

Fokker-Planck Eqⁿ:

$$\frac{\partial n}{\partial t} = D \cdot \frac{\partial^2 n}{\partial x^2} + \frac{1}{\mathcal{I}} \frac{\partial}{\partial x} \left(n \frac{\partial u}{\partial x} \right)$$

with $U = -f(x)$, we obtain

$$\frac{\partial n}{\partial t} = D \cdot \frac{\partial^2 n}{\partial x^2} - \frac{f}{\mathcal{I}} \frac{\partial n}{\partial x}$$

for reflecting boundary condition, with $x \in [0, L]$
domain, we have

$$\frac{\partial n}{\partial x} \Big|_{x=0} = \frac{\partial n}{\partial x} \Big|_{x=L} = 0$$

$$n(x, 0) = \mathcal{S}(x)$$

Just shifting
the domain
to have
earlier B/c

Now, let. $n = X(x) \cdot T(t)$

It gives us,

$$X \frac{\partial T}{\partial t} = DT \frac{\partial^2 X}{\partial x^2} - \frac{fT}{\mathcal{I}} \cdot \frac{\partial X}{\partial x}$$

$$\Rightarrow \frac{1}{D} \cdot \frac{\dot{T}}{T} = \frac{X''}{X} - \frac{f}{D\mathcal{I}} \cdot \frac{X'}{X} = -K$$

Therefore, $\dot{T} = -KDT$

$$\therefore T(t) = e^{-KDt}$$

$$\text{And, } \frac{\partial^2 X}{\partial x^2} - \frac{f}{D\mathcal{I}} \frac{\partial X}{\partial x} = -KX$$

So, the eigenvalue problem becomes,

$$\lambda^2 - \frac{f}{D\mathcal{I}} \lambda + K = 0$$

$$\Rightarrow \lambda_1 = \frac{f + \sqrt{f^2 - 4KD\mathcal{I}^2}}{2D\mathcal{I}} ; \lambda_2 = \frac{f - \sqrt{f^2 - 4KD\mathcal{I}^2}}{2D\mathcal{I}}$$

And, the general solution takes the form.

$$X = e^{\frac{t}{2D\zeta}} \left[A \cos \left(\frac{\sqrt{f^r - 4KD^r\zeta^2}}{2D\zeta} x \right) + B \sin \left(\frac{\sqrt{f^r - 4KD^r\zeta^2}}{2D\zeta} x \right) \right]$$

$$\begin{aligned} \text{Now, } X' &= e^{\frac{t}{2D\zeta}} \frac{\sqrt{f^r - 4KD^r\zeta^2}}{2D\zeta} \\ &\cdot \left[-A \sin \left(\frac{\sqrt{f^r - 4KD^r\zeta^2}}{2D\zeta} x \right) + B \cos \left(\frac{\sqrt{f^r - 4KD^r\zeta^2}}{2D\zeta} x \right) \right] \end{aligned}$$

Applying reflecting B/c's @ $x=0$,

we find. $B=0$

and applying reflecting B/c's @ $x=L$

we find. $-A_p \sin pL = 0$ [taking $\frac{\sqrt{f^r - 4KD^r\zeta^2}}{2D\zeta} = p$]

$$\therefore p = \frac{n\pi}{L} = \frac{\sqrt{f^r - 4KD^r\zeta^2}}{2D\zeta}$$

$$\Rightarrow K = \frac{f}{4D^r\zeta^2} - \frac{n^2\pi^2}{L^2}$$

finally, we get,

$$\chi = e^{\frac{f}{2D\zeta}} \cdot \cos\left(\frac{\sqrt{f^2 - 4KD^2\zeta^2}}{2D\zeta}\right) x$$

$$\text{And, } n(x,t) = \chi \cdot T$$

$$= e^{\frac{f}{2D\zeta}} \left\{ \cos\left(\frac{\sqrt{f^2 - 4KD^2\zeta^2}}{2D\zeta}\right) x \right\} e^{-KDt};$$

$$\text{where, } K = \frac{f}{4D^2\zeta^2} - \frac{n\pi^2}{L^2}$$

(also, when $K=0$, $x_2=0$)

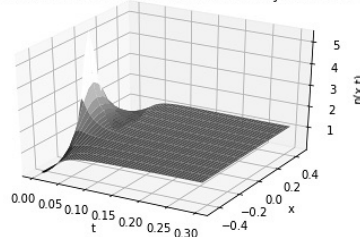
with $\lim_{t \rightarrow \infty}$ we get the steady-state

solution -

$$n_{\infty}(x) = e^{\frac{f}{2D\zeta}} \cos\left(\frac{\sqrt{f^2 - 4KD^2\zeta^2}}{2D\zeta}\right) x$$

And numerically we get the following plot:

Fokker-Plank Diffusion for Reflective Boundary Conditions



and it depicts that, as time passes by the probability of finding the particle at any point of the domain is equally distributed