

# Parallel Transport

To set up the idea of parallel transport, we need to understand curvature on manifolds.

Covariant Derivative: The operator  $\nabla$  that performs the functions of the partial derivative, independent of the coordinates.

$\nabla$  is a map from  $(k, l)$  tensor fields to  $(k, l+1)$  tensor fields that satisfies:

$$1. \text{ Linearity: } \nabla(T+S) = \nabla T + \nabla S$$

$$2. \text{ Product Rule: } \nabla(T \otimes S) = (\nabla T) \otimes S + T \otimes (\nabla S)$$

↳ thus it can always be written as the partial derivative plus a linear transformation.

Consider a vector  $V^r$ . For each direction  $\mu$ , the covariant derivative  $\nabla_\mu$  is the partial derivative  $\partial_\mu$  plus a correction given by a matrix  $(\Gamma_\mu^r)_\nu^\lambda$ .

This matrix is called the connection coefficient:  $\Gamma_{\mu\nu}^r$

$$\text{Thus, } \nabla_\mu V^r = \partial_\mu V^r + \Gamma_{\mu\nu}^r V^\lambda$$

We can define the transformation law:

$$\nabla_\mu V^{r'} = \frac{\partial x^\mu}{\partial x^{m'}} \frac{\partial x^{r'}}{\partial x^\nu} \nabla_\nu V^r$$

Thus, the transformation property of  $\Gamma_{\mu\nu}^r$ :

$$\Gamma_{\mu\nu}^{r'} = \frac{\partial x^{m'}}{\partial x^\mu} \frac{\partial x^\lambda}{\partial x^\nu} \frac{\partial x^{r'}}{\partial x^\lambda} \Gamma_{\mu\nu}^r - \frac{\partial x^\mu}{\partial x^{m'}} \frac{\partial x^\lambda}{\partial x^\nu} \frac{\partial^2 x^{r'}}{\partial x^\mu \partial x^\lambda}$$

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The connection coefficients are not the components of a tensor

We introduce 2 more properties of  $\nabla$

3. commutes with contractions:  $\nabla_\mu(T^\lambda_{\alpha\beta}) = (\nabla T)^\lambda_{\mu\alpha\beta}$

4. reduces to  $\partial$  on scalars:  $\nabla_\mu \phi = \partial_\mu \phi$

To take the covariant derivative of an arbitrary tensor:

$$\nabla_\mu T^{\mu_1 \mu_2 \dots \mu_k}_{\nu_1 \dots \nu_l}$$

Partial:

$$= \partial_\mu T^{\mu_1 \dots \mu_k}_{\nu_1 \dots \nu_l}$$

Sum for upper indices  
 $+ \Gamma^\lambda_{\mu_2} T^{\mu_1 \mu_2 \dots \mu_k}_{\nu_1 \dots \nu_l} + \dots + \Gamma^\lambda_{\mu_k} T^{\mu_1 \mu_2 \dots \mu_k}_{\nu_1 \dots \nu_l}$

Subtract for lower indices  
 $- \Gamma^\lambda_{\nu_1} T^{\mu_1 \dots \mu_k}_{\nu_1 \nu_2 \dots \nu_l} - \dots - \Gamma^\lambda_{\nu_k} T^{\mu_1 \dots \mu_k}_{\nu_1 \dots \nu_{k-1}}$

Now we move on to parallel transport

The machinery of connections is vital to compare vectors at different points in the manifold.

Parallel Transport: the concept of moving a vector along a path, while keeping it constant.

The result of parallel transport on a curved space is dependent on the path

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Let's set up the mathematics.

