D}{2} - \frac{3}{2}}}{(4\pi)^{\frac{D}{2}} \Gamma(\frac{3}{2})} \quad / \quad D \rightarrow 3 - 2\varepsilon \\ &= \frac{1}{4} \mu^{2\varepsilon} \frac{\Gamma(\varepsilon) m^{2\varepsilon}}{(4\pi)^{\frac{3}{2} - \varepsilon} \Gamma(\frac{3}{2})} \end{aligned}$$

Expand around  $\varepsilon = 0$

$$= \frac{1}{16\pi^2} \left( \frac{1}{\epsilon} - \gamma_E + \log(4\pi) - 2 \log\left(\frac{m}{\mu}\right) \right) \textcircled{D}$$

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Regularization of  $F_3$ :

$$\begin{aligned} & \frac{2}{q^3} \int \frac{d^3 q}{(2\pi)^3} \cos^2 \theta F_3 \\ &= \int \frac{d^3 q}{(2\pi)^3} \frac{2}{q^3} \cos^2 \theta \frac{3q^5}{16(q^2+m^2)^{\frac{5}{2}}} \\ &= \frac{3}{8} \int \frac{d^D q}{(2\pi)^D} \frac{q^2 \cos^2 \theta}{(q^2+m^2)^{\frac{5}{2}}} \\ &= \frac{3}{8} \int \frac{q^{D-1} dq}{(2\pi)^D} \frac{q^2}{(q^2+m^2)^{\frac{5}{2}}} \int d\Omega_D \cos^2 \theta \\ &= \frac{3}{8} \cdot \frac{4\pi}{3} \int \frac{q^{D-1} dq}{(2\pi)^D} \frac{q^2}{(q^2+m^2)^{\frac{5}{2}}} \end{aligned}$$

$$= \frac{\pi}{2} \int \frac{q^{D-1} dq}{(2\pi)^D} \left[ \frac{q^2+m^2}{(q^2+m^2)^{\frac{5}{2}}} - \frac{m^2}{(q^2+m^2)^{\frac{5}{2}}} \right]$$

$$= \frac{\pi}{2} \left[ \mu^{2\epsilon} \int \frac{q^{D-1} dq}{(2\pi)^D} \frac{1}{(q^2+m^2)^{\frac{5}{2}}} - m^2 \mu^{2\epsilon} \int \frac{q^{D-1} dq}{(2\pi)^D} \frac{1}{(q^2+m^2)^{\frac{5}{2}}} \right]$$

$$= \frac{\pi}{2} \left[ \mu^{2\epsilon} \frac{\Gamma(\frac{D}{2}) \Gamma(\frac{3}{2}-\frac{D}{2}) (m^2)^{\frac{D}{2}-\frac{3}{2}}}{2 (2\pi)^D \Gamma(\frac{3}{2})} - m^2 \mu^{2\epsilon} \frac{\Gamma(\frac{D}{2}) \Gamma(\frac{5}{2}-\frac{D}{2}) (m^2)^{\frac{D}{2}-\frac{5}{2}}}{2 (2\pi)^D \Gamma(\frac{5}{2})} \right]$$

$$D \rightarrow 3-2\epsilon$$

$$= \frac{\pi}{2} \left[ \mu^{2\epsilon} \frac{\Gamma(\frac{3}{2}-\epsilon) \Gamma(\epsilon) m^{-2\epsilon}}{2 (2\pi)^{3-2\epsilon} \Gamma(\frac{3}{2})} - m^2 \mu^{2\epsilon} \frac{\Gamma(\frac{3}{2}-\epsilon) \Gamma(1+\epsilon) m^{-2-2\epsilon}}{2 (2\pi)^{3-2\epsilon} \Gamma(\frac{5}{2})} \right]$$

