Hows on the Line - (tint-Order Systems) One-Dinentional System: > [i=f(n)] 2 = dn/dt => n = x(t) Antonomous System ii) $\chi(t)$ is a real rated function of t (time). ii) f(x) is a smooth real-rated function of χ . Existence and Unique ness Theorem: for the initial-rathe problem, [x = f(n)], [x = x(t)] and [x(o) = xo] (initial dition). il. f(n) and f(n) = df/dx are finite-valued and I mooth. ii/. Continuous on an open interval R of the x-axis (in the one-dimensional system) iii/. Xo is a point in R (on the x-axis). Fulfilling the foregoing conditions, the Some time interval (-7,7) about t=0. This solution exists and is unique. Note: 1/2 Effectively the solution is single-valued. 21. The solution may not exist forever.

Phase Partraits: Plotting is versus x in, Infinite norts for Example: | i=fh) = sinx Opposite in (attractor) n(0) => left areson (repeller)

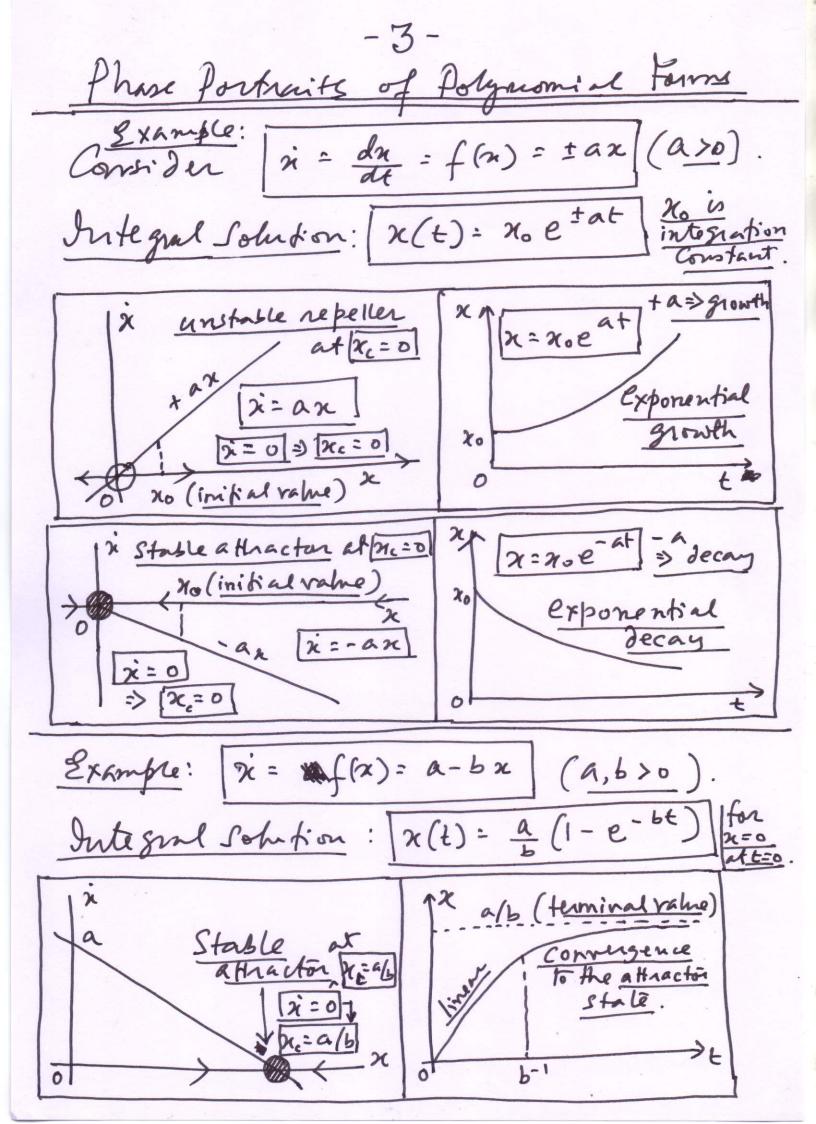
(nestable)