Imagine a curve  $x^\mu(\lambda)$ . Now imagine we want to parallel transport a tensor  $T$ . We need  $T$  to not change along the path, so in flat space we say:

$$\frac{dT}{d\lambda} = \frac{dx^\mu}{d\lambda} \frac{\partial T}{\partial x^\mu} = 0$$

In curved space we use the covariant derivative.  
The covariant derivative along the path is:

$$\frac{D}{d\lambda} = \frac{dx^\mu}{d\lambda} \nabla_\mu \quad \text{replacing the partial from flat space.}$$

Thus for parallel transport:

$$\left( \frac{D}{d\lambda} T \right)_{v_1, \dots, v_K}^{u_1, \dots, u_K} = \frac{dx^\sigma}{d\lambda} \nabla_\sigma T_{u_1, \dots, u_K}^{v_1, \dots, v_K} = 0$$

This equation is well defined, and for a single index vector:

$$\frac{d}{d\lambda} V^\mu + \Gamma_{\sigma\mu}^\nu \frac{dx^\sigma}{d\lambda} V^\nu = 0$$

Proof.

$$\frac{dx^\sigma}{d\lambda} \nabla_\sigma V^\mu = \frac{dx^\sigma}{d\lambda} \left[ \frac{\partial}{\partial x^\sigma} V^\mu + \Gamma_{\sigma\mu}^\nu V^\nu \right]$$

$$= \frac{d}{d\lambda} V^\mu + \Gamma_{\sigma\mu}^\nu V^\nu \cdot \frac{dx^\sigma}{d\lambda}$$

↑  
const so  
commutes

## Parallel Transport

We can find an explicit general solution to parallel transport.

It amounts to finding a matrix  $P_p^{\mu}(\lambda, \lambda_0)$  that relates the vector at  $\lambda_0$  to its value further down the path:

$$v^{\mu}(\lambda) = P_p^{\mu}(\lambda, \lambda_0) v^{\mu}(\lambda_0)$$

$P_p^{\mu}(\lambda, \lambda_0)$  is called the parallel propagator and it depends on the chosen path  $\gamma$ .

$$\text{Defin } A_p^{\mu}(\lambda) = -\Gamma_{\alpha\beta}^{\mu} \frac{dx^{\beta}}{d\lambda}$$

$$\text{Then } \frac{d}{d\lambda} v^{\mu} - A_p^{\mu}(\lambda) v^{\mu} = 0, \text{ thus}$$

$$\frac{d}{d\lambda} v^{\mu} = A_p^{\mu}(\lambda) v^{\mu} = A_0^{\mu} v^{\mu}$$

Substitute the P equation into this for:

$$\frac{d}{d\lambda} P_p^{\mu}(\lambda, \lambda_0) v^{\mu}(\lambda_0) = A_0^{\mu}(\lambda) P_p^{\mu}(\lambda, \lambda_0) v^{\mu}(\lambda_0)$$

$$\hookrightarrow \frac{d}{d\lambda} P_p^{\mu}(\lambda, \lambda_0) = A_0^{\mu}(\lambda) P_p^{\mu}(\lambda, \lambda_0)$$

Now we must solve this equation by integrating.

$$P_p^{\mu}(\lambda, \lambda_0) = \delta_p^{\mu} + \int_{\lambda_0}^{\lambda} A_0^{\mu}(n) P_p^{\mu}(n, \lambda_0) dn$$

Now we solve by iteration, plugging the right side into the original equation repeatedly:

The  $\delta_p^{\mu}$  gives the identity matrix when  $\lambda = \lambda_0$ .

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$$P_p^\sigma(\lambda, \lambda_0) = \delta_p^\sigma + \int_{\lambda_0}^{\lambda} A_\sigma^\sigma(n) P_p^\sigma(n, \lambda_0) dn$$

$$\frac{d}{d\lambda} P_p^\mu(\lambda, \lambda_0) = \delta_p^\sigma A_\sigma^\mu(\lambda) + A_\sigma^\mu \int_{\lambda_0}^{\lambda} A_\sigma^\sigma(n) P_p^\sigma(n, \lambda_0) dn$$

$$= A_\rho^\mu(\lambda) + \int_{\lambda_0}^{\lambda} A_\sigma^\mu \int_{\lambda_0}^{\lambda} A_\sigma^\sigma(n) P_p^\sigma(n, \lambda_0) dn$$

