# Estimates on the convergence of expansions at finite baryon chemical potentials

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Based on arXiv:2403.06770









## Outline

- Introduction
  - QCD phase structure
  - Extrapolation to finite density
- EoS from fRG
  - Polyakov-Quark-Meson (PQM) Model
  - Comparison of different Expansion methods
- Summary



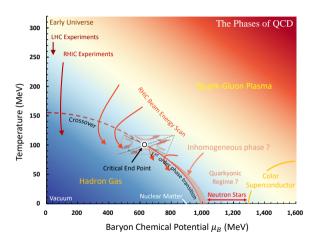
## QCD phase structure

#### Experiments:

- LHC Experiments at low density
- RHIC Beam Energy Scan at finite density
- ...

#### Goals:

- Location of Critical End Point (CEP)
- New phases at high density region
- EoS
- ...



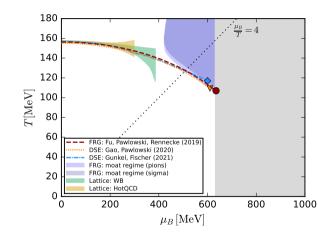
Nucl. Tech. 46 (2023) 04, 040002-040002



## Theoretical predictions

- Lattice QCD (at vanishing chemical potential)
- Dyson-Schwinger Equations (DSE)
- Functional renormalization group (fRG)

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QCD phase diagram (fQCD 2025)



## Well known Advantages & Disadvantages

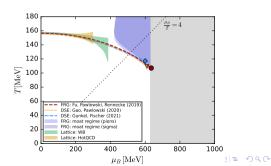
#### Lattice:

- Reliable at 0 chemical potential
- Expansion methods are needed to reach finite density

- Making reasonable use of reliable vanishing chemical potential results
- Attempt to reach higher chemical potentials

#### Functional methods:

- Can compute at finite density
- Have to use truncations to prevent an infinite number of loop diagrams

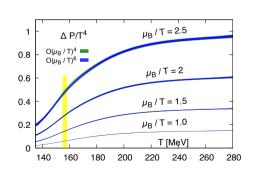


## Extrapolation to finite density

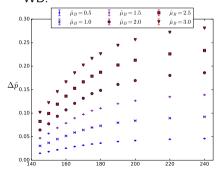
#### Taylor Expansion of the pressure:

$$\frac{p(T,\hat{\mu}_B) - p(T,0)}{T^4} = \sum_{n=1}^{\infty} \frac{1}{(2n)!} \chi_{2n}^B(T,0) \,\hat{\mu}_B^{2n} \tag{1}$$

#### HotQCD:



#### WB:



180 200 220 240

S. Yin (JLU) Extrapolation to finite density

## Extrapolation to finite density

#### Padé approximation:

$$P[m,n] = \frac{p(T,\hat{\mu}_B) - p(T,0)}{T^4} = \frac{\sum_{i=1}^{n/2} a_i \; \hat{\mu}_B^{2i}}{1 + \sum_{i=1}^{m/2} b_i \; \hat{\mu}_B^{2j}}$$
(2)

Here the coefficients  $a_i$  and  $b_i$  are determined by

$$\frac{\partial^{i} P[m, n]}{\partial \hat{\mu}_{B}^{i}} = \chi_{i}^{B} \tag{3}$$

- When m = 0 the Padé approximation will go back to Taylor expansion
- The poles of Padé approximation can be used to estimate the convergence radius of Taylor expansion

S. Yin (JLU) Extrapolation to finite density

## Extrapolation to finite density

#### Ratio estimator: (the pole of P[2,n])

$$r_{c,2n}^{\text{ratio}} = \left| \frac{(2n+1)(2n+2)\chi_{2n}^B}{\chi_{2n+2}^B} \right|^{\frac{1}{2}}$$
 (4)