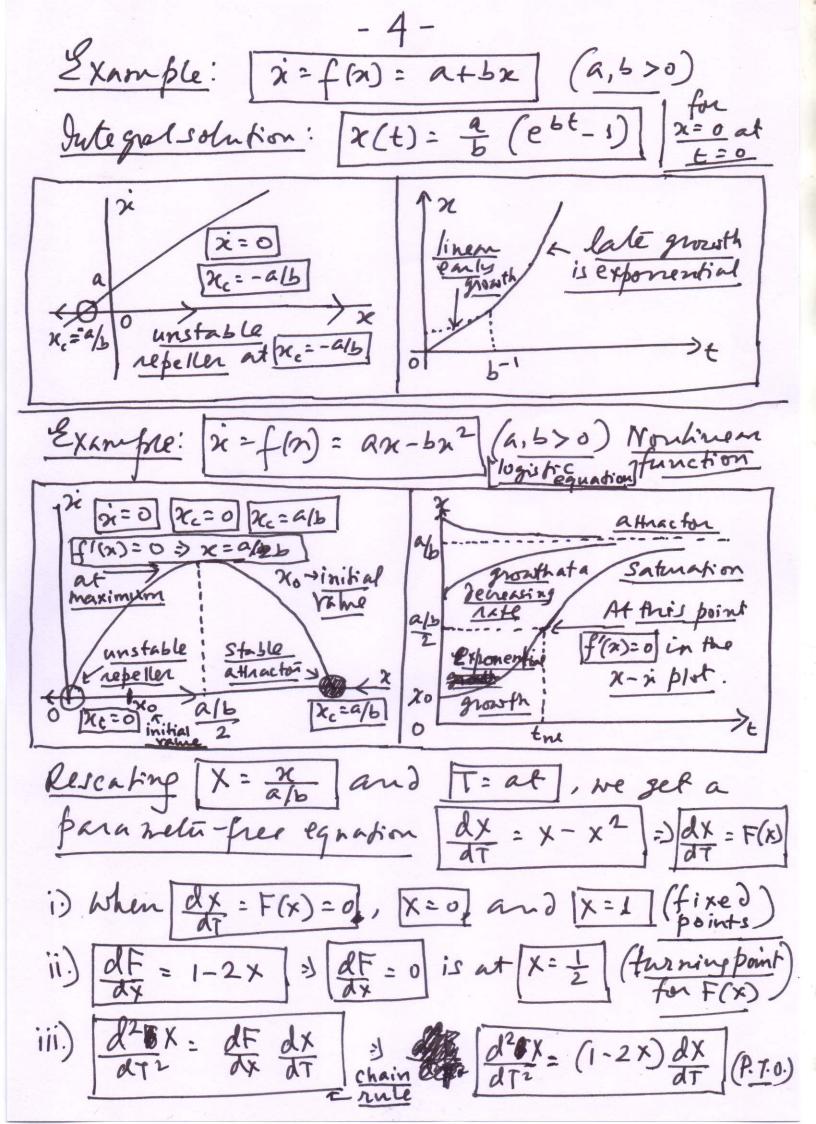
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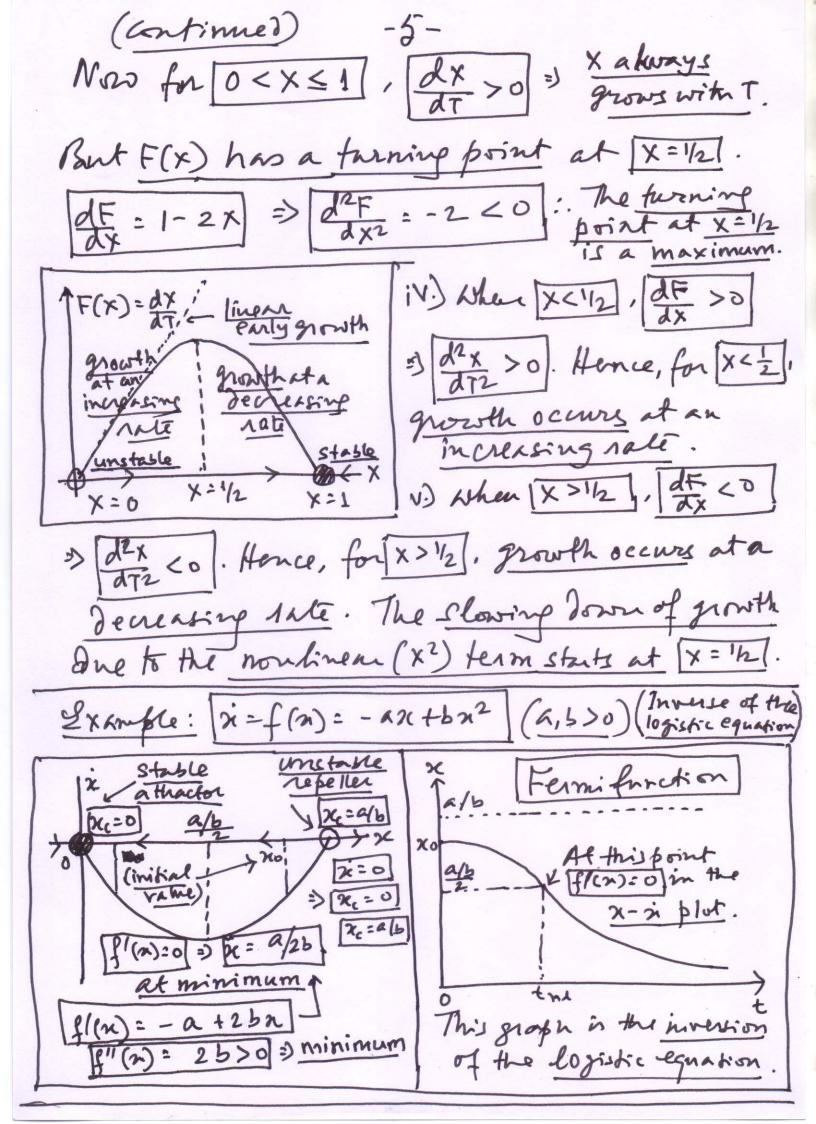
(nestable)