Then we integrate again:

$$P_p^\mu(\lambda, \lambda_0) = \delta_p^\mu + \int_{\lambda_0}^{\lambda} A_\rho^\mu(\eta) d\eta + \int_{\lambda_0}^{\lambda} A_\sigma^\mu(n) dn \int_{\lambda_0}^{\lambda} A_\sigma^\sigma(n') dn'$$

$$= \delta_p^\mu + \int_{\lambda_0}^{\lambda} A_\rho^\mu(\eta) d\eta + \int_{\lambda_0}^{\lambda} \int_{\lambda_0}^{\lambda} A_\sigma^\mu(n) A_\sigma^\sigma(n') P_p^\sigma(n', \lambda_0) dn' dn$$

Then:

$$P_p^\sigma(\lambda, \lambda_0) = \delta_p^\sigma + \int_{\lambda_0}^{\lambda} A_\rho^\sigma(n) dn + \int_{\lambda_0}^{\lambda} \int_{\lambda_0}^{\lambda} A_\sigma^\sigma(n) A_\sigma^\sigma(n') P_p^\sigma(n', \lambda_0) dn' dn$$

Subbing in:

$$\frac{d}{d\lambda} P_p^\mu(\lambda, \lambda_0) = \delta_p^\mu A_\rho^\mu(\lambda) + A_\sigma^\mu \int_{\lambda_0}^{\lambda} A_\rho^\sigma(n) dn$$

$$+ A_\sigma^\mu \int_{\lambda_0}^{\lambda} \int_{\lambda_0}^{\lambda} A_\sigma^\sigma(n) A_\sigma^\sigma(n') P_p^\sigma(n', \lambda_0) dn' dn$$

Integrating again:

$$P_p^\mu(\lambda, \lambda_0) = \delta_p^\mu + \int_{\lambda_0}^{\lambda} A_\rho^\mu(n) dn + \int_{\lambda_0}^{\lambda} \int_{\lambda_0}^{\lambda} A_\sigma^\mu(n) A_\rho^\sigma(n') dn' dn + \dots$$

The pattern continues for infinite integrals.

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The  $n$ th term in the series is an integral over an  $n$ -dimensional right triangle, or  $n$ -simplex.

The  $n$ th order term is then:

$$\int_{\lambda_0}^{\lambda} \int_{\lambda_0}^{\eta_n} \cdots \int_{\lambda_0}^{\eta_2} A(\eta_n) \cdots A(\eta_1) d^n \eta$$

$$= \frac{1}{n!} \int_{\lambda_0}^{\lambda} \cdots \int_{\lambda_0}^{\lambda} \mathcal{D}[A(\eta_n) \cdots A(\eta_1)] d^n \eta$$

$\mathcal{D}$  is the path ordering symbol. The expression

$$\mathcal{D}[A(\eta_n) \cdots A(\eta_1)]$$

stands for the product of  $n$  matrices  $A(\eta_i)$ , ordered in such a way that the largest value of  $\eta_i$  is on the left and each subsequent  $\eta_i$  is less than or equal to the previous one.

Thus we have the solution:

$$P(\lambda, \lambda_0) = 1 + \sum_{n=1}^{\infty} \frac{1}{n!} \int_{\lambda_0}^{\lambda} \cdots \int_{\lambda_0}^{\lambda} \mathcal{D}[A(\eta_n) \cdots A(\eta_1)] d^n \eta$$

Which is the Taylor expansion for the exponential:

$$P(\lambda, \lambda_0) = \mathcal{D} \exp \left[ \int_{\lambda_0}^{\lambda} A(\eta) d^n \eta \right]$$

Or more explicitly:

$$P_{\nu}^{\mu} (\lambda, \lambda_0) = \mathcal{D} \exp \left[ - \int_{\lambda_0}^{\lambda} \Gamma_{\nu}^{\mu} \frac{d\lambda_0}{d\eta} d\eta \right]$$