#### Mercer-Roberts estimator: (the pole of P[4,n])

$$r_{c,2n}^{\text{MR}} = \left| \left[ \frac{\chi_{2n+2}^B \chi_{2n-2}^B}{(2n+2)!(2n-2)!} - \left( \frac{\chi_{2n}^B}{(2n)!} \right)^2 \right] \right|^{\frac{1}{4}} \left| \left[ \frac{\chi_{2n}^B \chi_{2n+4}^B}{(2n)!(2n+4)!} - \left( \frac{\chi_{2n+2}^B}{(2n+2)!} \right)^2 \right] \right|^{-\frac{1}{4}}$$
(5)

- The estimators are given by the poles of the Padé approximation
- They can give a approximation convergence radius of Taylor expansion

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S. Yin (JLU) Extrapolation to finite density

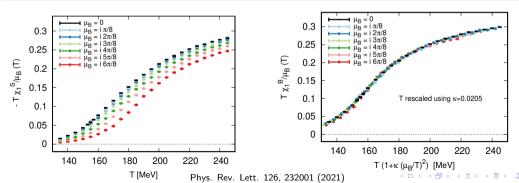
## Extrapolation to finite density

#### expansion:

$$\frac{\chi_1^B(T,\hat{\mu}_B)}{\hat{\mu}_B} = \chi_2^B(T',0) \tag{6}$$

$$\frac{\chi_1^B(T,\hat{\mu}_B)}{\hat{\mu}_B} = \chi_2^B(T',0)$$

$$T'(T,\hat{\mu}_B) = T\left(1 + \kappa_2(T)\,\hat{\mu}_B^2 + \kappa_4(T)\,\hat{\mu}_B^4 + \mathcal{O}(\hat{\mu}_B^6)\right)$$
(6)



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## Polyakov-Quark-Meson (PQM) Model

#### Effective action:

$$\Gamma_{k} = \int_{x} \left\{ Z_{q} \,\bar{q} \left[ \gamma_{\mu} \partial_{\mu} - \gamma_{0} (\mu + igA_{0}) \right] q + \frac{1}{2} Z_{\phi} \left( \partial_{\mu} \phi \right)^{2} + h \,\bar{q} (T^{0} \sigma + i \gamma_{5} \,\vec{T} \cdot \vec{\pi}) + V_{k}(\rho) - c\sigma + V_{\text{glue}}(L, \bar{L}) \right\}$$
(8)

Here we use Local potential approximation (LPA):

$$\partial_t Z_{q/\phi} = 0 \tag{9}$$

$$\partial_t h = 0 \tag{10}$$

We only consider a simple computation of the Grand potential.



S. Yin (JLU) EoS from fRG Polyakov-Quark-Meson (PQM) Model

## Effective potential

#### Flow equation of effective potential:

$$\partial_t V_k(\rho) = \frac{k^4}{4\pi^2} \left[ 3 l_0^{(B)}(m_\pi^2; T) + l_0^{(B)}(m_\sigma^2; T) - 4N_c N_f l_0^{(F)}(m_f^2; \mu, T) \right]$$
(11)

The fermion loop for real and imaginary chemical potential

$$I_0^{(F)}(m_f^2; \mu, T) = \frac{k}{3\sqrt{k^2 + m_f^2}} \left( 1 - n_F(m_f^2; \mu, T; L, \bar{L}) - \bar{n}_F(m_f^2; -\mu, T; L, \bar{L}) \right)$$
(12)

$$=\frac{k}{3\sqrt{k^2+m_e^2}}\left(1-2\operatorname{Re}\left(n_F(m_f^2;\mu,T;L,\bar{L})\right)\right) \tag{13}$$



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## **Equation of State**

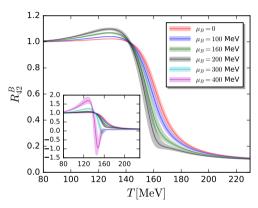
#### Pressure and Baryon number fluctuations:

$$p(T, \mu) = -\Omega(\mu, T) \tag{14}$$

$$\chi_n^B = \frac{\partial^n}{\partial \hat{\mu}_B^n} \frac{p}{T^4}, \qquad \hat{\mu}_B = \frac{\mu_B}{T}$$
 (15)

$$R_{m,n}^B = \frac{\chi_m^B}{\chi_n^B} \tag{16}$$

 The baryon number fluctuations at vanishing chemical potential can be used as Taylor coefficients