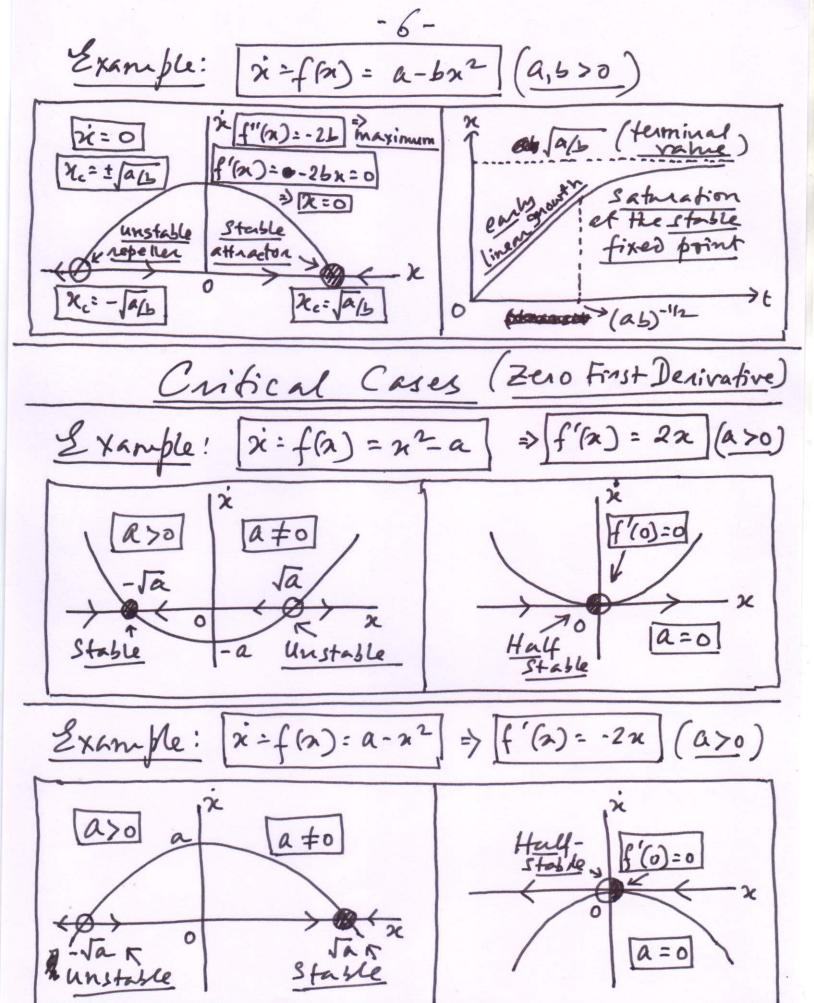
(nestable) 211 Lyton x=x(t), an antonomous first-order System is given by [i=fm). The fixed point (or equilibrium point) of such a System is obtained when [i=f(n)=6. Hence, ffre)=0 gives the fixed point (equilibrium point) on the line is = of at [x = xc) (in which x is the fixed point). This is not a tworning point for [x = x(t)]. den = de (de) = de = i = de de (apphying chain rule) When $|\dot{x}=0|$ = $|\dot{x}=\frac{df}{dn}\dot{n}=0|$. Hence both [i=0] and [i=0] at the same fine.