For dimensional regularization:

$$\begin{aligned} \int \frac{d^D q}{(2\pi)^D} \frac{1}{(q^2+L)^a} &= \int \frac{q^{D-1} dq d\Omega_D}{(2\pi)^D} \frac{1}{(q^2+L)^a} \\ &= \frac{\Gamma(\frac{D}{2}) L^{\frac{D}{2}-a}}{(4\pi)^{\frac{D}{2}} \Gamma(a)} \end{aligned}$$

$$\text{and } \int d\Omega_D = \frac{2\pi^{\frac{D}{2}}}{\Gamma(\frac{D}{2})}$$

$$\begin{aligned} \text{So } \int \frac{q^{D-1} dq}{(2\pi)^D} &= \frac{\Gamma(\frac{D}{2}) L^{\frac{D}{2}-a} \Gamma(\frac{D}{2})}{(4\pi)^{\frac{D}{2}} \Gamma(a) 2\pi^{\frac{D}{2}}} \\ &= \frac{\Gamma(\frac{D}{2}) \Gamma(\frac{D}{2}) L^{\frac{D}{2}-a}}{2 (2\pi)^D \Gamma(a)} \end{aligned}$$

finite.

Expand around  $\varepsilon=0$

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$$= \frac{\pi}{2} \left[ \frac{1}{16\pi^2} \left( \frac{1}{\varepsilon} - \gamma_E + 2\text{Log}(2\pi) - \text{Poly}\Gamma(0, \frac{3}{2}) - 2\text{Log}(\frac{m}{\mu}) \right) - \frac{1}{24\pi^3} \right] \textcircled{2}$$

$S_0 \quad Z_\pi^S$ :

$$Z_\pi^S(0) = -2N_c h^2 \frac{1}{q^3} \int \frac{d^3 q}{(2\pi)^3} (-F_2 + 2\cos\theta F_3)$$

$$= -2N_c h^2 (-\textcircled{1} + \textcircled{2})$$

$$= -2N_c h^2 \left\{ -\frac{1}{16\pi^2} \left( \frac{1}{\varepsilon} - \gamma_E + \text{Log}(4\pi) - 2\text{Log}(\frac{m}{\mu}) \right) + \frac{\pi}{2} \left[ \frac{1}{16\pi^2} \left( \frac{1}{\varepsilon} - \gamma_E + 2\text{Log}(2\pi) - \text{Poly}\Gamma(0, \frac{3}{2}) - 2\text{Log}(\frac{m}{\mu}) \right) - \frac{1}{24\pi^3} \right] \right\}$$

$$= -2N_c h^2 \left\{ -\frac{1}{16\pi^2} \left( \frac{1}{\varepsilon} - \gamma_E + \text{Log}(4\pi) - 2\text{Log}(\frac{m}{\mu}) \right) + \frac{1}{16\pi^2} \left( \frac{1}{2} \frac{1}{\varepsilon} - \frac{1}{2} \gamma_E + \text{Log}(2\pi) - \frac{1}{2} \text{Poly}\Gamma(0, \frac{3}{2}) - \text{Log}(\frac{m}{\mu}) \right) - \frac{1}{48\pi^3} \right\}$$

$$= -\frac{2N_c h^2}{16\pi^2} \left\{ -\frac{1}{\varepsilon} + \gamma_E - \text{Log}(4\pi) + 2\text{Log}(\frac{m}{\mu}) + \frac{1}{2} \frac{1}{\varepsilon} - \frac{1}{2} \gamma_E + \text{Log}(2\pi) - \frac{1}{2} \text{Poly}\Gamma(0, \frac{3}{2}) - \text{Log}(\frac{m}{\mu}) - \frac{1}{3\pi^2} \right\}$$



$$= -\frac{N_c h^2}{8\pi^2} \left\{ -\frac{1}{2} \frac{1}{2} + \frac{1}{2} \gamma_E - \text{Log}(4\pi) + \text{Log}(2\pi) - \frac{1}{2} \text{Poly}\Gamma(0, \frac{3}{2}) - \frac{1}{3\pi^2} \right. \\ \left. + \text{Log}\left(\frac{m}{\mu}\right) \right\} \quad \langle 9 \rangle$$

Here we use  $\overline{MS}$  then

$$Z_{\pi(0)}^S = -\frac{N_c h^2}{8\pi^2} \text{Log}\left(\frac{m}{\mu}\right)$$