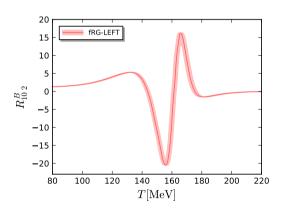


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S. Yin (JLU) EoS from fRG Polyakov-Quark-Meson (PQM) Model

## **Equation of State**

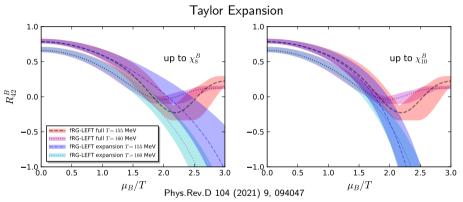
• The baryon number fluctuations can be computed up to (at least) the 10th order



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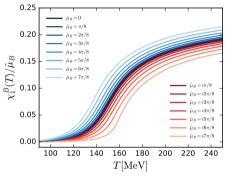
## Comparison of different Expansion methods

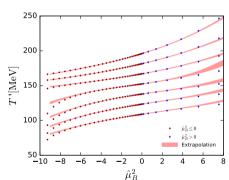


- Direct calculation vs. Taylor expansion of  $R_{42}^B$
- The  $R_{42}^B$  around  $T_{pc}$  exhibits strong fluctuations at high chemical potential, which are difficult to capture with a finite-order Taylor expansion.

## Comparison of different Expansion methods





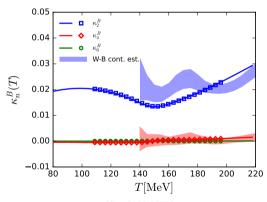


- arXiv: 2403.06770
- Directly compute at imaginary chemical potential
- Apply the T' expansion to perform extrapolation

## Comparison of different Expansion methods

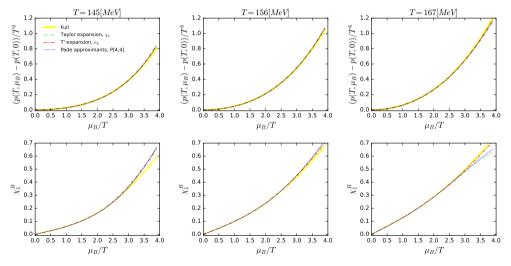
#### T' Expansion

$$T' = T \left( 1 + \kappa_2^B(T) \,\hat{\mu}_B^2 + \kappa_4^B(T) \,\hat{\mu}_B^4 + \kappa_6^B(T) \,\hat{\mu}_B^6 + \cdots \right)$$
(17)

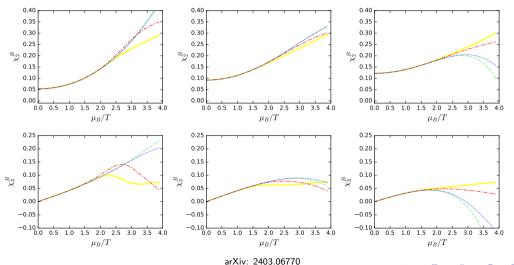


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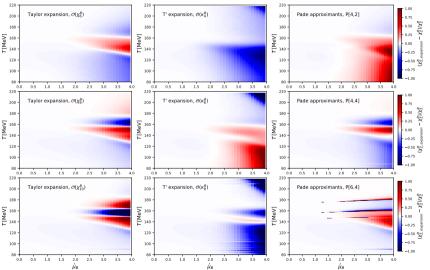
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FoS from fRG

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Happy TEXing!

Mention your findings here.

## Happy TEXing!



## References I

## Outline

Appendix I

6 Appendix II

## Appendix I: Title

#### Remark

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#### Example

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## Outline

4 Appendix I

6 Appendix II

## Appendix II: Title

#### Relevant Title

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#### Example of Appendix II

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