Turning points usually have non-zero seand derivatives.)

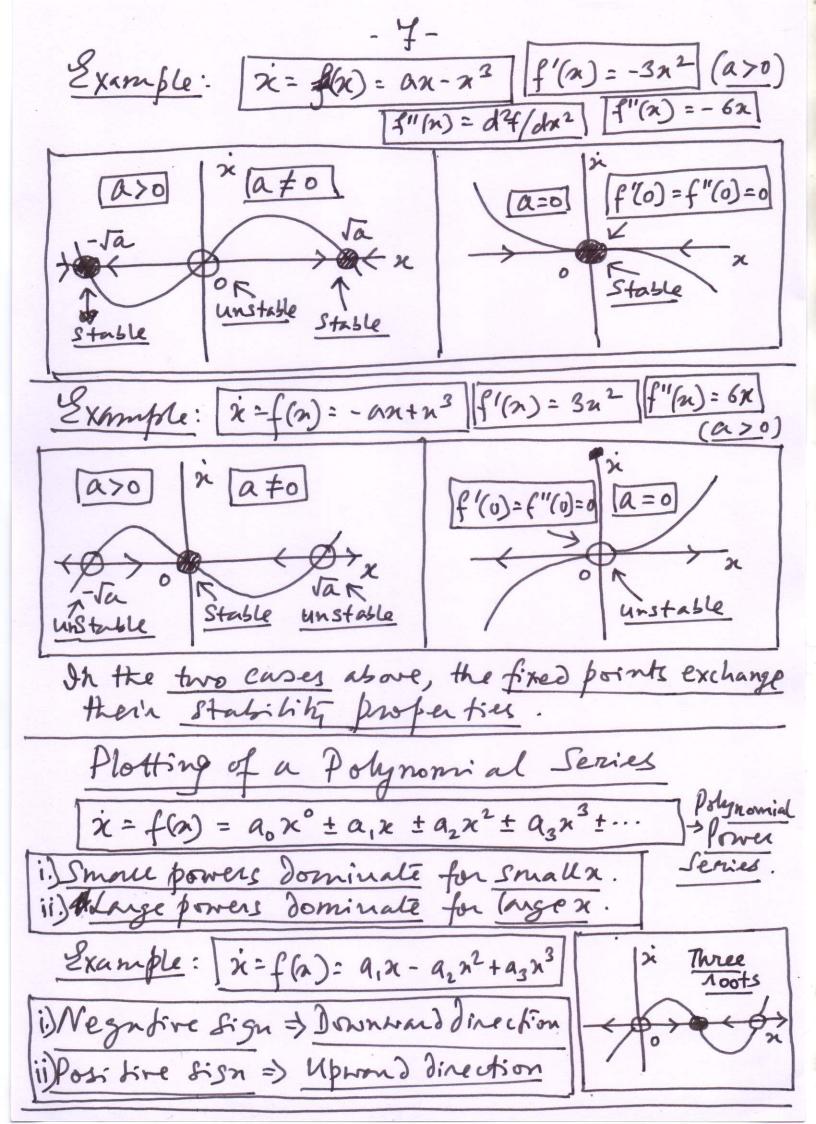


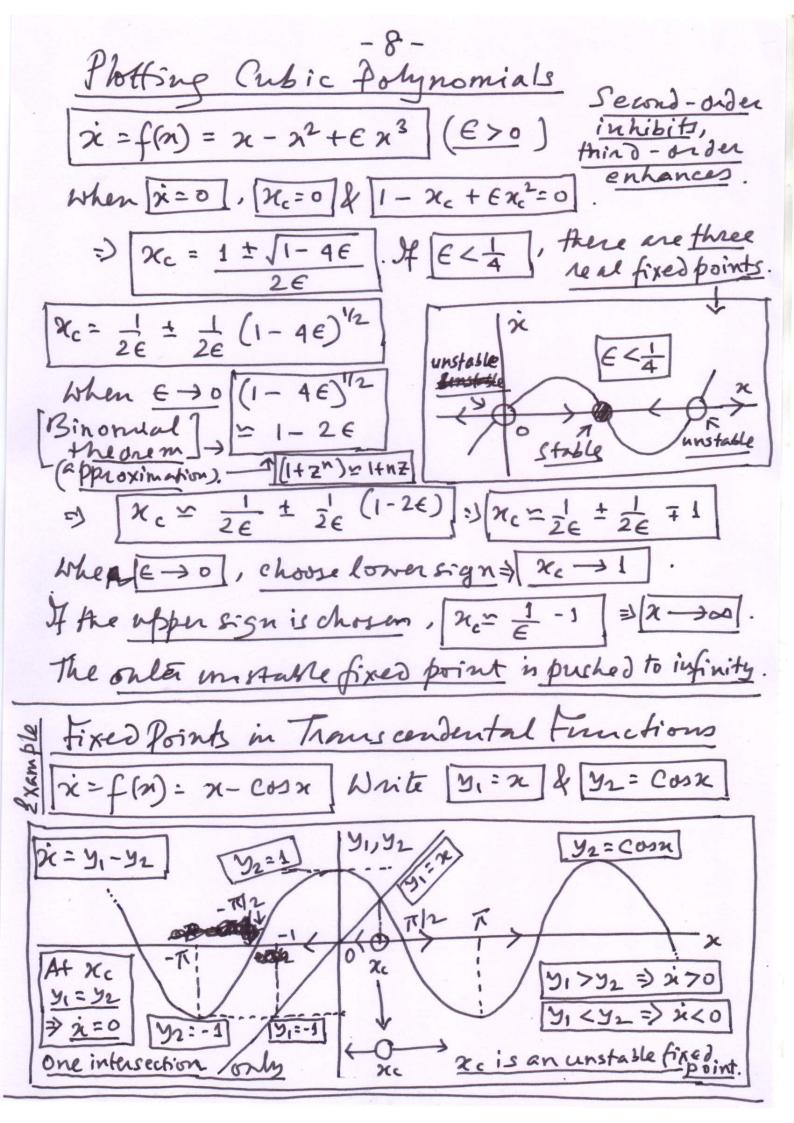






In the two cases whove, the stuble advactor and the unstuble repeller exchange positions.





Stability Analysis of fixed Points For a first-order autonomous dynamical System, | i = f (m) , the fixed point condition is [x=0] at x=xc =) f(xc)=0 at x=xc. The fixed point wordinale ne is pertented by a small amount E, i.e. EKXc. House, he write [x= xc+ E]. Using this in [x=f(x). (x) = xc + \(\epsi = \epsilon(x) = \(\epsilon(x_e + \epsilon)\) Now \(\int_{c=0}\). :. | e = f (nc+E) = f(nc) + f'(nc) E + f'(nc) E2+... by a Taylor expansion. Now far =0. and truncating the Taylor expansion at the linear order (i.e. ignoring E2 and all higher order terms, due to the smallness of E). We get $\dot{\epsilon} = f'(x_i) \epsilon$, a direm differential equation in ϵ . =) $d\epsilon = f'(n\epsilon) \epsilon$ => $\int d\epsilon = f'(n\epsilon) d\epsilon$ = $\int \frac{d\epsilon}{\epsilon} = f'(n\epsilon) d\epsilon$ > ln e = ln A + f'(nc) 4t => E = A e f'(nc) 4t Hence | x = xc + A ef'(ne) + (A is an integration). The foregoing result in due to a LINEAR STABILITY ANALYSIS when fmi)=0, f'(n) =0. (P.T.O.)

(continued) - 10-1/. For a stable fixed point, as t >00, E >0. :. x -> xc, i.e. there is a couvergence towards This happens only when f'(no) <0]. 2/ For an unstable fired point, as t-300. € → 00, j. e. a divergence away from Xc occurs. This happens only when [f'(n)>0]. => i) If f'me) < 0, the fixed point is stable. ii) If [f'(ni) >0]. The fixed point is unstable. 3/ Now x=xc+Aef(xe)t)=> E=Aef(xe)t $\Rightarrow t = \frac{1}{f'(nc)} ln\left(\frac{\epsilon}{A}\right) = \frac{1}{f'(nc)} ln\left(\frac{\lambda - \lambda c}{A}\right)$ For a stuble fixed point, [f'(a) <0] and X-3 xc or [= 30. Hence [t-)00, for x->xc The convergence to x takes infinitely long. Honce, for a first-order system, there is No overshoot of the fixed point, and no oscillation about the fixed point is possible. Decillations we only possible when f'(xc) is imaginary, but since | i=f(x) is real, this is not allowed in a first-order system.

Crifical Condition in the Stability Analysis

Siven [x=f(x)], the fixed point is at [x=0], i.e. f(xe)=0. Perturbing $[x=x_c+E]$, we get

x = E = f(xx) + f'(xx) E + 1 f"(xx) E2 + ... (xe=0).

At the fixed point [f(m)=0]. In addition if

f(xi)=0, then we have a critical condition.
(no longer linear).

 $= \frac{\int (no \log x) (no \log x)}{E^2} \left(\frac{E^2 (no \log x)}{\log x} \right) \frac{E^2}{\log x} \left(\frac{E^2 (no \log x)}{\log x} \right) \frac{E^2}{\log x} \left(\frac{E^2 (no \log x)}{\log x} \right)$

=) $d\epsilon = f''(\pi_c) \epsilon^2$ =) $\int d\epsilon = f''(\pi_c) \int dt$

=) [e-1 = f"(ne) (t-A)] -> A is an integration constant [f"(ne)is]

 $\Rightarrow \boxed{E = \frac{-2}{f''(\pi c)} \frac{1}{t - A}} \Rightarrow \boxed{\chi = \chi_c - \frac{2}{f'''(\pi c)} \frac{1}{t - A}}$

When [t > 00), [x -> xc] (slow power-law Conversence).

1/. Where f'(xc) <0, x = xc + e f'(ne)t. As t >0,

exponential convergence.

21. When [f(nc)=0] (critical condition),

 $N = 2c + \frac{B}{t-A}$ (B= $\frac{-2}{f''(nc)}$), the convergence

towards Xean + >00, is a slow power-law convergence.

